

Review of particle properties

Particle Data Group

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This review of the properties of leptons, mesons, and baryons is an updating of Review of Particle Properties, Particle Data Group [Phys. Letters **50B**, No.1 (1974), and Supplement, Rev. Mod. Phys. **47** (1975) 535]. Data are evaluated, listed, averaged, and summarized in tables. Numerous tables, figures, and formulae of interest to particle physicists are also included. A data booklet is available.

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I. INTRODUCTION, CREDITS, CONSULTANTS

This review is an updating through December 1975 of our previous review (Particle Data Group, 1974, and Supplement, 1975). In this version, we have attempted to make the text as complete and self-contained as possible.

As usual, the results of our compilation are presented in two sections. The Tables of Particle Properties, usually referred to as simply the Tables, contain a summary of the properties of only those peaks or resonances which in our judgment have a large probability of standing the test of time. This is a conservative judgment, and surely means that some genuine resonances are (temporarily) omitted. See Sec. V below.

The Data Card Listings, on the other hand, are an attempt to give up-to-date information, along with references, on all the reported particles whose existence is either considered confirmed (in which case they are also included in the Tables) or not yet confirmed. The Listings also contain mini-reviews pertaining to questions of interest.

A short survey of the history of some of the constants we compile, as well as some of those compiled by others, has been added to the text as a new feature (Sec. VIII). In general, the reliability of the data has been pretty good; that is, the percentage of the time that our best estimates of various constants have changed by more than one standard deviation over the years is roughly what would be expected from statistical considerations. A history of the Particle Data Group, with a discussion of procedures and problems, has been given by Rosenfeld (1975).

This year we are experimenting with a new statistical procedure in our data averaging. Given in the Tables are our best estimates (and their errors), calculated in the old way. In the Listings, we give simultaneously the old (labeled "AVG") and new (labeled "STUDENT") values. In most cases there is little difference. Details may be found in Sec. VII. User comments are solicited.

A pocket-sized data booklet, containing the Tables and a reprint of the figures and formulae from the first part of the book, is available on request.

For North and South America, Australia, and the Far East, write to Technical Information Division, Lawrence Berkeley Laboratory, Berkeley, California 94720, USA.

For all other areas, write to CERN Scientific Information Service, CH-1211 Genève 23, Switzerland.

As usual, we wish to emphasize that we compile the experimental results of others. It is inappropriate to give us the credit for their countless hours of effort. We urge that references be given directly to the original data, and we provide complete references in the Data Card Listings for that purpose.

The responsibilities for the various sections can be broken down as follows:

(1) *Stable particles*: A. Barbaro-Galtieri, N. Barash-Schmidt, and T. G. Trippe.

(2) *Meson resonances*: V. Chaloupka, R. J. Hemingway, M. J. Losty, and M. Roos.

(3) *Baryon resonances*: A. Barbaro-Galtieri, C. Bricman, and R. L. Kelly.

(4) *General, including text*: All authors.

Consultants: To overcome unavoidable gaps in our coverage, both intellectual and geographical, we have solicited the help of consultants:

- Ugo Amaldi (CERN)
- Stanley J. Brodsky (Stanford Linear Accelerator Center)
- Denyse M. Chew (Lawrence Berkeley Laboratory)
- Ronald Crawford (University of Glasgow)
- J. Engler (CERN)
- Anatoli Kuznetsov (JINR, Dubna)
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- F. Mönig (CERN)
- R. Gordon Moorhouse (University of Glasgow)
- David R. Nygren (Lawrence Berkeley Laboratory)
- Oliver E. Overseth (University of Michigan)
- Sherwood I. Parker (Lawrence Berkeley Laboratory)
- Bernard Sadoulet (Lawrence Berkeley Laboratory)
- Paul Söding (DESY)
- Fumiyo Uchiyama (Lawrence Berkeley Laboratory)

The usefulness of this compilation depends in large part on the interaction between the users and the authors and consultants. We appreciate comments, criticisms, and suggestions for improvements of all stages of data retrieval, processing, and presentation.

II. SELECTION OF DATA

All particles are considered to fall into one of the three groups:

- (1) Stable particles, immune to decay via the strong interaction, including the η .
- (2) Meson resonances (including all the new high mass particles, whether "stable" or not).
- (3) Baryon resonances.

These groups are maintained within the two main parts of the compilation:

- (1) Tables of Particle Properties.
- (2) Data Card Listings.

The Data Card Listings contain the original information (data, references etc.), weighted averages, comments, and "mini-reviews". Immediately preceding the Data Card Listings is an Illustrative Key thereto. We attempt to give complete Data Card Listings up to our closing date (January 1, 1976) for all journals listed in the Illustrative Key. We also include preprints and unpublished conference reports which have come to our attention, but make no attempt at completeness.

Roughly 40% of our encoded results are not used for averaging. They are set off in parentheses: our reasoning is then often given in a footnote below the data. If the reason is not given, it is one of the following:

- The quantity was presented with no error stated.
- The result comes from a preprint or conference report. It is our experience that such results (and particularly the errors) often change before final publication. Accordingly we keep these new results in parentheses until they are published (or explicitly verified to us by the authors).
- It involves some assumptions that we do not wish to incorporate.

- It is of poor quality, e.g., bad signal-to-noise-ratio.
- The result is inconsistent with others, e.g., because of different methods employed, rendering averaging meaningless.
- It is not independent of other results, e.g., a result from one of several partial-wave analyses all using the same data, again rendering averaging meaningless.

When the data for a particle have received special treatment or when they present special problems, this is noted in a mini-review in the Data Card Listings.

The Tables of Particle Properties represent the output of weighted averages and some critical judgment. The extent to which "blind" averaging has been tempered with judgment is explained in footnotes to the Tables. In general, however, the footnotes are less complete than is the collection of notes and mini-reviews in the Data Card Listings. The reader is therefore encouraged to familiarize himself with the Data Card Listings and, ultimately, with the original experiments.

III. NOMENCLATURE

A. Quantum numbers

The symbols $I^C(J^P)C$ represent:

I = isospin

G = G -parity

J = spin

P = space parity

C = charge conjugation parity .

1. Mesons

The charge conjugation operator C turns particle into antiparticle and has eigenvalues ± 1 only for neutral states; so it is useful to define an extension G which has eigenvalues for charged states too. It is usually¹ defined by

$$G = C \exp(i\pi I_y) . \quad (1)$$

A neutral nonstrange state is an eigenstate of $\exp(i\pi I_y)$ with eigenvalue $(-1)^I$. Then we can write the eigenvalue equation for the whole multiplet as

$$G = C_n (-1)^I , \quad (2)$$

where C_n (n for neutral) is the eigenvalue C would have if applied to the neutral member of the multiplet. Thus, for a π^0 , C has the eigenvalue $+1$, and since $I=1$, $G = -1$. For the charged pion there are no eigenvalues corresponding to C and to the isospin rotation, but Eqs. (1) and (2) still give $G = -1$.

Consider a meson as a bound state of fermion-antifermion, e.g., $\bar{q}q$, with orbital angular momentum l , and with the two fermion spins coupling to give a spin S . Then one can show that the charge-conjugation eigenvalue [defined in Eq. (2)] is

$$C_n = (-1)^{l+S} . \quad (3)$$

Equations (2) and (3) combine to give

$$G = (-1)^{l+S+I} . \quad (4)$$

The parity is

$$P = -(-1)^l . \quad (5)$$

Equations (3) and (5) combine to give

$$C_n P = -(-1)^S \quad (6)$$

so all singlet ($^1S_0, ^1P_1, \dots$) have $C_n P = -1$, and all triplet ($^3S_1, \dots$) have $C_n P = +1$. For proofs of the above, see our 1969 text (Particle Data Group, 1969) and Appendix by C. Zemach.

If, instead of $\bar{q}q$, we consider the meson as a state of *boson-antiboson* (e.g., $A_2 \rightarrow \bar{K}K$), it turns out that some signs cancel, and Eqs. (3) and (4) [not (5)!] apply *unchanged*. Of course the mesons are often spinless so S is zero, but the equations are more general. Eqs. (3) and (4) can be considered as selection rules forbidding many decays.

We now use Eqs. (3) and (4) to introduce the concept of "Abnormal- C " mesons, i.e., mesons that cannot be composed of $\bar{q}q$.

The unitary triplet of quarks is of course defined to have isospin and hypercharge properties such that $\bar{q}q$ can combine (according to the SU(3) relations $\{3\} \otimes \{3\} = \{8\} \oplus \{1\}$) so as to form only unitary octets and singlets. The non-observation of "exotic" mesons (i.e., mesons in more complicated supermultiplets) is of course one of the bases of the naive quark model. But it is slightly less obvious that even some *octets* are forbidden by the model, namely those with $(J^P)C_n = (0^*)-, (1^-)+, (2^*)-, \dots$. Such states are also not observed, and this is an additional success of the naive quark model classification scheme.

In what follows, do not confuse "Abnormal- C " with "Normal" or "Abnormal" J^P , both of which are allowed by the quark model. The series $J^P = 0^*, 1^-, 2^*, \dots$ is called Normal because $P = (-1)^J$ as for normal spherical harmonics, and $J^P = 0^-, 1^+, \dots$ is called Abnormal.

The top part of Table I shows all the low angular momentum states that can be formed from $\bar{q}q$. Note that half of the J^P states can be formed by both a triplet and a singlet $\bar{q}q$ state, e.g., $^3P_1, ^1P_1$ or $^3D_2, ^1D_2$. Eq. (3) shows that 3P_1 and 1P_1 have opposite C_n , so the $\bar{q}q$ model allows both. But the states 3P_0 and 3P_2 have no 1P counterparts. According to Eq. (6) they have $C_n P = +1$, and with the $\bar{q}q$ model there is no way to form a state with a J^P of $^3P_{0,2}$ (i.e., $J^P = \text{Normal}$) and with $C_n P = -1$. As mentioned, such octets have not shown up. With the help of Table I one can also see that the special state $^1S_0, C_n P = +1$, cannot be formed, so has Abnormal C .

2. General remarks

Well-established quantum numbers are underlined in the Tables of Particle Properties (except for stable particles, where most of the quantum numbers are established). We have used flimsy evidence to guess many of the remaining ones, and we have indicated with "?" ones (in the baryon table) for which there is almost no evi-

¹Most texts define it as in Eq. (1); see, e.g., Gasiorowicz (1966); however, sometimes the rotation is taken about I_x . The difference between the two conventions is mentioned in a footnote in Källén (1964).

TABLE I. $I^G(J^P)$ of non-strange mesons from $\bar{q}q$ model. For the distinction between abnormal J^P and abnormal C, see text following Eq. (6). K mesons share the same values of J^P as the $I=0$ and 1 states shown, but are not eigenstates of G. The middle column, which gathers together $(J^P)_{\text{Normal or abnormal}}$ or CP , is a redundant intermediate step intended to make the table easier to read.

Parity	$\bar{q}q$ State		$(J^P)_{\text{Normal or abnormal}}$ CP	$I^G(J^P)C_n$	Examples and comments
	CP	CP			
	-	+			
Parity -	$1S_0$		$(0^-)_{A^-}$	$\begin{cases} 0^+(0^-)+ \\ 1^-(0^-)+ \end{cases}$	$\eta, \eta', E ?$ π
	$3S_1$		$(1^-)_{N^+}$	$\begin{cases} 0^-(1^-)- \\ 1^+(1^-)- \end{cases}$	ω, ϕ ρ
Parity +	$1P_1$		$(1^+)_{A^-}$	$\begin{cases} 0^-(1^+)- \\ 1^+(1^+)- \end{cases}$	B
	$3P_0$		$(0^+)_{N^+}$	$\begin{cases} 0^+(0^+)+ \\ 1^-(0^+)+ \end{cases}$	ϵ, S^* $\delta ?$
	$3P_1$		$(1^+)_{A^+}$	$\begin{cases} 0^+(1^+)+ \\ 1^-(1^+)+ \end{cases}$	D ? A1
	$3P_2$		$(2^+)_{N^+}$	$\begin{cases} 0^+(2^+)+ \\ 1^-(2^+)+ \end{cases}$	f, f' A2
Parity -	$1D_2$		$(2^-)_{A^-}$	$\begin{cases} 0^+(2^-)+ \\ 1^-(2^-)+ \end{cases}$	A3
	$3D_1$		$(1^-)_{N^+}$	same as $3S_1$	ρ'
	$3D_2$		$(2^-)_{A^+}$	$\begin{cases} 0^-(2^-)- \\ 1^+(2^-)- \end{cases}$	Regge recurrence of the abnormal-C state $(J^P)_{C_n} = (0^-)-$
	$3D_3$		$(3^-)_{N^+}$	$\begin{cases} 0^-(3^-)- \\ 1^+(3^-)- \end{cases}$	g
Parity +	$1F_3$		$(3^+)_{A^-}$	$\{ J > 2 \}$	
	$3F_2$		$(2^+)_{N^+}$	same as $3P_2$	
	$3F_3$		$(3^+)_{A^+}$	$\{ J > 2 \}$	
	$3F_4$		$(4^+)_{N^+}$	etc.	

ABNORMAL C STATES THAT CANNOT COME FROM $\bar{q}q$ MODEL

Abnormal C states Have no $\bar{q}q$ model	$(0^-)_{A^+}$	$\begin{cases} 0^-(0^-)- \\ 1^+(0^-)- \end{cases}$	All except $J^P = 0^-$ are $J^P = \text{normal},$ $CP = -1$
	$(1^-)_{N^-}$	$\begin{cases} 0^+(1^-)+ \\ 1^-(1^-)+ \end{cases}$	
	$(0^+)_{N^-}$	$\begin{cases} 0^-(0^+)- \\ 1^+(0^+)- \end{cases}$	
	$(2^+)_{N^-}$	$\begin{cases} 0^-(2^+)- \\ 1^+(2^+)- \end{cases}$	
	$(3^-)_{N^-}$	$\begin{cases} 0^+(3^-)+ \\ 1^-(3^-)+ \end{cases}$	

dence.

As is customary, we define antiparticles as the result of operating with CPT on particles, so both share the same spins, masses, and mean lives. Whenever there is a particularly interesting test of CPT invariance we include it in the Stable Particles Table.

B. Particle names

If a *meson* has a well-accepted colloquial name, we use it. If not, we name it by a single symbol which specifies its baryon number B ($=0$ for mesons), its isospin I , its hypercharge Y , and, for a nonstrange meson its G parity. For convenience, we also list the strangeness S , which is related to Y and B by

$$S = Y - B.$$

The name conventions for mesons are given in the first part of Table II.

To crowd even more information onto the symbol, we sometimes add a subscript giving J^P . If J^P is not known, but must be "Normal" ($0^+, 1^-, 2^+, \dots$), e.g., because $K\pi$ decays are seen, we use the subscript N . If such modes are *not* seen (and are not otherwise forbidden), we *guess* that it is because J is "Abnormal", and we use the subscript A .

For *some* pairs of *mesons* with supposedly identical quantum numbers, we also use primes; e.g., η, η' ; f, f' .

For *baryons* no attempt has been made to attach a subscript about J and P . The name conventions are given in the second part of Table II. For stable baryons of each I and Y we use the symbol standing alone; for resonances, the mass is in parentheses [i.e., $N(1688)$, $\Lambda(1405)$, $\Sigma(1765)$, etc.]. The J^P assignments are reported in the Baryon Table as $\frac{1}{2}^+, \frac{3}{2}^-, \frac{5}{2}^+$, etc., and also by the symbols P_{11}, D_{13}, F_{15} , which refer to the πp or Kp partial-wave amplitude in which the resonant state occurs (the first subscript refers to the isospin state:

TABLE II. Particle name conventions.

Name	I	Y	S	G
Mesons				
η	0	0	0	+
ω or ϕ^a	0	0	0	-
ρ	1	0	0	+
π	1	0	0	-
K^+, K^0	$\frac{1}{2}$	+1	+1	
K^-, \bar{K}^0	$\frac{1}{2}$	-1	-1	
Baryons				
N	$\frac{1}{2}$	+1	0	
Δ	$\frac{3}{2}$	+1	0	
Z_0, Z_1	0, 1	+2	+1	
Λ	0	0	-1	
Σ	1	0	-1	
Ξ	$\frac{1}{2}$	-1	-2	
Ω	0	-2	-3	

^a Since 1973, we have used the symbol ω for those $I^G=0^-$ mesons that decay mainly into 3π [$\omega(783)$, $\omega(1670)$]; we reserve the symbol ϕ for $\phi(1020)$ and possible future higher-mass $I^G=0^-$ mesons that decay mainly into $K\bar{K}$.

$2 \times I$ for N and Δ and just I for Z , Λ , and Σ). When two or more baryons have identical quantum numbers we warn the reader by adding primes to the spectroscopic symbol as explained in footnote (a) of the Baryon Table.

IV. CONVENTIONS AND PARAMETERS FOR STRONG INTERACTIONS

A. Partial-wave amplitudes and resonance parameters

The vast majority of information concerning baryon resonances comes in the form of partial-wave analyses. In addition data concerning meson resonances ($\pi\pi, K\pi, \pi\pi\pi$) are, with increasing frequency, being subjected to partial-wave analyses. We thus find it natural to introduce the resonance parameters which we compile in terms of a Breit-Wigner approximation for the partial-wave amplitude.

In general the elastic amplitude for a given angular momentum l may be written as

$$T_{11} = \frac{\eta \exp(2i\delta) - 1}{2i}, \quad (1)$$

where η is referred to as the absorption parameter ($0 \leq \eta \leq 1$) and δ , as the phase shift. The subscripts 11 on T denote scattering from channel 1 to channel 1 (i.e., $\pi\pi \rightarrow \pi\pi$ or $\bar{K}N \rightarrow \bar{K}N$).

In Fig. 1 we show an Argand plot of the elastic partial wave amplitude T_{11} . It illustrates geometrically how the real parameters η and δ are related to the real and imaginary parts of T_{11} . Many examples of such Argand plots may be found in the Baryon Data Card Listings.

Consider the so-called non-relativistic Breit-Wigner approximation for T_{11} :

$$T_{11} = \frac{1}{2} \Gamma_1 / (M - E - \frac{1}{2} i\Gamma) \quad (2)$$

where E is the c.m. energy of invariant mass, Γ_1 and Γ are the *elastic* and *total* widths, and M is called the *resonance mass*. Equation (2) is, of course, not the only possible description of a resonant amplitude; it suffices to illustrate the properties of partial-wave amplitudes which we associate with resonance behavior in the absence of any background in the same partial wave (see, e.g., the $\pi N D_{15}$ and F_{15} waves in the Baryon Data Card Listings). Usually the widths contain barrier-

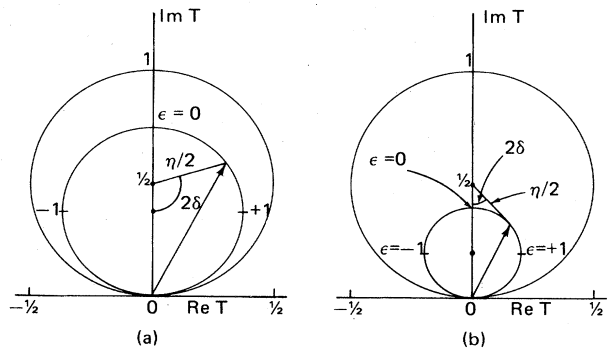


FIG. 1. Argand plots for the partial wave amplitude T_{11} . The outer circles are the unitarity bound ($\eta=1$). The inner circles correspond to the Breit-Wigner approximation of Eq. (2) for (a) $x_1 = \Gamma_1/\Gamma = 0.75$ and (b) $x_1 = 0.4$. Note: $\epsilon = 2(M - E)/\Gamma$.

penetration factors which can vary rapidly with energy. Near threshold, $\Gamma_1(E)$ should start up as q^{2l+1} (also true for the inelastic width Γ_β). Various E dependences are then used for Γ_1 , mostly of the form

$$\Gamma_1(E) \propto \frac{(qR)^{2l+1}}{\text{const.} + \dots + (qR)^{2l}}, \quad (3)$$

see Jackson (1964), Pišút and Roos (1968), and Barbaro-Galtieri (1968).

The BW-approximation to the amplitude for an inelastic process leading from channel 1 to channel β ($\pi\pi \rightarrow \bar{K}K$ or $KN \rightarrow \Sigma\pi$, for example), is

$$T_{1\beta} = \frac{1}{2}(\Gamma_1\Gamma_\beta)^{1/2}/(M-E - \frac{1}{2}i\Gamma) \\ = (x_1x_\beta)^{1/2}[\frac{1}{2}\Gamma/(M-E - \frac{1}{2}i\Gamma)], \quad (4)$$

where

$$\Gamma = \sum_1^N \Gamma_\beta, \quad x_\beta = \Gamma_\beta/\Gamma, \quad (5)$$

and x_1 (called the elasticity) is often written x_e . (Note that in the Data Card Listings we use the symbol P_β to denote x_β .) The channel cross-section $\sigma_{1\beta}$ for the reaction $1 \rightarrow \beta$, for spin $0 \rightarrow$ spin $\frac{1}{2}$ scattering, is

$$\sigma_{1\beta} = 4\pi\lambda^2(J + \frac{1}{2}) |T_{1\beta}|^2, \quad (6)$$

where $J = l \pm \frac{1}{2}$.

The important features of Eq. (4) which characterize resonant behavior in the Argand diagram ($\text{Im}T_{1\beta}$ versus $\text{Re}T_{1\beta}$) are:

(1) Energy variation given by circles with diameter $(x_1x_\beta)^{1/2}$ and maximum amplitude at $E=M$ of

$$T_{1\beta}^{\text{max}} = i(x_1x_\beta)^{1/2}. \quad (7)$$

(2) A maximum in the speed near resonance, given approximately by

$$\text{“Speed” (res)} = |dT_{1\beta}/dE|_{E=M} = \frac{2(x_1x_\beta)^{1/2}}{\Gamma(E)}, \quad (8)$$

for slowly varying $\Gamma(E)$. These features may be related to the η, δ representation of T_{11} . Thus when $E=M$, δ is either $90^\circ (x_1 > \frac{1}{2})$ or $0^\circ (x_1 < \frac{1}{2})$ and η dips to its minimum

value.

These simple properties can be used to judge the presence or absence of resonance behavior in an Argand plot. However, it must be kept in mind that Eqs. (2) and (4) are only approximations to the “true” amplitude. The simple picture given above can be distorted by various effects:

- the presence of “background” in the same partial wave as the resonance
- two resonances in the same partial wave overlapping in energy
- the resonant energy M being close to an inelastic channel threshold, in which case a K -matrix-like parametrization is more appropriate.

B. SU(3) sign conventions for Λ and Σ resonances

Consider the partial width Γ_β of a resonance decaying into the channel β . We can always define a coupling constant such that

$$\Gamma_\beta \propto G_\beta^2.$$

In this case the inelastic amplitude in the Breit-Wigner approximation, Eq. (4) will go as

$$T_{1\beta} \propto G_1G_\beta/(M-E - \frac{1}{2}i\Gamma),$$

where G_1 is the coupling constant for the elastic channel. In the context of exact SU(3) symmetry the relative signs of the product G_1G_β for different resonances are often useful as a consistency check on SU(3) assignment of Λ and Σ resonances. See Appendix II for further details.

In the Data Card Listings for Λ and Σ resonances, we tabulate measured values for $(x_1x_\beta)^{1/2} \propto G_1G_\beta$. Whenever there is an explicit sign, it will be according to the convention advocated by Levi-Setti (1969) and used in the table of SU(3) Isoscalar Factors presented in this review. Thus the signs multiplying the Breit-Wigner amplitudes for $\bar{K}N \rightarrow \Sigma(1385) \rightarrow \Sigma\pi$, $\Lambda\pi$ and $\bar{K}N \rightarrow \Lambda(1405) \rightarrow \Sigma\pi$ are simply the product of the phases of the appropriate isoscalar factors. This convention is shown in Fig. 2 from Levi-Setti (1969).

SU(3) RELATIVE SIGN OF RESONANT AMPLITUDES

$$T_{\text{RES}} \sim \alpha (G_{N\bar{K}Y} \cdot G_{Y\pi Y^*}) / (M-E - i\Gamma/2)$$

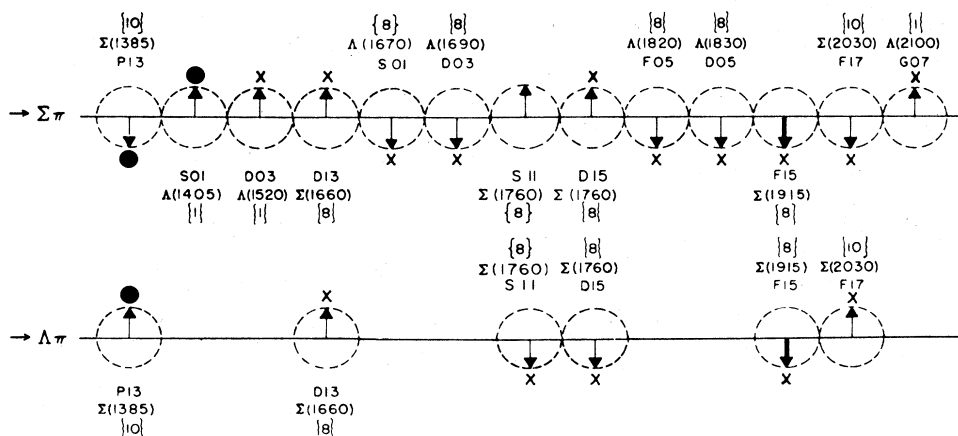


FIG. 2. Plot adapted from Levi-Setti (1969) showing the sign convention adopted here for the $\Sigma\pi$ and $\Lambda\pi$ amplitudes. Once the signs of one $I=0$ and one $I=1$ amplitude are fixed, the others can be measured relative to these two. Arrows here indicate signs predicted by SU(3); × marks indicate the observed phases; ● indicates phase chosen according to sign convention described in text. The $\Sigma(1915)$ predictions have been changed from Levi-Setti's original figure.

C. Types of partial-wave analyses

Partial-wave analyses (PWA) are classified into three categories in the Data Card Listings: energy-independent partial-wave analyses (IPWA), energy-dependent partial-wave analyses (DPWA), and model-dependent partial-wave analyses (MPWA), in increasing order of the number of explicit supplementary hypotheses that are used to extract the amplitudes from experimental data.

In an IPWA, data at different energies are analyzed separately. Usually each partial wave included in the fit is allowed to vary freely (subject to unitarity constraints) over some large region, and waves whose angular momenta are above some cutoff value are assumed to be negligible. The sharp cutoff in angular momentum resolves continuum ambiguities in the solution (such as the overall phase ambiguity), but there remains a finite number of indistinguishable "best" solutions (i.e., solutions corresponding to identical physical observables) which have been codified by Barrelet (1972). In addition, there are generally some nearby solutions (and their associated Barrelet ambiguities) which have chi-squared values close to the minimum one.

At the end of the analysis a choice is made among these many solutions, usually on the basis of energy continuity. A popular criterion for making this choice is the shortest path technique in which the total "length" of the preferred solution is chosen to be a minimum. The definition of "length" used here is not universal but is usually closely related to the total geometrical length of the lines representing the various partial-wave amplitudes in Argand plots (see the Baryon section of the Data Card Listings for examples of Argand plots). Various other criteria which are also used in some analyses are, e.g., matching with known solutions at low energies, the presence of known resonances in the final results, and limited inelasticity in high partial waves.

In a DPWA, data at different energies are fit simultaneously by using an energy dependent parametrization of the partial-wave amplitudes. The parametrization is usually chosen to include both resonances and nonresonant background of some sort and an attempt is made to keep it as "model independent" as possible. Often the data are grouped into several energy bins which are fit separately rather than trying to fit the whole energy range under consideration simultaneously. One of the main advantages of DPWA over IPWA is that sparse data spread over many different energies can be analyzed, e.g., nearly all $S=-1$ analyses are DPWA. In addition, the built-in energy continuity helps to resolve the ambiguities that plague IPWA and eases the problems associated with resonance parameter extraction. The price one pays for these advantages lies in the danger of systematic error in the amplitudes and poor fits to the data if the parametrization is poorly chosen or insufficiently flexible.

An MPWA also uses an energy-dependent parametrization, but one based on explicit model-dependent theoretical assumptions such as Regge exchanges. This technique is usually applied to reactions where the data are incomplete. There is, of course, no sharp distinction between DPWA and MPWA, and a well chosen MPWA

parametrization may actually be less biased than a model-independent but poorly chosen DPWA parametrization.

D. Production of resonances

Hereby, we mean the observation of statistically significant peaks in invariant mass plots or, loosely, in integrated cross sections. Many meson resonances are of this type. We expect most of these peaks to be associated with Breit-Wigner behavior in appropriate Argand plots; thus the ρ meson peak in $\pi\pi$ mass plots is firmly related to the $I=1, l=1$ $\pi\pi$ phase shift passing through 90° .

From mass plots we can determine M , Γ , and the approximate branching ratios

$$x_\alpha/x_\beta = \Gamma_\alpha/\Gamma_\beta . \quad (9)$$

In the case of total cross sections, the peak above background gives us, using the optical theorem, the product $(J+\frac{1}{2})x_e$:

$$\sigma^{\text{tot}}(E=M) = 4\pi\chi^2(J+\frac{1}{2})x_e . \quad (10)$$

V. CRITERIA FOR RESONANCES

An experimentalist who sees indications of a resonance in some energy (or mass) region will of course want to know what has been seen in that region in the past; hence, we strive to have the Data Card Listings serve as an archive for all substantial claims or evidences for resonances.

For the Tables of Particle Properties, on the other hand, we wish to be more conservative and to include only those peaks or resonances which we feel have a large chance of survival. An arrow (→) at the left of the Tables of Particle Properties indicates that a questionable candidate has been omitted from the Table, but that it can be found in the corresponding part of the Data Card Listings. One's betting odds for survival are of course subjective; therefore no precise criteria can be defined. In what follows we shall attempt to specify some guide lines.

(a) When energy-independent partial-wave analyses are available (mostly for N^* 's), approximate Breit-Wigner behavior of the amplitude appears to us to be the most satisfactory test for a resonance. We can check that the Argand plot follows roughly a left-hand circle, and that the "speed" of the amplitude also shows a maximum near the resonance energy; further, there should be data well above the resonance, showing that the speed again decreases. Indeed proper behavior of the partial-wave amplitude could accredit a resonance even if its elasticity is too small to make a noticeable peak in the cross section.

Of course even if Argand plots are available, it may still be a matter of opinion as to what behavior constitutes a resonance. Such an example is the $Z_0(1780)$ state seen in KN total cross-section experiments and in partial-wave analysis. The recent partial-wave analyses of Giacomelli (1974) and Martin (1975) find preferred solutions which exhibit a resonance-like loop in the P_{01} wave near 1740 MeV (see Fig. 5 of the $S=+1$

mini-review in the Baryon Data Card Listings). However, Giacomelli *et al.* and Martin point out that, despite the resonant-like *appearance* of the loop, the evidence for resonant *energy dependence* is inconclusive. Thus we omit the $Z_0(1780)$ from the Baryon Table. A similar quandary has existed for some time concerning the $Z_1(1900)$, and it too has been omitted from the Tables.

(b) When there are insufficient data to perform energy-independent analyses, one often resorts to energy-dependent partial-wave analyses (mostly for Y^* 's). In this case Breit-Wigner behavior is an input. We therefore require that resonance solutions be found by several different analyses, preferably in different channels ($\bar{K}N \rightarrow \bar{K}N, \pi\Sigma$, etc.), before putting the claim in the table.

(c) Partial-wave analyses of three-body final states ($\pi N \rightarrow \pi\pi N$) are becoming available. While these analyses are based on the isobar model ($\pi N \rightarrow \rho N, \pi\Delta$, etc.) and are subject to theoretical objections of varying importance, they provide increasingly reliable information on inelastic decay modes of otherwise established resonances.

(d) Most mesons, Ξ^* peaks, and high mass N^* and Y^* peaks fall into a category for which no partial-wave analyses exist. In general we accept such peaks if they are experimentally reliable, of high statistical significance or if they are observed in several different production processes.

(e) A special category of "diffractive mesons" consists of statistically significant peaks like A_1, A_3 or Q , which are not far above the $\rho\pi, f\pi$, or $K^*\pi$ thresholds. The question of a resonance interpretation for these states is complicated, because the behavior near threshold in these channels may be described by the Deck effect. Modern partial-wave analyses can shed considerable light on these problems. See the mini-reviews for details.

Thus, we enter into the tables of Particle Properties only states for which there is experimentally convincing evidence, and we expect that most of these will be confirmed as resonances.

VI. CONVENTIONS AND PARAMETERS FOR WEAK AND ELECTROMAGNETIC DECAYS

A. Muon-decay parameters

The μ -decay parameters describe the momentum spectrum (ρ and η), the asymmetry (ξ and δ), and the helicity (h) of the electron in the process $\mu^+ \rightarrow e^+ + \nu + \bar{\nu}$. Assuming a local and lepton-conserving interaction, the matrix element may be written as

$$\sum_i \langle \bar{e} | \Gamma_i | \mu \rangle \langle \bar{\nu} | \Gamma_i (C_i + C'_i \gamma_5) | \nu \rangle,$$

where the summation is taken over $i = S, V, T, A, P$. Using the definitions and sign conventions of Kinoshita and Sirlin (1957), we have for the momentum parameters

$$\rho = [3g_A^2 + 3g_V^2 + 6g_T^2] / D,$$

$$\eta = [g_S^2 - g_P^2 + 2g_A^2 - 2g_V^2] / D,$$

for the asymmetry parameters:

$$\xi = \frac{6g_S g_P \cos\phi_{SP} - 8g_A g_V \cos\phi_{AV} + 14g_T^2 \cos\phi_{TT}}{D},$$

$$\delta = [-6g_A g_V \cos\phi_{AV} + 6g_T^2 \cos\phi_{TT}] / D\xi,$$

and for the parameter describing the helicity of the electron:

$$h = \frac{2g_S g_P \cos\phi_{SP} - 8g_A g_V \cos\phi_{AV} - 6g_T^2 \cos\phi_{TT}}{D}.$$

Here

$$D = g_S^2 + g_P^2 + 4g_V^2 = 6g_T^2 + 4g_A^2,$$

$$g_i^2 = |C_i|^2 + |C'_i|^2,$$

and

$$\cos\phi_{ij} = \text{Re}(C_i^* C'_j + C'_i C_j^*).$$

The quantities g_i are defined to be real non-negative numbers, and the ϕ_{ij} are phase angles between the i -type and j -type interactions. Under the assumption of two-component neutrinos $C'_i = -C_i$ and $C'_j = -C_j$, the S, P , and T terms vanish, and ϕ_{AV} is the phase angle between C_A and C_V in the complex plane.

By using the above equations and the experimental determinations of ρ, η, ξ, δ , and h , limits can be placed on $g_S/g_V, g_A/g_V, g_T/g_V, g_P/g_V$, and ϕ_{AV} . The results, given in the Data Card Listings assume neither two-component neutrinos nor time-reversal invariance. If, however, two-component neutrinos are assumed, then $\sin\phi_{AV}$ is the amplitude of time-reversal violation. Note that most experiments study only the upper end of the spectrum where ρ and η are highly correlated, so they can only report ρ for $\eta=0$ and η for $\rho=\frac{3}{4}$. The values for ρ and η we use here were obtained by combining measurements of both upper and lower ends of the spectrum and turn out to be nearly uncorrelated.

Note also that the radiative corrections are unambiguous only when $g_S = g_T = g_P = 0$. The same limits on g_A/g_V and ϕ_{AV} are obtained, however, as when g_S, g_T , and g_P are left free.

Current values for the asymmetry parameters as well as $|g_A/g_V|$ and ϕ_{AV} are given in the Addendum to the Stable Particle Table. In addition, upper limits on $|g_S/g_V|, |g_T/g_V|$ and $|g_P/g_V|$ are given in the μ section of the Stable Particle Data Card Listings.

B. K -decay parameters

1. Dalitz plot for $K \rightarrow 3\pi$ decays

The small deviation from uniformity of the Dalitz plot for the 3π decay of the K meson is usually described by a "slope parameter" (Dalitz, 1956). For the τ and τ' decays of the charged K 's, and the τ^0 decay mode of the K_L^0 , we parametrize the Dalitz plot distribution by the expression

$$|M|^2 \propto 1 + g \frac{S_3 - S_0}{m_{\pi^+}^2} + h \left(\frac{S_3 - S_0}{m_{\pi^+}^2} \right)^2 + j \frac{S_2 - S_1}{m_{\pi^+}^2} + k \left(\frac{S_2 - S_1}{m_{\pi^+}^2} \right)^2 + \dots, \quad (1)$$

where $m_{\pi^+}^2$ has been introduced so as to make the coeffi-

cients g , h , j , and k dimensionless, and

$$s_i = (P_K - P_i)^2 = (m_K - m_i)^2 - 2m_K T_i, \quad i = 1, 2, 3$$

$$s_0 = \frac{1}{3} \sum_i s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2).$$

Here the P_i are 4-vectors, m_i and T_i are the mass and kinetic energy of the i th pion, and the index 3 is used for the odd pion.

The coefficient g is a measure of the slope in the variable s_3 (or T_3) of the Dalitz plot, while h and k measure the quadratic dependence on s_3 and $(s_2 - s_1)$, respectively. The coefficient j is related to the asymmetry of the plot and must be zero if CP invariance holds. Note also that if CP is good, g must be the same for τ^+ and τ^- , and similarly for h and k .

At present there is no compelling experimental evidence for the h , k , or j term (for upper limits on the j term, see Sec. B.3(b) below). Thus we stop the above expansion at the first term and list only g . Since different experiments use different forms for $|M|^2$, in order to compare the experiments we have converted to g whatever coefficients have been measured. See the mini-review in the K^* section of the Stable Particle Data Card Listings for details on this point. The results are given in the Addendum to the Stable Particle Table and in the K^* and K_L^0 sections of the Stable Particle Data Card Listings.

Relations among τ^+ , τ^{*+} , and τ^0 are predicted by the $\Delta I = \frac{1}{2}$ rule. See Appendix I for these relations and a discussion of this rule.

2. Form factors in K_{l3} leptonic decays

Assuming that only the vector current contributes to these decays, we write the matrix element as

$$M \propto f_+(t) [(P_K + P_\pi)_\mu \bar{u}_l \gamma_\mu (1 + \gamma_5) u_\nu] + f_-(t) [m_l \bar{u}_l (1 + \gamma_5) \mu_\nu], \quad (2)$$

where P_K and P_π are the four momenta of K and π mesons; m_l is the lepton mass; f_+ and f_- are dimensionless form factors which can depend only on $t = (P_K - P_\pi)^2$, the square of the four-momentum transfer to the leptons. f_+ and f_- are relatively real if time reversal invariance holds for these decays. $K_{\mu 3}$ experiments measure f_+ and f_- , while K_{e3} experiments are sensitive only to f_+ because the presence of the lepton mass makes the f_- term negligible.

(a) $K_{\mu 3}$ experiments

Analyses of $K_{\mu 3}$ data frequently assume a linear dependence of f_+ and f_- on t , i.e.

$$f_\pm(t) = f_\pm(0) [1 + \lambda_\pm (t/m_\pi^2)]. \quad (3)$$

Most $K_{\mu 3}$ data are adequately described by Eq. (3) for f_+ and a constant f_- (i.e. $\lambda_- = 0$). There are two equivalent parametrizations commonly used in these analyses:

(1) $\lambda, \xi(0)$ *Parametrization*. Analyses of $K_{\mu 3}$ data often introduce the ratio of the two form factors

$$\xi(t) = f_-(t)/f_+(t).$$

The $K_{\mu 3}$ decay distribution is then described by the two parameters λ_+ and $\xi(0)$ (assuming time reversal invariance and $\lambda_- = 0$). These parameters can be determined by three different methods:

Method A: By studying the Dalitz plot or the pion spectrum of $K_{\mu 3}$ decay. The Dalitz plot density is (See, e.g. Chounet *et al.*, 1972):

$$\rho(E_\pi, E_\mu) \propto f_+^2(t) [A + B\xi(t) + C\xi(t)^2],$$

where

$$A = m_K (2E_\mu E_\nu - m_K E'_\pi) + m_\mu^2 (\frac{1}{4} E'_\pi - E_\nu),$$

$$B = m_\mu^2 (E_\nu - \frac{1}{2} E'_\pi),$$

$$C = \frac{1}{4} m_\mu^2 E'_\pi,$$

$$E'_\pi = E_\pi^{\max} - E_\pi = \frac{m_K^2 + m_\pi^2 - m_\mu^2}{2m_K} - E_\pi.$$

Here E_π , E_μ , and E_ν are respectively the pion, muon, and neutrino energies in the kaon center of mass. The density ρ is fit to the data to determine the values of λ_+ , $\xi(0)$, and their correlation.

Method B: By measuring the $K_{\mu 3}/K_{e3}$ branching ratio and comparing it with the theoretical ratio (See, e.g., Fearing *et al.*, 1970) as given in terms of λ^+ and $\xi(0)$, assuming μ - e universality.

$$\Gamma(K_{\mu 3}^\pm)/\Gamma(K_{e3}^\pm) = 0.6457 + 1.4115\lambda_+ + 0.1264\xi(0) + 0.0192\xi(0)^2 + 0.0080\lambda_+\xi(0),$$

$$\Gamma(K_{\mu 3}^0)/\Gamma(K_{e3}^0) = 0.6452 + 1.3162\lambda_+ + 0.1246\xi(0) + 0.0186\xi(0)^2 + 0.0064\lambda_+\xi(0).$$

This cannot determine λ_+ and $\xi(0)$ simultaneously but simply fixes a relationship between them.

Method C: By measuring the muon polarization in $K_{\mu 3}$ decay. In the rest frame of the K , the μ is expected to be polarized in the direction \vec{A} with $\vec{P} = \vec{A}/|\vec{A}|$, where \vec{A} is given (Cabibbo and Maksymowicz, 1964) by

$$\vec{A} = a_1(\xi) \vec{p}_\mu - a_2(\xi) \left\{ \frac{\vec{p}_\mu}{m_\mu} \left[m_K - E_\pi + \frac{\vec{p}_\pi \cdot \vec{p}_\mu}{|\vec{p}_\mu|^2} (E_\mu - m_\mu) \right] + \vec{p}_\pi \right\} + m_K \text{Im} \xi(t) (\vec{p}_\pi \times \vec{p}_\mu).$$

If time-reversal invariance holds, ξ is real, and thus there is no polarization perpendicular to the K -decay plane. Polarization experiments measure the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$.

(2) λ, λ_0 *Parametrization*. Some of the more recent $K_{\mu 3}$ analyses have parametrized in terms of the form factors f_+ and f_0 which are associated with vector and scalar exchange respectively to the lepton pair. f_0 is related to f_+ and f_- by

$$f_0(t) = f_+(t) + [t/(m_K^2 - m_\pi^2)] f_-(t).$$

Here $f_0(0)$ must equal $f_+(0)$ unless $f_-(t)$ diverges at $t=0$. The earlier assumption that f_+ is linear in t and f_- is constant leads to f_0 linear in t

$$f_0(t) = f_0(0) [1 + \lambda_0 (t/m_\pi^2)].$$

With the assumption that $f_0(0) = f_+(0)$, the two parametrizations, $(\lambda_+, \xi(0))$ and (λ_+, λ_0) are equivalent as long as correlation information is retained. (λ_+, λ_0) correlations tend to be less strong than $(\lambda_+, \xi(0))$ correlations.

The experimental results for $\xi(0)$ and its correlation with λ_+ are listed in the K^+ and K_L^0 sections of the Stable Particle Data Card Listings in subsection XIA, XIB, or XIC depending on whether Method A, B or C discussed above was used. The corresponding values of λ^+ are listed in subsection $L+M$.

Because current experiments tend to use the (λ_+, λ_0) parametrization, we have added a subsection $L0$ for λ_0 results. Wherever possible we have converted $\xi(0)$ results into λ_0 results and vice versa.

(b) K_{e3} experiments

Analysis of K_{e3} data is simpler than that of $K_{\mu 3}$ because the second term of the matrix element assuming a pure vector current (Eq. 2 above) can be neglected. Here f_+ is usually assumed to be linear in t , and the linear coefficient λ_+ of Eq. (3) is determined.

If we remove the assumption of a pure vector current, then the matrix element for the decay, in addition to the terms in Eq. (2), would contain

$$+ 2m_K(f_S \bar{u}_1(1 + \gamma_5)u_\nu + (2f_T/m_K)(P_K)_\lambda (P_\pi)_\mu \bar{u}_1 \sigma_{\lambda\mu}(1 + \gamma_5)u_\nu),$$

where f_S is the scalar form factor, and f_T is the tensor form factor. In the case of the K_{e3} decays where the f_- term can be neglected, experiments have yielded limits on $|f_S/f_+|$ and $|f_T/f_+|$.

The K_{e3} results for λ_+ , $|f_S/f_+|$ and $|f_T/f_+|$ are listed in the subsections $L+M$, FS , and FT respectively of the K^+ and K_L^0 sections of the Stable Particle Data Card Listings.

See also the *Note on K_{13}^+ and K_{13}^0 Form Factors* in the K^+ section of the Stable Particle Data Card Listings for additional discussion of the $K_{\mu 3}$ parameters, correlations, and conversion between parametrization and also for a comparison of the experimental results.

3. CP violation in K^0 decays

We list parameters for four different reactions in which CP can be tested [for details, see Okun and Rubbia (1967), Steinberger (1969), and Wolfenstein (1969)].

(a) $K_S \rightarrow \pi^+ \pi^- \pi^0$. The quantity measured here is the ratio of amplitudes

$$A_S(K_S \rightarrow \pi^+ \pi^- \pi^0) / A_L(K_L \rightarrow \pi^+ \pi^- \pi^0) = x + iy. \quad (4)$$

If CPT invariance holds and there is no $I=3$ state present, then x can be neglected and CP violation would be observed as a nonzero y . We give the result for Eq. (4) in the K_L^0 section of the Stable Particle Table and under Branching Ratio $R4$ in the K_S^0 section of the Stable Particle Data Card Listings. Our procedure is to assume that $x=0$, and to list $(A_S/A_L)^2$ in the form of a branching ratio.

(b) *Charge asymmetry in $K_L \rightarrow 3\pi$ decays.* As mentioned above, the presence of a term in $(s_2 - s_1)$ in expression (1) describing the Dalitz plot distribution for τ^\pm , τ^0 decays of K mesons would be an indication of CP violation. Rather than listing values of the $(s_2 - s_1)$

coefficient j in Eq. (1), we choose to list σ_\pm from the equivalent expression

$$|M|^2 \propto 1 + \sigma_\pm (2/\sqrt{3})(T_+ - T_-)/T_{\pm \max} + (CP \text{ nonviolating terms}), \quad (5)$$

where T_\pm are the kinetic energies of the charged pions. We have momentarily abandoned the form involving the Mandelstam variables s_i in favor of Eq. (5) because the latter has been consistently used by experimenters searching for CP violation. We list σ_\pm among the CP -violating parameters at the back of the K_L^0 section of the Stable Particle Data Card Listings. Note that only upper limits have been reported for this quantity.

(c) *Asymmetry in the $K_L - \pi^\mp l^\pm \nu$ decays.* The quantity measured and compiled here is

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^- l^+ \nu) - \Gamma(K_L \rightarrow \pi^+ l^- \nu)}{\Gamma(K_L \rightarrow \pi^- l^+ \nu) + \Gamma(K_L \rightarrow \pi^+ l^- \nu)}.$$

This asymmetry violates CP invariance. If CPT is good, for a pure K_L^0 beam, δ can be written as

$$\delta = 2[(1 - |x|^2)/(1 + |x|^2)]\text{Re}\epsilon,$$

where x is the $\Delta S = \Delta Q$ -violating parameters defined in section B4, and ϵ is the parameter of the expansion

$$|K_L\rangle = [(1 + \epsilon)|K\rangle - (1 - \epsilon)|\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}, \quad (6a)$$

$$|K_S\rangle = [(1 + \epsilon)|K\rangle + (1 - \epsilon)|\bar{K}\rangle] / [2(1 + |\epsilon|^2)]^{1/2}. \quad (6b)$$

We give δ in the Addendum to the Stable Particle Table. In addition, in the K_L^0 CP -violation section of the Stable Particle Data Card Listings, we list δ separately for $K_L^0 \rightarrow \pi \mu \nu$ and $K_L^0 \rightarrow \pi e \nu$.

(d) $K_L \rightarrow 2\pi$ decay. The relevant parameters are

$$\eta_{+-} = A(K_L \rightarrow \pi^+ \pi^-) / A(K_S \rightarrow \pi^+ \pi^-) = |\eta_{+-}| \exp(i\phi_{+-}),$$

$$\eta_{00} = A(K_L \rightarrow \pi^0 \pi^0) / A(K_S \rightarrow \pi^0 \pi^0) = |\eta_{00}| \exp(i\phi_{00}),$$

ϵ , defined in Eqs. (6) above, and

$$\epsilon' = \frac{1}{2} i \sqrt{2} [\exp(i\delta_2 - \delta_0)] \text{Im}(A_2/A_0).$$

Here, A_i and δ_i are the amplitude and phase of $\pi\pi$ scattering at the K mass, defined by

$$\langle I=0 | T | K \rangle = \exp(i\delta_0) A_0,$$

$$\langle I=2 | T | K \rangle = \exp(i\delta_2) A_2.$$

Wu and Yang (1964) have derived the relationships

$$\eta_{+-} = \epsilon + \epsilon',$$

$$\eta_{00} = \epsilon - 2\epsilon'.$$

We give η_{+-} , η_{00} , ϕ_{+-} , and ϕ_{00} in the Addendum to the Stable Particle Table. The phases are measured directly, whereas the magnitudes η_{+-} and η_{00} are derived parameters. We use, as far as we can, the directly measured quantities as input and calculate η_{+-} and η_{00} from the values given by our constrained fits. Therefore, if one looks at the Data Card Listings, most of the $|\eta|$ measurements appear in the form of branching ratios, with appropriate comments. We then give the values of η_{+-} and $|\eta_{00}|^2$ in a separate list at the end of the CP -violating parameters section of the K_L^0 section of the Stable Particle Data Card Listings.

4. $\Delta S = \Delta Q$ rule in K^0 decays

The relative amount of $\Delta S \neq \Delta Q$ component present is measured by the parameter x , defined as

$$x = A(\bar{K}^0 \rightarrow \pi^- l^+ \nu) / A(K^0 \rightarrow \pi^- l^+ \nu).$$

We list $\text{Re}\{x\}$ and $\text{Im}\{x\}$ for both K_{e3} and $K_{\mu 3}$ at the end of the Stable Particle Data Card Listings and give values in the Addendum to the Stable Particle Table.

C. η -decay parameters

1. C -violation in η -decays

As a test of possible C -violation in electromagnetic interactions, a number of experiments have looked for possible charge asymmetries in the decays $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$. We list the following parameters:

(a) The left-right asymmetry

$$A = (N^+ - N^-) / (N^+ + N^-),$$

where $N^{(\pm)}$ means the number of events with the $\pi^{(\pm)}$ energy greater than the $\pi^{(\mp)}$ energy in the η rest frame.

(b) The sextant asymmetry

$$A_s = \frac{N_1 + N_3 + N_5 - N_2 - N_4 - N_6}{N_1 + N_2 + N_3 + N_4 + N_5 + N_6}$$

for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$. The numbers refer to the sextants of the Dalitz plot [see, for example, Layter (1972)]. A_s is sensitive to an $I=0$ C -violating asymmetry.

(c) The quadrant asymmetry A_q , defined in a similar way as A_s , but with each sector of the Dalitz plot now containing $\pi/2$ rather than $\pi/3$ radians. A_q is sensitive to an $I=2$ C -violating final state.

(d) The d -wave contribution to the C -violating amplitude in the decay $\eta \rightarrow \pi^+ \pi^- \gamma$. The upper limit for this contribution is measured by the parameter β , defined by

$$dN/d|\cos\theta| \propto \sin^2\theta(1 + \beta \cos^2\theta),$$

where θ is the angle between the π^+ and the γ in the dipion center of mass. A term proportional to $\cos^2\theta$ could also be due to p - and f -wave interference.

We list A for the decay modes $\eta \rightarrow \pi^+ \pi^- \pi^0$ and $\eta \rightarrow \pi^+ \pi^- \gamma$, A_s and A_q for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$, and β for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ in the η section of the Stable Particle Data Card Listings.

2. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \pi^0$

The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \pi^0$ may be fit by the distribution

$$|M(x, y)|^2 \propto 1 + ay + by^2 + cx + dx^2 + exy.$$

Here,

$$x = \sqrt{3} (T_+ - T_-) / Q,$$

$$y = (3T_0 / Q) - 1,$$

T_+ , T_- , T_0 are the kinetic energies of the π^+ , π^- , and π^0 in the η rest system, and $Q = m_\eta - m_{\pi^+} - m_{\pi^-} - m_{\pi^0}$. The coefficient of the term linear in x is sensitive to C -violation due to an $I=0$ or $I=2$ final state. We list papers presenting determinations of the parameters a , b , c , and d in the η section of the Stable Particle

Data Card Listings. However, we do not tabulate values of these parameters because the assumptions made by different authors are not compatible and do not allow comparison of the numerical values.

3. Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \gamma$

The Dalitz plot for the decay $\eta \rightarrow \pi^+ \pi^- \gamma$ may be fit to the expression

$$|M|^2 \propto 1 + 2\alpha z,$$

where

$$z = \frac{2}{3} \sum_{i=1}^3 \left[\frac{3}{m_\eta - 3m_\pi} \left(E_i - \frac{1}{3} m_\eta \right) \right]^2 = \frac{\rho^2}{\rho_{\max}^2}.$$

Here E_i is the energy of the i th pion in the η rest frame, and ρ is the distance to the center of the Dalitz plot. We list the parameter α in the η section of the Stable Particle Data Card Listings.

D. Baryon-decay parameters

1. A/V ratio for baryon leptonic decays

Consider the decay

$$B_i \rightarrow B_f + l + \nu.$$

Assuming V, A theory, neglecting "induced" scalar, "induced" pseudoscalar, and axial weak-magnetism terms, and neglecting the q^2 dependence of the form factors, the baryon part of the matrix element for these decays may be written (Goldberger and Treiman, 1958) as

$$\langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) + (g_W / m_{B_i}) \sigma^{\lambda\nu} q_\nu | B_i \rangle,$$

where B_i and B_f represent initial and final baryons, g_A and g_V the axial and vector coupling constant, g_W the weak magnetism coupling constant, and q_ν the sum of the lepton momenta. Here the Pauli representation is used for the γ matrices. The definition of g_A/g_V is

$$g_A/g_V = |g_A/g_V| \exp(i\delta),$$

where δ is $0 + n\pi$ if time-reversal invariance holds (see Jackson *et al.*, 1957).

In neutron beta decay the measurements are consistent with time reversal, so g_A/g_V is nearly real and has been considered to be such in all the baryon leptonic decays. Notice that by using the above definition of the matrix element with the Pauli representations, the value of g_A/g_V in neutron beta decay is negative.

Due to statistical limitation the weak magnetism form factor g_W is usually assumed from CVC and SU(3), so only g_A and g_V are determined experimentally. This determination is accomplished in a variety of ways:

(a) The lepton-neutrino angular correlation provides a measure of the absolute value of g_A/g_V (for relevant formulas, see, e.g., Albright, 1959).

(b) The up-down asymmetry of the lepton from polarized baryon decays provides a measure of g_A/g_V with its sign (for relevant formulas, see, e.g. Albright, 1959).

(c) The lepton spectrum, given enough statistics, provides a measure of g_A/g_V with its sign (for relevant formulas, see, e.g., Bender, 1968).

(d) The polarization of the decay baryon, from polar-

ized or unpolarized initial baryon, also provides g_A/g_V with its sign (for formulas, see, e.g., Willis and Thompson, 1968).

We compile the ratio g_A/g_V with its sign, for those decays for which it has been measured. For the neutron beta decay we compile also the phase δ .

All the coupling constants and decay rates for baryon leptonic decays are related by Cabbibo's theory (Cabbibo, 1964). A recent fit to this theory has been done by Roos (1974).

2. Asymmetry parameters in nonleptonic hyperon decays

The transition matrix for the hyperon decay may be written as

$$M = s + p(\vec{\sigma} \cdot \vec{q}), \quad (7)$$

where s and p are the parity-changing and the parity conserving amplitudes, respectively; $\vec{\sigma}$ is the Pauli spin operator, and \vec{q} is a unit vector along the direction of the decay baryon in the hyperon rest frame.

The asymmetry parameters are defined by the relations

$$\begin{aligned} \alpha &= 2 \operatorname{Re}(s^*p) / (|s|^2 + |p|^2), \\ \beta &= 2 \operatorname{Im}(s^*p) / (|s|^2 + |p|^2), \\ \gamma &= (|s|^2 - |p|^2) / (|s|^2 + |p|^2). \end{aligned}$$

With the transition matrix (7), the angular distribution of the decay baryon, in the hyperon rest system, is of the form

$$I = 1 + \alpha \vec{P}_Y \cdot \vec{q},$$

where $\vec{P}_Y = \langle Y | \sigma | Y \rangle$ is the hyperon polarization.

In the notation of Lee and Yang (1957) the polarization \vec{P}_B of the decay baryon is²

$$\vec{P}_B = \frac{(\alpha + \vec{P}_Y \cdot \vec{q})\vec{q} + \beta(\vec{P}_Y \times \vec{q}) + \gamma\vec{q} \times (\vec{P}_Y \times \vec{q})}{1 + \alpha \vec{P}_Y \cdot \vec{q}}$$

where \vec{P}_B is defined in that rest system of the baryon obtained by a Lorentz transformation along \vec{q} from the hyperon rest system in which \vec{q} and \vec{P}_Y are defined. Note that α is the helicity of the decay baryon for unpolarized hyperons.

The three parameters α , β , and γ satisfy the relation

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

It is then convenient to describe hyperon nonleptonic decays in terms of the two independent parameters α and the angle ϕ defined by

$$\begin{aligned} \beta &= (1 - \alpha^2)^{1/2} \sin \phi, \\ \gamma &= (1 - \alpha^2)^{1/2} \cos \phi, \end{aligned}$$

which has a more nearly Gaussian distribution than β or γ . Evidently

²Note that Lee and Yang (1957) contains a misprint. The minus sign in the definition of β should be replaced by a 2. In addition, our unit vector \vec{q} is the direction of the baryon, whereas their unit vector \vec{p} is the direction of the pion.

$$-\frac{1}{2}\pi \leq \phi \leq \frac{1}{2}\pi \quad \text{for } \gamma > 0.$$

$$+\frac{1}{2}\pi \leq \phi \leq \frac{3}{2}\pi \quad \text{for } \gamma < 0.$$

In discussing time-reversal invariance, the quantity of interest is Δ , defined by

$$\begin{aligned} \alpha &= 2 |s| |p| \cos \Delta / (|s|^2 + |p|^2), \\ \beta &= -2 |s| |p| \sin \Delta / (|s|^2 + |p|^2); \end{aligned}$$

that is, Δ is the phase angle of s relative to p . Evidently

$$-\frac{1}{2}\pi \leq \Delta \leq \frac{1}{2}\pi \quad \text{for } \alpha > 0,$$

$$+\frac{1}{2}\pi \leq \Delta \leq \frac{3}{2}\pi \quad \text{for } \alpha < 0.$$

Under the assumption of time-reversal invariance, the angle Δ must satisfy the relation

$$\Delta = \delta_s - \delta_p,$$

modulo π , where δ_s and δ_p are the pion-baryon scattering phase shifts at the appropriate energy and for the appropriate isospin state. For Λ decay, assuming the validity of the $|\Delta I| = \frac{1}{2}$ rule,

$$\Delta = \delta_s - \delta_p = (6.8 \pm 2.0) \text{ deg.}^3$$

In the Stable Particle Data Card Listings we give α and ϕ for each decay since they are the most closely related to the experiments and are essentially uncorrelated. Whenever necessary we have changed the signs of the reported values, so as to agree with our conventions. In the Stable Particle Table we give α , ϕ , and Δ with errors; and for convenience we also give the central value of γ , without an error.

VII. STATISTICAL PROCEDURES

We divide this discussion on obtaining averages and errors into two sections:

- A. the unconstrained case, or "simple averaging", and
- B. the constrained case.

In what follows, the term "error" means one standard deviation (1σ); that is, for central value \bar{x} and error $\delta\bar{x}$, the range $\bar{x} \pm \delta\bar{x}$ constitutes a 68.3% confidence interval.

A. Unconstrained averaging

We first describe the standard procedure which we have used for several years to determine averages and errors. We will then discuss a second method, newly proposed, which we feel offers a less conservative, and possibly more accurate, estimate of errors.

1. Standard procedure—Gaussian distribution with scale factor

We begin by assuming that measurements of a given quantity obey a Gaussian distribution, and thus we calculate a weighted average and error

³This value for $\delta_s - \delta_p$ is derived from the phase-shift analyses by Roper *et al.* (1965). The error is our estimation of the uncertainty.

$$\bar{x} \pm \delta\bar{x} = \left(\frac{\sum_i w_i x_i}{\sum_i w_i} \right) \pm \left[1 / \left(\sum_i w_i \right) \right]^{-1/2}, \quad (1)$$

$$w_i = [1/(\delta x_i)^2],$$

where x_i and δx_i are the value and error, respectively, reported by the i th experiment, and the sums run over N experiments. We also calculate χ^2 and compare it with its expectation value of $N - 1$.

If $\chi^2/(N - 1)$ is less than or equal to 1, and there are no known problems with the data, we accept the above results.

If $\chi^2/(N - 1)$ is ridiculously large, or if there is prior knowledge of extremely large inconsistencies between experiments, we may choose not to average the data at all. Or, in some cases, we may quote the calculated average, but then give an "educated guess" as to the error; such a guess is a generally quite conservative estimate designed to take into account known problems with the data.

Finally, if $\chi^2/(N - 1)$ is greater than 1, but not to such a large extent, we still average the data, but then try to make up for this fact in two ways:

(i) We plot an ideogram to guide the reader in deciding which data he might reject before making his own selected average. An example of such an ideogram is given in Fig. 3 below. Each experiment appearing in the plot is represented by a Gaussian with central value x_i , error δx_i , and area proportional to $1/\delta x_i$. The choice of area is a somewhat arbitrary one; it is based on the assumption that an experimenter will work to reduce his systematic errors until they are slightly smaller (but seldom much smaller) than his statistical

errors. Thus as a bubble-chamber physicist gets more events, he will use them both to reduce his statistical errors and to study his biases. Our confidence that a significant systematic error has not been made in his experiment, as compared with other contradictory experiments, then tends to go up as $1/\delta x_i$.

But why not assign a weight $1/\delta x_i^2$, as is done when computing a weighted average? We feel that this is equivalent to assuming that large systematic errors are as infrequent as large statistical fluctuations, and that this is unrealistic.

We want to emphasize the difference between least-squares averaging (where the weighting factor is the inverse square of the error) and the ideograms prepared for visual display. The former arithmetic is of course best if one has statistically distributed input, and yields a narrow Gaussian distribution centered at the weighted mean. The ideogram (often multi-peaked and certainly not Gaussian) is based on the opposite hypothesis that some of the input is systematically in error. The idea behind least-squares averaging is that experiments 1, 2, 3, etc., are *all* valid (so we should multiply their probabilities); our *ideograms* are based on the assumption that 1 *or* 2 *or* 3, etc., is valid, "hedged" with $1/\delta x_i$ betting odds; we then add their probabilities. Both approaches cannot simultaneously be right; we leave it to the reader to choose. A glance at the ideogram will show, however, that the discrepancy is often not severe for reasonably distributed input.

(ii) The second way in which we try to take account of $\chi^2/(N - 1)$ being greater than 1 is to scale up our quoted error $\delta\bar{x}$ in Eq. (1) by a factor

$$\text{SCALE} = [\chi^2/(N - 1)]^{1/2}. \quad (2)$$

Our reasoning is as follows. Since we don't know which one or more of the experiments are wrong, we assume that all experimentalists underestimated their errors by the same scale factor (2). If we scale up all input errors by this factor, χ^2 returns to $N - 1$, and of course the output error scales up by the same factor.

If all the experiments have errors of about the same size, the above (straightforward) procedure for calculating SCALE is carried out. If, however, we are to combine experiments with widely varying errors, we must modify the procedure slightly. This is because it is the more precise experiments that most influence not only the average value \bar{x} , but also the error $\delta\bar{x}$. Now, on the average, the low-precision experiments each contribute about unity to *both* the numerator and the denominator of SCALE, hence the χ^2 contribution of the sensitive experiments is diluted, i.e., reduced. Therefore, we evaluate SCALE by using *only* experiments for which the error are not much greater than those of the more precise experiments. Explicitly, to calculate SCALE we use only the most sensitive experiments, i.e., those with errors less than δ_0 , where the ceiling δ_0 is (arbitrarily) chosen to be

$$\delta_0 = 3N^{1/2}\delta\bar{x}.$$

Here $\delta\bar{x}$ is the unscaled error of the mean of all the experiments. Note that if each experiment had the same error δx_i , then $\delta\bar{x}$ would be $\delta x_i/N^{1/2}$, so each individual

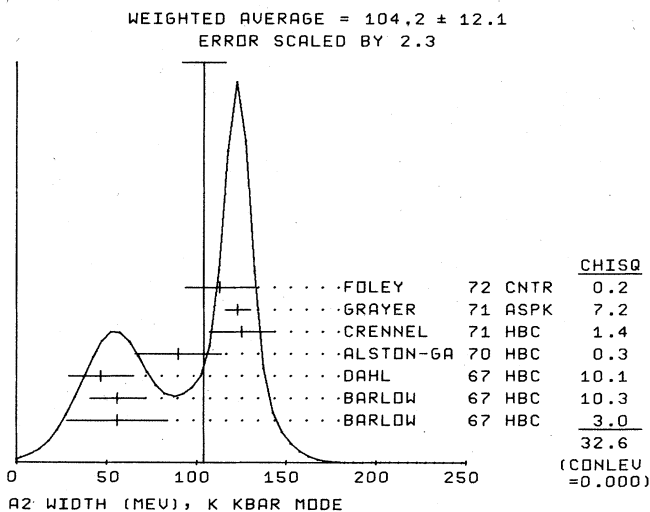


FIG. 3. Ideogram of measurements of the A_2 width, as determined from the $K\bar{K}$ mode. The vertical line indicates the position of the weighted average, while the horizontal bar atop the line gives the error in the average after scaling by the SCALE factor. Only those experiments indicated by + error flags were precise enough to be accepted in the calculation of the SCALE factor; the column on the far right gives the χ^2 contribution of each of these experiments. The less precise experiments were included in the calculation of the weighted average, but not SCALE; they have \pm error flags.

experiment would be well under the ceiling on SCALE.

This scaling approach has the property that if there are two values with comparable errors separated by much more than their stated errors (with or without a number of other experiments of lower accuracy), the error on the mean value $\delta\bar{x}$ is increased so that it is approximately half the interval between the two discrepant values.

We wish to emphasize the fact that our scaling procedures for errors in no way affect central values. In addition, if one wishes to recover the unscaled error $\delta\bar{x}$, he need only divide the given error by the SCALE factor for that error.

2. A second procedure—Student's distribution

The newly proposed method of averaging data, described in detail in Roos *et al.* (1975), relies upon an empirical determination of the distribution of the residuals for the ensemble of data appearing in the Review. The residual for the i th measurement of a quantity with average value \bar{x} is defined as

$$h_i = (x_i - \bar{x}) / \delta x_i.$$

Roos *et al.* select several different subsamples of the data, and show that the residuals for each subsample have approximately the same properties; in particular, their first few even moments are similar. Since the distributions have longer tails than a Gaussian, the authors choose to represent them by a distribution function having such a property, namely the Student distribution

$$S_n(h/c) = K \left[1 + \frac{(h/c)^2}{n} \right]^{-(n+1)/2}. \quad (3)$$

Here K is a normalization constant, and n and c are parameters which the authors then fit to the combined sample of data. The resulting empirical distribution is

$$S_{10}(h/1.11) = 0.351 \left[1 + \frac{(h/1.11)^2}{10} \right]^{-11/2}. \quad (4)$$

Note that the shape of S_{10} is somewhere between that of a Gaussian ($=S_\infty$) and that of a Breit-Wigner ($=S_1$).

The proposed method of averaging the data for a given quantity then consists of finding the value of \bar{x} which maximizes the log-likelihood function

$$\text{Log} \mathcal{L}(\{x_i\} | \bar{x}) = \sum_i S_{10} \left(\frac{x_i - \bar{x}}{1.11 \delta x_i} \right); \quad (5)$$

the sum here is again taken over all N measurements of x . The error $\delta\bar{x}$ is determined by finding the variation in \bar{x} needed to decrease the log-likelihood by 1/2:

$$\text{log} \mathcal{L}(\{x_i\} | \bar{x}) - \text{log} \mathcal{L}(\{x_i\} | \bar{x} \pm \delta\bar{x}) = \frac{1}{2}. \quad (6)$$

3. Comparison of procedures

Both of the procedures described above adopt a partially empirical approach to the problem that measured values for the quantities tabulated in this Review do not exhibit the Gaussian behavior naively expected. (This problem, it should be noted, persists even when careful attempts are made to resolve difficulties and inconsistencies in the data prior to averaging.)

The first approach operates on a quantity-by-quantity basis and adjusts the error in each case so that no scaled $\chi^2/(N-1)$ is greater than 1. This is obviously rather conservative, since even if the data obeyed a Gaussian distribution, about half of the quantities would be expected to have $\chi^2/(N-1) > 1$.

The second approach, on the other hand, assumes that (provided we first eliminate quantities with obvious, known problems) all quantities have the same theoretical distribution function, namely the fairly long-tailed $S_{10}(h/1.11)$. With this supposition, if a particular quantity has a large χ^2 , it is assumed to be just a happenstance, occasioned by a random fluctuation into the long tails, and no special scaling for this quantity is done. This procedure thus results in generally smaller, or less conservative, error estimates for quantities having $\chi^2/(N-1) > 1$. (However, it should be noted that, because of the overall scale of 1.11 appearing in the empirical Student's distribution, the errors for quantities with $\chi^2/(N-1) \leq 1$ are actually increased by about 10%.) Table III shows some comparisons of sample results from the two procedures, using data from the 1974 edition of the Review. Shifts in both \bar{x} and $\delta\bar{x}$ can be observed, especially where $\text{SCALE} > 1$.

Since the new procedure is a significant departure from the past, we have adopted the following approach for this year: in the Data Card Listings we give the average-and-error for each quantity calculated both ways; the standard way is labeled at the left with the code "AVG," while the proposed newer way is labeled

TABLE III. Comparison of procedures.

Particle property	Pure Gaussian $\bar{x} \pm \delta\bar{x}$	Standard method: Gaussian + scale factor $\bar{x} \pm \delta\bar{x}$	SCALE	Proposed method: Student's distribution $\bar{x} \pm \delta\bar{x}$
ρ^0 mass (MeV)	770.32 ± 0.65	770.32 ± 0.91	1.4	770.37 ± 0.82
η' mass (MeV)	957.59 ± 0.24	957.59 ± 0.24	1.0	957.58 ± 0.28
ϕ mass (MeV)	1019.69 ± 0.15	1019.69 ± 0.28	1.9	1019.83 ± 0.20
$K^*(1420)$ mass (MeV)	1421.3 ± 2.3	1421.3 ± 2.3	1.0	1421.3 ± 2.6
K_S^0 mean life (10^{-8} sec)	5.158 ± 0.042	5.158 ± 0.042	1.0	5.158 ± 0.046
Σ^+ mean life (10^{-10} sec)	0.8004 ± 0.0058	0.8004 ± 0.0058	1.0	0.8004 ± 0.0064
Σ^- mean life (10^{-10} sec)	1.482 ± 0.011	1.482 ± 0.017	1.5	1.479 ± 0.013
$K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ (%)	5.521 ± 0.075	5.521 ± 0.098	1.3	5.533 ± 0.089
$\Lambda \rightarrow p \pi^-$ (%)	63.99 ± 0.49	63.99 ± 0.49	1.0	63.98 ± 0.55

“STUDENT.” In the Tables of Particle Properties, we continue to use the standard procedure—Gaussian with SCALE factor. As in the past, a SCALE factor greater than 1 is indicated by the appearance of “S=...” next to the value and error.

We heartily invite your comments on the proposed Student’s distribution method. They will assist us in deciding on procedures for future editions.

B. Constrained fits

Except for trivial cases, all branching ratios and rate measurements are analyzed by computer program AHR. This program makes a simultaneous least-squares fit to all the data, and outputs the partial-decay fractions \bar{P}_i , width Γ , partial widths Γ_i , and their error matrix.

The original version of AHR was written by J. Peter Berge. It is documented separately, and we wish here only to give the simplest nontrivial example that permits us to comment on the error matrix and the scale factor.

Assume that a state has only three partial-decay fractions, P_1 , P_2 , and P_3 ($\sum P_i = 1$), which have been measured in four different ratios, R_1, \dots, R_4 , where, e.g., $R_1 = P_1/P_2, R_2 = P_1/P_3$, etc.⁴ Further assume that *each* ratio has been measured by N experiments (we designate each experiment with a subscript x , e.g., R_{1x}). Then AHR finds the best values of P_1 , P_2 , and P_3 by minimizing χ^2 , namely

$$\chi^2 = \sum_{r=1}^4 \left[\sum_{x=1}^N \left(\frac{R_{rx} - R_r(P_1, P_2, P_3)}{\delta R_{rx}} \right)^2 \right]. \quad (7)$$

In addition to the fitted values \bar{P}_i , the program calculates an error matrix $\langle \delta \bar{P}_i \delta \bar{P}_j \rangle$. We tabulate the diagonal elements $\delta \bar{P}_i = \langle \delta \bar{P}_i \delta \bar{P}_i \rangle^{1/2}$ [except that some errors are scaled according to Eq. (2) as discussed below]. In the listings we give the complete error matrix; we also calculate the fitted value of each ratio, for comparison with the input data, and list it below the relevant input, along with a simple unconstrained average of the same input.

Two further comments on the example above:

(1) There was no connection between measurements of the width and the branching ratios. But often we also have information on partial widths Γ_i as well as total width Γ . In this case AHR must introduce Γ as a parameter into the fit, along with the relations $\Gamma_i = \Gamma P_i$, $\sum \Gamma_i = \Gamma$. When appropriate, we tabulate the Γ_i along with the P_i , and give error matrices in the listings.

(2) Note that we do *not* allow for correlations between input data. We *do* try to pick those ratios and widths which are as independent and as close to the original data as possible.

In *asymmetric* errors, we use a continuous function of $\delta(P)^+$ and $\delta(P)^-$ in the fitting. When no errors are reported, we merely list the data for inspection.

Hyperon-decay parameters

The program AHR handles any type of input, α , Φ , Δ , β , or γ , according to the definitions of Sec. VI. If for a

⁴We can handle any R of the form $R = \sum \alpha_i P_i / \sum \beta_i P_i$, where α_i and β_i are constants, usually 1 or 0.

particular hyperon decay there are data for more than two of the decay parameters, they are analyzed by using the constraint

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

Inconsistent constrained data

According to our simple example, which led to Eq. (7), the double sum for χ^2 is summed over experiments $x = 1$ to N , leaving a single sum over ratios

$$\chi^2 = \sum_r \chi_r^2.$$

Even before fitting, some of the χ_r^2 may be too large. But if we scaled them before fitting, then the scaling would move the central value, contrary to our policy. So we do not scale until after the first fit; then, knowing the fitted χ_r^2 and its expectation value $\langle \chi_r^2 \rangle$ we form SCALE factors (just as before), i.e.,

$$(\text{SCALE})_r^2 = \chi_r^2 / \langle \chi_r^2 \rangle,$$

and if any $(\text{SCALE})_r$ is greater than 1, all N of the measurements of that particular ratio are equally penalized by having their errors increased by SCALE. Program AHR then recycles on all the data, those with errors unchanged as well as those with errors increased. We then get new values, $\delta \bar{P}'_i$ for the errors in the partial decay modes.

Because of the constraint ($\sum P_i = 1$) some SCALE factors may still be greater than 1 even after this second pass. If this is so, the whole procedure (i.e., increasing errors by the new SCALE factors and recycling through AHR) is repeated.

At the end of AHR’s final pass we have *two* measures of the errors for the \bar{P}_i . One is, of course, the $\delta \bar{P}'_i$, i.e., the errors in the final fitted values \bar{P}'_i which include the effects of scaling the input errors. The other measure of the errors is $(\bar{P}_i - \bar{P}'_i)$, i.e., the *shift* in the central values of the i th mode between the first (unscaled) fit and the final (scaled) fit. In practice we find that on the average these two measures of the uncertainty are about equal. Rather than selecting just one or the other, our tabulated errors are given by the combination

$$(\delta \bar{P}_i)_{\text{tab}} = [\delta \bar{P}'_i{}^2 + (\bar{P}_i - \bar{P}'_i)^2]^{1/2},$$

where \bar{P}_i is the fitted value of the i th partial-decay mode before scaling, \bar{P}'_i is its value after scaling, and $\delta \bar{P}'_i$ is the error in \bar{P}'_i . The SCALE factors we finally list in such cases are defined by

$$(\text{SCALE})_i = (\delta \bar{P}_i)_{\text{tab}} / \delta \bar{P}_i.$$

However, in line with our policy of not letting SCALE affect the central values, we give the values of \bar{P}_i obtained from the original (unscaled) fits. [The differences between the \bar{P}_i calculated with either the scaled or the unscaled errors are, of course, always within the tabulated errors, $(\delta \bar{P}_i)_{\text{tab}}$.]

VIII. A LOOK AT HISTORY

It may be said that one can estimate the age of a high energy physicist by asking him or her the mass of the

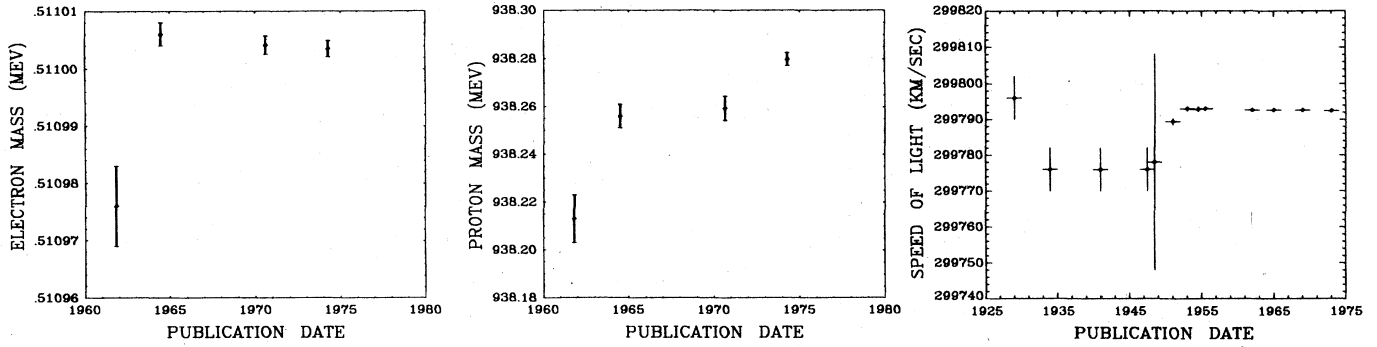


FIG. 4. The "generally accepted values" of the proton mass, the electron mass, and the speed of light, as a function of the publication data of the compilation used (not done by Particle Data Group). Data for the speed of light plot courtesy of E. R. Cohen, Rockwell International Science Center. See the Stable Particle Data Card Listings for references on proton and electron masses.

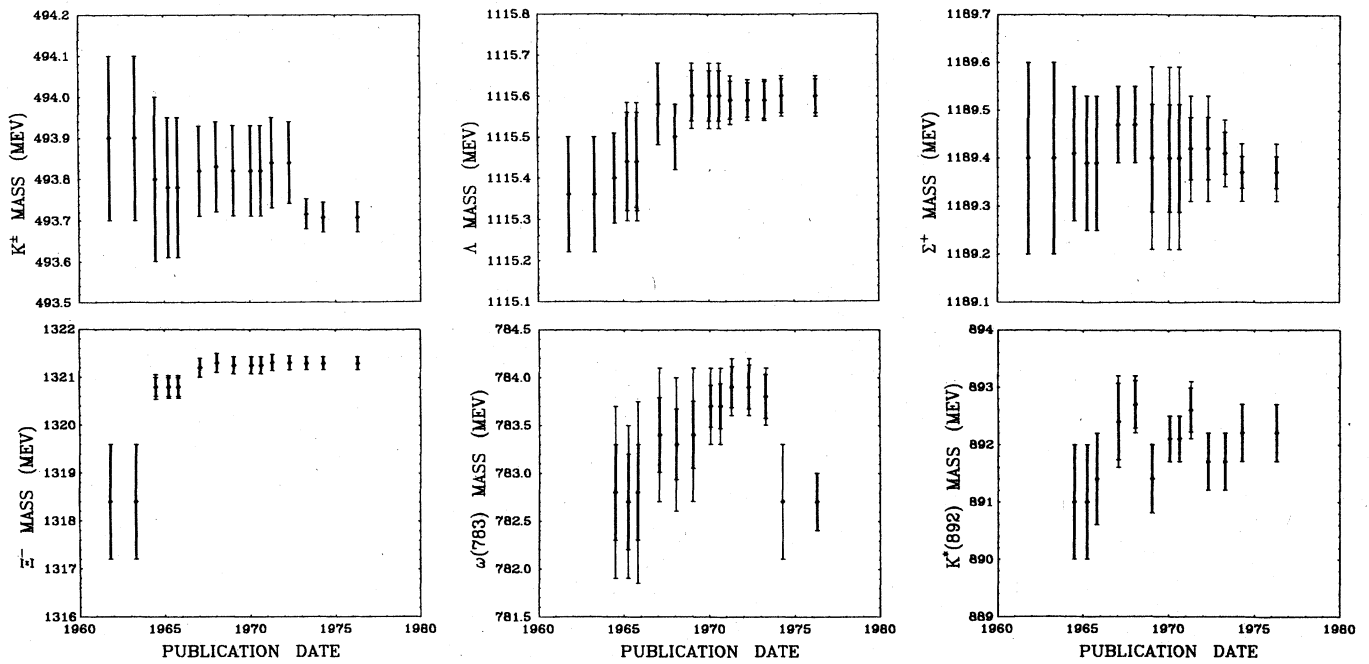


FIG. 5. Particle Data Group averages of the masses of various particles, as a function of date of publication of Review of Particle Properties (Adapted, with permission, from *Annual Review of Nuclear Science*, Volume 25. Copyright © 1975 by Annual Reviews, Inc. All rights reserved).

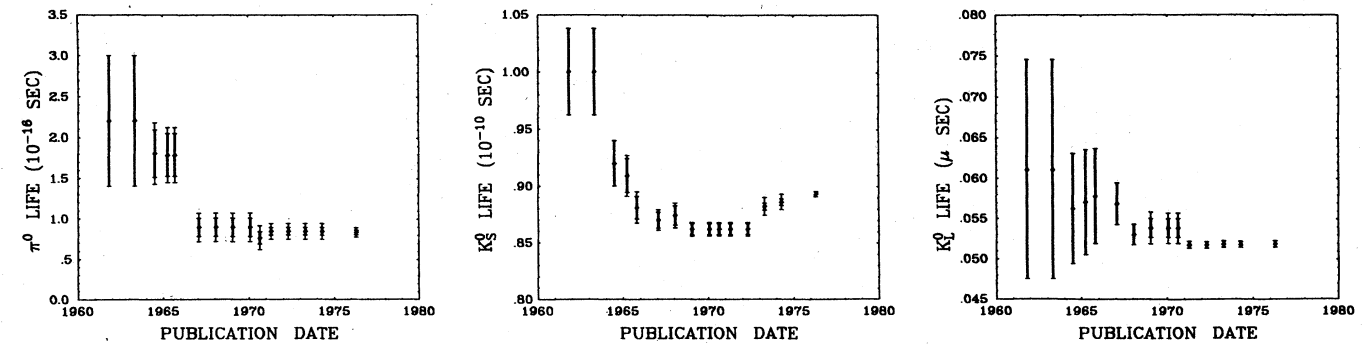


FIG. 6. Particle Data Group averages of the lifetimes of various particles, as a function of publication date of RPP.

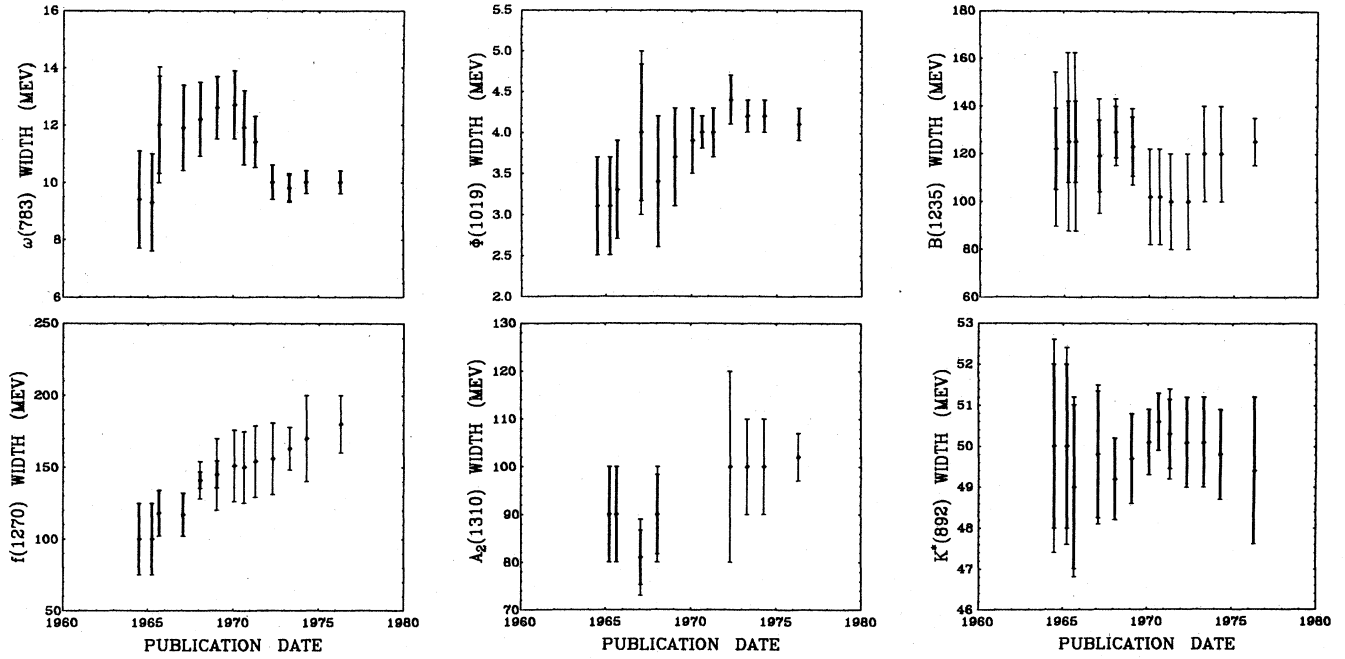


FIG. 7. Particle Data Group averages of the widths of various resonances, as a function of date of publication of RPP. The gap in the A_2 data indicates the years when the A_2 was thought to be split.

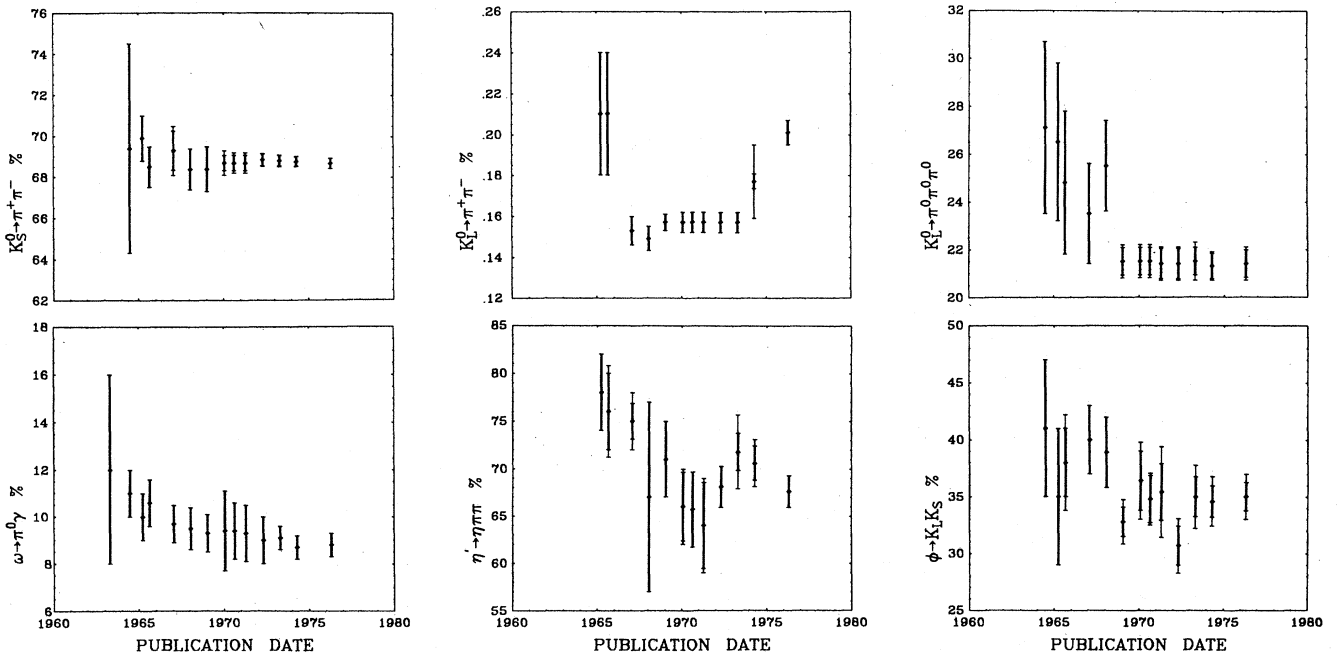


FIG. 8. Particle Data Group averages of various branching fractions, as a function of date of publication of RPP.

Λ^0 . If the answer is 1115.44 MeV, he probably was deep into his graduate training in 1965.

A history of the values of physical constants has more than whimsical value. In Fig. 4 we show how the generally accepted values for the speed of light and a couple of other constants have changed with time. The "generally accepted value" is usually an average over several experiments, performed by a compiler (in Fig. 4, the compiler is other than the Particle Data Group in all cases, although we do of course quote the compiled results). The x axis on all these figures is the date of publication of the value shown. Clearly there is a general progression toward better understanding—at least as measured by the size of the error bars. However, the size of the error bars do not tell the full story, as we can see by the frequency with which the "best" value has changed by more than one standard deviation. Changes in these values can come from several sources—a new experimental measurement, re-evaluation of an old measurement (which can come about if a previously unrecognized source of bias is discovered and corrected, or if a new value for one of the input constants, e.g. the electric charge, is available), or a change in the averaging procedure.

In Fig. 5 we show the history of some masses (including the Λ^0 , for radioactive Λ^0 dating of your colleagues), based on averages which we ourselves have performed. All of these were originally presented in Rosenfeld (1975). The publication date refers to the publication of the Review of Particle Properties.

In Fig. 6 we show the best estimates for the lifetimes of some of the particles stable against strong decay. These and subsequent figures have been compiled since publication of Rosenfeld (1975). In Fig. 7 we show the widths of some of the resonances, and in Fig. 8, the values of some of the branching fractions. All values are taken from the Tables. Before 1964, very few branching fractions were listed in the Tables. In all cases, a representative sample is chosen. In each figure, the heavy inner error bar represents the statistical error computed in the averaging procedure, and the thin outer error bars, when present, indicate the increase in the error due to the scale factor. The scale factor is described in the preceding section. It represents an attempt to quantify the increase in the uncertainty which is present in the case of experiments which disagree by more than a certain amount. In the case where the error represents an "educated guess," rather than a calculation, the inner error bar is absent.

On the whole, the number of times the values have changed by more than one standard deviation over the years is remarkably few. Even those branching fractions which involve rare decays and which are therefore presumably difficult to measure (Fig. 8) are, for the most part, within one or two standard deviations in 1974 of their value in any year since 1960. This is in spite of the vast amount of new experimental input, and indicates the general reliability of the results.

Of course, the data points for the different years are hardly independent of each other, but those differing by several years frequently have quite different experimental input. The relative lack of change is a comment both

on the experiments and on the averaging procedures. We, of course, are responsible only for the averages (but not on Fig. 4). These averages entail considerable exercise of judgment: there are conflicting experiments, experiments with impossibly small errors, "preliminary" results, and so forth. Statistical procedures will tell us that two experiments do not agree; they do not give a clue as to which (if either) is a good representation of the truth. Major decisions, and their motivations, are usually discussed on a case-by-case basis in the Data Card Listings; general comments may be found in Sec. II and in Rosenfeld (1975). Note that, occasionally, the error bars increase from one publication to the next, in these figures. This is usually the result of decision making by the compiler, e.g., to cease using a particular result.

We show these figures not only to demonstrate that there is not much change in these averages in the usual case, but also to show that there exist cases with relatively large changes. There is a psychological danger in preparing tables of "right" answers. The old joke about the experimenter who fights the systematics until he or she get the "right" answer (read "agrees with previous experiments"), and then publishes, contains a germ of truth (presumably, those who compile and average experimental results are also not immune to this disease). A result can disagree with the average of all previous experiments by five standard deviations, and still be right! Hence, perhaps it is of value to show that large changes can (and do) sometimes occur.

In summary, with the addition of Figs. 7–8, not available at the time of publication of Rosenfeld (1975), we find we can reiterate his conclusions. Namely, that the combination of careful work by experimenters and by compilers (which involved excluding around 40% of the data from the averages, adjusting impossibly small errors, etc.), and the frequent use of the conservative scale factor, has produced averages whose reliability is remarkably good.

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We would like to take this opportunity to thank all those who have assisted in the many phases of preparing this Review. In particular, we want to acknowledge the usefulness of the feed-back from the physics community, especially those who have taken the trouble to make suggestions or point out errors. The European members of the Particle Data Group wish to acknowledge the generous support of CERN; in particular through Dr. A. Günther and his services.

REFERENCES

- Albright, C. H., 1959, *Phys. Rev.* **115**, 750.
- Barbaro-Galtieri, A., 1968, "Baryon resonances", in: *Advances in Particle Physics*, edited by R. L. Cool and R. E. Marshak (Wiley, New York), Vol. 2. See specifically, Table IV and Figs. 10 and 12.
- Barrelet, E., 1972, *Nuovo Cim.* **8A**, 331.
- Bender, I., V. Linke, and H. J. Rothe, 1968, *Z. Physik* **221**, 190.
- Cabibbo, N., 1963, *Phys. Rev. Lett.* **10**, 531.

- Cabibbo, N., and A. Maksymowicz, 1961, *Phys. Lett.* **9**, 352.
- Chounet, L. M., J. M. Gaillard, and M. K. Gaillard, 1972, *Phys. Rep.* **4C**, 199.
- Dalitz, R. H., 1956, *Proc. Phys. Soc. (London)* **69A**, 527.
- Fearing, H. W., E. Fischbach, and J. Smith, 1970, *Phys. Rev. D* **2**, 542.
- Gasiorowicz, S., 1966, *Elementary Particle Physics* (Wiley, New York).
- Giacomelli, G., *et al.*, 1974, *Nucl. Phys.* **B71**, 138.
- Goldberger, M. L., and S. B. Treiman, 1958, *Phys. Rev.* **11**, 354.
- Herndon, D. J., A. Barbaro-Galtieri, and A. H. Rosenfeld, Feb. 1972, " πN Partial-wave Amplitudes—a Compilation, Lawrence Radiation Laboratory Report UCRL-20030 πN ," p. 4, Eq. (9).
- Jackson, J. D., S. D. Treiman, and H. W. Wyld, Jr., 1957, *Phys. Rev.* **106**, 517.
- Jackson, J. D., 1964, *Nuovo Cim.* **34**, 1644.
- Källén, G., 1964, *Elementary Particle Physics* (Addison-Wesley Publ. Co., Reading, Mass.)
- Kinoshita, T., and A. Sirlin, 1957, *Phys. Rev.* **108**, 844.
- Layter, J. G., J. A. Appel, A. Kotlewski, W. Lee, S. Stein, and J. J. Thaler, 1972, *Phys. Rev. Lett.* **29**, 316.
- Lee, T. D., and C. N. Yang, 1957, *Phys. Rev.* **108**, 1615.
- Levi-Setti, R., June 1969, Rapporteur talk at the Lund Inter. Conf. on Particle Physics (Lund).
- Martin, B. R., 1975, *Nucl. Phys.* **B94**, 413.
- Okun, L. B., and C. Rubbia, 1967, *Proc. Heidelberg Conf. on Elementary Particles*, p. 301.
- Particle Data Group (1969): N. Barash-Schmidt, A. Barbaro-Galtieri, L. R. Price, A. H. Rosenfeld, P. Söding, C. G. Wohl, M. Roos, and G. Conforto, *Rev. Mod. Phys.* **41**, 109.
- Particle Data Group (1974): V. Chaloupka, C. Bricman, A. Barbaro-Galtieri, D. M. Chew, R. L. Kelly, T. A. Lasinski, A. Rittenberg, A. H. Rosenfeld, T. G. Trippe, F. Uchiyama, N. Barash-Schmidt, P. Söding, and M. Roos, *Phys. Lett.* **50B**, No. 1.
- Particle Data Group (1975): V. Chaloupka, C. Bricman, A. Barbaro-Galtieri, D. M. Chew, R. L. Kelly, T. A. Lasinski, A. Rittenberg, A. H. Rosenfeld, T. G. Trippe, F. Uchiyama, G. P. Yost, N. Barash-Schmidt, and M. Roos, *Rev. of Mod. Phys.* **47**, 535.
- Pišút, J., and M. Roos, 1968, *Nucl. Phys.* **B6**, 325.
- Roper, L. D., R. M. Wright, and B. T. Feld, 1965, *Phys. Rev.* **138**, B190.
- Roos, M., 1974, *Nucl. Phys.* **B77**, 420; see also Table VIII in M. Kleinknecht, *Proc. 17th Intern. Conf. on High Energy Physics*, London, 1974 (Science Research Council, Chilton, 1974), p. III-23.
- Roos, M., M. Hietanen, and J. Luoma, 1975, *Physica Fennica* **10**, 21.
- Rosenfeld, A. H., 1975, *Ann. Rev. Nucl. Sci.* **25**, 555.
- Steinberger, J., 1969, CERN Topical Conference on Weak Interactions, CERN 69-7, p. 291.
- Willis, W., and J. Thompson, 1968, "Leptonic Decays of Elementary Particles," in *Advances in Particle Physics*, edited by R. L. Cool and R. E. Marshak (Wiley, New York), Vol. 1, p. 295.
- Wolfenstein, L., 1969, in: *Theory and Phenomenology in Particle Physics*, edited by A. Zichichi (Academic, New York), p. 218.
- Wu, T. T., and C. N. Yang, 1964, *Phys. Rev. Lett.* **12**, 380.

TABLES OF PARTICLE PROPERTIES

April 1976

N. Barash-Schmidt, A. Barbaro-Galtieri, C. Bricman, V. Chaloupka,
 R. J. Hemingway, R. L. Kelly, M. J. Losty, A. Rittenberg,
 M. Roos, A. H. Rosenfeld, T. G. Trippe, G. P. Yost

(Closing date for data: Jan. 1, 1976)

Stable Particle Table

For additional parameters, see Addendum to this table.

Quantities in italics have changed by more than one (old) standard deviation since April 1974.

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean Life (sec) $c\tau$ (cm)	Partial decay mode		
				Mode	Fraction ^a	p or P _{max} ^b (MeV/c)
γ	0,1(1 ⁻) ⁻	0(<7×10 ⁻²²)	stable	stable		
ν	J= $\frac{1}{2}$ ν_e : 0(<0.00006) ν_μ : 0(<0.65)		stable	stable		
e	J= $\frac{1}{2}$	0.5110034 ±.0000014	stable (>5×10 ²¹ y)	stable		
μ	J= $\frac{1}{2}$ $m^2=0.01116$ $m_\mu-m_\pi\pm=-33.909$ ±.006	105.65948 ±.00035 	<i>2.197134</i> ×10 ⁻⁶ <i>±.000077</i> <i>c\tau=6.5868</i> ×10 ⁴	$e\nu\bar{\nu}$ $e\gamma\gamma$ $3e$ $e\gamma$ $e^+\nu_e\nu_\mu$	100 (<4)×10 ⁻⁶ (<6)×10 ⁻⁹ (<2.2)×10 ⁻⁸ (<25)%	53 53 53 53
π^\pm	1 ⁻ (0 ⁻)	139.5688 ±.0064 $m^2=0.0195$	2.6030×10 ⁻⁸ ±.0023 <i>c\tau=780.4</i> <i>(\tau^+-\tau^-)/\tau=</i> <i>(0.05±0.07)%</i> <i>(test of CPT)</i>	$\mu\nu$ $e\nu$ $\mu\nu\gamma$ $\pi^0e\nu$ $e\nu\gamma$ $e\nu e^+e^-$	100 (1.267±0.023)×10 ⁻⁴ ^c (1.24±0.25)×10 ⁻⁴ (1.02±0.07)×10 ⁻⁸ ^c (3.0±0.5)×10 ⁻⁸ (<3.4)×10 ⁻⁸	30 70 30 5 70 70
π^0	1 ⁻ (0 ⁻) ⁺	134.9645 ±.0074 $m^2=0.182$ $m_\pi\pm-m_\pi^0=4.6043$ ±.0037	0.828×10 ⁻¹⁶ ±.057 S=1.8* <i>c\tau=2.5</i> ×10 ⁻⁶	$\gamma\gamma$ γe^+e^- $\gamma\gamma\gamma$ $e^+e^-e^+e^-$ ^d $\gamma\gamma\gamma$ e^+e^-	(98.85±0.05)% (1.15±0.05)% (<5)×10 ⁻⁶ (3.32)×10 ⁻⁵ (<6)×10 ⁻⁵ (<2)×10 ⁻⁶	67 67 67 67 67 67

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean life (sec) $c\tau$ (cm)	Partial decay mode		p or Pmax ^b (MeV/c)			
				Mode	Fraction ^a				
K[±]	$\frac{1}{2}(0^-)$	493.707 ±0.037 m ² =0.244 m _{K[±]} -m _{K⁰} =-3.99 ±0.13 S=1.1*	1.2371×10 ⁻⁸ ±0.0026 S=1.9* cτ=370.9 (τ ⁺ -τ ⁻)/τ ⁻ = (.11±.09)% (test of CPT) S=1.2*	μν	(63.61±0.16)%	236			
				ππ ⁰	(21.05±0.14)%	205			
				ππ ⁻ π ⁺	(5.59±0.03)% S=1.1*	125			
				ππ ⁰ π ⁰	(1.73±0.05)% S=1.4*	133			
				μπ ⁰ ν	(3.20±0.09)% S=1.7*	215			
				επ ⁰ ν	(4.82±0.05)% S=1.1*	228			
				μνγ	^c (5.8 ±3.5)×10 ⁻³	236			
				επ ⁰ μ ⁰ ν	(1.8 ±2.4)×10 ⁻⁵	207			
				ππ [±] e [±] ν	(3.7 ±0.5)×10 ⁻⁵	203			
				ππ [±] e [±] ν	(<5)×10 ⁻⁷	203			
				ππ [±] μ [±] ν	(0.9 ±0.4)×10 ⁻⁵	151			
				ππ [±] μ [±] ν	(<3.0)×10 ⁻⁶	151			
				εν	(1.54±0.09)×10 ⁻⁵	247			
				ενγ	^c (1.62±0.47)×10 ⁻⁵	247			
				ππ ⁰ γ	^{e,c} (2.71±0.19)×10 ⁻⁴	205			
				ππ ⁺ π ⁻ γ	^c (1.0 ±0.4)×10 ⁻⁴	125			
				μπ ⁰ νγ	^c <6)×10 ⁻⁵	215			
				επ ⁰ νγ	^c (3.7 ±1.4)×10 ⁻⁴	228			
				πe ⁺ e ⁻	(2.6 ±0.5)×10 ⁻⁷	227			
				π [±] e [±] e [±]	(<1.5)×10 ⁻⁵	227			
				πμ ⁺ μ ⁻	(<2.4)×10 ⁻⁶	172			
				πγγ	^c <3.5)×10 ⁻⁵	227			
				πγγγ	^c <3.0)×10 ⁻⁴	227			
				πνν	(<0.6)×10 ⁻⁶	227			
				πγ	(<4)×10 ⁻⁶	227			
				επ [±] μ [±]	(<2.8)×10 ⁻⁸	214			
				επ [±] μ [±]	(<1.4)×10 ⁻⁸	214			
				μννν	(<6)×10 ⁻⁶	236			
				K⁰	$\frac{1}{2}(0^-)$	497.70 ±0.13 S=1.1* m ² =0.248	50% K _{Short} , 50% K _{Long}		
				K_S⁰	$\frac{1}{2}(0^-)$	0.8930×10 ⁻¹⁰ (f) ±0.0023 cτ=2.68	π ⁺ π ⁻	(68.67)%	206
							π ⁰ π ⁰	(31.33±0.25)% S=1.1*	209
							μ ⁺ μ ⁻	<3.2)×10 ⁻⁷	225
e ⁺ e ⁻	<3.4)×10 ⁻⁴	249							
π ⁺ π ⁻ γ	^c (2.0 ±0.4)×10 ⁻³	206							
γγ	<0.4)×10 ⁻³	249							
K_L⁰	$\frac{1}{2}(0^-)$	5.181×10 ⁻⁸ ±0.040 cτ=1553 m _{K_L} -m _{K_S} = 0.5349×10 ¹⁰ ħ sec ⁻¹ ±0.0022	π ⁰ π ⁰ π ⁰	(21.4 ±0.7)% S=1.2*	139				
			π ⁺ π ⁻ π ⁰	(12.25±0.18)% S=1.1*	133				
			πμν	(27.1 ±0.5)%	216				
			πeν	^e (39.0 ±0.5)% S=1.1*	229				
			πeνγ	^{e,c} (1.3 ±0.8)%	229				
			π ⁺ π ⁻	^f (0.201±0.006)%	206				
			π ⁰ π ⁰	^c (0.094±0.019)% S=1.5*	209				
			π ⁺ π ⁻ γ	(6.0 ±2.0)×10 ⁻⁵	206				
			π ⁰ γγ	<2.4)×10 ⁻⁴	231				
			γγ	(4.9 ±0.5)×10 ⁻⁴	249				
			εμ	<2.0)×10 ⁻⁹	238				
			μ ⁺ μ ⁻	^h (1.0 ±0.3)×10 ⁻⁸	225				
			μ ⁺ μ ⁻ γ	<7.8)×10 ⁻⁶	225				
			μ ⁺ μ ⁻ π ⁰	<5.7)×10 ⁻⁵	177				
			e ⁺ e ⁻	<2.0)×10 ⁻⁹	249				
e ⁺ e ⁻ γ	<2.8)×10 ⁻⁵	249							
π ⁺ π ⁻ e ⁺ e ⁻	<7.2)×10 ⁻⁶	206							
π ⁰ π [±] e [±] ν	<2.2)×10 ⁻³	207							
η	0 ⁺ (0 ⁻) ⁺	548.8 ±0.6 S=1.4* m ² =0.301 eΓ=(0.85±0.12)keV ⁽¹⁾ Neutral decays (71.0±0.7)% S=1.1* Charged decays (29.0±0.7)% S=1.1*	γγ	(38.0 ±1.0)% S=1.2*	274				
			π ⁰ γγ	ⁱ (3.1 ±1.1)% S=1.2*	258				
			3π ⁰	(29.9 ±1.1)% S=1.1*	180				
			π ⁺ π ⁻ π ⁰	(23.6 ±0.6)% S=1.1*	175				
			π ⁺ π ⁻ γ	(4.89±0.13)% S=1.1*	236				
			e ⁺ e ⁻ γ	(0.50±0.12)%	274				
			π ⁰ e ⁺ e ⁻	<0.04)%	258				
			π ⁺ π ⁻	<0.15)%	236				
			π ⁺ π ⁻ e ⁺ e ⁻	(0.1 ±0.1)%	236				
			π ⁺ π ⁻ π ⁰ γ	<6)×10 ⁻⁴	175				
			π ⁺ π ⁻ γγ	<0.2)%	236				
μ ⁺ μ ⁻	(2.2 ±0.8)×10 ⁻⁵	253							
μ ⁺ μ ⁻ π ⁰	<5)×10 ⁻⁴	211							

Stable Particle Table (cont'd)

Particle	$I^G(J^P)C_n$	Mass (MeV) Mass ² (GeV) ²	Mean Life (sec) $c\tau$ (cm)	Partial decay mode		
				Mode	Fraction ^a	p or p_{max}^b (MeV/c)
p	$\frac{1}{2}(\frac{1}{2}^+)$	938.2796 ± 0.0027 $m^2 = 0.8804$	stable ($>2 \times 10^{30}y$)			
n	$\frac{1}{2}(\frac{1}{2}^+)$	939.5731 ± 0.0027 $m^2 = 0.8828$ $m_p - m_n = -1.29343$ ± 0.00004	918 \pm 14 $c\tau = 2.75 \times 10^{13}$	$pe^- \nu$	100 %	1
Λ	$0(\frac{1}{2}^+)$	1115.60 ± 0.05 $S = 1.2^*$ $m^2 = 1.245$	2.578×10^{-10} ± 0.021 $S = 1.6^*$ $c\tau = 7.73$	$p\pi^-$ $n\pi^0$ $pe^- \nu$ $p\mu^- \nu$ $p\pi^- \gamma$	(64.2 \pm 0.5)% (35.8 \pm 0.5)% (8.13 \pm 0.29) $\times 10^{-4}$ (1.57 \pm 0.35) $\times 10^{-4}$ ^c (0.85 \pm 0.14) $\times 10^{-3}$	100 104 163 131 100
Σ^+	$1(\frac{1}{2}^+)$	1189.37 ± 0.06 $S = 1.8^*$ $m^2 = 1.415$ $m_{\Sigma^+} - m_{\Sigma^-} = -7.98$ ± 0.08 $S = 1.2^*$	0.800×10^{-10} ± 0.006 $c\tau = 2.40$ $\frac{\Gamma(\Sigma^+ \rightarrow p^+ \pi^0 \nu)}{\Gamma(\Sigma^+ \rightarrow p^+ \pi^- \nu)} < .043$	$p\pi^0$ $n\pi^+$ $p\gamma$ $n\pi^+ \gamma$ $\Lambda e^+ \nu$ $\left\{ \begin{array}{l} n\mu^+ \nu \\ ne^+ \nu \end{array} \right.$ $pe^+ e^-$	(51.6 \pm 0.7)% (48.4 \pm 0.7)% ^c (1.24 \pm 0.18) $\times 10^{-3}$ $S = 1.4^*$ ^c (0.93 \pm 0.10) $\times 10^{-3}$ (2.02 \pm 0.47) $\times 10^{-5}$ (<3.0) $\times 10^{-5}$ (<0.5) $\times 10^{-5}$ (<7) $\times 10^{-6}$	189 185 225 185 71 202 224 225
Σ^0	$1(\frac{1}{2}^+)$	1192.47 ± 0.08 $m^2 = 1.422$	$< 1.0 \times 10^{-14}$ $c\tau < 3 \times 10^{-4}$	$\Lambda\gamma$ $\Lambda e^+ e^-$ $\Lambda\gamma\gamma$	100 % ^d (5.45) $\times 10^{-3}$ (<3)%	74 74 74
Σ^-	$1(\frac{1}{2}^+)$	1197.35 ± 0.06 $m^2 = 1.434$ $m_{\Sigma^0} - m_{\Sigma^-} = -4.88$ ± 0.06	1.482×10^{-10} ± 0.017 $S = 1.5^*$ $c\tau = 4.44$	$n\pi^-$ $ne^- \nu$ $n\mu^- \nu$ $\Lambda e^- \nu$ $n\pi^- \gamma$	100 % (1.08 \pm 0.04) $\times 10^{-3}$ (0.45 \pm 0.04) $\times 10^{-3}$ (0.60 \pm 0.06) $\times 10^{-4}$ ^c (4.6 \pm 0.6) $\times 10^{-4}$	193 230 210 79 193
Ξ^0	$\frac{1}{2}(\frac{1}{2}^+)(j)$	1314.9 ± 0.6 $m^2 = 1.729$ $m_{\Xi^0} - m_{\Xi^-} = -6.4$ ± 0.6	2.96×10^{-10} ± 0.12 $c\tau = 8.87$	$\Lambda\pi^0$ $\Lambda\gamma$ $\Sigma^0\gamma$ $p\pi^-$ $pe^- \nu$ $\Sigma^+ e^- \nu$ $\Sigma^- e^+ \nu$ $\Sigma^+ \mu^- \nu$ $\Sigma^- \mu^+ \nu$ $p\mu^- \nu$	100 % (0.5 \pm 0.5)% (<7)% (<3.6) $\times 10^{-5}$ (<1.3) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<0.9) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<0.9) $\times 10^{-3}$ (<1.3) $\times 10^{-3}$	135 184 117 299 323 120 112 64 49 309
Ξ^-	$\frac{1}{2}(\frac{1}{2}^+)(j)$	1321.29 ± 0.14 $m^2 = 1.746$	1.652×10^{-10} ± 0.023 $S = 1.1^*$ $c\tau = 4.95$	$\Lambda\pi^-$ $\Lambda e^- \nu$ $\Sigma^0 e^- \nu$ $\Lambda\mu^- \nu$ $\Sigma^0 \mu^- \nu$ $n\pi^-$ $ne^- \nu$ $n\mu^- \nu$ $\Sigma^- \gamma$ $p\pi^- \pi^-$ $p\pi^- e^- \nu$ $p\pi^- \mu^- \nu$ $\Xi^0 e^- \nu$	100 % ^k (0.69 \pm 0.18) $\times 10^{-3}$ (<0.5) $\times 10^{-3}$ (3.5 \pm 3.5) $\times 10^{-4}$ (<0.8) $\times 10^{-3}$ (<1.1) $\times 10^{-3}$ (<3.2) $\times 10^{-3}$ (<1.5)% (<1.2) $\times 10^{-3}$ (<4) $\times 10^{-4}$ (<4) $\times 10^{-4}$ (<4) $\times 10^{-4}$ (<2.3) $\times 10^{-4}$	139 190 123 163 70 303 327 313 118 223 304 250 6
Ω^-	$0(\frac{3}{2}^+)(j)$	1672.2 ± 0.4 $m^2 = 2.796$	$1.3^{+0.3}_{-0.2} \times 10^{-10}$ $c\tau = 4.0$	$\Xi^0 \pi^-$ $\Xi^- \pi^0$ ΛK^-	Total of 43 events seen	293 290 211

ADDENDUM TO
Stable Particle Table

e $1.001\ 159\ 6567 \frac{e\hbar}{2m_e c}$ $\pm .000\ 000\ 0035$		μ Decay parameters (4)			
μ $1.001\ 166\ 897 \frac{e\hbar}{2m_\mu c}$ $\pm .000\ 000\ 027$		$\rho = 0.752 \pm 0.003$ $\xi = 0.972 \pm 0.013$ $ g_A/g_V = 0.86^{+0.33}_{-0.11}$	$\eta = -0.12 \pm 0.21$ $\delta = 0.755 \pm 0.009$ $\phi = 180^\circ \pm 15^\circ$	$h = 1.00 \pm 0.13$	
K^\pm Mode	Partial rate (sec ⁻¹)		$\Delta I = \frac{1}{2}$ rule for $K \rightarrow 3\pi$ (m)		
	$\mu\nu$ (51.42±0.17)×10 ⁶ S=1.2* $\pi\pi^0$ (17.02±0.12)×10 ⁶ S=1.1* $\pi\pi^+\pi^-$ (4.52±0.02)×10 ⁶ S=1.1* $\pi\pi^0\pi^0$ (1.40±0.04)×10 ⁶ S=1.4* $\mu\pi^0\nu$ (2.58±0.07)×10 ⁶ S=1.7* $e\pi^0\nu$ (3.90±0.04)×10 ⁶ S=1.1*	$K^+ \rightarrow \pi^+\pi^+\pi^-$ $g = -0.214 \pm 0.005$ S=1.7* $K^- \rightarrow \pi^-\pi^-\pi^+$ $g = -0.214 \pm 0.007$ S=2.7* $K^\pm \rightarrow \pi^0\pi^0\pi^\pm$ $g = 0.550 \pm 0.020$ S=1.6* $K_L^0 \rightarrow \pi^+\pi^-\pi^0$ $g = 0.646 \pm 0.014$ S=2.5*			
K_S^0	$\pi^+\pi^-\pi^0$ n(0.7689±.0035)×10 ¹⁰ $\pi^0\pi^0$ n(0.3509±.0030)×10 ¹⁰ S=1.1*	Form factors for K_{13} decays (For ξ , λ_{13}^\pm , and λ_0^\pm see Data Card Listings, especially note in K^\pm section.) K_{e3}^+ $\lambda_{13}^+ = 0.029 \pm 0.004$ K_{e3}^0 $\lambda_{13}^0 = 0.0288 \pm 0.0028$ S=1.4*			
K_L^0	$\pi^0\pi^0\pi^0$ (4.13±0.14)×10 ⁶ S=1.2* $\pi^+\pi^-\pi^0$ (2.36±0.04)×10 ⁶ S=1.1* $\pi\mu\nu$ (5.23±0.10)×10 ⁶ $\pi e\nu$ (7.52±0.11)×10 ⁶ $\pi^+\pi^0$ f,n (3.88±0.11)×10 ⁴ $\pi^0\pi^0$ n(1.81±0.37)×10 ⁴ S=1.5*	CP violation parameters (o,n,t) $ \eta_{+-} = (2.272 \pm 0.023) \times 10^{-3}$ $ \eta_{00} = (2.32 \pm 0.09) \times 10^{-3}$ S=1.1* $\phi_{+-} = (45.0 \pm 1.2)^\circ$ $\phi_{00} = (48 \pm 13)^\circ$ $ \eta_{+-} ^2 < 0.12$ $ \eta_{00} ^2 < 0.28$ $\delta = (0.330 \pm 0.012) \times 10^{-2}$ $\Delta S = -\Delta Q$ $\text{Re } x = 0.008 \pm 0.020$ S=1.4* $\text{Im } x = -0.003 \pm 0.027$ S=1.2*			
	η Mode	Left-right asymmetry $\pi^+\pi^-\pi^0$ (0.12±.17)% $\pi^+\pi^-\gamma$ (0.88±.40)%	Sextant asymmetry (0.19±0.16)%	Quadrant asymmetry (-0.17±0.17)% $\beta = 0.047 \pm 0.062$	
Magnetic moment $(e\hbar/2m_p c)$	Decay parameters (p)				
	P 2.7928456 $\pm .0000011$	Measured α	Derived ϕ (degree)	Derived γ	Derived Δ (degree)
P 2.7928456 $\pm .0000011$					
π -1.913148 $\pm .000066$	$p e^- \nu$				-1.250 ± 0.009 $\delta = (181.1 \pm 1.3)^\circ$
Λ -0.67 $\pm .06$	$p\pi^-$ 0.647±0.013 (-6.5±3.5)° $n\pi^0$ 0.651±0.045 $p e^- \nu$		0.76 (7.6 ^{+4.0} _{-4.1})°		-0.66 ± 0.05 S=1.2*
Σ^+ 2.62 $\pm .41$	$p\pi^0$ -0.979±0.016 (36±34)° $n\pi^+$ +0.066±0.016 (167±20)° $p\gamma$ -1.03 ⁺⁵² ₋₄₂ S=1.1*		0.17 (187±6)° -0.97 (-73 ⁺¹³⁶ ₋₁₀)°		
Σ^- -1.48 $\pm .37$	$n\pi^-$ -0.069±0.008 (10±15)° $n e^- \nu$ $\Lambda e^- \nu$		0.98 (249 ⁺¹² ₋₁₁₅)°		$\pm(0.435 \pm 0.035)$ 0.24 ± 0.23 S=1.3*
Ξ^0	$\Lambda\pi^0$ -0.44±0.08 (21±12)° S=1.3*		0.84 (216 ⁺¹³ ₋₁₉)°		
Ξ^- -1.85 $\pm .75$	$\Lambda\pi^-$ -0.392±0.021 (2±6)° S=1.1*		0.92 (185±13)°		
Ω^-	ΛK^- -0.66 ^{+0.36} _{-0.30}				

Stable Particle Table (cont'd)

*S = Scale factor = $\sqrt{\chi^2/(N-1)}$, where $N \approx$ number of experiments. S should be ≈ 1 . If $S > 1$, we have enlarged the error of the mean, $\delta\bar{x}$; i.e., $\delta\bar{x} \rightarrow S\delta\bar{x}$. This convention is still inadequate, since if $S \gg 1$ the experiments are probably inconsistent, and therefore the real uncertainty is probably even greater than $S\delta\bar{x}$. See text, and ideograms in Stable Particle Data Card Listings.

- a. Quoted upper limits correspond to a 90% confidence level.
- b. In decays with more than two bodies, p_{\max} is the maximum momentum that any particle can have.
- c. See Stable Particle Data Card Listings for energy limits used in this measurement.
- d. Theoretical value; see also Stable Particle Data Card Listings.
- e. The direct emission branching fraction is $(1.56 \pm .35) \times 10^{-5}$.
- f. The $\tau(K_S^0)$ and $|\eta_{+-}|$ averages (and the related $K_L^0 \rightarrow \pi^+\pi^-$ branching fraction and rate averages) contain only post-1971 results. The pre-1971 averages were $|\eta_{+-}| = (1.95 \pm 0.03) \times 10^{-3}$ and $\tau(K_S^0) = (0.862 \pm 0.006) \times 10^{-10}$ sec. See notes on $|\eta_{+-}|$ and $\tau(K_S^0)$ discrepancies in Stable Particle Data Card Listings.
- g. The branching fraction for $K_L^0 \rightarrow \pi\nu$ includes the radiative events $K_L^0 \rightarrow \pi\nu\gamma$.
- h. This is above the contradictory result of Clark et al. ($< 0.3 \times 10^{-8}$). See note in Stable Particle Data Card Listings.
- i. See note in Stable Particle Data Card Listings.
- j. P for Ξ and J^P for Ω^- not yet measured. Values reported are SU(3) predictions.
- k. Assumes rate for $\Xi^- \rightarrow \Sigma^0 e^- \nu$ small compared with $\Xi^- \rightarrow \Lambda e^- \nu$.
- l. $|g_A/g_V|$ defined by $g_A^2 = |C_A|^2 + |C'_A|^2$, $g_V^2 = |C_V|^2 + |C'_V|^2$, and $\Sigma(\epsilon|\Gamma_i|\mu)(\bar{\nu})\Gamma_i(C_1+C'_1\gamma_5)|\nu\rangle$; ϕ defined by $\cos \phi = -\text{Re}(C_A^* C'_V + C'_A C_V)/g_A g_V$ [for more details, see text Section VI A].
- m. The definition of the slope parameter of the Dalitz plot is as follows [see also text Section VI B.1]:

$$|M|^2 = 1 + g \left(\frac{s_3 - s_0}{m_{\pi^+}^2} \right)$$
- n. The $K_S^0 \rightarrow \pi\pi$ and $K_L^0 \rightarrow \pi\pi$ rates (and branching fractions) are from independent fits and do not include results of K_L^0 - K_S^0 interference experiments. The $|\eta_{+-}|$ and $|\eta_{00}|$ values given in the addendum are these rates combined with the $|\eta_{+-}|$ and $|\eta_{00}|$ results from interference experiments.
- o. The definition for the CP violation parameters is as follows [see also text Section VI B.3]:

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L^0 \rightarrow \pi^+\pi^-)}{A(K_S^0 \rightarrow \pi^+\pi^-)} \quad \eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L^0 \rightarrow \pi^0\pi^0)}{A(K_S^0 \rightarrow \pi^0\pi^0)}$$

$$\delta = \frac{\Gamma(K_L^0 \rightarrow \pi^+\pi^-) - \Gamma(K_S^0 \rightarrow \pi^+\pi^-)}{\Gamma(K_L^0 \rightarrow \pi^+\pi^-) + \Gamma(K_S^0 \rightarrow \pi^+\pi^-)}, \quad |\eta_{+-}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^+\pi^- \pi^0)}{\Gamma(K_L^0 \rightarrow \pi^+\pi^- \pi^0)}, \quad |\eta_{00}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^0\pi^0 \pi^0)}{\Gamma(K_L^0 \rightarrow \pi^0\pi^0 \pi^0)}$$

- p. The definition of these quantities is as follows [for more details on sign convention, see text Section VI B]:

$$\alpha = \frac{2|s||p|\cos\Delta}{|s|^2 + |p|^2} \quad \left| \quad \beta = \sqrt{1 - \alpha^2} \sin\phi \quad \left| \quad g_A/g_V \text{ defined by } \langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) | B_i \rangle \right.$$

$$\beta = \frac{-2|s||p|\sin\Delta}{|s|^2 + |p|^2} \quad \left| \quad \gamma = \sqrt{1 - \alpha^2} \cos\phi \quad \left| \quad \delta \text{ defined by } g_A^f/g_V^f = |g_A/g_V| e^{i\delta} \right.$$

Meson Table

April 1976

In addition to the entries in the Meson Table, the Meson Data Card Listings contain all substantial claims for meson resonances. See Contents of Meson Data Card Listings⁽¹⁾.

Quantities in italics have changed by more than one (old) standard deviation since April 1974.

Name	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode			
					Mode	Fraction (%) [Upper limits are 1σ (%)]	p or P _{max} ^(b) (MeV/c)	
π^\pm π^0	$1^-(0^-)_+$	139.57 134.96	0.0 7.8 eV ± 1.9 eV	0.019479 0.018215	See Stable Particle Table			
η	$0^+(0^-)_+$	548.8 ± 0.6	2.63 keV ± 0.58 keV	0.301 ± 0.000	Neutral Charged	71.1 28.9	See Stable Particle Table	
$\rho(770)$	$1^+(1^-)_-$	773 ± 3 [§]	152 ± 3 [§]	0.598 ± 0.117	$\pi\pi$ $\pi\gamma$ e^+e^- $\mu^+\mu^-$	≈ 100 0.024 ± 0.007 0.0043 ± 0.0005 (d) 0.0067 ± 0.0012 (d)	360 374 386 372	
M and Γ from neutral mode.								
$\omega(783)$	$0^-(1^-)_-$	782.7 ± 0.3	10.0 ± 0.4	0.613 ± 0.008	$\pi^+\pi^-\pi^0$ $\pi^0\pi^-$ $\pi^0\gamma$ e^+e^-	89.9 ± 0.6 1.3 ± 0.3 8.8 ± 0.5 0.0076 ± 0.0017	S=1.2* S=1.5* S=1.9*	327 366 380 391
η' (958)	$0^+(0^-)_+$	957.6 ± 0.3	< 1	0.917 < .001	$\eta\pi\pi$ $\rho^0\gamma$ $\gamma\gamma$	67.6 ± 1.7 30.4 ± 1.7 2.0 ± 0.3	S=1.1*	231 167 479
For upper limits, see footnote (g)								
$\delta(970)$	$1^-(0^+)_+$	976 ^(h) ± 10 [§]	50 ^(h) ± 20 [§]	0.953 ± 0.049	$\eta\pi$	seen	315	
Possibly coupled to the $I = 1$ $K\bar{K}$ system [¶] .								
S^* (993)	$0^+(0^+)_+$	~ 993 ^(c) ± 5	40 ^(c) ± 8	0.986 ± 0.040	$K\bar{K}$ $\pi\pi$	near threshold	53 476	
See note on $\pi\pi$ S wave [¶] .								
$\Phi(1020)$	$0^-(1^-)_-$	1019.7 ± 0.3 S=1.6*	4.1 ± 0.2	1.040 ± 0.004	K^+K^- $K_L^0K_S^0$ $\pi^+\pi^-\pi^0$ (incl. $\rho\pi$) $\eta\gamma$ $\pi^0\gamma$ e^+e^- $\mu^+\mu^-$	46.6 ± 2.3 35.0 ± 2.0 16.4 ± 1.5 2.0 ± 0.4 0.14 ± 0.05 .032 ± 0.002 .025 ± 0.003	S=1.6* S=1.6* S=1.1* S=1.4*	128 111 462 362 501 510 499
For upper limits, see footnote (i)								
$A_1(1100)$	$1^-(1^+)_+$	~ 1100	~ 300	1.21 ± 0.33	$\rho\pi$	~ 100	251	
Broad enhancement in the $J^P=1^+$ $\rho\pi$ partial wave; not an established resonance [¶] .								
$\epsilon(1200)$	$0^+(0^+)_+$	1100 to 1300	~ 600		$\pi\pi$			
Existence of pole not established. See note on $\pi\pi$ S wave [¶] .								
$B(1235)$	$1^+(1^+)_-$	1228 [§] ± 10 [§]	125 [§] ± 10 [§]	1.51 ± 0.15	$\omega\pi$ [D/S amplitude ratio = .25 ± 0.06] For upper limits, see footnote (j)	only mode seen	345	
$f(1270)$	$0^+(2^+)_+$	1271 [§] ± 5 [§]	180 [§] ± 20 [§]	1.62 ± 0.23	$\pi\pi$ $2\pi^+2\pi^-$ $K\bar{K}$ $\pi^+\pi^-2\pi^0$	81 ± 1 [§] 2.8 ± 0.3 2.7 ± 0.6 seen	S=1.1* 620 557 395 560	
For upper limits, see footnote (l)								
$D(1285)$	$0^+(A)_+$	1286 [§] ± 10 [§]	30 [§] ± 20 [§]	1.65 ± 0.04	$K\bar{K}\pi$ $\eta\pi\pi$ $[\delta\pi]$ $2\pi^+2\pi^-$ (prob. $\rho^0\pi^+\pi^-$)	seen seen seen seen	305 484 245 565	
$J^P = 0^-, 1^+, 2^-,$ with 1^+ favoured								

Meson Table (cont'd)

Name	$\frac{G}{\omega/\phi} \frac{0}{\eta} \frac{1}{\rho} \frac{1}{\pi}$	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		p or P _{max} (MeV/c) ^(b)
						Mode	Fraction (%) [Upper limits are 1 σ (%)]	
A ₂ (1310)	$\frac{1}{\rho}$	$1^-(2^+)_{+}$	1310 _{±5}	102 _{±5}	1.72 ±.13	$\rho\pi$ $\eta\pi$ $\omega\pi\pi$ $K\bar{K}$ $\eta'\pi$	70.9±1.8 15.0±1.2 9.3±1.9 4.7±0.5 <1	S=1.1* S=1.2* 411 529 354 428 279
E(1420)		$0^+(A)_{+}$	1416 _{±10}	60 _{±20}	2.01 ±.08	$K\bar{K}\pi$ $[K^*\bar{K} + \bar{K}^*K]$ $\eta\pi\pi$ $[\delta\pi]$	~ 40 ~ 20 ~ 60 possibly seen]	421 130 564 352
Not a well established resonance.								
f'(1514)		$0^+(2^+)_{+}$	1516 ±3	40 ±10	2.30 ±.06	$K\bar{K}$	only mode seen	572
For upper limits, see footnote (k)								
F ₁ (1540)		$1(A)_{-}$	1540 ±5	40 ±15	2.37 ±.06	$K^*\bar{K} + \bar{K}^*K$ 3π	seen possibly seen	321 737
Not a well established resonance.								
ρ' (1600)		$1^+(1^-)_{-}$	~ 1600	200-800	2.56	4π $[0\pi^+\pi^-]$ $\pi\pi\pi$ $K\bar{K}$	seen with $\pi^+\pi^-$ in S-wave] dominant possibly seen < 8	738 573 788 629
Not a well established resonance. ^{fl}								
A ₃ (1640)		$1^-(2^-)_{+}$	~ 1640	~ 300	2.69 ±.49	$f\pi$		304
Broad enhancement in the $J^P = 2^- f\pi$ partial wave; not a well established resonance. ^{fl}								
ω (1675)		$0^-(3^-)_{-}$	1667 _{±10}	150 _{±20}	2.78 ±.25	$\rho\pi$ 3π 5π $[\omega\pi\pi]$	seen possibly seen possibly seen possibly seen]	646 806 778 615
g(1680) ^{fl}		$1^+(3^-)_{-}$	1690 _{±20}	180 _{±30}	2.86 ±.30	2π 4π (incl. $\pi\pi\rho, \rho\rho, A_2\pi, \omega\pi$) $K\bar{K}$ $K\bar{K}\pi$ (incl. $K^*\bar{K}$)	24±1 large small small	833 787 683 624
J^P, M and Γ from the 2π mode.								
h(2040)		$0^+(4^+)_{+}$	2040 ±20	193 ±50	4.16 ±.39	$\pi\pi$ $K\bar{K}$	seen seen	1010 890
See note (1) for possible heavier states.								
K ⁺ K ⁰		$1/2(0^-)_{-}$	493.71 497.70		0.244 0.248		See Stable Particle Table	
K [*] (892)		$1/2(1^-)_{-}$	892.2 ±0.5	49.4 ±1.8	0.796 ±.044	$K\pi$ $K\pi\pi$ $K\gamma$	~ 100 < 0.2 0.15±0.07	288 216 309
M and Γ from charged mode; $m^0 - m^\pm = 4.1 \pm 0.6$ MeV.								
κ (1250)		$1/2(0^+)_{+}$	1250 _{±100}	~ 450	1.56 ±.56	$K\pi$		
See note on $K\pi$ S wave ^{fl} .								
Q region		$1/2(A)_{-}$	1200 to 1400			$K\pi\pi$ $[K^*\pi]$ $[K\rho]$ $[K(\pi\pi)_{\ell=0}]$	only mode seen large] seen] possibly seen]	
$J^P = 1^+$ is dominant contribution; not a well established resonance ^{fl} .								
K [*] (1420)		$1/2(2^+)_{+}$	1421 _{±3}	108 _{±10}	2.02 ±.15	$K\pi$ $K^*\pi$ $K\rho$ $K\omega$ $K\eta$	56.1±2.6 30.9±2.1 6.6±1.7 4.5±1.7 2.0±2.0	616 415 316 305 482
See note (m).								

Meson Table (cont'd)

Name	$I^G(J^P)C_n$	Mass M (MeV)	Full Width Γ (MeV)	M^2 $\pm \Gamma M^{(a)}$ (GeV) ²	Partial decay mode		p or Pmax (MeV/c) ^(b)
					Mode	Fraction (%) [Upper limits are 1 σ (%)]	
L(1770)	$1/2(A)$	1765 _s ± 10	140 _s ± 50	3.11 $\pm .25$	$K\pi\pi$ $K\pi\pi\pi$	dominant seen	788 757
Not a well established resonance [¶] . †[K*(1420) π and other subreactions [¶]]							
See note (1) for possible heavier states.							
J/ ψ (3100)	$0^-(1^-)$	3098 ± 3	0.067 $\pm .012$	9.6 $\pm .0$	e^+e^- $\mu^+\mu^-$ hadrons	7 \pm 1 7 \pm 1 86 \pm 2	1549 1545
†[identified hadron modes ~ 12] [¶]							
†[γ X(2750) possibly seen] [¶]							
ψ (3700)	$0^-(1^-)$	3684 ± 4	0.228 $\pm .056$	13.6 $\pm .0$	e^+e^- $\mu^+\mu^-$ hadrons	0.9 \pm .2 0.9 \pm .2 98.1 \pm .3	1842 1839
†[J/ ψ $\pi^+\pi^-$ 33 \pm 3] 474							
†[J/ ψ $\pi^0\pi^0$ 17 \pm 2] 478							
†[J/ ψ η 4.2 \pm .7] 189							
†[$\gamma P_c, P_c \rightarrow J/\psi \gamma$ 3.6 \pm .7] [¶]							
†[$\gamma\chi$ (3410) seen] [¶] 264							
†[$\gamma\chi$ (3530) seen] [¶] 151							
†[other identified hadron modes ~ 0.5] [¶]							
ψ (4100)	(1^-)	~ 4100	~ 200	16.8 $\pm .8$			
Broad enhancement in the e^+e^- total cross section; probably not a single resonance. [¶]							
ψ (4400)	(1^-)	4414 ± 7	33 ± 10	19.5 $\pm .1$	e^+e^-	.0013 \pm .0003	2207
†X(2750) †P _c (3300 or 3500) †X(3410) †X(3530)							
States observed in radiative decays of J/ ψ (3100) and ψ (3700). See Meson Data Card Listings for a compilation and discussion of the experimental data.							

(1) Contents of Meson Data Card Listings

Non-strange (Y = 0)				Strange (Y = 1)	
entry	$I^G(J^P)C_n$	entry	$I^G(J^P)C_n$	entry	I (J ^P)
π	$1^-(0^-)+$	A_1 (1100)	$1^-(1^+)+$	ρ' (1600)	$1^+(1^-)-$
η	$0^+(0^-)+$	$\rightarrow M$ (1150)		A_3 (1640)	$1^-(2^-)+$
ρ (770)	$1^+(1^-)-$	$\rightarrow A_{1,s}$ (1170)		ω (1675)	$0^-(3^-)-$
ω (783)	$0^-(1^-)-$	ϵ (1200)	$0^+(0^+)+$	g (1680)	$1^+(3^-)-$
$\rightarrow M$ (940)		B (1235)	$1^+(1^+)-$	$\rightarrow X$ (1690)	-
$\rightarrow M$ (953)		$\rightarrow \rho'$ (1250)	$1^+(1^-)-$	$\rightarrow X$ (1795)	1
η' (958)	$0^+(0^-)+$	f (1270)	$0^+(2^+)+$	$\rightarrow A_4$ (1900)	1^-
δ (970)	$1^-(0^+)+$	D (1285)	$0^+(A)+$	$\rightarrow S$ (1930)	1
$\rightarrow H$ (990)		A_2 (1310)	$1^-(2^+)+$	h (2040)	$0^+(4^+)+$
S^* (993)	$0^+(0^+)+$	E (1420)	$0^+(A)+$	$\rightarrow \rho$ (2100)	
ϕ (1020)	$0^-(1^-)-$	$\rightarrow X$ (1430)	0	$\rightarrow T$ (2200)	1
$\rightarrow M$ (1033)		$\rightarrow X$ (1440)	1	$\rightarrow U$ (2360)	1
$\rightarrow B_1$ (1040)		f' (1514)	$0^+(2^+)+$	$\rightarrow N\bar{N}$ (2375)	0
$\rightarrow \eta_N$ (1080)	$0^+(N)+$	F_1 (1540)	1 (A)	$\rightarrow X$ (2500-3600)	

\rightarrow Exotics

New heavy mesons

J/ ψ (3100) ψ (3700) ψ (4100) ψ (4400) $\rightarrow X$ (2750) $\rightarrow P_c$ (3300 or 3500) $\rightarrow X$ (3410) $\rightarrow X$ (3530)

Meson Table (cont'd)

- + Indicates an entry in Meson Data Card Listings not entered in the Meson Table. We do not regard these as established resonances.
- ¶ See Meson Data Card Listings.
- * Quoted error includes scale factor $S = \sqrt{\chi^2/(N-1)}$. See footnote to Stable Particle Table.
- † Square brackets indicate a subreaction of the previous (unbracketed) decay mode(s).
- § This is only an educated guess; the error given is larger than the error of the average of the published values. (See Meson Data Card Listings for the latter.)
- (a) ΓM is approximately the half-width of the resonance when plotted against M^2 .
- (b) For decay modes into ≥ 3 particles, p_{\max} is the maximum momentum that any of the particles in the final state can have. The momenta have been calculated by using the averaged central mass values, without taking into account the widths of the resonances.
- (c) From pole position $(M - i\Gamma/2)$.
- (d) The e^+e^- branching ratio is from $e^+e^- \rightarrow \pi^+\pi^-$ experiments only. The $\omega\rho$ interference is then due to $\omega\rho$ mixing only, and is expected to be small. See note in Meson Data Card Listings. The $\mu^+\mu^-$ branching ratio is compiled from 3 experiments; each possibly with substantial $\omega\rho$ interference. The error reflects this uncertainty; see notes in Meson Data Card Listings. If $e\mu$ universality holds, $\Gamma(\rho^0 \rightarrow \mu^+\mu^-) = \Gamma(\rho^0 \rightarrow e^+e^-) \times 0.99785$.
- (e) Empirical limits on fractions for other decay modes of $\rho(770)$ are $\pi^+\eta < 0.8\%$, $\pi^+\pi^+\pi^-\pi^- < 0.15\%$, $\pi^+\pi^+\pi^-\pi^0 < 0.2\%$.
- (f) Empirical limits on fractions for other decay modes of $\omega(783)$ are $\pi^+\pi^-\gamma < 5\%$, $\pi^0\pi^0\gamma < 1\%$, $\eta + \text{neutral}(s) < 1.5\%$, $\mu^+\mu^- < 0.02\%$, $\pi^0\mu^+\mu^- < 0.2\%$, $\eta\gamma < 0.5\%$.
- (g) Empirical limits on fractions for other decay modes of $\eta'(958)$: $\pi^+\pi^- < 2\%$, $\pi^+\pi^-\pi^0 < 5\%$, $\pi^+\pi^+\pi^-\pi^- < 1\%$, $\pi^+\pi^+\pi^-\pi^0 < 1\%$, $6\pi < 1\%$, $\pi^+\pi^+e^+e^- < 0.6\%$, $\pi^0e^+e^- < 1.3\%$, $\eta e^+e^- < 1.1\%$, $\pi^0\rho^0 < 4\%$, $\gamma\omega < 5\%$.
- (h) The mass and width are from the $\eta\pi$ mode only. If the $K\bar{K}$ channel is strongly coupled, the width may be 120 MeV or more.
- (i) Empirical limits on fractions for other decay modes of $\phi(1020)$ are $\pi^+\pi^- < 0.03\%$, $\pi^+\pi^-\gamma < 0.7\%$, $\omega\gamma < 5\%$, $\rho\gamma < 2\%$, $2\pi^+2\pi^-\pi^0 < 1\%$.
- (j) Empirical limits on fractions for other decay modes of $B(1235)$: $\pi\pi < 15\%$, $K\bar{K} < 2\%$, $4\pi < 50\%$, $\phi\pi < 1.5\%$, $\eta\pi < 25\%$, $(\bar{K}K)\pi^0 < 8\%$, $K_S K_S \pi^\pm < 2\%$, $K_S K_L \pi^\pm < 6\%$.
- (k) Empirical limits on fractions for other decay modes of $f'(1514)$ are $\pi^+\pi^- < 20\%$, $\eta\pi < 50\%$, $\eta\pi\pi < 30\%$, $K\bar{K}\pi + K^*\bar{K} < 35\%$, $2\pi^+2\pi^- < 32\%$.
- (l) Empirical limits on fractions for other decay modes of $f(1270)$ are $\eta\pi\pi < 1\%$, $K^0K^-\pi^+ + \text{c.c.} < 1\%$, $\eta\eta < 2\%$.
- (m) The tabulated mass of 1421 MeV comes from the $K\pi\pi$ mode; the $K\pi\pi$ mode can be contaminated with diffractively produced Q^\pm .

Established Nonets, and octet-singlet mixing angles from Appendix IIB, Eq. (2'). Of the two isosinglets, the "mainly octet" one is written first, followed by a semicolon.

$(J^P)C_n$	Nonet members	$\theta_{\text{lin.}}$	$\theta_{\text{quadr.}}$
$(0^-)^+$	$\pi, K, \eta; \eta'$	$-24 \pm 1^\circ$	$-11 \pm 1^\circ$
$(1^-)^-$	$\rho, K^*, \phi; \omega$	$37 \pm 1^\circ$	$40 \pm 1^\circ$
$(2^+)^+$	$A_2, K^*(1420), f'; f$	$29 \pm 2^\circ$	$31 \pm 2^\circ$

Baryon Table

April 1976

The following short list gives the status of all the Baryon States in the Data Card Listings. In addition to the status, the name, the nominal mass, and the quantum numbers (where known) are shown. States with three- or four-star status are included in the main Baryon Table; the others have been omitted because the evidence for the existence of the effect and/or for its interpretation as a resonance is open to considerable question.

N(939)	P11	****	$\Delta(1232)$	P33	****	$\Lambda(1116)$	P01	****	$\Sigma(1193)$	P11	****	$\Xi(1317)$	P11	****
N(1470)	P11	****	$\Delta(1650)$	S31	****	$\Lambda(1330)$	Dead		$\Sigma(1385)$	P13	****	$\Xi(1530)$	P13	****
N(1520)	D13	****	$\Delta(1670)$	D33	****	$\Lambda(1405)$	S01	****	$\Sigma(1440)$	Dead		$\Xi(1630)$		**
N(1535)	S11	****	$\Delta(1690)$	P33	*	$\Lambda(1520)$	D03	****	$\Sigma(1480)$	*		$\Xi(1820)$		***
N(1670)	D45	****	$\Delta(1890)$	F35	**	$\Lambda(1600)$	F01	*	$\Sigma(1580)$	D13	**	$\Xi(1940)$		***
N(1688)	F15	****	$\Delta(1900)$	S31	*	$\Lambda(1670)$	S01	****	$\Sigma(1620)$	S41	**	$\Xi(2030)$		**
N(1700)	S11	****	$\Delta(1910)$	P31	****	$\Lambda(1690)$	D03	****	$\Sigma(1660)$	P11	**	$\Xi(2250)$		*
N(1700)	D13	**	$\Delta(1950)$	F37	****	$\Lambda(1800)$	P01	**	$\Sigma(1670)$	D13	****	$\Xi(2500)$		**
N(1780)	P11	**	$\Delta(1960)$	D35	**	$\Lambda(1800)$	G09	*	$\Sigma(1670)$	**				
N(1810)	P13	**	$\Delta(2160)$	**	**	$\Lambda(1815)$	F05	****	$\Sigma(1690)$	**		$\Omega(1672)$	P03	****
N(1990)	F17	**	$\Delta(2420)$	H311	***	$\Lambda(1830)$	D05	****	$\Sigma(1750)$	S11	***			
N(2000)	F15	*	$\Delta(2850)$	***	***	$\Lambda(1860)$	P03	***	$\Sigma(1765)$	D15	****			
N(2040)	D13	**	$\Delta(3230)$	***	***	$\Lambda(1870)$	S01	**	$\Sigma(1770)$	P11	*			
N(2100)	S11	*				$\Lambda(2010)$	**	**	$\Sigma(1840)$	P13	*			
N(2100)	D45	*	Z0(1780)	P01	*	$\Lambda(2020)$	F07	*	$\Sigma(1880)$	P11	**			
N(2190)	G17	***	Z0(1865)	D03	*	$\Lambda(2100)$	G07	****	$\Sigma(1915)$	F15	****			
N(2220)	H19	***	Z1(1900)	P13	*	$\Lambda(2110)$	F05	**	$\Sigma(1940)$	D13	***			
N(2650)		***	Z1(2150)	*	*	$\Lambda(2350)$	****	****	$\Sigma(2000)$	S11	*			
N(3030)		***	Z1(2500)	*	*	$\Lambda(2585)$	***	***	$\Sigma(2030)$	F17	****			
N(3245)		*							$\Sigma(2070)$	F15	*			
N(3690)		*							$\Sigma(2080)$	P13	**			
N(3755)		*							$\Sigma(2100)$	G17	*			
									$\Sigma(2250)$	****				
									$\Sigma(2455)$	***				
									$\Sigma(2620)$	***				
									$\Sigma(3000)$	**				

 **** Good, clear, and unmistakable. *** Good, but in need of clarification or not absolutely certain.
 ** Needs confirmation. * Weak.

[See notes on N's and Δ 's, on possible Z's, and on Y's and Ξ 's at the beginning of those sections in the Baryon Data Card Listings; also see notes on individual resonances in the Baryon Data Card Listings.]

Particle ^a	I (J ^P) ^a estab.	π or K Beam ^b p _{beam} (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass ^c M (MeV)	Full Width ^c Γ (MeV)	M ² $\pm\Gamma M^b$ (GeV ²)	Partial decay mode		
						Mode	Fraction %	p or d P _{max} (MeV/c)
p	$1/2(1/2^+)$		938.3		0.880	See Stable Particle Table		
n	$1/2(1/2^+)$		939.6		0.883			
N(1470) ^g	$1/2(1/2^+)$ P' ₁₁	p = 0.66 $\sigma = 27.8$	1390 to 1470	180 to 220 (200)	2.16 ± 0.29	N π N η N $\pi\pi$ [N ϵ [$\Delta\pi$ [N ρ p γ^f n γ^f	~60 ~18 ~25 ~ 7] ^e ~19] ^e < 9] ^e 0.07-0.14 < 0.05	420 d 368 d 177 d 435 435
N(1520) ^g	$1/2(3/2^-)$ D' ₁₃	p = 0.74 $\sigma = 23.5$	1510 to 1530	110 to 150 (125)	2.31 ± 0.19	N π N $\pi\pi$ [N ϵ [N ρ [$\Delta\pi$ N η p γ^f n γ^f	~55 ~45 < 5] ^e ~15] ^e ~25] ^e < 1 0.4-0.7 0.3-0.6	456 410 d d 228 d 471 471
N(1535) ^g	$1/2(1/2^-)$ S' ₁₁	p = 0.76 $\sigma = 22.5$	1500 to 1530	50 to 120 (100)	2.36 ± 0.15	N π N η N $\pi\pi$ [N ρ [N ϵ [$\Delta\pi$ p γ^f n γ^f	~30 ~65 ~ 5 ~ 3] ^e ~ 2] ^e ~ 1] ^e < 0.4 < 0.4	467 182 422 d d 243 481 481

Baryon Table (cont'd)

Particle ^a	I (J ^P) ^a estab.	$\frac{\pi \text{ or } K \text{ Beam}^b}{p_{\text{beam}}(\text{GeV}/c)}$ $\sigma = 4\pi\lambda^2$ (mb)	Mass M ^c (MeV)	Full Width Γ^c (MeV)	M ² $\pm\Gamma M^b$ (GeV ²)	Partial decay mode		
						Mode	Fraction %	p or P _{max} ^d (MeV/c)
N(1670) ^g	<u>1/2(5/2⁻)D₁₅</u>	p = 1.00 $\sigma = 15.6$	1660 to 1685	145 to 165 (155)	2.79 ± 0.26	N π N $\pi\pi$ [$\Delta\pi$ ΔK N η p γ^f n γ^f	~45 ~55 ~50] ^e <0.3 <0.5 <0.03 <0.14	560 525 360 200 368 572 572
N(1688) ^g	<u>1/2(5/2⁺)F₁₅</u>	p = 1.03 $\sigma = 14.9$	1670 to 1690	120 to 145 (140)	2.85 ± 0.24	N π N $\pi\pi$ [N η N ρ $\Delta\pi$ N η p γ^f n γ^f	~60 ~40 ~14] ^e ~14] ^e ~11] ^e <0.3 0.1-0.4 <0.03	572 538 340 d 375 388 583 583
N(1700) ^g	<u>1/2(1/2⁻)S₁₁</u>	p = 1.05 $\sigma = 14.3$	1660 to 1690	100 to 200 (150)	2.89 ± 0.26	N π N $\pi\pi$ [N η N ρ $\Delta\pi$ ΔK ΣK p γ^f n γ^f	~55 ~30 ~10] ^e ~7] ^e ~4] ^e ~4 ~2 <0.1 <0.15	580 547 355 d 385 250 109 591 591
N(1780)	<u>1/2(1/2⁺)P₁₁</u>	p = 1.20 $\sigma = 12.2$	1700 to 1800	100 to 250 (200)	3.17 ± 0.36	N π N $\pi\pi$ [N η N ρ $\Delta\pi$ ΔK ΣK N η p γ^f n γ^f	~20 >40 15-40] ^e 20-50] ^e 10-20] ^e ~7 ~10 2-20 ^h <0.15 <0.13	633 603 440 249 448 353 267 476 643 643
N(1810)	<u>1/2(3/2⁺)P₁₃</u>	p = 1.26 $\sigma = 11.5$	1700 to 1850	100 to 300 (200)	3.28 ± 0.36	N π N $\pi\pi$ [N ρ ΔK ΣK N η p γ^f n γ^f	~20 ~70 ~70] ^e ~5 ~2 <5 <0.2 <0.2	652 624 297 386 307 503 661 661
N(2190)	<u>1/2(7/2⁻)G₁₇</u>	p = 2.07 $\sigma = 6.21$	2100 to 2250	150 to 300 (250)	4.80 ± 0.55	N π ΔK ΣK	15-35 <0.2 <0.2	888 710 664
N(2220)	<u>1/2(9/2⁺)H₁₉</u>	p = 2.14 $\sigma = 5.97$	2200 to 2250	250 to 350 (300)	4.93 ± 0.67	N π	~20	905
N(2650)	<u>1/2(?)</u>	p = 3.26 $\sigma = 3.67$	~2650	~350 (350)	7.02 ± 0.93	N π	(J+1/2) _x <0.4] ^j	1154
N(3030)	<u>1/2(?)</u>	p = 4.41 $\sigma = 2.62$	~3030	~400 (400)	9.18 ± 1.21	N π	(J+1/2) _x <0.1] ^j	1366
$\Delta(1232)^g$	<u>3/2(3/2⁺)P₃₃</u>	p = 0.30 $\sigma = 94.3$	1230 to 1234	110 to 120 (115)	1.52 ± 0.14	N π N $\pi^+\pi^-$ p γ^f	~99.4 ~0 0.58-0.66	227 80 259
		$\Delta(++)$ Pole position: ^k M-i $\Gamma/2$ = (1211.0 \pm 0.8) -i(49.9 \pm 0.6)						
		$\Delta(0)$ Pole position: ^k M-i $\Gamma/2$ = (1210.9 \pm 1.0) -i(53.1 \pm 1.0)						
$\Delta(1650)^g$	<u>3/2(1/2⁻)S₃₁</u>	p=0.96 $\sigma=16.4$	1615 to 1695	140 to 200 (140)	2.72 ± 0.23	N π N $\pi\pi$ [N ρ $\Delta\pi$ p γ^f	~35 ~65 10-25] ^e ~50] ^e <0.25	547 511 d 344 558

Baryon Table (cont'd)

Particle ^a	I	(J ^P) ^a estab.	π or K Beam ^b $\frac{p_{beam}}{\sigma} (\text{GeV}/c)$ $\sigma = 4\pi \lambda^2$ (mb)	Mass M ^c (MeV)	Full Width Γ^c (MeV)	M ² $\pm \Gamma M^b$ (GeV ²)	Partial decay mode		
							Mode	Fraction %	p or P _{max} ^d (MeV/c)
$\Delta(1670)^g$	$3/2(3/2^-)D_{33}$		p = 1.00 $\sigma = 15.6$	1650 to 1720	190 to 260 (200)	2.79 ± 0.33	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ^f p γ ^f	~ 15 ~ 85 30-60] ^e ~ 45] ^e 0.05-0.3	560 525 d 361 572
$\Delta(1890)^g$	$3/2(5/2^+)F_{35}$		p = 1.42 $\sigma = 9.88$	1860 to 1900	150 to 300 (250)	3.57 ± 0.47	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ^f ΣK p γ ^f	~ 15 ~ 80 ~ 60] ^e 10-30] ^e <3 <0.1	704 677 403 531 400 712
$\Delta(1910)^g$	$3/2(1/2^+)P_{31}$		p = 1.46 $\sigma = 9.54$	1780 to 1950	160 to 230 (200)	3.65 ± 0.38	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ^f ΣK p γ ^f	15-35 ? small] ^e small] ^e 2-20 <0.1	716 691 429 545 420 725
$\Delta(1950)^g$	$3/2(7/2^+)F_{37}$		p = 1.54 $\sigma = 8.90$	1910 to 1940	200 to 240 (220)	3.80 ± 0.43	N π N $\pi\pi$ [N ρ [$\Delta\pi$ ^f ΣK p γ ^f	~ 40 >25 ~ 10] ^e ~ 20] ^e <1 0.09-0.15	741 716 471 574 460 749
$\Delta(2420)^g$	$3/2(11/2^+)H_{311}$		p = 2.64 $\sigma = 4.68$	2380 to 2450	300 to 500 (300)	5.86 ± 0.73	N π	10-15	1023
$\Delta(2850)$	$3/2(?^+)$		p = 3.85 $\sigma = 3.05$	2800 to 2900	~ 400 (400)	8.12 ± 1.14	N π	(J+1/2) _x ~ 0.25]	1266
$\Delta(3230)$	$3/2(?)$		p = 5.08 $\sigma = 2.25$	3200 to 3350	~ 440 (440)	10.43 ± 1.42	N π	(J+1/2) _x ~ 0.05]	1475
Z* Evidence for states with strangeness +1 is controversial. See the Baryon Data Card listings for discussion and display of data.									
Λ	$0(1/2^+)$			1115.6		1.245	See Stable Particle Table		
$\Lambda(1405)$	$0(1/2^-)S'_{01}$		below K ⁻ p threshold	1405 $\pm 5^l$	40 ± 10^l (40)	1.97 ± 0.06	$\Sigma\pi$	100	142
$\Lambda(1520)$	$0(3/2^-)D'_{03}$		p = 0.389 $\sigma = 84.5$	1519 $\pm 2^l$	15 $\pm 2^l$ (15)	2.31 ± 0.02	N \bar{K} $\Sigma\pi$ $\Delta\pi\pi$ $\Sigma\pi\pi$	46 ± 1 42 ± 1 10 ± 1 0.9 ± 0.1	234 258 250 140
$\Lambda(1670)$	$0(1/2^-)S''_{01}$		p = 0.74 $\sigma = 28.5$	1660 to 1680	20 to 60 (40)	2.79 ± 0.07	N \bar{K} $\Delta\eta$ $\Sigma\pi$	15-35 15-35 20-60	410 64 393
$\Lambda(1690)$	$0(3/2^-)D''_{03}$		p = 0.78 $\sigma = 26.1$	1690 $\pm 10^l$	30 to 80 (60)	2.86 ± 0.10	N \bar{K} $\Sigma\pi$ $\Delta\pi\pi$ $\Sigma\pi\pi$	20-30 15-40 ~ 25 ~ 20	429 409 415 352
$\Lambda(1815)$	$0(5/2^+)F'_{05}$		p = 1.05 $\sigma = 16.7$	1820 $\pm 5^l$	70 to 100 (85)	3.29 ± 0.15	N \bar{K} $\Sigma\pi$ $\Sigma(1385)\pi$	~ 60 ~ 12 15-20	542 508 362
$\Lambda(1830)$	$0(5/2^-)D_{05}$		p = 1.09 $\sigma = 15.8$	1810 to 1840	60 to 110 (95)	3.35 ± 0.17	N \bar{K} $\Sigma\pi$ $\Delta\eta$	<10 35-75 <4	554 519 367
$\Lambda(1860)$	$0(1/2^+)P_{03}$		p = 1.14 $\sigma = 14.7$	1860 to 1910	40 to 110 (80)	3.46 ± 0.15	N \bar{K} $\Sigma\pi$	15-35 5-10	576 534
$\Lambda(2100)$	$0(7/2^-)G_{07}$		p = 1.68 $\sigma = 8.68$	2100 to 2120	150 to 300 (250)	4.41 ± 0.53	N \bar{K} $\Sigma\pi$ $\Delta\eta$ ΞK $\Lambda\omega$	~ 30 ~ 5 <3 <3 <8	748 699 617 483 443

Baryon Table (cont'd)

Particle ^a	I	(J ^P) ^a estab.	π or K Beam ^b		Mass M ^c (MeV)	Full Width Γ ^c (MeV)	M ² ±ΓM ^b (GeV ²)	Partial decay mode		
			p _{beam} (GeV/c)	σ = 4πλ ² (mb)				Mode	Fraction %	p or d p _{max} (MeV/c)
Λ(2350)	0	(?)	p = 2.29 σ = 5.85		2340 to 2360	100 to 200 (120)	5.52 ±0.28	N \bar{K} Σπ	(J+1/2) _x ~0.9 _j seen	913 865
Λ(2585)	0	(?)	p = 2.91 σ = 4.37		~2585	~300 (300)	6.68 ±0.78	N \bar{K}	(J+1/2) _x ~1.0 _j	1058
Σ		1(1/2 ⁺)			(+)1189.4 (0)1192.5 (-)1197.4		1.415 1.422 1.434	See Stable Particle Table		
Σ(1385)		1(3/2 ⁺)P ₁₃		below K ⁻ p threshold	(+)1382.5±0.5 S=1.2 ^m (-)1386.6±1.2 S=2.3 ^m	(+)35±2 S=1.9 ^m (-)42±4 S=3.2 ^m (35)	1.92 ±0.05	Λπ Σπ	88±2 12±2	208 117
Σ(1670) ⁿ		1(3/2 ⁻)D ₁₃	p = 0.74 σ = 28.5		1670 ±10 ^l	35 to 70 (50)	2.79 ±0.08	N \bar{K} Σπ Λπ	10-25 20-60 <20	410 387 447
Σ(1750)		1(1/2 ⁻)S ₁₁	p = 0.91 σ = 20.7		1700 to 1790	50 to 120 (75)	3.06 ±0.13	N \bar{K} Λπ Σπ Ση	10-40 5-20 <18 15-55	483 507 450 54
Σ(1765)		1(5/2 ⁻)D ₁₅	p = 0.94 σ = 19.6		1723 ±7 ^l	110 to 150 (130)	3.12 ±0.23	N \bar{K} Λπ Λ(1520)π Σ(1385)π Σπ	~41 ~14 ~16 ~10 ~1	496 518 187 315 461
Σ(1915) ^g		1(5/2 ⁺)F ₁₅	p = 1.25 σ = 13.0		1905 to 1930	70 to 140 (100)	3.67 ±0.19	N \bar{K} Λπ Σπ	5-15 20 ?	612 619 568
Σ(1940) ⁱ		1(3/2 ⁻)D ₁₃	p = 1.32 σ = 12.0		1900 to 1960	110 to 280 (220)	3.76 ±0.43	N \bar{K} Λπ Σπ	<20 ~4 ~7	678 680 589
Σ(2030) ^g		1(7/2 ⁺)F ₁₇	p = 1.52 σ = 9.93		2020 to 2040	120 to 200 (180)	4.12 ±0.37	N \bar{K} Λπ Σπ ΞK	~20 ~20 5-10 <2	700 700 652 412
Σ(2250)		1(?)	p = 2.04 σ = 6.76		2200 to 2300	50 to 200 (150)	5.06 ±0.34	N \bar{K} Λπ Σπ	(J+1/2) _x ~0.3 _j seen seen	849 841 801
Σ(2455)		1(?)	p = 2.57 σ = 5.09		~2455	~120 (120)	6.03 ±0.29	N \bar{K}	(J+1/2) _x ~0.2 _j	979
Σ(2620)		1(?)	p = 2.95 σ = 4.30		~2600	~200 (200)	6.86 ±0.52	N \bar{K}	(J+1/2) _x ~0.3 _j	1064
Ξ		1/2(1/2 ⁺)			(0)1314.9 (-)1321.3		1.729 1.746	See Stable Particle Table		
Ξ(1530) ^o		1/2(3/2 ⁺)P ₁₃			(0)1531.8±0.3 S=1.3 ^m (-)1535.1±0.6	(0) 9.1±0.5 (-) 10.1±1.9 (10)	2.34 ±0.02	Ξπ	100	144
Ξ(1820) ^{o,P}		1/2(?)			1800 to 1850	12 to 100 (60)	3.31 ±0.11	Λ \bar{K} Σ \bar{K} Ξπ Ξ(1530)π	seen seen seen seen	396 306 413 234
Ξ(1940) ^{o,q}		1/2(?)			1900 to 1970	30 to 140 (90)	3.76 ±0.17	Ξπ Ξ(1530)π	seen seen	499 336
Ω ⁻		0(3/2 ⁺)			1672.2		2.796	See Stable Particle Table		

Baryon Table (*cont'd*)

- For convenience all Baryon States for which information exists in the Baryon Data Card Listings are listed at the beginning of the Baryon Table. States with only a one or two star (*) rating in that list have been omitted from the main Baryon Table; each omitted state is indicated by an arrow in the left-hand margin of the Table. In the Listings there is an arrow under the name of each state omitted from the Table.
- a. The names of the Baryon States in Col. 1 [such as N(1470)] contain a nominal mass which is a rounded average of the reported values in the Data Card Listings. The convention for using primes in the spectroscopic notation for the quantum numbers in Col. 2 [such as P_{11}^1] is as follows: no prime is attached when the Data Card Listings include only one resonance in the given partial wave; when there is more than one resonance the first has been designated with a prime, the second with a double prime, etc. The name and the quantum numbers for each state are also given in large print at the beginning of the Data Card Listings for that state.
- b. The numbers in Col. 3 and Col. 6 are calculated using the nominal mass (see a. above) for M and the nominal width (see c. below) for Γ .
- c. For M and Γ of most baryons we report here an interval instead of an average. Averages are appropriate if each result is based on independent measurements, but inappropriate where the spread in parameters arises because different models or procedures have been applied to a common set of data. A single value with an approximation sign (~) indicates that there is not enough data to give a meaningful interval. A nominal width is included in parentheses in Col. 5; this nominal width is used to calculate the value of ΓM given in Col. 6.
- d. For two body decay modes we give the momentum, p, of the decay products in the decaying baryon rest frame. For decay modes into ≥ 3 particles we give the maximum momentum, p_{\max} , that any of the particles in the final state can have in this frame. The momenta are calculated using the nominal mass (see a. above) of the decaying baryon, and of any isobars in the final state. Some decays which would be energetically forbidden for the nominal masses actually occur because of the finite widths of the decaying Baryon and/or isobars in the final state. In these cases, the decay momentum is omitted from Col. 9 and replaced with a reference to this footnote.
- e. Square brackets around an isobar decay mode indicate that it is a sub-reaction of the previous unbracketed decay mode. In the case of N^* and Δ decays into isobar modes we have used the isobar model results of LONGACRE 75 in addition to other data from the listings (where available) to estimate the branching fractions.
- f. The tabulated radiative fractions involve a sum over two helicities (1/2, 3/2). In the case of $I = 1/2$ resonances, there are two distinct isospin couplings, whence γ_p and γ_n . For conventions and further details, see the Mini-Review preceding the Baryon Data Card Listings.
- g. Only information coming from partial-wave analyses has been used here. For the production experiments results see the Baryon Data Card Listings.
- h. The range given here does not include the branching ratio of approximately 80% reported by FELTESSE 75.
- i. There may be more than one state in this region. The only analysis which reports an elastic coupling (LEA 73) also finds unusually low mass and width values. The inelastic branching fractions quoted here are based on an elasticity of 10%, which is a compromise between LEA 73 and RLIC 76.
- j. This state has been seen only in an energy-dependent fit to total, channel, or fixed angle cross-section data. J is not known; x is Γ_{e1}/Γ .
- k. See note on determination of resonance parameters in the Baryon Data Card Listings. Values of mass and width are dependent upon resonance shape used to fit the data. The pole position is much less dependent upon the parametrization used. The pole positions given here are taken from results (in the Data Card Listings) of fits to the phase shifts of CARTER 73 without Coulomb corrections.
- l. The error given here is only an educated guess; it is larger than the error of the average of the published values (see the Baryon Data Card Listings for the latter).
- m. Quoted error includes an S (scale) factor. See first footnote to Stable Particle Table.
- n. In this energy region the situation is still confused. In addition to the effect at ~ 1670 MeV seen in both production and formation experiments, recent formation experiments have found evidence for fairly narrow S_{11} and/or P_{11} states at 1620-1660 MeV. A narrow bump in the $I = 1$ $\bar{K}N$ total cross section has also been seen recently at ~ 1590 MeV. It is not clear how many states really exist here. No one has reported a strong coupling of any of these states to $\bar{K}N$ but there is much disagreement about branching ratios into $\pi\Lambda$ and $\pi\Sigma$. See the mini-reviews preceding the $\Sigma(1620)$ and $\Sigma(1670)$ Data Card Listings for more information.
- o. Only $\Xi(1530)$ is firmly established; information on the other states comes from experiments that have poor statistics due to the fact that the cross sections for $S = -2$ states are very low. For Ξ states, because of the meager statistics, we lower our standards and tabulate resonant effects if they have at least a four-standard-deviation statistical significance and if they are seen by more than one group. See the Baryon Data Card Listings for the other states.
- p. All four decay modes shown have been seen. Branching ratios are not quoted because there may be more than one state here.
- q. This bump has been seen in both final states shown; it is not clear if one, or more, states are present.

PHYSICAL AND NUMERICAL CONSTANTS*

PHYSICAL CONSTANTS

		Uncert. (ppm)
N	= 6.0220943(63)×10 ²³ mole ⁻¹	1.05
V _m	= 22413.83(70) cm ³ mole ⁻¹ = molar volume of ideal gas at STP	31
c	= 2.99792458(1.2)×10 ¹⁰ cm sec ⁻¹	0.004
e	= 4.803242(14)×10 ⁻¹⁰ esu = 1.6021892(46)×10 ⁻¹⁹ coulomb	2.9; 2.9
1 MeV	= 1.6021892(46)×10 ⁻⁶ erg	2.9
ħ=h/2π	= 6.582173(17)×10 ⁻²² MeV sec = 1.0545887(57)×10 ⁻²⁷ erg sec	2.6; 5.4
ħc	= 1.9732858(51)×10 ⁻¹¹ MeV cm = 197.32858(51) MeV Fermi	2.6; 2.6
	= 0.6240078(16) GeV mb ^{1/2}	2.6
α	= e ² /ħc = 1/137.035982(30)	0.22
k _{Boltzmann}	= 1.380662(44)×10 ⁻¹⁶ erg °K ⁻¹	32
	= 8.61735(28)×10 ⁻¹¹ MeV °K ⁻¹ = 1 eV/11604.50(36) °K	32; 32
m _e	= 0.5110034(14) MeV = 9.109534(47)×10 ⁻³¹ kg	2.8; 5.1
m _p	= 938.2796(27) MeV = 1836.15152(70) m _e = 6.72270(31) m _{π±}	2.8; 0.38; 46
	= 1.007276470(11) amu	0.011
1 amu	= 1/12 m _{C12} = 931.5016(26) MeV	2.8
m _d	= 1875.628(5) MeV	3
r _e	= e ² /m _e c ² = 2.8179380(70) fermi (1 fermi = 10 ⁻¹³ cm)	2.5
λ _e	= ħ/m _e c = r _e α ⁻¹ = 3.8615905(64)×10 ⁻¹¹ cm	1.6
a _{∞Bohr}	= ħ ² /m _e e ² = r _e α ⁻² = 0.52917706(44)Å (1Å = 10 ⁻⁸ cm)	0.82
σ _{Thomson}	= (8/3)πr _e ² = 0.6652448(33)×10 ⁻²⁴ cm ² (10 ⁻²⁴ cm ² = 1 barn)	4.9
μ _{Bohr}	= eħ/2m _e c = 0.57883785(95)×10 ⁻¹⁴ MeV gauss ⁻¹	1.6
μ _p	= eħ/2m _p c = 3.1524515(53)×10 ⁻¹⁸ MeV gauss ⁻¹	1.7
μ _p /μ _{Bohr}	= 1.520993136(21)	0.014
1/2ω _{cyclotron}	= e/2m _e c = 8.794023(25)×10 ⁶ rad sec ⁻¹ gauss ⁻¹	2.8
1/2ω _p	= e/2m _p c = 4.789378(13)×10 ³ rad sec ⁻¹ gauss ⁻¹	2.8
Hydrogen-like atom (nonrelativistic, μ = reduced mass):		
	$\frac{v}{c}_{rms} = \frac{ze^2}{n\hbar c}$, $E_n = \frac{\mu v^2}{2} = \frac{\mu z^2 e^4}{2(n\hbar)^2}$, $a_n = \frac{n^2 \hbar^2}{\mu z e^2}$	
R _∞ = m _e e ⁴ /2ħ ²	= m _e c ² α ² /2 = 13.605804(36) eV (Rydberg)	2.6
	= m _e cα ² /2ħ = 109737.3143(10) cm ⁻¹	0.009
pc = 0.3 H _p (MeV, kilogauss, cm)		
1 year (sidereal)	= 365.256 days = 3.1558×10 ⁷ sec (≈π×10 ⁷ sec)	
density of dry air	= 1.205 mg cm ⁻³ (at 20°C, 760 mm)	
acceleration by gravity	= 980.62 cm sec ⁻² (sea level, 45°)	
gravitational constant	= 6.6732(31)×10 ⁻⁸ cm ³ g ⁻¹ sec ⁻²	
1 calorie (thermochemical)	= 4.184 joules	
1 atmosphere	= 1033.2275 dynes cm ⁻² = 1.01325 bar	
1 eV per particle	= 11604.50(36) °K (from E = kT)	

NUMERICAL CONSTANTS

π	= 3.1415927	1 rad	= 57.2957795 deg	√π	= 1.7724539
e	= 2.7182818	1/e	= 0.3678794	√2	= 1.4142136
ln2	= 0.6931472	ln10	= 2.3025851	√3	= 1.7320508
log ₁₀ 2	= 0.3010300	log ₁₀ e	= 0.4342945	√10	= 3.1622777

*Prepared by Stanley J. Brodsky, based mainly on the adjustment of the fundamental physical constants by E. R. Cohen and B. N. Taylor, J. Phys. Chem. Ref. Data 2, 663 (1973), plus current values for N, α, μ_p/μ_{Bohr}, R_∞ [see B. N. Taylor and E. R. Cohen, Proceedings of the Fifth International Conference on Atom Masses and Fundamental Constants (AMCO-5), Paris, 1975]. The figures in parentheses correspond to the one-standard-deviation uncertainty in the last digits of the main number. (Updated April 1976.)

CLEBSCH-GORDAN COEFFICIENTS, SPHERICAL HARMONICS, AND d FUNCTIONS

Note: A $\sqrt{\quad}$ is to be understood over every coefficient; e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation:

J	J	...
M	M	...
m_1	m_2	
m_1	m_2	Coefficients
...	...	

$1/2 \times 1/2$	1	0
$+1/2 + 1/2$	1	0
$+1/2 - 1/2$	$1/2$	$1/2$
$-1/2 + 1/2$	$1/2$	$-1/2$
$-1/2 - 1/2$	1	0

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

$$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$$

$$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$$

$2 \times 1/2$	$5/2$	$3/2$
$+2$	$1/2$	1
$+2 - 1/2$	$1/5$	$4/5$
$+1 + 1/2$	$4/5 - 1/5$	$5/2$
$0 + 1/2$	$2/5$	$3/5$
$0 + 1/2$	$3/5 - 2/5$	$5/2$
$0 - 1/2$	$3/5$	$2/5$
$-1 + 1/2$	$2/5 - 3/5$	$5/2$
$-1 - 1/2$	$3/5$	$3/2$

$1 \times 1/2$	$3/2$	$1/2$
$+1 + 1/2$	1	$1/2 + 1/2$
$+1 - 1/2$	$1/3$	$2/3$
$0 + 1/2$	$2/3 - 1/3$	$3/2$
$0 + 1/2$	$1/3$	$1/3$
$0 - 1/2$	$2/3$	$1/3$
$-1 + 1/2$	$1/3 - 2/3$	$3/2$
$-1 - 1/2$	1	$3/2$

2×1	3	2
$+2 + 1$	1	2
$+2 + 0$	$1/3$	$2/3$
$+1 + 1$	$2/3 - 1/3$	3
$+1 + 1$	1	2
$+1 + 1$	1	1
$+1 + 1$	1	1
$+1 + 1$	1	1

$3/2 \times 1$	$5/2$	$3/2$
$+3/2 + 1$	1	$3/2 + 3/2$
$+3/2 + 0$	$2/5$	$3/5$
$+1/2 + 1$	$3/5 - 2/5$	$5/2$
$+1/2 + 1$	$3/5$	$2/5$
$+1/2 + 1$	$3/5$	$1/2$
$+1/2 + 1$	$3/5$	$1/2$
$+1/2 + 1$	$3/5$	$1/2$
$+1/2 + 1$	$3/5$	$1/2$

$3/2 \times 1/2$	2	1
$+3/2 + 1/2$	1	1
$+3/2 - 1/2$	$1/4$	$3/4$
$+1/2 + 1/2$	$3/4 - 1/4$	2
$+1/2 + 1/2$	$3/4$	$1/4$
$+1/2 + 1/2$	$3/4$	$1/4$
$+1/2 + 1/2$	$3/4$	$1/4$
$+1/2 + 1/2$	$3/4$	$1/4$
$+1/2 + 1/2$	$3/4$	$1/4$

1×1	2	1
$+1 + 1$	1	1
$+1 + 0$	$1/2$	$1/2$
$0 + 1$	$1/2 - 1/2$	2
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$

1×1	2	1
$+1 + 1$	1	1
$+1 + 0$	$1/2$	$1/2$
$0 + 1$	$1/2 - 1/2$	2
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$
$0 + 1$	$1/2$	$1/2$

$$Y_\ell^{-m} = (-1)^m Y_\ell^m$$

$$d_{m,0}^\ell = \sqrt{\frac{4\pi}{2\ell+1}} Y_\ell^m e^{-im\phi}$$

$$\langle j_1 j_2 m_1 m_2 | j_1 j_2 J M \rangle = (-1)^{J-j_1-j_2} \langle j_2 j_1 m_2 m_1 | j_2 j_1 J M \rangle$$

$$d_{m',m}^j = (-1)^{m-m'} d_{m,m'}^j = d_{-m,-m'}^j$$

$3/2 \times 3/2$	3	2
$+3/2 + 3/2$	1	2
$+3/2 + 1/2$	$1/2$	$1/2$
$+1/2 + 3/2$	$1/2 - 1/2$	3
$+1/2 + 3/2$	$1/2$	2
$+1/2 + 3/2$	$1/2$	1
$+1/2 + 3/2$	$1/2$	1
$+1/2 + 3/2$	$1/2$	1

$$d_{1/2,1/2}^{1/2} = \cos \frac{\theta}{2}$$

$$d_{1/2,-1/2}^{1/2} = -\sin \frac{\theta}{2}$$

$2 \times 3/2$	$7/2$	$5/2$
$+2 + 3/2$	1	$5/2 + 5/2$
$+2 + 1/2$	$3/7$	$4/7$
$+1 + 3/2$	$4/7 - 3/7$	$7/2$
$+1 + 3/2$	$4/7$	$5/2$
$+1 + 3/2$	$4/7$	$5/2$
$+1 + 3/2$	$4/7$	$5/2$
$+1 + 3/2$	$4/7$	$5/2$

$$d_{1,1}^1 = \frac{1+\cos\theta}{2}$$

$$d_{1,0}^1 = -\frac{\sin\theta}{\sqrt{2}}$$

$$d_{1,-1}^1 = \frac{1-\cos\theta}{2}$$

$$d_{0,0}^1 = \cos\theta$$

2×2	4	3
$+2 + 2$	1	3
$+2 + 1$	$1/2$	$1/2$
$+1 + 2$	$1/2 - 1/2$	4
$+1 + 2$	$1/2$	3
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2

2×2	4	3
$+2 + 2$	1	3
$+2 + 1$	$1/2$	$1/2$
$+1 + 2$	$1/2 - 1/2$	4
$+1 + 2$	$1/2$	3
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2

2×2	4	3
$+2 + 2$	1	3
$+2 + 1$	$1/2$	$1/2$
$+1 + 2$	$1/2 - 1/2$	4
$+1 + 2$	$1/2$	3
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2

2×2	4	3
$+2 + 2$	1	3
$+2 + 1$	$1/2$	$1/2$
$+1 + 2$	$1/2 - 1/2$	4
$+1 + 2$	$1/2$	3
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2

2×2	4	3
$+2 + 2$	1	3
$+2 + 1$	$1/2$	$1/2$
$+1 + 2$	$1/2 - 1/2$	4
$+1 + 2$	$1/2$	3
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2
$+1 + 2$	$1/2$	2

$$d_{3/2,3/2}^{3/2} = \frac{1+\cos\theta}{2} \cos \frac{\theta}{2}$$

$$d_{3/2,1/2}^{3/2} = -\sqrt{3} \frac{1+\cos\theta}{2} \sin \frac{\theta}{2}$$

$$d_{3/2,-1/2}^{3/2} = \sqrt{3} \frac{1-\cos\theta}{2} \cos \frac{\theta}{2}$$

$$d_{3/2,-3/2}^{3/2} = -\frac{1-\cos\theta}{2} \sin \frac{\theta}{2}$$

$$d_{1/2,1/2}^{3/2} = \frac{3\cos\theta-1}{2} \cos \frac{\theta}{2}$$

$$d_{1/2,-1/2}^{3/2} = -\frac{3\cos\theta+1}{2} \sin \frac{\theta}{2}$$

$$d_{2,2}^2 = \left(\frac{1+\cos\theta}{2} \right)^2$$

$$d_{2,1}^2 = -\frac{1+\cos\theta}{2} \sin \theta$$

$$d_{2,0}^2 = \frac{\sqrt{6}}{4} \sin^2 \theta$$

$$d_{2,-1}^2 = -\frac{1-\cos\theta}{2} \sin \theta$$

$$d_{2,-2}^2 = \left(\frac{1-\cos\theta}{2} \right)^2$$

$$d_{1,1}^2 = \frac{1+\cos\theta}{2} (2\cos\theta-1)$$

$$d_{1,0}^2 = -\sqrt{\frac{3}{2}} \sin \theta \cos \theta$$

$$d_{1,-1}^2 = \frac{1-\cos\theta}{2} (2\cos\theta+1)$$

$$d_{1,-2}^2 = \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

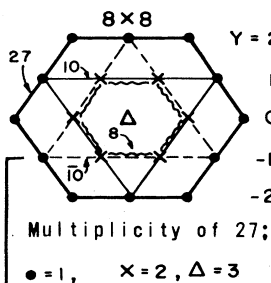
Sign convention is that of Wigner (Group Theory, Academic Press, New York, 1959), also used by Condon and Shortley (The Theory of Atomic Spectra, Cambridge Univ. Press, New York, 1953), Rose (Elementary Theory of Angular Momentum, Wiley, New York, 1957), and Cohen (Tables of the Clebsch-Gordan Coefficients, North American Rockwell Science Center, Thousand Oaks, Calif., 1974). The signs and numbers in the current tables have been calculated by computer programs written independently by Cohen and at LBL. (Table extended April 1974.)

SU(3) ISOSCALAR FACTORS

Adapted from J. J. de Swart, Rev. Mod. Phys. 35, 916 (1963)

The convention used here is: baryon first, meson second.

$$\{8\} \otimes \{8\} = \{27\} \oplus \{10\} \oplus \{10^*\} \oplus \{8\}_1 \oplus \{8\}_2 \oplus \{1\}.$$



* Five single-coefficient tables are omitted. The one involving a $\{10^*\}$ has a negative coefficient, i.e. $(NK|10^*) = -1$. The others, involving $\{27\}$ and $\{10\}$, are all +1.

		$Y=1, I=1/2, N$				$Y=1, I=3/2, \Delta$	
$\xi_1 \downarrow$		27	8 _D	8 _F	10*	27	10
N π		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	-1/2	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	-1/2	$\sqrt{2}/2$	$\sqrt{2}/2$
N η		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	1/2		
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	-1/2		

		$Y=0, I=0, \Lambda$				$Y=0, I=1, \Sigma$				
$\xi_1 \downarrow$		27	8 _D	1	8 _F	27	8 _D	8 _F	10	10*
N K		$\sqrt{15}/10$	$\sqrt{10}/10$	1/2	$\sqrt{2}/2$	$\sqrt{5}/5$	$-\sqrt{30}/10$	$\sqrt{6}/6$	$-\sqrt{6}/6$	$-\sqrt{6}/6$
ΣK		$-\sqrt{15}/10$	$-\sqrt{10}/10$	-1/2	$\sqrt{2}/2$	$\sqrt{5}/5$	$-\sqrt{30}/10$	$-\sqrt{6}/6$	$\sqrt{6}/6$	$-\sqrt{6}/6$
$\Sigma \pi$		$-\sqrt{10}/20$	$-\sqrt{15}/5$	$\sqrt{6}/4$	0	0	0	$\sqrt{6}/3$	$\sqrt{6}/6$	$-\sqrt{6}/6$
$\Delta \eta$		$3\sqrt{30}/20$	$-\sqrt{5}/5$	$-\sqrt{2}/4$	0	$\sqrt{30}/10$	$\sqrt{5}/5$	0	1/2	1/2
$\Lambda \pi$						$\sqrt{30}/10$	$\sqrt{5}/5$	0	-1/2	-1/2

		$Y=-1, I=1/2, \Xi$				$Y=-1, I=3/2$	
$\xi_1 \downarrow$		27	8 _D	8 _F	10	27	10*
$\Xi \pi$		$-\sqrt{5}/10$	$-3\sqrt{5}/10$	1/2	1/2	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		$\sqrt{5}/10$	$3\sqrt{5}/10$	1/2	1/2	$\sqrt{2}/2$	$\sqrt{2}/2$
$\Xi \eta$		$3\sqrt{5}/10$	$-\sqrt{5}/10$	-1/2	1/2		
ΔK		$3\sqrt{5}/10$	$-\sqrt{5}/10$	1/2	-1/2		

The phase factor $\xi_1 = \pm 1$, from de Swart's Table I, enters in his symmetry formula (14. 3):

$$(\mu_1 \mu_2 | \mu) = \xi_1 (-1)^{I_1 + I_2 - I} (\mu_2 \mu_1 | \mu).$$

This factor is irrelevant if you are doing your own self-consistent calculations; it enters when you try to check someone else who chose $\mu_2 \otimes \mu_1$ instead of $\mu_1 \otimes \mu_2$.

$$\{10\} \otimes \{8\} = \{35\} \oplus \{27\} \oplus \{10\} \oplus \{8\}.$$

* Four single coefficient tables are omitted; only the $\{27\}$ is -1; the three with $\{35\}$ are +1.

Multiplicity of 35;
●=1, x=2

		$Y=1, I=1/2, N$		$Y=1, I=3/2, \Delta$		
$\xi_1 \downarrow$		27	8	35	27	10
$\Delta \pi$		$-\sqrt{5}/5$	$-2\sqrt{5}/5$	1/4	$-\sqrt{5}/4$	$\sqrt{10}/4$
ΣK		$2\sqrt{5}/5$	$\sqrt{5}/5$	$\sqrt{5}/4$	3/4	$\sqrt{2}/4$
				$\sqrt{10}/4$	$-\sqrt{2}/4$	-1/2

		$Y=0, I=0, \Lambda$		$Y=0, I=1, \Sigma$				$Y=0, I=2$	
$\xi_1 \downarrow$		27	8	35	27	10	8	35	27
$\Sigma \pi$		$-\sqrt{10}/5$	$-\sqrt{15}/5$	$\sqrt{3}/6$	$-3\sqrt{5}/10$	$\sqrt{3}/3$	$-\sqrt{30}/15$	$\sqrt{3}/2$	-1/2
ΞK		$-\sqrt{15}/5$	$\sqrt{10}/5$	$\sqrt{2}/2$	$\sqrt{30}/10$	0	$-\sqrt{5}/5$	-1/2	$\sqrt{3}/2$
$\Sigma \eta$				$\sqrt{3}/3$	$-\sqrt{5}/5$	$-\sqrt{3}/3$	$\sqrt{30}/15$		
ΔK				$\sqrt{3}/6$	$\sqrt{5}/10$	$\sqrt{3}/3$	$2\sqrt{30}/15$		

		$Y=-1, I=1/2, \Xi$				$Y=-1, I=3/2$	
$\xi_1 \downarrow$		35	27	10	8	35	27
$\Xi \pi$		1/4	$-7\sqrt{5}/20$	$\sqrt{2}/4$	$-\sqrt{5}/5$	$\sqrt{2}/2$	$-\sqrt{2}/2$
$\Xi \eta$		3/4	$3\sqrt{5}/20$	$-\sqrt{2}/4$	$-\sqrt{5}/5$	$\sqrt{2}/2$	$\sqrt{2}/2$
ΩK		$\sqrt{2}/4$	$-3\sqrt{10}/20$	-1/2	$\sqrt{10}/5$	$\sqrt{2}/2$	$-\sqrt{2}/2$
ΣK		1/2	$\sqrt{5}/10$	$\sqrt{2}/2$	$\sqrt{5}/5$	$\sqrt{2}/2$	$\sqrt{2}/2$

		$Y=-2, I=0, \Omega^-$		$Y=-2, I=1$	
$\xi_1 \downarrow$		35	10	35	27
$\Omega \eta$		$\sqrt{2}/2$	$-\sqrt{2}/2$	1/2	$-\sqrt{3}/2$
ΞK		$\sqrt{2}/2$	$\sqrt{2}/2$	$\sqrt{3}/2$	1/2

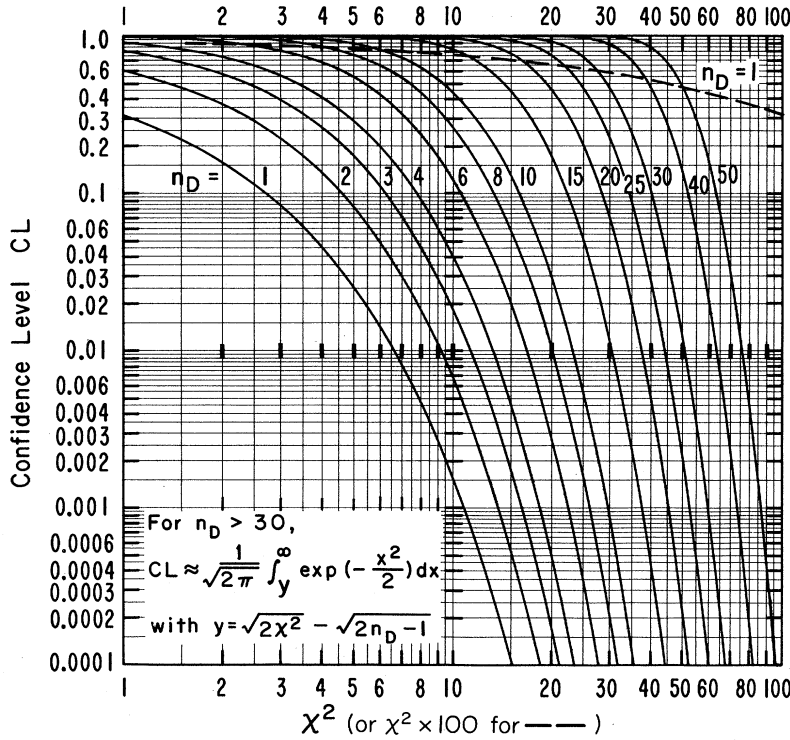
PROBABILITY AND STATISTICS

A. PROBABILITY DISTRIBUTIONS AND CONFIDENCE LEVELS

We give here properties of the three probability distributions most commonly used in high energy physics: Normal (or Gaussian), Chi-squared, and Poisson. We warn the reader that there is no universal convention for the term "confidence level"

as used by physicists; thus, explicit definitions are given for each distribution, and we have attempted to choose definitions that correspond to common usage. It is explained below how confidence levels for all three distributions can be extracted from the following figure.

χ^2 Confidence Level vs. χ^2 for n_D Degrees of Freedom

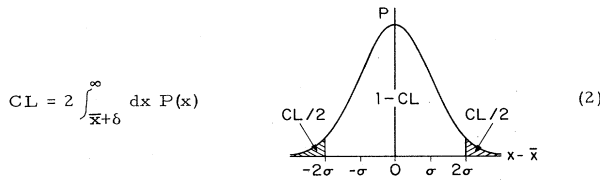


A.1. Normal Distribution

The normal distribution with mean \bar{x} and standard deviation σ (variance σ^2) is:

$$P(x)dx = \frac{1}{\sqrt{2\pi}\sigma} e^{-(x-\bar{x})^2/2\sigma^2} dx. \quad (1)$$

The confidence level associated with an observed deviation from the mean, δ , is the probability that $|x-\bar{x}| > \delta$, i.e.,



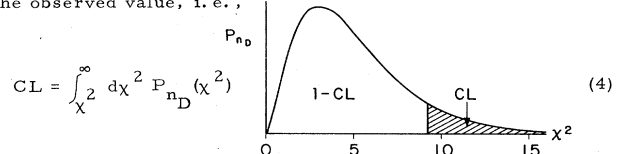
[The small figure in Eq. (2) is drawn with $\delta = 2\sigma$.] CL is given by the ordinate of the $n_D = 1$ curve in the figure at $\chi^2 = (\delta/\sigma)^2$. The confidence level for $\delta = 1\sigma$ is 31.7%; 2σ , 4.6%; 3σ , 0.3%. The central confidence interval, $1-CL$, (which is also sometimes called confidence level) for $\delta = 1\sigma$ is 68.3%; 2σ , 95.4%; 3σ , 99.7%. The odds against exceeding δ , $(1-CL)/CL$, for $\delta = 1\sigma$ are 2.15:1; 2σ , 21:1; 3σ , 370:1; 4σ , 16,000:1; 5σ , 1,700,000:1. Relations between σ and other measures of the width: probable error (CL = 0.5 deviation) = 0.67σ ; mean absolute deviation = 0.80σ ; RMS deviation = σ ; half width at half maximum = 1.18σ .

A.2. Chi-squared Distribution

The chi-squared distribution for n_D degrees of freedom is:

$$P_{n_D}(\chi^2) d\chi^2 = \frac{1}{2^h \Gamma(h)} (\chi^2)^{h-1} e^{-\chi^2/2} d\chi^2 \quad (\chi^2 \geq 0), \quad (3)$$

where h (for "half") = $n_D/2$. The mean and variance are n_D and $2n_D$, respectively. In evaluating Eq. (3) one may use Stirling's approximation: $\Gamma(h) = (h-1)! \approx 2.507 e^{-h} h^{h-1/2} \times (1 + 0.0833/h)$ which is accurate to $\pm 0.1\%$ for all $h \geq 1/2$. The confidence level associated with a given value of n_D and an observed value of χ^2 is the probability of chi-squared exceeding the observed value, i.e.,



[The small figure in Eq. (4) is drawn with $n_D = 5$ and $CL = 10\%$.] CL is plotted as a function of χ^2 for several values of n_D in the above figure. For large n_D , χ^2 becomes normally distributed about n_D . Thus,

$$y_1 = (\chi^2 - n_D) / \sqrt{2n_D} \quad (5)$$

becomes normally distributed with unit standard deviation. A better approximation, due to Fisher,¹ is that χ , not χ^2 , becomes normally distributed, specifically

$$y_2 = \sqrt{2\chi^2} - \sqrt{2n_D - 1} \quad (6)$$

approaches normality with unit standard deviation. For small CL's in particular, y_2 is much more accurate than y_1 . Thus, for $n_D = 50$ and $\chi^2 = 80$, the true CL = 0.45%, but y_1 is 3.0 corresponding to a CL of 0.13%, while y_2 is 2.7 corresponding to a CL of 0.35%.

PROBABILITY AND STATISTICS (Cont'd)

A.3. Poisson Distribution

The Poisson distribution with mean \bar{n} is:

$$P_{\bar{n}}(n) = \frac{e^{-\bar{n}} \bar{n}^n}{n!} \quad (n = 0, 1, 2, \dots) \quad (7)$$

The variance is equal to the mean. Confidence levels for Poisson distributions are usually defined in terms of quantities called "upper limits" as follows: The confidence level associated with a given upper limit N and an observed value n of n is the probability that $n > n_0$ if $\bar{n} = N$, i. e.,

$$CL = \sum_{n=n_0+1}^{\infty} P_N(n) = 1 - \sum_{n=0}^{n_0} P_N(n) \quad (8)$$

[The small figure in Eq. (8) is drawn with $n_0 = 2$ and $CL = 90\%$.] A useful relation between Poisson and chi-squared confidence levels allows one to look up this quantity on the above figure. Specifically, the quantity $1-CL$ is given by the ordinate of the $n_D = 2(n_0+1)$ curve at $\chi^2 = 2N$. Thus, 90% confidence level upper limits for $n_0 = 0, 1, \text{ and } 2$ are given by half the χ^2 value corresponding to an ordinate of 0.1 on the $n_D = 2, 4, \text{ and } 6$ curves, respectively; the values are $N = 2.3, 3.9, \text{ and } 5.3$.

Tables of confidence levels for all three of these distributions, the relation between Poisson and chi-squared confidence levels, and numerous other useful tables and relations may be found in Ref. 2.

B. STATISTICS

We consider here the situation in which one is presented with N independent data, $y_n \pm \sigma_n$, and it is desired to make some inference about the "true" value of the quantity represented by these data. For this purpose we interpret each datum y_n as a single sample point drawn randomly (and independently of the other data) from a distribution having mean \bar{y}_n (which we wish to estimate) and variance σ_n^2 . (Identification of the true σ_n with the σ_n datum is an approximation which may become seriously inaccurate when σ_n is an appreciable fraction of y_n .) Some methods of estimation commonly used in high energy physics are given below; see Ref. 3 for numerous applications. Section B.1. deals with the case in which all \bar{y}_n are the same, e. g., several different measurements of the same quantity; Sec. B.2. deals with the case in which $\bar{y}_n = \bar{y}(x_n)$, where x_n represents some set of independent variables, e. g., cross-section measurements at various values of energy and angle, $x_n = \{\theta_n, \theta_n\}$.

B.1. Single Mean and Variance Estimates

(1) If the y_n represent a set of values all supposedly drawn from a single distribution with mean \bar{y} and variance σ^2 (i. e., the σ_n are all the same, but their common value is unknown) then

$$\bar{y}_e = \frac{1}{N} \sum y_n \quad \text{and} \quad (9)$$

$$\sigma_e^2 = \frac{1}{N-1} \sum (y_n - \bar{y}_e)^2 \quad (10)$$

are unbiased estimates of \bar{y} and σ^2 . The variance of \bar{y}_e is σ^2/N . If the parent distribution is normal and N is large, the variance of σ_e^2 is $2\sigma^4/N$.

(2) If the \bar{y}_n all have the common value \bar{y} and the σ_n are known, then the weighted average

$$\bar{y}_e = \frac{1}{w} \sum w_n y_n \quad (11)$$

where $w_n = 1/\sigma_n^2$ and $w = \sum w_n$, is an appropriate unbiased estimate of \bar{y} . This choice of weighting factors in Eq. (11) minimizes the variance of the estimate; the variance is $1/w$.

B.2. Linear Least Squares Fit

A least squares fit of the function $y(x) = \sum_i a_i f_i(x)$ to independent data $y_n \pm \sigma_n$ at points x_n (e. g., a Legendre fit in which the f_i are Legendre polynomials and the a_i are Legendre coefficients) gives the following estimates of the parameters a_i :

$$a_{e,i} = \sum_{j=1}^n V_{ij} f_j(x_n) y_n / \sigma_n^2 \quad (12)$$

Here V is the covariance matrix of the fitted parameters

$$V_{ij} = \overline{(a_{e,i} - \bar{a}_{e,i})(a_{e,j} - \bar{a}_{e,j})} \quad (13)$$

which is given by

$$(V^{-1})_{ij} = \sum f_i(x_n) f_j(x_n) / \sigma_n^2 \quad (14)$$

The variance of an interpolated or extrapolated value of y at point x, $y_e = \sum a_{e,i} f_i(x)$, is:

$$(y_e - \bar{y}_e)^2 = \sum_{ij} V_{ij} f_i(x) f_j(x) \quad (15)$$

For the case of a straight line fit, $y(x) = a + bx$, one obtains the following estimates of a and b,

$$a_e = (S_y S_{xx} - S_x S_{xy}) / D, \quad (16)$$

$$b_e = (S_1 S_{xy} - S_x S_y) / D,$$

where

$$S_1, S_x, S_y, S_{xx}, S_{xy} = \sum (1, x_n, y_n, x_n^2, x_n y_n) / \sigma_n^2, \quad (17)$$

$$D = S_1 S_{xx} - S_x^2$$

The covariance matrix of the fitted parameters is:

$$\begin{pmatrix} V_{aa} & V_{ab} \\ V_{ab} & V_{bb} \end{pmatrix} = \frac{1}{D} \begin{pmatrix} S_{xx} & -S_x \\ -S_x & S_1 \end{pmatrix} \quad (18)$$

The variance of an interpolated or extrapolated value of y at point x is:

$$(y_e - \bar{y}_e)^2 = \frac{1}{S_1} + \frac{S_1}{D} \left(x - \frac{S_x}{S_1} \right)^2 \quad (19)$$

C. ERROR PROPAGATION

We consider here the situation in which one wishes to calculate the value and error of a function of some other quantities with errors, e. g., in a Monte Carlo program. Let {y} be a set of random variables with means $\{\bar{y}\}$ and covariance matrix V. Then the mean and variance of a function of these variables are approximately (to second order in $\{y - \bar{y}\}$):

$$\bar{f} \approx f(\{\bar{y}\}) + \frac{1}{2} \sum_{mn} V_{mn} \left(\frac{\partial^2 f}{\partial y_m \partial y_n} \right) \{\bar{y}\} = \{f\} \quad (20)$$

$$(f - \bar{f})^2 = \sum_{mn} V_{mn} \left(\frac{\partial f}{\partial y_m} \right) \{\bar{y}\} \left(\frac{\partial f}{\partial y_n} \right) \{\bar{y}\} = \{f\} \quad (21)$$

E. g., the mean and variance of a function of a single variable with mean \bar{y} and variance σ^2 are:

$$\bar{f} \approx f(\bar{y}) + \frac{1}{2} \sigma^2 f''(\bar{y}), \quad (22)$$

$$(f - \bar{f})^2 = \sigma^2 f'(\bar{y})^2. \quad (23)$$

Note that these equations will usually be applied by substituting some measured quantities, $\{\tilde{y}\}$ say, for the true means, $\{\bar{y}\}$. If, as is often the case, $\tilde{y}_n - \bar{y}_n$ is of order $\sqrt{V_{nn}}$, then there is no point in keeping the second order terms in Eq. (20) or (22) since the substitution itself introduces first order errors.

1. R. A. Fisher, Statistical Methods for Research Workers (Oliver and Boyd, Edinburgh and London, 1958).
2. M. Abramovitz and I. Stegun, eds., Handbook of Mathematical Functions (National Bureau of Standards, Applied Mathematics Series, Vol. 55, Washington, 1964).
3. W. T. Eadie, D. Drijard, F. E. James, M. Roos, and B. Sadoulet, Statistical Methods in Experimental Physics (North-Holland, Amsterdam and London, 1971).

Revised and expanded April 1974.

RELATIVISTIC KINEMATICS (Cont'd)

Application to the reaction $a + b \rightarrow c + d$:

1. General formulae

$$t \equiv t_{ac} = t_{bd} = t_{ac}^{\min} - 4p_a p_c \sin^2\left(\frac{\theta_{ac}}{2}\right),$$

with $t_{ac}^{\max} = (E_a - E_c)^2 - (p_a \mp p_c)^2$ (II-20)

which, after expansion in powers of $(1/s)$, gives

$$t_{ac}^{\min} = -\frac{(m_a^2 - m_c^2)(m_b^2 - m_d^2)}{s} - \frac{(m_a^2 + m_b^2 - m_c^2 - m_d^2)(m_a^2 m_b^2 - m_c^2 m_d^2)}{s^2} + O\left(\frac{1}{s^3}\right)$$
 (II-22)

In a similar way, defining $u_{ad} = (p_a - p_d)^2$, one finds

$$u \equiv u_{ad} = t_{bc} = u_{ad}^{\min} - 4p_a p_d \sin^2\left(\frac{\theta_{ad}}{2}\right).$$

$$u_{ad}^{\min} = -\frac{(m_a^2 - m_d^2)(m_b^2 - m_c^2)}{s} - \frac{(m_a^2 + m_b^2 - m_c^2 - m_d^2)(m_a^2 m_b^2 - m_c^2 m_d^2)}{s^2} + O\left(\frac{1}{s^3}\right)$$
 (II-23)

A general relation between the invariants is

$$s + t + u = m_a^2 + m_b^2 + m_c^2 + m_d^2.$$

IV. RELATIONS BETWEEN THE (j) PARTICLE AND (ij) PARTICLE REST FRAMES: LORENTZ TRANSFORMATION

The general Lorentz transformation has the matrix form

$$\begin{pmatrix} E \\ p_{\parallel} \\ p_{\perp} \end{pmatrix}^{(ij)} = \begin{pmatrix} \gamma & -\eta & 0 \\ -\eta & \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} E \\ p_{\parallel} \\ p_{\perp} \end{pmatrix}^{(j)}$$
 (IV-1)

If we define, in any frame, the quantity: $P_{\pm} = E \pm p_{\parallel}$,

$$\begin{pmatrix} P_{+} \\ P_{-} \\ p_{\perp} \end{pmatrix}^{(ij)} = \begin{pmatrix} (\gamma - \eta) & 0 & 0 \\ 0 & (\gamma + \eta) & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_{+} \\ P_{-} \\ p_{\perp} \end{pmatrix}^{(j)}$$
 (IV-3)

$$\gamma \mp \eta = \frac{s_{ij} + m_j^2 - m_i^2 \mp \sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{2m_j \sqrt{s_{ij}}} = e^{\mp \xi_{ij}}$$
 (IV-4)

$$\left\{ \begin{aligned} |\beta| &= \frac{|\eta|}{\gamma} = \tanh \xi_{ij} = \frac{\sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{s_{ij} + m_j^2 - m_i^2} = \frac{p_i^{(j)}}{E_i^{(j)} + m_j} \end{aligned} \right.$$
 (IV-5)

$$\left\{ \begin{aligned} \gamma &= \frac{1}{\sqrt{1 - \beta^2}} = \cosh \xi_{ij} = \frac{s_{ij} + m_j^2 - m_i^2}{2m_j \sqrt{s_{ij}}} = \frac{E_i^{(j)} + m_j}{\sqrt{s_{ij}}} \end{aligned} \right.$$
 (IV-6)

$$\left\{ \begin{aligned} |\eta| &= |\beta|\gamma = \sinh \xi_{ij} = \frac{\sqrt{\Delta(s_{ij}, m_i^2, m_j^2)}}{2m_j \sqrt{s_{ij}}} = \frac{p_i^{(j)}}{\sqrt{s_{ij}}} \end{aligned} \right.$$
 (IV-7)

The spatial part of the matrix in Eq. (IV-1) above can be written with the classical vector form

$$\vec{p}^{(ij)} = \vec{p}^{(j)} + \vec{\eta} \left(\frac{\vec{\eta} \cdot \vec{p}^{(j)}}{\gamma + 1} - E^{(j)} \right)$$
 (IV-9)

A convenient way to write the latter expression is

$$\vec{p}^{(ij)} = \vec{p}^{(j)} - \vec{\eta} \left[\frac{E^{(ij)} + E^{(j)}}{\gamma + 1} \right]$$
 (IV-10)

2. Particular case

In the c.m.: $dt = 2p_a p_c d(\cos \theta_{ac})$.

i) $m_a = m_c, m_b \neq m_d$:

$$t_{ac}^{\min} \approx -\frac{m_a^2(m_b^2 - m_d^2)^2}{s^2}$$
 (II-24)

$$u_{ad}^{\min} \approx -\frac{(m_a^2 - m_d^2)(m_b^2 - m_c^2)}{s}$$
 (II-25)

ii) For elastic scattering ($a + b \rightarrow a' + b'$),

$$t_{aa'}^{\min} = 0,$$

and $t_{aa'} = -4p_a^2 \sin^2\left(\frac{\theta_{aa'}}{2}\right)$,

$$dt_{aa'} = 2p_a^2 d(\cos \theta_{aa'});$$
 (II-26)

also $u_{ab'}^{\min} = \frac{(m_a^2 - m_b^2)^2}{s}$ (II-27)

At high energy:

$$e^{\xi_{ij}} \approx \frac{\sqrt{s_{ij}}}{m_j} \quad \text{and} \quad \xi_{ij} \approx \ln \frac{\sqrt{s_{ij}}}{m_j}.$$

ξ_{ij} is called the *boost parameter* of the Lorentz transformation that connects the particle ij rest frame to the particle j rest frame.

Application

1. To reaction a (beam) + b (target)

The transformation from the lab frame into the c.m. frame is given by:

$$\gamma = \frac{E_a^{\text{lab}} + m_b}{\sqrt{s}} \quad \text{and} \quad \eta = \frac{p_{\text{inc}}}{\sqrt{s}};$$

$$p_{\pm}^{\text{cm}} = e^{\mp \xi} p_{\pm}^{\text{lab}}$$

with

$$e^{\pm \xi} = \frac{s + m_b^2 - m_a^2 \mp \sqrt{\Delta(s, m_a^2, m_b^2)}}{2m_b \sqrt{s}}$$
 (IV-11)

2. The Lorentz transform of any vector

$\vec{p}_1 = (E_1, \vec{p}_1)$, given in a frame containing another vector $\vec{p}_2 = (E_2, \vec{p}_2)$, into the p_2 rest frame — where $p_2 = (m_2, 0)$, and $p_1 = (E_1, \vec{p}_1)$ — may be calculated using:

$$\gamma = \frac{E_2}{m_2} \quad \text{and} \quad \vec{\eta} = \frac{\vec{p}_2}{m_2};$$

$$\left. \begin{aligned} E_1' &= \gamma E_1 - \vec{\eta} \cdot \vec{p}_1 = \frac{E_1 E_2 - \vec{p}_1 \cdot \vec{p}_2}{m_2}, \\ \vec{p}_1' &= \vec{p}_1 - \frac{\vec{p}_2}{m_2} \left(\frac{E_1 + E_1'}{\gamma + 1} \right) = \vec{p}_1 - \vec{p}_2 \left(\frac{E_1 + E_1'}{E_2 + m_2} \right) \end{aligned} \right\}$$
 (IV-12)

RELATIVISTIC KINEMATICS (Cont'd)

V. RAPIDITY VARIABLE: DEFINITION AND KINEMATIC RELATIONS

Definition: in any system, for a particle of energy E and longitudinal momentum P_{\parallel} , the rapidity is defined by

$$y \equiv \frac{1}{2} \ln \left(\frac{E + P_{\parallel}}{E - P_{\parallel}} \right) = \tanh^{-1} \left(\frac{P_{\parallel}}{E} \right) \tag{V-1}$$

$$= \frac{1}{2} \ln \left(\frac{P_+}{P_-} \right) = \ln \left(\frac{P_+}{m_1} \right) = -\ln \left(\frac{P_-}{m_1} \right)$$

with the transverse mass $m_1 \equiv (m^2 + p_{\perp}^2)^{1/2}$.

From the last of the two equations above, one gets

$$P_+ = m_1 e^y \tag{V-2}$$

$$P_- = m_1 e^{-y} \tag{V-3}$$

so that $\begin{cases} E = m_1 \cosh y, \\ P_{\parallel} = m_1 \sinh y. \end{cases} \tag{V-4, V-5}$

A. Relations between the rapidities of different frames:

$$y^{(ij)} = y^{(j)} - \xi_{ij} \tag{V-7}$$

B. Application to a reaction $a+b \rightarrow c+d$ (or $a+b \rightarrow c+X$):

In the lab frame,

$$\left. \begin{aligned} p_a^{\text{lab}} &= (m_a \cosh y_a^{\text{lab}}, m_a \sinh y_a^{\text{lab}}, 0) \\ p_b^{\text{lab}} &= (m_b, 0, 0) \\ p_c^{\text{lab}} &= (m_c \cosh y_c^{\text{lab}}, m_c \sinh y_c^{\text{lab}}, \vec{p}_{1c}) \end{aligned} \right\} \tag{V-8}$$

(and a similar formula for particle d or X).

In the c.m. frame,

$$\left. \begin{aligned} p_a^{\text{cm}} &= (m_a \cosh y_a^{\text{cm}}, m_a \sinh y_a^{\text{cm}}, 0) \\ p_b^{\text{cm}} &= (m_b \cosh y_b^{\text{cm}}, m_b \sinh y_b^{\text{cm}}, 0) \\ p_c^{\text{cm}} &= (m_c \cosh y_c^{\text{cm}}, m_c \sinh y_c^{\text{cm}}, \vec{p}_{1c}) \end{aligned} \right\} \tag{V-9}$$

(and a similar formula for particle d or X).

$$y_a^{\text{lab}} = \cosh^{-1} \left(\frac{s - m_a^2 - m_b^2}{2m_a m_b} \right) \xrightarrow{\text{at large } s} y_a^{\text{lab}} \approx \ln \frac{s}{m_a m_b} \tag{V-10}$$

$$y_a^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_a^2 - m_b^2}{2m_a \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_a^{\text{cm}} \approx \ln \frac{\sqrt{s}}{m_a} \tag{V-11}$$

$$y_b^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_b^2 - m_a^2}{2m_b \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_b^{\text{cm}} \approx -\ln \frac{\sqrt{s}}{m_b} \tag{V-12}$$

$$y_c^{\text{cm}} = \cosh^{-1} \left(\frac{s + m_c^2 - m_d^2}{2m_c \sqrt{s}} \right) \xrightarrow{\text{at large } s} y_c^{\text{cm}} \approx \ln \frac{\sqrt{s}}{m_c} \tag{V-13}$$

and m_d^2 (or m_X^2) $\ll s$

With $\xi \approx \ln \frac{\sqrt{s}}{m_b}$, one gets, at large s:

$$y_c^{\text{lab}} \approx \ln \frac{s}{m_b m_c} \tag{V-14}$$

$$y^{\text{lab}} \approx \frac{1}{2} \ln \left(\frac{P^{\text{lab}} + P_{\parallel}^{\text{lab}}}{P^{\text{lab}} - P_{\parallel}^{\text{lab}}} \right) \approx -\ln \left(\tan \frac{\theta}{2} \right)$$

The maximum rapidity gap for the reaction $a+b \rightarrow c+d$ at large s is:

i) for the incoming particles (a,b):

$$y_{ab} = y_a^{\text{lab}} = y_a^{\text{cm}} - y_b^{\text{cm}} \approx \ln \frac{s}{m_a m_b} \tag{V-15}$$

ii) for the outgoing particles (c,d):

$$y_{cd} = y_c^{\text{lab}} - y_d^{\text{lab}} = y_c^{\text{cm}} - y_d^{\text{cm}} \approx \ln \frac{s}{m_c m_d} \tag{V-16}$$

Application to reaction $a+b \rightarrow 1+2 \dots +N$, with the N outgoing particles being ordered by increasing rapidity.

1. The maximum rapidity gap between the outgoing particles has an approximate value

$$y_{1N} \approx \ln \frac{s}{m_{11} m_{1N}} \tag{V-17}$$

2. The rapidity gap $(y_i - y_j)$ between the rapidities y_i and y_j of particles i and j, respectively, is related to the invariant mass squared of particles i and j

$$s_{ij} = m_i^2 + m_j^2 + 2m_i \cdot m_j \cosh(y_i - y_j) - 2\vec{p}_{1i} \cdot \vec{p}_{1j} \tag{V-18}$$

$$y_i - y_j = \cosh^{-1} \frac{s_{ij} - m_i^2 - m_j^2 + 2\vec{p}_{1i} \cdot \vec{p}_{1j}}{2m_i m_j} \tag{V-19}$$

3. The rapidity gap between either of the extreme particles (1 or N) and its adjacent neighbor [2 or (N-1) respectively] is, on a statistical basis, and in first approximation

$$|y_1 - y_2| \approx \ln \frac{s}{M_1^2} \tag{V-20}$$

and

$$|y_{N-1} - y_N| \approx \ln \frac{s}{M_N^2} \tag{V-21}$$

M_1 (M_N) being the missing mass with respect to particle 1 (N) in the reaction $a + b \rightarrow 1 + M_1$ ($a + b \rightarrow N + M_N$).

4. The rapidity gap y_{a1} between the incoming particle a and the outgoing particle closest to a, 1, is related to the transverse momentum $t_{a1} = (p_a - p_1)^2$:

$$y_a - y_1 = \cosh^{-1} \left(\frac{t_{a1} - m_a^2 - m_1^2}{2m_a m_{11}} \right) \tag{V-23}$$

RELATIVISTIC KINEMATICS (Cont'd)

VI. INCLUSIVE REACTIONS AND SCALING VARIABLES: DEFINITIONS AND KINEMATIC RELATIONS

A. Scaling Variable: Definitions

• For the *inclusive reaction* $a + b \rightarrow c + X$

1) Feynman's definition:

$$x = p_{\parallel}^{cm} / p_{\parallel \max}^{cm} \tag{VI-3}$$

where
$$p_{\parallel \max}^{cm} = \frac{\sqrt{\Delta(s, m_c^2, m_{X, \min}^2)}}{2\sqrt{s}} \tag{VI-4}$$

is the center-of-mass momentum of particle c and the lightest possible particle(s) $m_{X, \min}$ which, consistent with conservation laws, could recoil against particle c.

A relationship between x and m_X^2 which becomes useful at high energy [see (VI-6)] derives from

$$m_X^2 = (p_a + p_b - p_c)^2 = s + m_c^2 - 2p_c \cdot (p_a + p_b)$$

2) An alternative definition:

$$x' = p_{\parallel}^{cm} / p_{\parallel \max}^{cm} \tag{VI-7}$$

with
$$p_{\parallel \max}^{cm} = \sqrt{(p_{\max}^{cm})^2 - p_{\perp}^2}$$

• For *electro- and photoproduction*:

1) Bjorken's definition:

$$\omega = \frac{2m_p v}{q^2}$$

where
$$\left\{ \begin{array}{l} v = e - e' = \text{the difference between the incident and the outgoing energies} \\ m_p = \text{proton mass} \\ q^2 = \text{momentum transfer squared} \end{array} \right.$$

(NOTE: ω^{-1} is sometimes also called x .)

2) Bloom-Gilman's definition:

$$\omega' = \omega + \frac{m_p^2 - q^2}{q^2}$$

B. For the inclusive reaction $a + b \rightarrow c + X$, the *invariant differential cross section* $E(d^3\sigma/dp_c^3)$ may be written in the following forms using various expressions of the phase space volume d^3p_c/E_c (and omitting the subscript c):

Longitudinal variable of particle c	Transverse variable of particle c	Corresponding invariant (differential) cross section $E(d^3\sigma/dp^3)$
p_{lab}	θ^{lab}	$\frac{E^{\text{lab}}}{(p^{\text{lab}})^2} \frac{d^3\sigma}{dp^{\text{lab}} d\Omega^{\text{lab}}} \tag{VI-10}$
p_{\parallel}	p_{\perp}	$\frac{E}{\pi} \frac{d^2\sigma}{dp_{\parallel} dp_{\perp}^2}$ (averaged over azimuth) $\tag{VI-11}$
x	p_{\perp}	$\frac{E^{cm}}{p_{\max}^{cm}} \cdot \frac{1}{\pi} \frac{d^2\sigma}{dx dp_{\perp}^2} \approx \frac{2E^{cm}}{\pi\sqrt{s}} \frac{d^2\sigma}{dx dp_{\perp}^2}$ (averaged over azimuth) $\tag{VI-12}$
y	p_{\perp} (or m_{\perp})	$\frac{1}{\pi} \frac{d^2\sigma}{dy dp_{\perp}^2} = \frac{1}{\pi} \frac{d^2\sigma}{dy dm_{\perp}^2}$ (averaged over azimuth) $\tag{VI-13}$
$m_X^2 = (p_a + p_b - p_c)^2$	$t = (p_a - p_c)^2$	$\frac{\sqrt{\Delta(s, m_a^2, m_b^2)}}{\pi} \frac{d^2\sigma}{dm_X^2 dt}$ (averaged over azimuth) $\tag{VI-14}$

• At high energy,

$$p_{\max}^{cm} \approx \frac{\sqrt{s}}{2} \quad \text{and} \quad x \approx \frac{2p_{\parallel}^{cm}}{\sqrt{s}} \tag{VI-5}$$

The relation between x and y is trivial:

$$x \approx 2 \frac{p_{\parallel}^{cm}}{\sqrt{s}} \approx \frac{2m_{\perp}}{\sqrt{s}} \sinh y^{cm} \tag{VI-8}$$

and, for y also large:

$$y^{cm} \approx \ln \frac{x\sqrt{s}}{m_{\perp}} \tag{VI-9}$$

$$x = 1 - \frac{m_X^2 - m_c^2}{s} \approx 1 - \frac{m_X^2}{s} \tag{VI-6}$$

(for $m_X^2 \gg m_c^2$).

For p_{\perp} fixed as $s \rightarrow \infty$,

$$x' \approx \frac{2p_{\parallel}^{cm}}{\sqrt{s}} \approx x$$

For q^2 and v large (Bjorken limit),

$$\omega' \approx \omega - 1$$

*See Denyse M. Chew, "Relativistic Kinematics", Particle Data Group Memo PDG-101, for details and references.

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PARTICLE DETECTORS, ABSORBERS, AND RANGES*

A. DETECTOR PARAMETERS

In this section we give various parameters for common detectors. The quoted numbers represent at best an order of magnitude useful only for preliminary design.

A.1 Scintillators: Photon yield $\approx 1\gamma/100$ eV in plastic scintillator¹ and $\approx 1\gamma/25$ eV in NaI.^{1,2}

A.2 Čerenkov³: Half angle θ_c of cone aperture in terms of velocity β and index of refraction n :

$$\theta_c = \arccos\left(\frac{1}{\beta n}\right) \sim \sqrt{2\left(1 - \frac{1}{\beta n}\right)}$$

Threshold velocity: $\beta_t = 1/n$; $\gamma_t = 1/\sqrt{1-\beta_t^2}$.

Therefore, $\beta_t \gamma_t = 1/\sqrt{2\delta + \delta^2}$, where $\delta = n-1$. Values of δ for various commonly used gases are given as a function of pressure and wavelength in Ref. 4; for values at atmospheric pressure, see the Table of Atomic and Nuclear Properties following.

Number of photons N per cm:

$$N = \frac{\alpha}{c} \int \left(1 - \frac{1}{\beta^2 n^2}\right) 2\pi \, d\nu = \frac{\alpha}{c} \beta_t^2 \int \left(\frac{1}{\beta_t^2 \gamma_t^2} - \frac{1}{\beta^2 \gamma^2}\right) 2\pi \, d\nu$$

$$\approx 500 \sin^2 \theta_c / \text{cm (visible spectrum)}$$

A.3 Photon Collection: In addition to the photon yield, one should take into account the light collection efficiency ($\approx 10\%$ for typical 1 cm thick scintillator), attenuation length (≈ 1 to 4 m for typical scintillators⁵), and quantum efficiency of the photomultiplier cathode ($\leq 25\%$).

A.4 Bubble, Streamer, Wire Chambers:

Chamber Type	Accuracy(rms)	Resolution Time	Dead Time
Bubble	$\pm 75\mu$	≈ 1 ms	$\approx 1/20$ s ^{a)}
Streamer	$\pm 300\mu$	≈ 2 μ s	≈ 100 ms
Optical Spark	$\pm 200\mu$ ^{b)}	≈ 2 μ s	≈ 10 ms
Magnetostrictive Spark	$\pm 500\mu$	≈ 2 μ s	≈ 10 ms
Proportional	$\approx \pm 300\mu$ ^{c, d)}	≈ 50 ns	≈ 200 ns
Drift	± 50 to 300μ	≈ 2 ns ^{e)}	≈ 100 ns

- a) Multiple pulsing time.
- b) 60μ for high pressure.
- c) 300μ is for 1 mm pitch.
- d) Delay line cathode readout can give $\pm 150\mu$ parallel to anode wire.
- e) For two chambers.

A.5 Shower Detectors: Typical energy resolution for incident electron in the 1 GeV range, E in GeV.

NaI (20 rad. lengths) ⁶ :	$\frac{2\%}{E^{1/4}}$	(FWHM)
Lead Glass (14 rad. lengths) ⁷ :	$\frac{8-12\%}{\sqrt{E}}$	(FWHM)
Lead Scintillator Sandwich (10 rad. lengths) ⁸ :	$\frac{22\%}{\sqrt{E}}$	(FWHM)

(10 lead plates, each 1 radiation length, and scintillators, each 1 cm in thickness).

A.6 Proportional Chamber Wire Instability: The limit on the voltage V for a wire tension T is given by⁹

$$V \leq \frac{sT^{1/2}}{lC}$$

where s , l , and C are the wire spacing, length, and capacitance per unit length. An approximation to C for a chamber half gap t and wire diameter d (good for $s \leq t$) gives¹⁰

$$V \leq 59T^{1/2} \left[\frac{t}{l} + \frac{s}{\pi l} \ln\left(\frac{s}{\pi d}\right) \right],$$

where V is in kV, and T is in grams.

B. COSMIC RAY FLUXES

The fluxes of particles of different types depend on the latitude, their energy, and on the conditions of measurement. Some typical sea-level values¹¹ are given below:

- I_v vertical flux
- J_1 total flux crossing a unit horizontal area
- J_2 total flux crossing a sphere of unit radius

	Total Intensity	Hard component	Soft component
I_v	1.14×10^{-2}	0.83×10^{-2}	0.31×10^{-2} cm ⁻² sec ⁻¹ sterad ⁻¹
J_1	1.79×10^{-2}	1.27×10^{-2}	0.52×10^{-2} cm ⁻² sec ⁻¹
J_2	2.41×10^{-2}	1.68×10^{-2}	0.73×10^{-2} cm ⁻² sec ⁻¹

Very approximately, about 75% of all particles at sea-level are penetrating, and are μ mesons. The absolute flux of protons at sea-level, in a momentum range 700-1100 MeV/c, is 1.5×10^{-5} cm⁻²sec⁻¹sterad⁻¹, or $\sim 0.1\%$ of all particles.

C. PASSAGE OF PARTICLES THROUGH MATTER

C.1 dE/dx ¹²: Ionization energy loss for particles heavier than an electron and having charge $z|e|$ (Bethe-Bloch equation):

$$\frac{1}{\beta} \frac{dE}{dx} = z^2 D \frac{Z}{A \beta^2} \left[\ln\left(\frac{2m_e c^2 \beta^2 T'_{\max}}{(1-\beta^2) I^2}\right) - 2\beta^2 - K \right],$$

where $D = 2\pi N_e^2 m_e c^2 = 0.1535$ MeV cm²/g (see Physical and Numerical Constants Table), Z and A are charge and mass number of the material, $\beta = v/c$ for the incident particle, T'_{\max} is the maximum-energy δ ray considered to be ionization loss, I is the mean ionization potential ($I \approx 20Z$ eV for $Z = 1$, tending toward $I \approx 10Z$ for $Z \geq 20$), and K includes the shell correction (maximum $\approx 10\%$ around $\beta = 0.1$ in the heaviest elements), density effect (important at high velocities and high densities), and the z^3 correction¹³ (important at low velocity, a few percent at $\beta \approx 0.1$).

A point of minimum ionization occurs in the range $\beta \approx 0.95$ to 0.97, with the dependence on absorber confined to that range. Energy loss increases approximately as $\ln \gamma$ thereafter (relativistic rise), until modified by the density effect (which reduces the rate of increase). Values of dE/dx at minimum are listed in the Table of Atomic and Nuclear Properties of Materials (following).

Figures are given below for ionization energy loss and range in lead (with scaling indicated for copper, aluminum, and carbon), and in liquid hydrogen, with T'_{\max} taken to be T'_{\max} as defined in Sec. C.2 below. Energy losses due to bremsstrahlung, nuclear collisions, etc., are not included. For thin samples of absorber, fluctuations are very important.¹⁴ A review of the problem of the penetration of charged particles into matter may be found in Fano¹⁵.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mixtures or compounds may be treated in first (but fairly accurate) approximation using Bragg's additivity rule:

$$\left(\frac{1}{\rho} \frac{dE}{dx}\right)_{\text{mix}} = \sum_i \left(\frac{1}{\rho} \frac{dE}{dx}\right)_i f_i,$$

where f_i is the fractional weight ($= \rho_i / \rho_{\text{mix}}$) of element i .

C.2 δ rays¹⁶: Number N per g/cm^2 per MeV of δ rays of kinetic energy T produced by an incident particle of momentum p and velocity β (Rutherford formula):

$$\frac{dN}{dTdx} = D \frac{Z}{A} \frac{1}{\beta^2 T^2},$$

where D is defined in the dE/dx section. This expression is valid for any unit charge incident particle if $T \ll T_{\text{max}}$. T_{max} is the maximum δ -ray energy, given by

$$T_{\text{max}} = \frac{2 m_e c^2 (\text{pc})^2}{s} = \frac{2 \beta^2 \gamma^2 m_e c^2}{1 + 2\gamma \frac{m_e}{M} + \frac{m_e^2}{M^2}}$$

where m_e and M are the electron and incident particle masses and s is their center-of-mass energy squared.

C.3 Multiple Coulomb Scattering^{17, 18, 19}: Sketched below are the multiple scattering angle θ , displacement y , and sagitta s , of a particle passing through a scatterer of thickness L . If the incident particle has mass M , charge $z|e|$, velocity β , and energy E , and the scatterer has atomic mass m and radiation length L_R , the projection of the $1/e$ multiple scattering angle (i.e., the value of the scattering angle at which the distribution falls to $1/e$ of its peak value) is given by

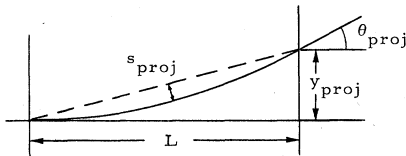
$$\theta_{\text{proj}}^{1/e} = z \frac{14 \text{MeV}/c}{\beta \gamma} \sqrt{\frac{L}{L_R}} \left[1 + \frac{1}{9} \log_{10} \left(\frac{L}{L_R} \right) \right] \left[1 + \frac{M^2}{E m_s} \right],$$

in the small angle approximation. The distribution is approximately Gaussian for $\theta < \theta^{1/e}$, but has long tails.²⁰ The accuracy of this expression is $\approx 5\%$ except for low z or β , where it is $\approx 10\text{-}20\%$.

For a given θ_{proj} ,

$$y_{\text{proj}} = \frac{L \theta_{\text{proj}}}{\sqrt{3}},$$

$$s_{\text{proj}} = \frac{L \theta_{\text{proj}}}{4 \sqrt{3}}.$$



C.4 Electron Range in Lead, Copper, Carbon and Hydrogen: See figure following.

C.5 Fractional Energy Loss for Electrons and Positrons in Lead: See figure following.

C.6 Contributions to Photon Cross Section in Lead: See figure following.

C.7 Photon Cross Section: See figure following.

D. ATOMIC AND NUCLEAR PROPERTIES OF MATTER

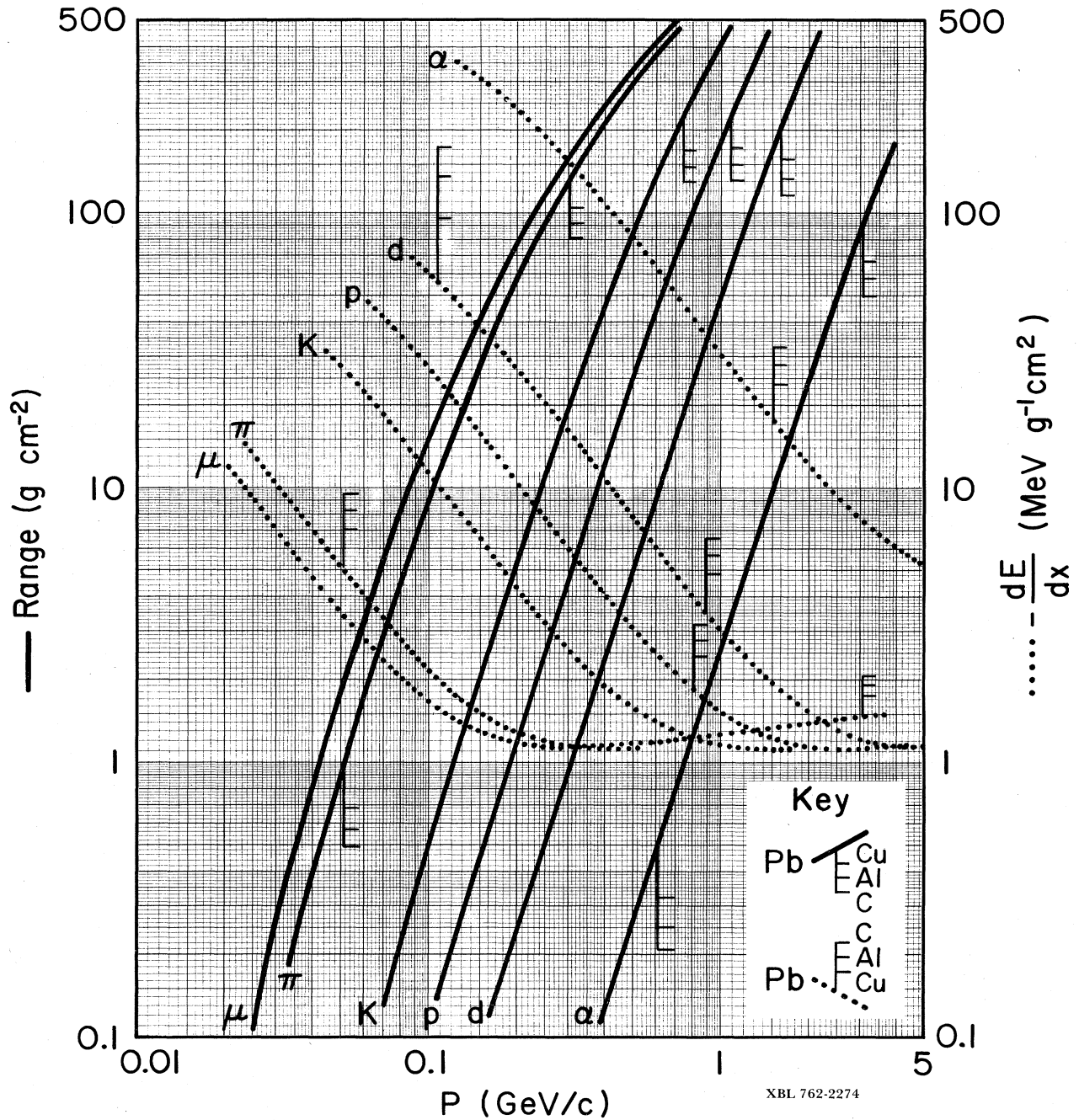
See table following.

*Prepared April 1974 by Sherwood Parker and Bernard Sadoulet. Revised April 1976.

1. Methods of Experimental Physics, L. C. L. Yuan and C.-S. Wu, editors, Academic Press, 1961, Vol. 5A, p. 127.
2. R. K. Swank, *Ann. Rev. Nuc. Sci.* **4**, 137 (1954) and G. T. Wright, *Proc. Phys. Soc.* **B68**, 929 (1955).
3. Methods of Experimental Physics, L. C. L. Yuan and C.-S. Wu, editors, Academic Press, 1961, Vol. 5A, p. 163.
4. E. R. Hayes, R. A. Schlutter, and A. Tamosaitis, "Index and Dispersion of Some Čerenkov Counter Gases," ANL-6916 (1964).
5. Nuclear Enterprises Catalogue.
6. E. B. Hughes et al., *IEEE Transactions on Nuclear Science*, **NS-19**, No. 3, 126 (1972).
7. M. Holder et al., *Phys. Letters* **40B**, 141 (1972) and J. S. Beale et al., "A Lead-Glass Čerenkov Detector for Electrons and Photons," CERN Writeup, Intl. Conf. on Instrumentation in H. E. P., Frascati (1973).
8. J. K. Walker and T. R. Knasel, *Rev. Sci. Instr.* **37**, 913 (1966).
9. T. Trippe, CERN NP Internal Report 69-18 (1969).
10. S. Parker and R. Jones, LBL-797 (1972), and A. Morse and B. Feshbach, Methods of Theoretical Physics, McGraw-Hill, New York, 1953, p. 1236.
11. B. Rossi, *Rev. Mod. Phys.* **20**, 537 (1948).
12. H. Bichsel, Passage of Charged Particles Through Matter, American Institute of Physics Handbook, McGraw-Hill, 1972, Sec. 8, p. 142.
13. H. H. Heckman et al., *Phys. Rev. Lett.* **22**, 871 (1969) and J. D. Jackson et al., *Phys. Rev. B* **6**, 4131 (1972).
14. John H. Cobb, "A Study of Some Electromagnetic Interactions of High Velocity Particles with Matter," (thesis), University of Oxford Report HEP/T/55 (1975).
15. U. Fano, *Ann. Rev. Nucl. Sci.* **13**, 1 (1963).
16. See B. Rossi, High Energy Particles (Prentice Hall, 1952), Section 2.3 for more accurate formulae.
17. Virgil L. Highland, *Nucl. Instr. and Meth.* **129**, 497 (1975).
18. J. B. Marion and B. A. Zimmerman, *Nucl. Instr. and Meth.* **51**, 93 (1967).
19. G. Z. Moliere, *Naturforsch.* **3(a)**, 78 (1948).
20. See, for example, the experimental work of A. D. Hansen, L. H. Lanzl, E. M. Lyman, and M. B. Scott, *Phys. Rev.* **84**, 634 (1951).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

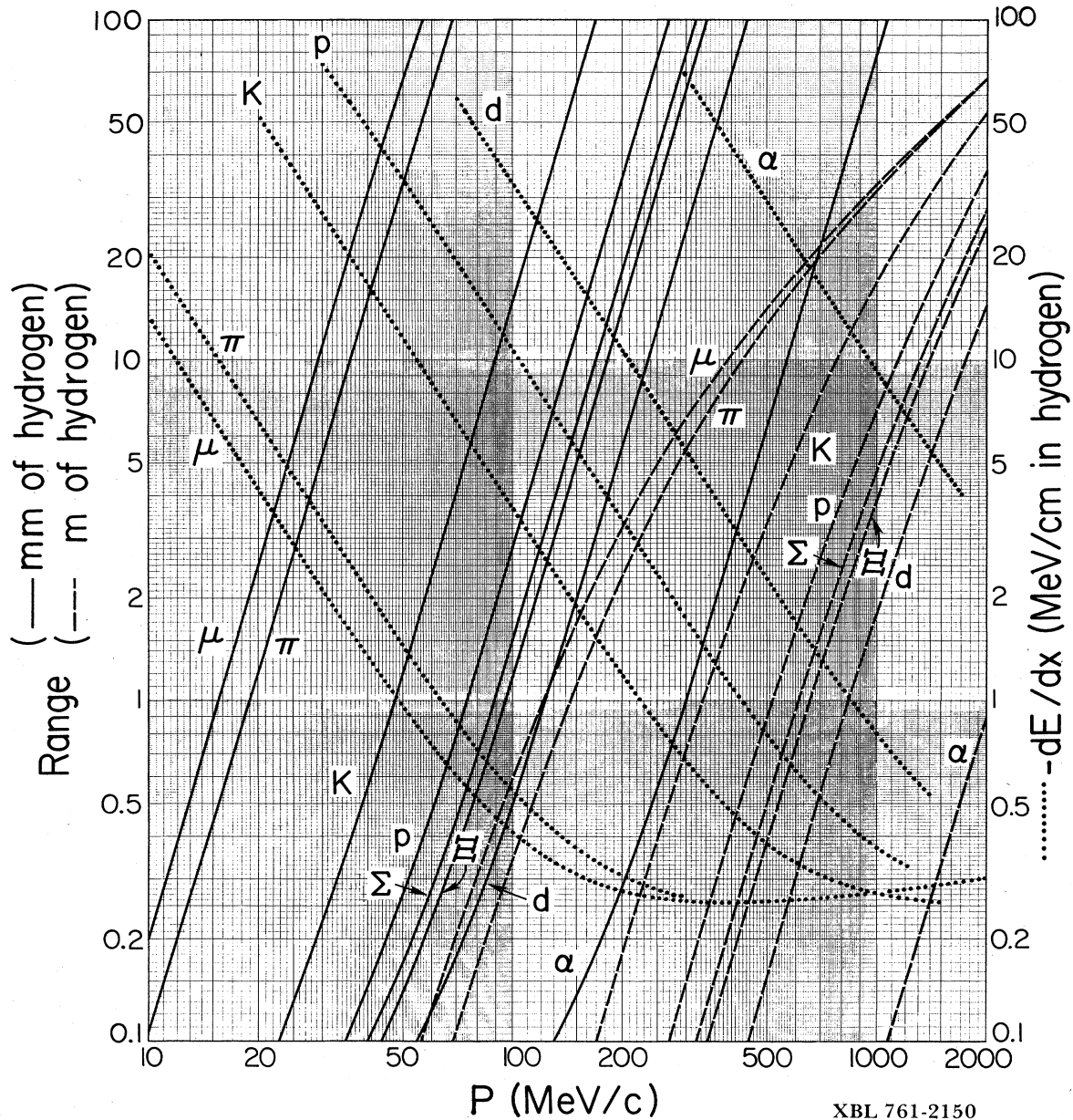
Mean Range and Energy Loss in Lead, Copper, Aluminum, and Carbon



Mean range and energy loss due to ionization for the indicated particles in Pb, with scaling to Cu, Al, and C indicated, using Bethe-Bloch equation (Section C.1 above) with corrections. Calculated using program of Hans Bichsel (UCRL-17538), with density correction added (Hans Bichsel, private communication). See also Joseph F. Janni [Air Force Weapons Laboratory Technical Report No. AFWL-TR-65-150 (1966)]. The average ionization potentials (I) assumed were: Pb (820 eV), Cu (320 eV), Al (166 eV), and C (77.5 eV). Figure indicates total path length; observed range may be smaller (by ~ 1% - 2% in heavy elements) due to multiple scattering, primarily from small energy-loss collisions with nuclei. The functional forms have not been experimentally verified to better than roughly ± 1%. For higher energies refer to discussion by Cobb ["A Study of Some Electromagnetic Interactions of High Velocity Particles with Matter," University of Oxford Report HEP/T/55 (1973)] and by Turner ["Penetration of Charged Particles in Matter: A Symposium", National Academy of Sciences, Washington D. C. (1970), p. 48]. Scaling to other beam particles is, to a good approximation, described by the expression on the next page.

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mean Range and Energy Loss in Liquid Hydrogen



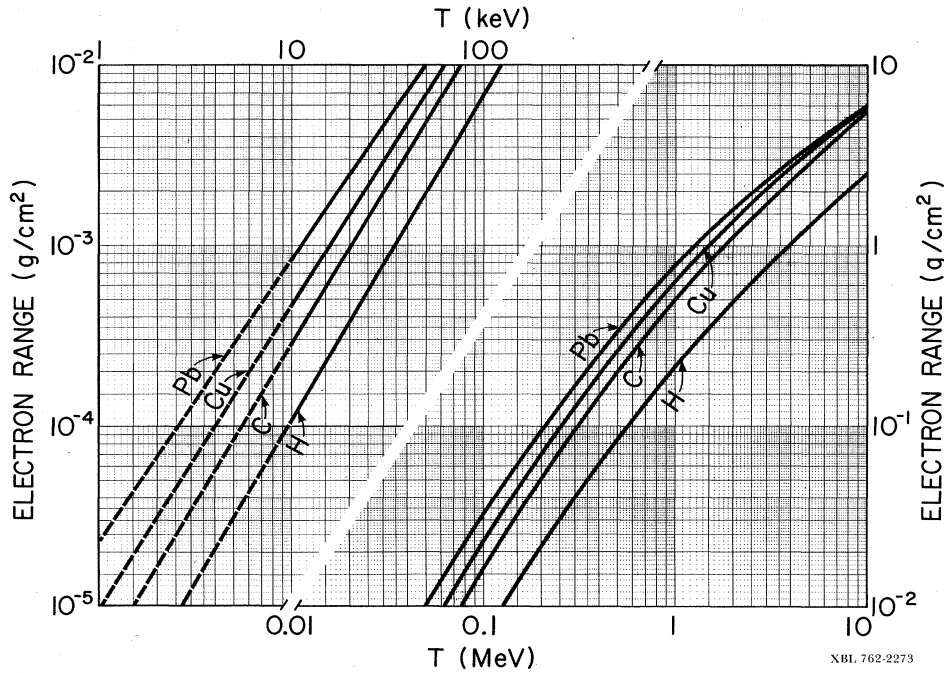
Range and energy loss in liquid hydrogen bubble chamber, based on Bethe-Bloch equation (Section C.1 above), using an average ionization potential for H_2 of $I = 20.0$ eV, which is an approximate average of the recent experimental result of Garbincius and Hyman [Phys. Rev. **A2**, 1834 (1970)] and the recent theoretical result of Ford and Browne [Phys. Rev. **A7**, 418 (1973)]: This is somewhat higher than used in the last edition. Bubble chamber conditions are chosen to be those of Garbincius and Hyman: parahydrogen of density = 0.0625 g/cm³ (note: range $\propto 1/\text{density}$), with vapor-pressure 60.8 lb/in² (absolute) and temperature 26.2°K . The functional dependence of the Bethe-Bloch equation is not experimentally verified to better than about $\pm 1\%$ over large momentum ranges. It should be noted that the number of bubbles per cm of a track in a bubble chamber is nearly proportional to $1/\beta^2$, not dE/dx . For the linear portions of the range curves, $R \propto p^{3.6}$.

Scaling law for particles of other mass or charge (except electrons): for a given medium, the range R_b of any beam particle with mass M_b , charge z_b , and momentum p_b is given in terms of the range R_a of any other particle with mass M_a , charge z_a , and momentum $p_a = p_b M_a / M_b$ (i. e., having the same velocity) by the expression

$$R_b(M_b, z_b, p_b) = \left[\frac{M_b/M_a}{z_b^2/z_a^2} \right] R_a(M_a, z_a, p_a = p_b M_a / M_b).$$

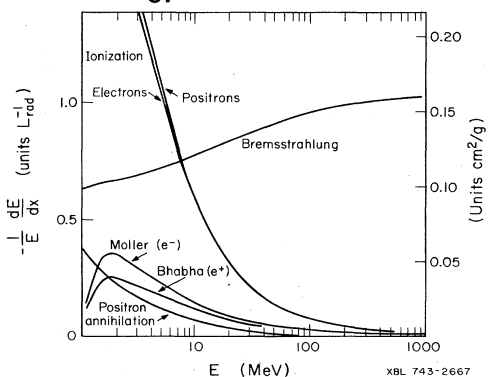
PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Mean Electron Range in Lead, Copper, Carbon, and Liquid Hydrogen



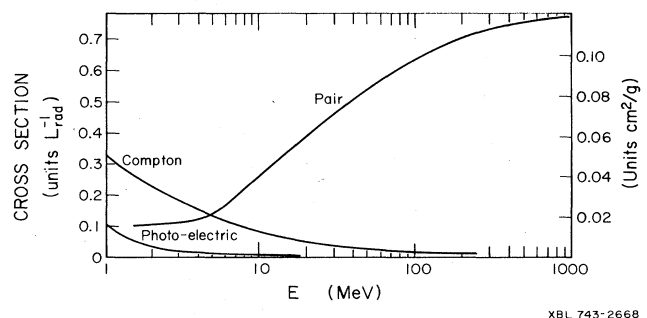
Mean range of electrons in the continuous-slowing-down approximation, taking into account energy loss by collisions with atomic electrons and by bremsstrahlung; strong fluctuations are to be expected for individual tracks. This range is the total path length; the practical range is shorter because of multiple Coulomb scattering, which becomes increasingly important as the electron slows down. E.g., for a fast electron the rms projected angle due to multiple Coulomb scattering reaches 1 radian by the time the electron has slowed to 0.4 MeV in hydrogen, 1.5 MeV in carbon, 9 MeV in copper, and 24 MeV (off scale) in lead. The "critical energy" (above which the energy loss due to bremsstrahlung exceeds that due to ionization and showering becomes important) is 400 MeV for hydrogen, 100 MeV for carbon, 25 MeV for copper, and 10 MeV for lead. The mean positron range may differ from the mean electron range by several percent. See Berger and Seltzer, NASA SP-3012 (1964) and SP-3036, and P. Trower, UCRL-2426, Vol. III Rev. (1966). 1-10 keV range was obtained by linear extrapolation.

Fractional Energy Loss for e⁺ and e⁻ in Lead



Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization when the energy loss per collision is below 0.255 MeV, and as Moller (Bhabha) scattering when it is above.

Contributions to Photon Cross Section in Lead

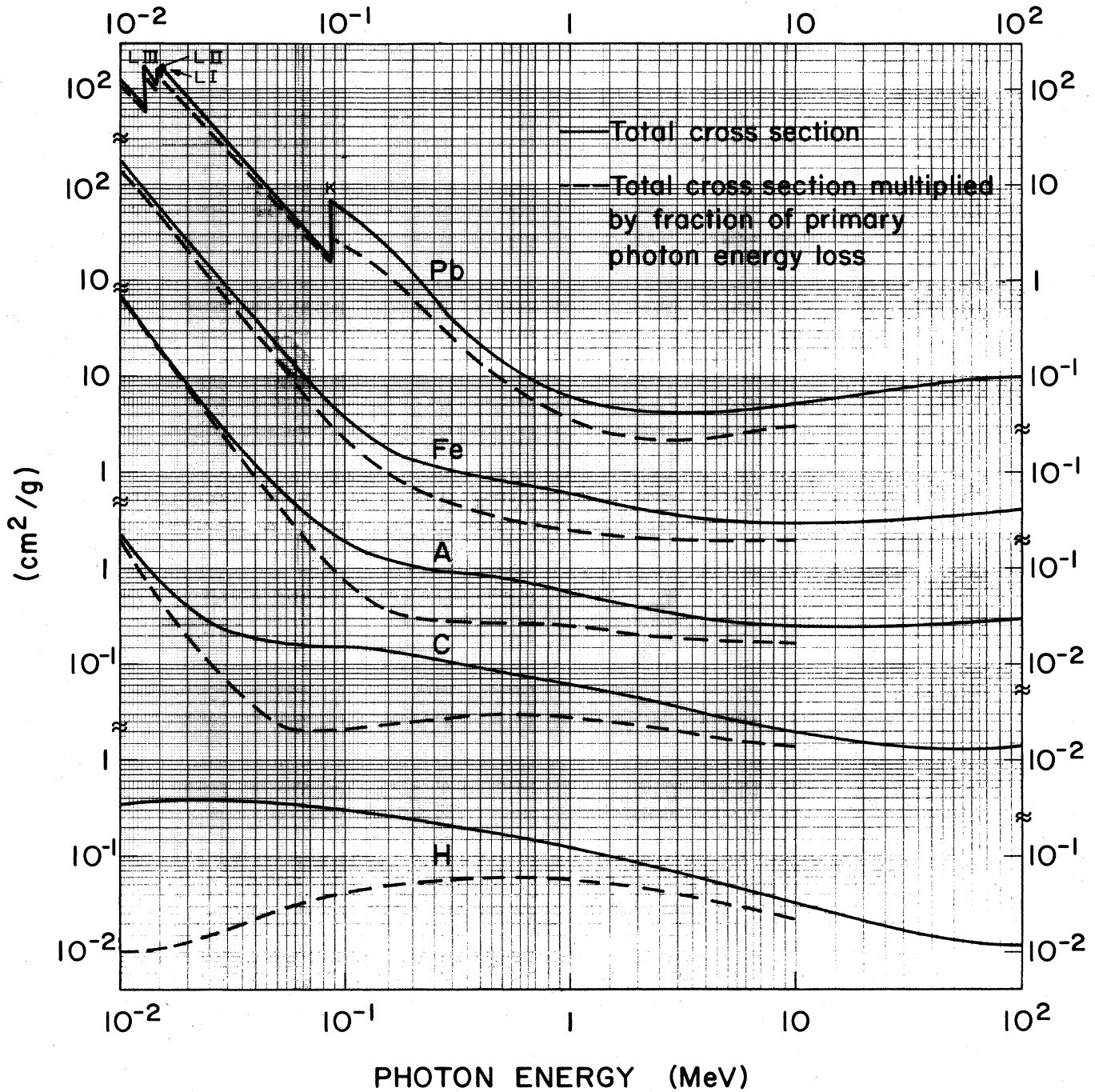


Photon cross section in lead in inverse radiation lengths as a function of photon energy. The intensity of photons can be expressed as $I = I_0 \exp(-\sigma x)$, where σ is read above and x is the path length in radiation lengths. See also figure following.

These figures are adapted from Fig. 3.2 and Fig. 3.3 from Messel and Crawford, *Electron-Photon Shower Distribution Function Tables for Lead, Copper and Air Absorbers*, Pergamon Press, 1970. Messel and Crawford use $L_r(\text{Pb}) = 5.82 \text{ g/cm}^2$, but we have modified the figures to reflect the value given in the Table of Atomic and Nuclear Properties of Materials (following), namely $L_r(\text{Pb}) = 6.4 \text{ g/cm}^2$. Note that the development of electron-photon cascades is approximately independent of absorber when the results are expressed in terms of inverse radiation lengths (i. e., scales on left of plots).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Photon Cross Sections



XBL 743-2662

The photon cross section as a function of photon energy for various substances. The solid curve for each substance gives the total cross section in (cm²/g). The dashed curve is the total cross section multiplied by the fraction of energy deposited by the photon (different from solid curve because of Compton scattering and bremsstrahlung of secondary electrons). See J. H. Hubbell, NSRDS-NBS 29 (1969).

PARTICLE DETECTORS, ABSORBERS, AND RANGES (Cont'd)

Atomic and Nuclear Properties of Materials*

Material Z	A	Nuclear cross section σ^a [barns]	Nuclear collision length L_{coll}^b [cm]	Absorption length λ^b [cm]	dE/dx min ^c [$\frac{MeV}{g/cm^2}$] [$\frac{MeV}{cm}$]	Radiation length L_{rad}^d [cm]	Density ^e [g/cm^3] [g/l]	Refractive index $n; n_e^e$ () is for gas () is $(n-1) \times 10^6$ for gas				
H ₂	1	1.01	0.039	43.0	607	790	4.12	0.292	63.05	890	{ 0.0708 (0.090)	{ 1.112 (140)
D ₂	1	2.01	0.074	45.1	273	342	2.07	0.342	126.1	764	0.163	1.128
He	2	4.00	0.134	49.6	397	478	1.94	0.243	94.32	755	{ 0.125 (0.178)	{ 1.024 (35)
Li	3	6.94	0.215	53.6	100.4	120.6	1.65	0.881	82.76	155	0.534	-
Be	4	9.01	0.270	55.4	30.0	36.7	1.61	2.97	65.19	35.3	1.848	-
C	6	12.01	0.340	58.7	≈37.8	49.9	1.78	≈2.76	42.70	≈27.5	≈1.55 ^f	-
N ₂	7	14.01	0.390	59.7	73.8	99.4	1.82	1.47	37.99	47.0	{ 0.808 (1.25)	{ 1.205 (300)
Ne	10	20.18	0.520	64.4	53.7	74.9	1.73	2.08	28.94	24.0	{ 1.207 (0.90)	{ 1.092 (67)
Al	13	26.98	0.650	68.9	25.5	37.2	1.62	4.37	24.01	8.9	2.70	-
A	18	39.95	0.890	74.5	53.2	80.9	1.51	2.11	19.55	14.0	{ 1.40 (1.78)	{ 1.233 (283)
Fe	26	55.85	1.160	79.9	10.2	17.1	1.48	11.6	13.84	1.76	7.87	-
Cu	29	63.54	1.270	83.1	9.3	14.8	1.44	12.9	12.86	1.43	8.96	-
Sn	50	118.69	2.040	96.6	13.2	22.8	1.28	9.4	8.82	1.21	7.31	-
W	74	183.85	2.810	108.6	5.6	10.3	1.17	22.6	6.76	0.35	19.3	-
Pb	82	207.19	3.080	111.7	9.8	18.5	1.13	12.8	6.37	0.56	11.35	-
U	92	238.03	3.380	116.9	≈6.2	12.0	1.09	≈20.7	6.00	≈0.32	≈18.95	-
Air				60.2	50000 ^g	67500 ^g	1.82	0.0022 ^g	36.20	30050 ^g	{ 0.001205 ^g (1.29)	{ 1.000273 ^g (293)
H ₂ O				58.3	58.3	78.8	2.03	2.03	36.08	36.1	1.00	1.33
H ₂ (bubble chamber 26°K) ^h				43.0	≈683	887	4.12	≈0.26	63.05	≈1000	≈0.063 ^h	1.112
D ₂ (bubble chamber 31°K) ^h				45.1	≈322	403	2.07	≈0.29	126.1	≈900	≈0.140 ^h	1.110
H-Ne mixture (50 mole percent) ⁱ				62.9	154.5	215	1.84	0.75	29.70	73.0	0.407	1.092
Propane (C ₃ H ₈) ^j				55.0	134	176	2.28	0.98	45.38	111	{ 0.41 ^j (2.0)	{ 1.25 ^j (1005)
Freon 13B1 (CF ₃ Br) ^j				74.3	≈49.5	73.5	1.52	≈2.3	16.53	≈11	{ ≈1.50 ^j (8.71)	{ 1.238 ^j (750)
Iford emulsion				79.5	23.6	39.1	1.44	5.49	11.02	2.94	3.815	-
NaI				91.9	25.0	41.3	1.32	4.84	9.49	2.59	3.67	1.775
LiF				61.1	23.1	30.7	1.69	4.46	39.25	14.9	2.64	1.394
Polyethylene (CH ₂)				55.7	≈59.6	78.4	2.09	≈1.95	44.78	≈48	0.92-0.95	-
Mylar (C ₅ H ₄ O ₂)				58.5	42.1	56.1	1.91	2.65	39.95	28.7	1.39	-
Polystyrene, scintillator (CH) ^k				57.0	55.2	68.5	1.97	2.03	43.8	42.9	1.032	1.581
Lucite, Plexiglas (C ₅ H ₈ O ₂)				57.7	≈48.9	65.0	1.97	≈2.32	40.55	≈34.5	1.16-1.20	≈1.49
Spark or proportional chamber ^l				0.05%		0.03%	-	0.073		2.7%	0.046	-
Shielding concrete ^m				64.9	26.0	32.2	1.70	4.25	26.7	10.7	2.5	-
CO ₂ ⁿ				60.4	33800	46000	1.82	0.0033	36.2	20210	(1.79) ⁿ	(410) ⁿ
Freon 12 (CCl ₂ F ₂) ⁿ				68.1	13800	20200	1.64	0.0081	23.7	4810	(4.93) ⁿ	(1080) ⁿ
Freon 13 (CClF ₃) ⁿ				66.0	15000	21400	1.70	0.0072	27.15	6380	(4.26) ⁿ	(720) ⁿ
Silica Aerogel ^o				62.3	≈311	430	1.82	≈0.36	30	≈150	0.1-0.3	1.0+0.25 ρ

* Table revised January 1976 by J. Engler and F. Mönning. For details and references, see CERN NP Internal Report 74-1.

a) σ of neutrons ($\approx \sigma$ of protons) at 20 GeV from Landolt-Börnstein, New Series I, Vol. 5. Energy dependence for all nuclei $\approx 1/2$ percent/GeV (from 5-25 GeV).

b) $L_{coll} = A/(N\sigma)$. In the absorption length the elastic scattering is subtracted.

c) For a minimum-ionizing, singly-charged particle in the material. From W.H. Barkas and M.J. Berger, Tables of Energy Losses and Ranges of Heavy Charged Particles, NASA-SP-3013 (1964).

d) From Y.S. Tsai, Pair Production and Bremsstrahlung of Charged Leptons, SLAC-PUB-1365 (1974), Table III.6.

e) Values for solids, or the liquid phase at boiling point, except where noted. Values in parentheses for gaseous phase STP (0°C, 1 atm.), except where noted.

f) Density variable.

g) Gas at 20°C.

h) Density may vary about $\pm 3\%$, depending on operating conditions.

i) Values for typical working condition with H₂ target: 50 mole percent, 29°K, 7 atm.

j) Values for typical chamber working conditions: Propane $\sim 57^\circ\text{C}$, 8-10 atm. Freon 13B1 $\sim 28^\circ\text{C}$, 8-10 atm.

k) Typical scintillator; e.g. PILOT B and NE 102A have an atomic ratio H/C = 1.10.

l) Values for typical construction: 2 layers 50 μm Cu/Be wires, 8 mm gap, 60% argon, 40% isobutane or CO₂; 2 layers 50 μm Mylar/Aclar foils.

m) Standard shielding blocks, typical composition O₂ 52%, Si 32.5%, Ca 6%, Na 1.5%, Fe 2%, Al 4% plus reinforcing iron bars. Attenuation length $l = 115 \pm 5$ g/cm², also valid for earth (typical $\rho = 2.15$) from CERN-LRL-RHEL Shielding exp. UCRL 17841 (1968).

n) Used in Čerenkov counters, value at 26°C and 1 atm. Indices of refraction from E.R. Hayes, R.A. Schluter, and A. Tamosaitis, ANL-6916 (1964).

o) $n(\text{SiO}_2) + 2n(\text{H}_2\text{O})$ used in Čerenkov counters, ρ = density in g/cm³. From M. Cantin et al., Nucl. Instr. Meth. 118, 177 (1974).

ELECTROMAGNETIC RELATIONS

Maxwell's Equations

Quantity	CGS (statcoul., statamp., sec cm ⁻¹)	MKSA (coul., amp., ohm)
Potentials:	$V = \frac{\sum \text{charges}}{r}$, $\vec{A} = \frac{1}{c} \sum \text{currents} \frac{\vec{I}}{r}$; c = speed of light in vacuum	$V = \frac{1}{4\pi\epsilon_0} \sum \text{charges} \frac{q}{r}$, $\vec{A} = \frac{\mu_0}{4\pi} \sum \text{currents} \frac{\vec{I}}{r}$; $\epsilon_0 = \frac{1}{36\pi} 10^{-9}$ MKSA, $\mu_0 = 4\pi 10^{-7}$ MKSA
Fields:	$\vec{E} = -\vec{\nabla}V, \vec{B} = \vec{\nabla} \times \vec{A}$	$\vec{E} = -\vec{\nabla}V, \vec{B} = \vec{\nabla} \times \vec{A}$
Materials:	$\vec{D} = \epsilon \vec{E}, \vec{B} = \mu \vec{H}$	$\vec{D} = \epsilon \vec{E}, \vec{B} = \mu \vec{H}$
Force:	$\vec{F} = q(\vec{E} + \frac{\vec{v}}{c} \times \vec{B})$	$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$
Maxwell:	$\vec{\nabla} \cdot \vec{E} = 4\pi\rho$, $\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$, $\vec{\nabla} \cdot \vec{B} = 0$, $\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$	$\vec{\nabla} \cdot \vec{D} = \rho$, $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$, $\vec{\nabla} \cdot \vec{B} = 0$, $\vec{\nabla} \times \vec{H} = \vec{j} + \frac{\partial \vec{D}}{\partial t}$

Impedances: Alternating Currents (MKSA)

Ohm's law: $V = ZI, V = V_0 e^{i\omega t}$

- Impedance of self-inductance of inductance L : $Z = i\omega L$
- Impedance of a capacitor of capacitance C : $Z = \frac{1}{i\omega C}$
- Impedance of a flat conductor of width w at high frequency:
 $Z = \frac{(1+i)\rho}{w\delta}$;
 $\rho = \text{resistivity} \begin{cases} \sim 1.7 \times 10^{-8} \Omega \text{m for Cu} \\ \sim 2.8 \times 10^{-8} \Omega \text{m for Al} \end{cases}$
 $\delta = \text{effective skin depth}$
 $= \sqrt{\frac{\rho}{\pi\nu\mu}} \approx \frac{6.6 \text{ cm}}{\sqrt{\nu(\text{sec}^{-1})}}$ for Cu

Capacitance C and Inductance L per Unit Length (MKSA)

- For flat plates of width w, separated by $d \ll w$:
 $C = \frac{\epsilon w}{d}$; $L = \mu \frac{d}{w}$
- For coax cable of interior and exterior radii r_1 and r_2 :
 $C = \frac{2\pi\epsilon}{\ln(r_2/r_1)}$; $L = \frac{\mu}{2\pi} \ln(r_2/r_1)$;
 $\epsilon = \text{dielectric constant} \begin{cases} 2 \text{ to } 6 \text{ for plastics} \\ 4 \text{ to } 8 \text{ for porcelain, glasses} \end{cases}$
 $\mu = \text{magnetic susceptibility}$

Transmission Lines (No Loss) (MKSA)

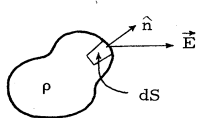
Velocity = $1/\sqrt{LC} = 1/\sqrt{\mu\epsilon}$
Impedance = $\sqrt{L/C}$
L, C are inductance and capacitance per unit length

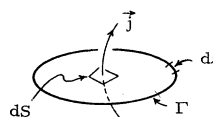
Synchrotron Radiation (CGS)

Energy loss/revolution = $\frac{4\pi}{3} \frac{e^2}{\rho} \beta^3 \gamma^4$, $\rho = \text{orbit radius}$.
For electrons ($\beta \approx 1$), $\frac{\Delta E}{\text{rev}} (\text{MeV}) = 0.0885 [E(\text{GeV})]^4 / \rho(\text{meter})$.
Critical frequency: $\omega_c = 3\gamma^3 c/\rho$

See J. D. Jackson, Classical Electrodynamics, 2nd Ed., John Wiley & Sons, New York, 1975, for more formulae and details. (Prepared April 1974; revised April 1976.)

Integral Forms (MKSA)

1. Gauss' theorem:
 $\int_{\text{surface}} \vec{E} \cdot \hat{n} dS = \int_{\text{volume}} \rho/\epsilon_0 dV$
($\rho = \text{charge/volume}$)

2. Ampere's law:
 $\int_{\Gamma} \vec{H} \cdot d\vec{l} = \int_{\text{surface}} \vec{j} \cdot \hat{n} dS$
($\vec{j} = \text{surface current density}$)

RADIOACTIVITY AND RADIATION PROTECTION

Unit of activity = Curie:
1 Ci = 3.7×10^{10} disintegrations/sec
Unit of exposure dose for x and γ radiation = Roentgen:
1 R = $1 \text{ esu/cm}^2 = 87.8 \text{ erg/g} (5.49 \times 10^7 \text{ MeV/g})$ of air
Unit of absorbed dose = rad:
1 rad = $100 \text{ erg/g} (6.25 \times 10^7 \text{ MeV/g})$ in any material
Unit of dose equivalent (for protection) = rem:
rems (Roentgen equivalents for man) = rads \times QF,
where QF (quality factor) depends upon the type of radiation and other factors. For γ rays and HE protons, QF \approx 1; for thermal neutrons, QF \approx 3; for fast neutrons, QF ranges up to 10; and for α particles and heavy ions, QF ranges up to 20.
Maximum permissible occupational dose for the whole body:
5 rem/year (or \approx 100 millirem/week)
Fluxes (per cm²) to liberate 1 rad in carbon:
 3.5×10^7 minimum ionizing singly charged particles
 1.0×10^9 photons of 1 MeV energy
(These fluxes are correct to within a factor of 2 for all materials.)
Natural background: 120 to 130 millirem/year
cosmic radiation (charged particles + neutrons) ~ 25
cosmic radiation (γ rays) ~ 25
radiation from rocks and air (γ rays) ~ 73 } mrem/yr
Cosmic ray background in counters: $\sim 1/\text{min/cm}^2/\text{ster}$

C.M. ENERGY AND MOMENTUM VS. BEAM MOMENTUM

$$E_{cm} dE_{cm} = m_p dT_{beam} = m_p v_{beam} dP_{beam} \approx m_p dP_{beam}$$

PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)				PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)				PBEAM (MEV/C)	C.M. ENERGY (MEV)				-MOMENTUM IN C.M.- (MEV/C)			
	YP ep	TP	Kp	PP	YP ep	TP	Kp	PP		YP ep	TP	Kp	PP	YP ep	TP	Kp	PP		YP ep	TP	Kp	PP	YP ep	TP	Kp	PP
0	939	1078	1432	1877	0	0	0	0	1500	1922	1930	2022	2254	732	729	696	624	3.0	2.56	2.61	2.77	1.10	1.08	1.02		
20	958	1078	1432	1877	20	17	13	10	1520	1932	1940	2031	2261	738	735	702	631	3.2	2.63	2.68	2.83	1.14	1.12	1.06		
40	977	1083	1433	1877	38	35	26	20	1540	1942	1950	2039	2268	744	741	709	637	3.4	2.70	2.75	2.89	1.18	1.16	1.10		
60	996	1089	1434	1877	56	52	39	30	1560	1951	1959	2048	2275	750	747	715	643	3.6	2.77	2.82	2.96	1.22	1.20	1.14		
80	1015	1096	1436	1878	74	68	52	40	1580	1961	1969	2057	2282	756	753	721	650	3.8	2.83	2.88	3.02	1.26	1.24	1.18		
					T(P1) = PBEAM - 59 MEV																					
100	1033	1105	1439	1879	91	85	65	50	1600	1970	1978	2065	2289	762	759	727	656	4.0	2.90	2.95	3.08	1.29	1.27	1.22		
120	1051	1116	1441	1880	107	101	78	60	1620	1980	1988	2074	2296	768	765	733	662	4.2	2.96	3.01	3.14	1.33	1.31	1.26		
140	1069	1127	1445	1882	123	117	91	70	1640	1989	1997	2083	2304	773	770	739	668	4.4	3.03	3.07	3.19	1.36	1.34	1.29		
160	1087	1139	1449	1883	138	132	104	80	1660	1999	2006	2091	2311	779	776	745	674	4.6	3.09	3.13	3.25	1.40	1.38	1.33		
180	1104	1152	1453	1885	153	147	116	90	1680	2008	2016	2100	2318	785	782	751	680	4.8	3.15	3.19	3.31	1.43	1.41	1.36		
					T(P1) = PBEAM - 92 MEV																					
200	1121	1165	1457	1887	167	161	129	99	1700	2018	2025	2109	2325	791	788	756	686	5.0	3.21	3.25	3.36	1.46	1.44	1.40		
220	1137	1178	1462	1889	182	175	141	109	1720	2027	2034	2117	2332	796	793	762	692	5.2	3.26	3.30	3.40	1.50	1.48	1.44		
240	1154	1192	1468	1892	195	189	153	119	1740	2036	2043	2126	2339	802	799	768	698	5.4	3.31	3.35	3.45	1.54	1.52	1.48		
260	1170	1206	1474	1894	209	202	166	129	1760	2045	2053	2134	2346	807	805	774	704	5.6	3.36	3.40	3.50	1.58	1.56	1.52		
280	1186	1219	1480	1897	222	215	178	138	1780	2054	2062	2143	2353	813	810	779	710	5.8	3.41	3.45	3.55	1.62	1.60	1.56		
					T(P1) = PBEAM - 107 MEV																					
300	1201	1233	1486	1900	234	228	189	148	1800	2064	2071	2151	2360	818	816	785	716	6.0	3.46	3.50	3.60	1.66	1.64	1.60		
320	1217	1247	1493	1903	247	241	201	158	1820	2073	2080	2159	2367	824	821	791	721	6.2	3.51	3.55	3.65	1.70	1.68	1.64		
340	1232	1261	1500	1906	259	253	213	167	1840	2082	2089	2168	2374	829	827	796	727	6.4	3.56	3.60	3.70	1.74	1.72	1.68		
360	1247	1274	1507	1910	271	265	224	177	1860	2091	2098	2176	2381	835	832	802	733	6.6	3.61	3.65	3.75	1.78	1.76	1.72		
380	1262	1288	1514	1913	282	277	235	186	1880	2100	2107	2184	2388	840	837	808	739	6.8	3.66	3.70	3.80	1.82	1.80	1.76		
					T(P1) = PBEAM - 115 MEV																					
400	1277	1302	1522	1917	294	288	247	196	1900	2108	2115	2193	2395	845	843	813	744	7.0	3.71	3.75	3.85	1.86	1.84	1.80		
420	1292	1315	1530	1921	305	300	258	205	1920	2117	2124	2201	2402	851	848	818	750	7.2	3.76	3.80	3.90	1.90	1.88	1.84		
440	1306	1329	1538	1925	316	311	268	214	1940	2126	2133	2209	2409	856	853	824	756	7.4	3.81	3.85	3.95	1.94	1.92	1.88		
460	1320	1342	1546	1929	327	322	279	224	1960	2135	2142	2217	2416	861	859	829	761	7.6	3.86	3.90	4.00	1.98	1.96	1.92		
480	1335	1356	1554	1933	337	332	290	233	1980	2144	2150	2226	2423	867	864	835	767	7.8	3.91	3.95	4.05	2.02	2.00	1.96		
					T(P1) = PBEAM - 120 MEV																					
500	1349	1369	1563	1938	348	343	300	242	2000	2153	2159	2234	2430	872	869	840	772	8.0	3.96	4.00	4.10	2.06	2.04	2.00		
520	1362	1382	1572	1943	358	353	310	251	2020	2161	2168	2242	2437	877	874	845	778	8.2	4.01	4.05	4.15	2.10	2.08	2.04		
540	1376	1395	1580	1947	368	363	321	260	2040	2170	2176	2250	2444	882	879	851	783	8.4	4.06	4.10	4.20	2.14	2.12	2.08		
560	1390	1408	1589	1952	378	373	331	269	2060	2179	2185	2258	2451	887	885	856	789	8.6	4.11	4.15	4.25	2.18	2.16	2.12		
580	1403	1421	1598	1957	388	383	341	278	2080	2187	2194	2266	2458	892	890	861	794	8.8	4.16	4.20	4.30	2.22	2.20	2.16		
					T(P1) = PBEAM - 123 MEV																					
600	1416	1434	1607	1962	397	393	350	287	2100	2196	2202	2274	2465	897	895	866	799	9.0	4.21	4.25	4.35	2.26	2.24	2.20		
620	1430	1446	1616	1967	407	402	360	296	2120	2204	2211	2282	2472	902	900	872	805	9.2	4.26	4.30	4.40	2.30	2.28	2.24		
640	1443	1459	1625	1973	416	412	370	304	2140	2213	2219	2290	2479	907	905	877	810	9.4	4.31	4.35	4.45	2.34	2.32	2.28		
660	1456	1472	1634	1978	425	421	379	313	2160	2221	2227	2298	2486	912	910	882	815	9.6	4.36	4.40	4.50	2.38	2.36	2.32		
680	1468	1484	1644	1984	434	430	388	322	2180	2230	2236	2306	2493	917	915	887	821	9.8	4.41	4.45	4.55	2.42	2.40	2.36		
					T(P1) = PBEAM - 125 MEV																					
700	1481	1496	1653	1989	443	439	397	330	2200	2238	2244	2314	2500	922	920	892	826	10.0	4.46	4.50	4.60	2.46	2.44	2.40		
720	1496	1509	1662	1995	452	448	406	339	2220	2246	2253	2322	2507	927	925	897	831	10.2	4.51	4.55	4.65	2.50	2.48	2.44		
740	1506	1521	1671	2001	461	457	415	347	2240	2255	2261	2330	2514	932	930	902	836	10.4	4.56	4.60	4.70	2.54	2.52	2.48		
760	1519	1533	1681	2007	470	465	424	355	2260	2263	2269	2338	2520	937	934	907	841	10.6	4.61	4.65	4.75	2.58	2.56	2.52		
780	1531	1545	1690	2013	478	474	433	364	2280	2271	2277	2346	2527	942	939	912	846	10.8	4.66	4.70	4.80	2.62	2.60	2.56		
					T(P1) = PBEAM - 127 MEV																					
800	1543	1557	1699	2019	486	482	442	372	2300	2280	2286	2353	2534	947	944	917	852	11.0	4.71	4.75	4.85	2.66	2.64	2.60		
820	1555	1569	1709	2025	495	490	450	380	2320	2288	2294	2361	2541	951	949	922	857	11.2	4.76	4.80	4.90	2.70	2.68	2.64		
840	1567	1580	1718	2031	503	499	459	388	2340	2296	2302	2369	2548	956	954	927	862	11.4	4.81	4.85	4.95	2.74	2.72	2		

PERIODIC TABLE OF THE ELEMENTS

IA		IIA				1 H 1.008						IIIA		IVA		VA		VIA		VIIA		2 He 4.0026	
3 Li 6.94	4 Be 9.0122											5 B 10.81	6 C 12.011	7 N 14.0067	8 O 15.9994	9 F 18.9984	10 Ne 20.17						
11 Na 22.9898	12 Mg 24.305											13 Al 26.9815	14 Si 28.086	15 P 30.9738	16 S 32.06	17 Cl 35.453	18 Ar 39.948						
19 K 39.09	20 Ca 40.08	21 Sc 44.956	22 Ti 47.90	23 V 50.941	24 Cr 51.996	25 Mn 54.9380	26 Fe 55.847	27 Co 58.9332	28 Ni 58.71	29 Cu 63.55	30 Zn 65.38	31 Ga 69.72	32 Ge 72.59	33 As 74.9216	34 Se 78.96	35 Br 79.904	36 Kr 83.80						
37 Rb 85.47	38 Sr 87.62	39 Y 88.906	40 Zr 91.22	41 Nb 92.906	42 Mo 95.94	43 Tc 98.9062	44 Ru 101.07	45 Rh 102.906	46 Pd 106.4	47 Ag 107.868	48 Cd 112.40	49 In 114.82	50 Sn 118.69	51 Sb 121.75	52 Te 127.60	53 I 126.9045	54 Xe 131.30						
55 Cs 132.905	56 Ba 137.34	57-71 Rare Earths	72 Hf 178.49	73 Ta 180.947	74 W 183.85	75 Re 186.2	76 Os 190.2	77 Ir 192.22	78 Pt 195.09	79 Au 196.967	80 Hg 200.59	81 Tl 204.37	82 Pb 207.2	83 Bi 208.981	84 Po (209)	85 At (210)	86 Rn (222)						
87 Fr (223)	88 Ra 226.025	89- Acti- nides	104 (261)	105 (262)	106 (263)																		

57 La 138.91	58 Ce 140.12	59 Pr 140.908	60 Nd 144.24	61 Pm (145)	62 Sm 150.4	63 Eu 151.96	64 Gd 157.25	65 Tb 158.925	66 Dy 162.50	67 Ho 164.930	68 Er 167.26	69 Tm 168.934	70 Yb 173.04	71 Lu 174.97
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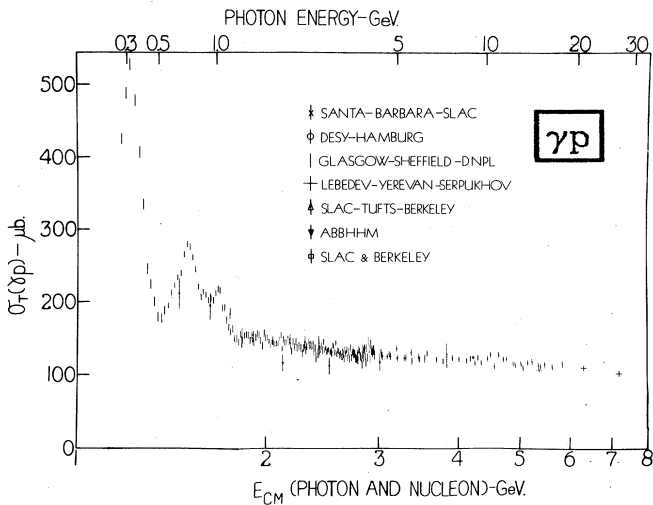
Rare earths (Lanthanide series).

89 Ac (227)	90 Th 232.038	91 Pa 231.036	92 U 238.03	93 Np 237.048	94 Pu (244)	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (254)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (260)
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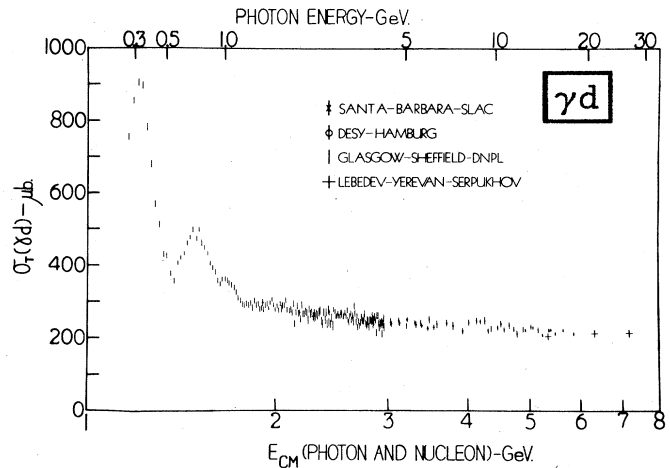
Actinide series

Numbers in parentheses are mass numbers of most stable isotope of that element. Adapted from the *Handbook of Chemistry and Physics, 1975-76.* (Particle Data Group update, April 1976.)

CROSS SECTION PLOTS

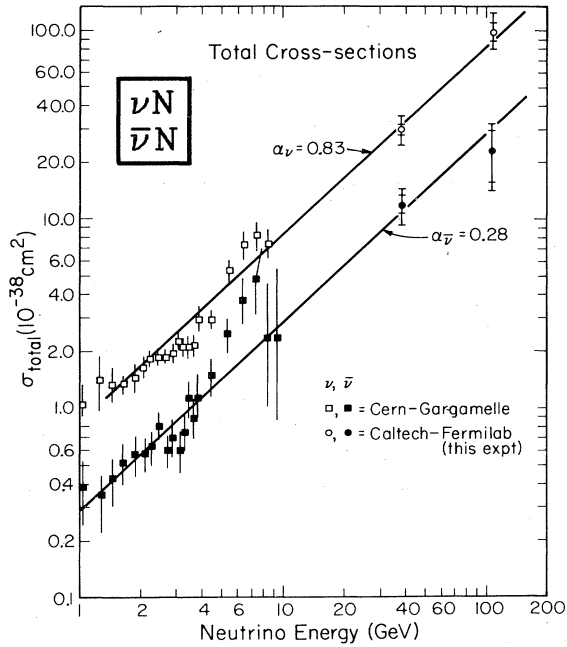


γp total cross section versus photon energy (top scale) and photon-plus-nucleon total center-of-mass energy (lower scale). Courtesy Gething M. Lewis, Glasgow.

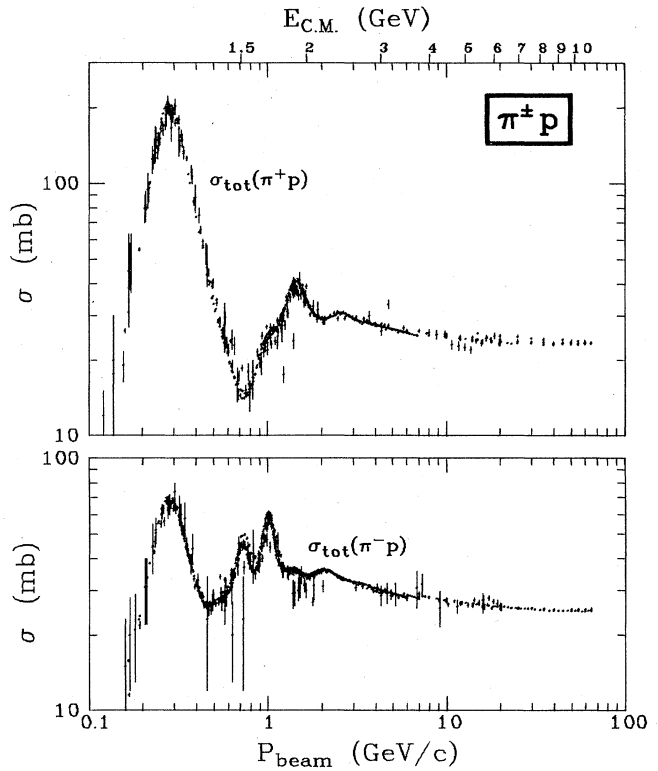


γd total cross section versus photon energy (top scale) and photon-plus-single-nucleon total center-of-mass energy (lower scale). Courtesy Gething M. Lewis, Glasgow.

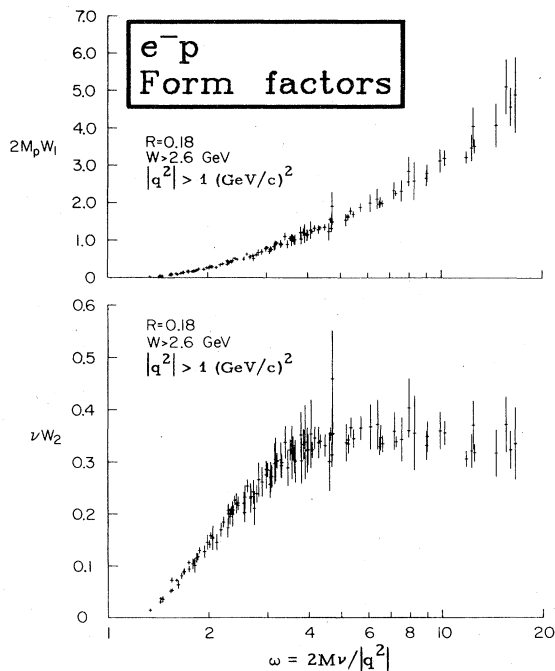
CROSS SECTION PLOTS (Cont'd)



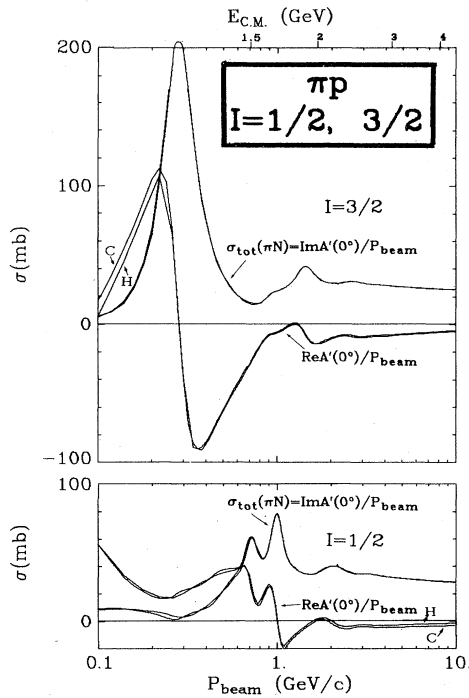
Neutrino and anti-neutrino total cross sections per nucleon [reproduced, with permission, from B.C. Barish et al., Phys. Rev. Lett. 35, 1316 (1975)]. The data are consistent with the form $\sigma_{tot} = \alpha E_{\nu, \bar{\nu}}$.



$\pi^\pm p$ total cross-section data from the Particle Data Group compilation "nN Two-Body Scattering Data," LBL-63 (1973).

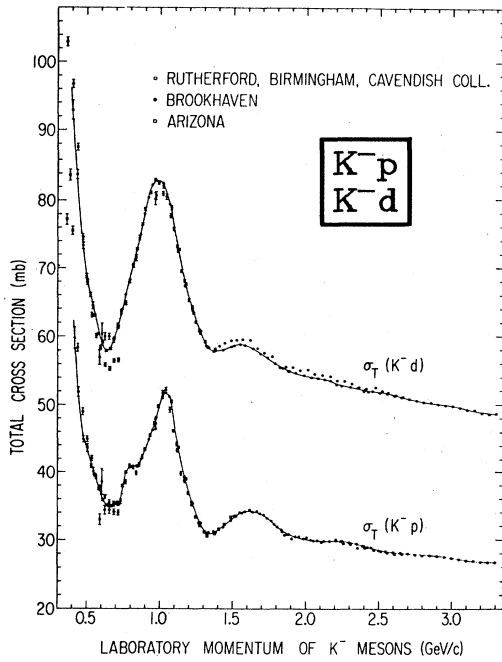


Inelastic electroproduction form factors for a virtual space-like photon scattering on a proton, from SLAC-MIT, Phys. Rev. D5, 528 (1972). $R = \sigma_S/\sigma_T$, the ratio of the longitudinal to transverse cross sections. See L. Hand, Phys. Rev. 129, 1834 (1963).

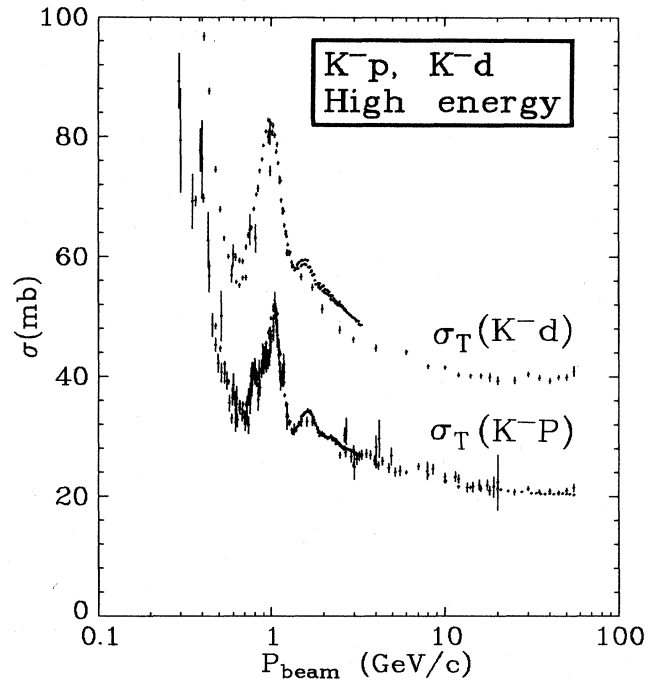


Interpolations of πN total cross sections for $I=3/2$ and $1/2$, and the corresponding real parts of the forward amplitudes as calculated from dispersion relations by A. A. Carter and J. R. Carter (RHEL ppt. RL-73-024, 1973; labeled C above) and by G. Hohler and H. P. Jakob (priv. comm., 1972; labeled H above). The normalization of the curves for each value of I is such that the sum of their squares divided by 19.6 gives $d\sigma/dt$ at 0° in $\text{mb}/(\text{GeV}/c)^2$.

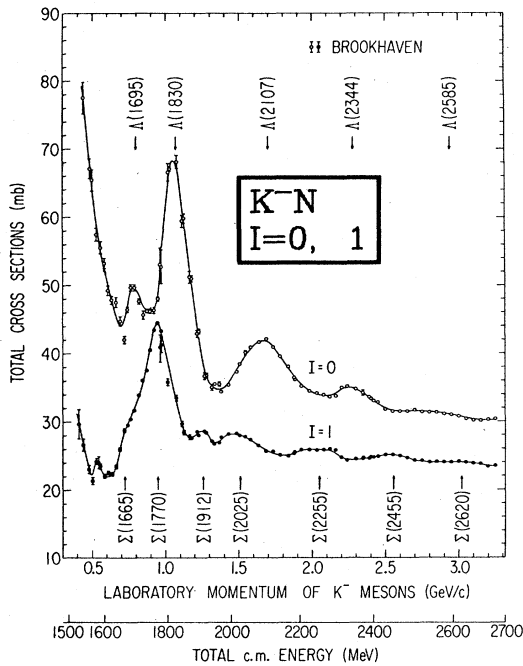
CROSS SECTION PLOTS (Cont'd)



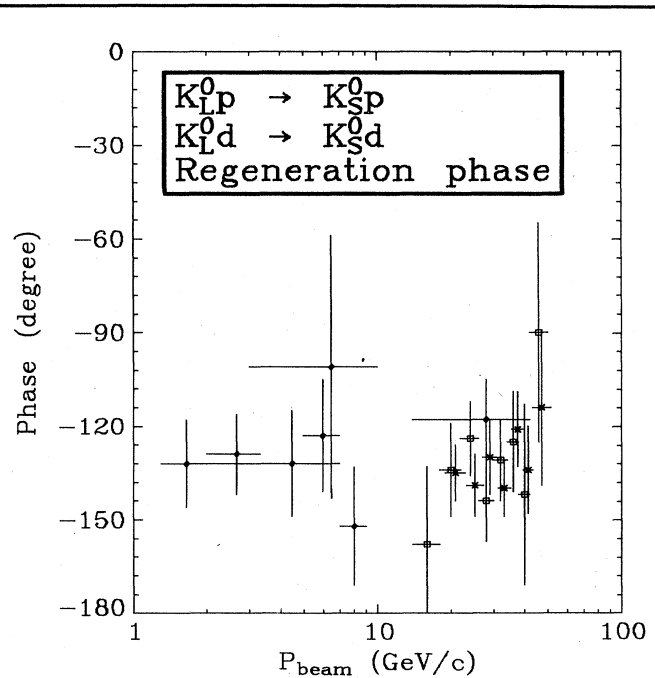
K^-p and K^-d total cross-section data compiled by Li et al., Proc. 1973 Purdue Conf. on Baryon Resonances. The solid curve passes through the Brookhaven data.



K^-p and K^-d total cross-section data. Compilation sources: E. Bracci et al., CERN/HERA 72-2, K^-p ; G.R. Lynch, K^-d (<3 GeV/c); Particle Data Group, K^-d (>3 GeV/c). The new BNL data below 1 GeV/c are not included.

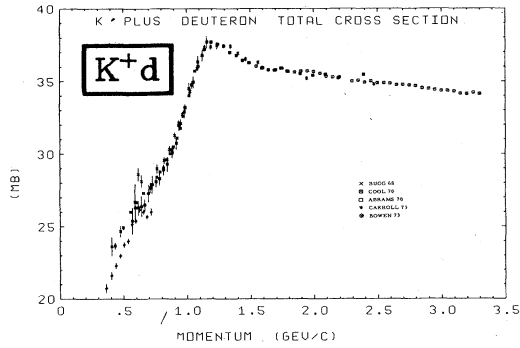
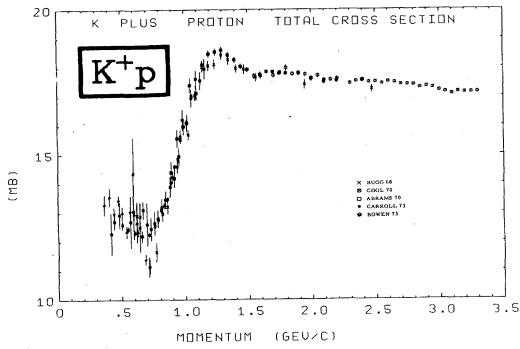


K^-N total cross sections for $I=0$ and $I=1$ below 3.3 GeV/c. Compiled and unfolded by Li et al., Proc. 1973 Purdue Conf. on Baryon Resonances.

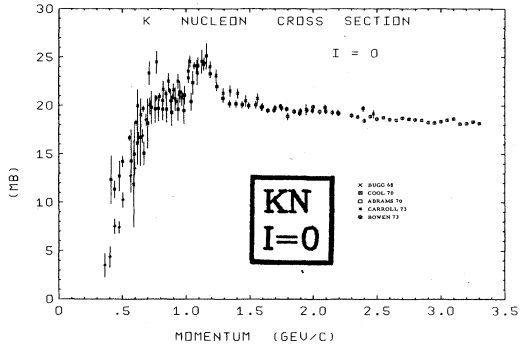


Phases of forward amplitudes for $K_{Lp}^0 \rightarrow K_{Sp}^0$ (\diamond, \square) and $K_{Ld}^0 \rightarrow K_{Sd}^0$ (*). The deuterium data are shifted 1 GeV/c higher to avoid overlap with the proton data. \diamond -p, compiled by Particle Data Group, LBL-55; \square -p, V. Birnlev et al; JINR-E1-6851 (1972); *-d, K.-F. Albrecht et al., Phys. Lett. 48B, 257 (1974).

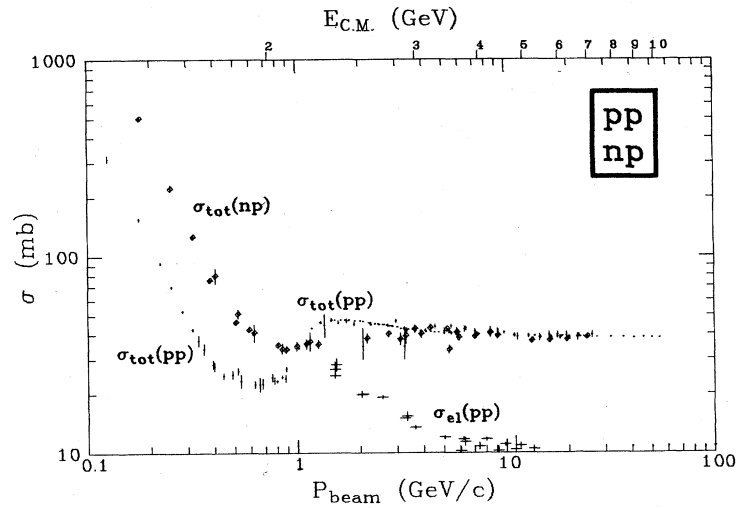
CROSS SECTION PLOTS (Cont'd)



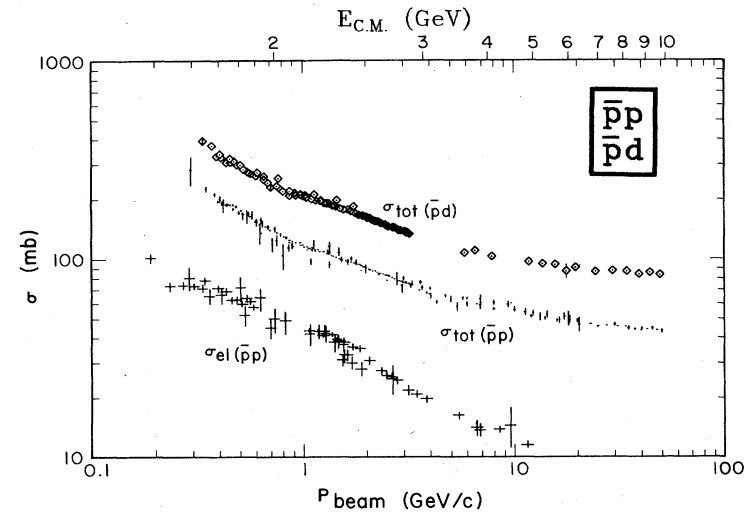
Compilation of recent K^+p and K^+d total cross-section measurements. References can be found in the Baryon Data Card Listings.



Total cross-section for isospin zero KN system. Unfolding of the BUGG 68 and BOWEN 70 and 73 data was done by G. R. Lynch (as in Proc. of 1970 Duke Conference). Tables of σ_0 were provided by the BNL authors. Lynch and BNL use the same method of unfolding; the BOWEN 73 unfolded distribution is obtained by a different method (see plot in Z^0 mini-review in the Baryon Data Card Listings).



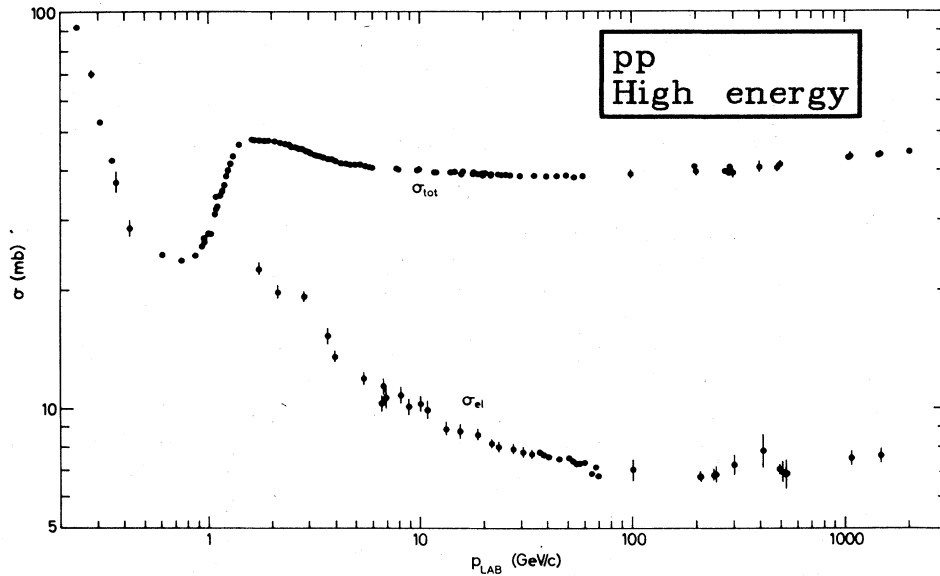
pp and np cross sections from Particle Data Group, "NN and ND Interactions -- A Compilation", UCRL-20 000 NN (August 1970); some points at higher energies added since original compilation.



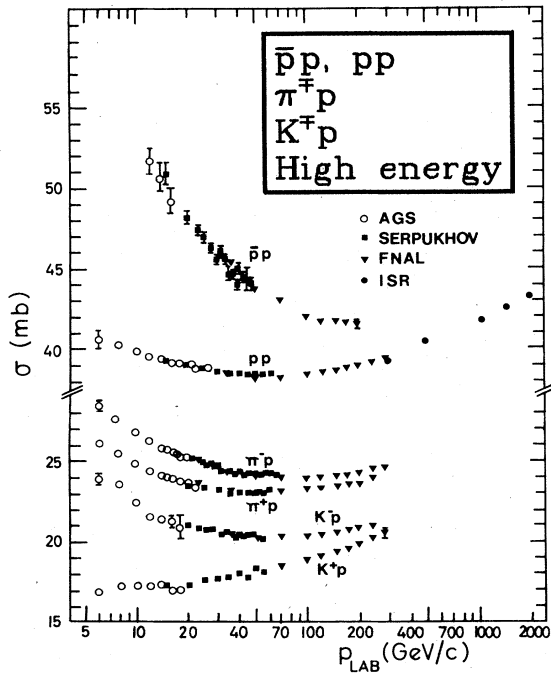
XBL 743-2661

$\bar{p}p$ and $\bar{p}d$ cross sections from Particle Data Group, "A Compilation of NN and ND Reactions", LBL-58 (1972); some points added since original compilation.

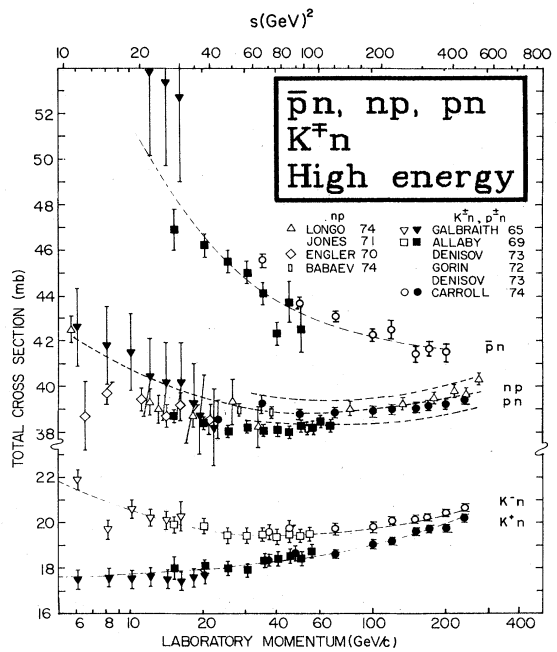
CROSS SECTION PLOTS (Cont'd)



Total and elastic pp cross-section data compiled by U. Amaldi, CERN.



$\bar{p}p$, pp , π^-p , π^+p , K^-p , and K^+p total cross sections versus $s(\approx 2m_p p_{lab})$, as compiled by U. Amaldi, CERN.



$\bar{p}n$, np , pn , K^-n , and K^+n total cross sections versus $s(\approx 2m_p p_{lab})$, as compiled by G. Giacomelli, CERN.

DATA CARD LISTINGS

Illustrative Key

Name of particle as it appears in table. **XX(1200)**

Arrow indicates this particle omitted from table. **74 XX MESON (1200, JPG= -) I=1**

Quantity tabulated below. **74 XX(1200) MASS (MEV)**

Code for quantity tabulated (M=mass, W=width, etc.)

Symbols used to key together data card and related comments.

Number of events above background.

Measured values (parentheses indicate value not used in average).

± Error in measured value (- field blank if error symmetric; parentheses on error only indicate data not used in average due to problems with error estimation).

Average value (and error) of quantity measured.

Vertical bar indicates average; width of horizontal bar on top is error (scaled) in average.

Value and error for each experiment.

Partial decay mode (labeled by P_i).

Branching ratio (labeled by R_j).

Value (and error) of quantity measured, as determined from constrained fit (using all measured branching ratios for this particle).

References listed by year, then author.

Abbreviated reference form used on data cards above.

Journal, report, preprint, etc. (see abbreviations on next page).

Particle name, and quantum numbers (if known).

Particle code (for internal use only).

General comments on particle.

Abbreviated reference for this result; full reference given below.

Measurement technique (see abbreviations on next page.)

Charge(s) of particle detected.

Reaction producing particle, or comments.

Date this result punched (asterisk indicates result added or changed since previous edition).

Scale factor > 1 indicates inconsistent data.

Ideogram to display inconsistent data; curve is sum of Gaussians, one for each experiment (area of Gaussian = 1/error; width of Gaussian = ± error).

Contribution of experiment to χ^2 (if no entry present, experiment not used in calculating χ^2 or scale factor because of large error).

Representative masses of decay products (used for calculating last column of Particle Property Tables).

Author(s).

Quantum number determinations in this reference.

Institution(s) of author(s) (see abbreviations on next page).

REFERENCES FOR XX(1200)

A. MERRILL (TORINO+CERN)/TJP
 B. LYNCH (BNL)
 N. PIERCE (LRL)
 D. FENNER, B. BEANE (NYSE+AMEX)
 J. SMITH (SLAC)

Illustrative Key (cont'd)

Abbreviations

Journals

APAH	Acta Phys. Acad. Hungarica
ADVP	Advances in Physics
ANP	Annals of Physics
ARNNS	Annual Reviews of Nuclear Science
BAPS	Bulletin of the Amer. Phys. Soc.
JETP	English Transl. of Soviet Physics JETP
JETPL	Letters of Soviet Physics JETP
LNC	Letters to Nuovo Cimento
NC	Nuovo Cimento
NP	Nuclear Physics
PL	Physics Letters
PN	Particles and Nuclei
PPSL	Proc. of the Phys. Soc. of London
PR	Physical Review
PRL	Physical Review Letters
PRML	Proc. of the Royal Soc. of London
RMI	Reviews of Modern Physics
SJNP	Soviet Journal of Nuclear Physics
ZPHV	Zeitschrift für Physik

Measurement techniques

ASPK	Automatic spark chambers	HYBR	Hybrid: BC + electronics
CC	Cloud chamber	IPWA	Energy-independent PWA
CNTR	Counters	MMS	Missing mass spectrometer
DASP	DESY double-arm spectrometer	MPWA	Model-dependent PWA
DBC	Deuterium bubble chamber	OSPK	Optical spark chamber
DPWA	Energy-dependent PWA	PBC	Propane bubble chamber
ELEC	Electronic combination	PLUT	DESY PLUTO detector
EMUL	Emulsions	PWA	Partial-wave analysis
PBC	Freon bubble chamber	RVUE	Review of previous data
FRAB	ADONE BB Group detector	SMAG	SPEAR magnetic detector
FRAG	ADONE YY Group detector	SPEC	Spectrometer
FRAM	ADONE MEA Group detector	SPRK	Spark chamber
HBC	Hydrogen bubble chamber	STRC	Streamer chamber
HERC	Helium bubble chamber	WIRE	Wire chamber
HLBC	Heavy liquid bubble chamber	XEBC	Xenon bubble chamber

Conferences

Conferences are referred to by the location in which they were held (e.g., DUBNA, BOULDER, LUND, etc.).

Institutions

AACH	TECHNISCHE UNIV. AACHEN	AACHEN, GERMANY	LUND	UNIV. I LUND	LUND, SWEDEN
AARH	AARHUS UNIV.	AARHUS, DENMARK	MADR	JUNTA DE ENERGIA NUCLEAR	MADRID, SPAIN
AERE	ATOMIC ENERGY RES. ESTAB.	HARWELL, BERKS., ENGLAND	MADU	UNIV. AUTONOME DE MADRID	MADRID, SPAIN
ALBA	STATE UNIV. OF NEW YORK AT ALBANY	ALBANY, N. Y., USA	MANH	MANHATTAN COLLEGE	NEW YORK, N. Y., USA
AMST	UNIV. OF AMSTERDAM	AMSTERDAM, NETHERLANDS	MANZ	UNIV. MAINZ	MAINZ, GERMANY
ANKA	MIDDLE EAST TECHNICAL UNIV.	ANKARA, TURKEY	MASA	UNIV. OF MASSACHUSETTS	AMHERST, MASS., USA
ANL	ARGONNE NAT. LAB.	ARGONNE, ILL., USA	MASH	UNIV. OF MASSACHUSETTS	BOSTON, MASS., USA
ARIZ	UNIV. OF ARIZONA	TUCSON, ARIZ., USA	MCGI	MCGILL UNIV.	MONTREAL, CANADA
ATEA	NUCLEAR RES. CENTRE DEMOKRITOS	ATHENS, GREECE	MCHS	UNIV. MANCHESTER	MANCHESTER, ENGLAND
ATHU	UNIV. OF ATHENS	ATHENS, GREECE	MCHS	MOUNT HOLYOKE COLL.	SOUTH HADLEY, MASS., USA
BARC	UNIV. DE BARCELONA	BARCELONA, SPAIN	MICH	UNIV. OF MICHIGAN	ANN ARBOR, MICH., USA
BARI	UNIV. DI BARI	BARI, ITALY	MILA	UNIV. DI MILANO	MILANO, ITALY
BELG	INST. INTERUNIV. DES SCI. NUC.	BRUXELLES, BELGIUM	MINN	UNIV. OF MINNESOTA	MINNEAPOLIS, MINN., USA
BERG	FYSISK INSTITUTT	BERGEN, NORWAY	MIDA	MICHIGAN STATE UNIV.	OXFORD, OHIO, USA
BERL	INST. HOCHENERGIEPHYS. CAW	ZELTEN/BERLIN, DDR	MIEI	MASSACHUSETTS INST. OF TECHNOLOGY	CAMBRIDGE, MASS., USA
BERN	UNIV. OF BERNE	BERN, SWITZERLAND	MODE	ISTITUTO DI FISICA DELLA UNIVERSITA	MONTREAL, CANADA
BGNA	UNIV. DI BOLOGNA	BOLOGNA, ITALY	MCNL	UNIV. DE MONTREAL	MONTREAL, CANADA
BING	STATE UNIV. OF NEW YORK AT BINGHAMTON	BINGHAMTON, N. Y., USA	MOND	UNIV. DE MONTPELLIER	MONTPELLIER, FRANCE
BIRM	SIRKINGHAM UNIV.	BIRMINGHAM, ENGLAND	MONT	UNIV. OF MONTPELLIER	MONTPELLIER, FRANCE
BNL	BROOKHAVEN NATIONAL LAB.	UPTON, L.I., N. Y., USA	MPH	MAX-PLANCK-INST. FUR PHYS.-ASTROPHYS.	MUNICH, GERMANY
BNHR	NIELS BOHR INSTITUTE	COPENHAGEN, DENMARK	MPHM	MAX-PLANCK-INST. FUR PHYS.-ASTROPHYS.	MUNICH, GERMANY
BONN	UNIV. OF BONN	BONN, GERMANY	MSNA	INS. DI FISICA DELL'UNIV.	MESSINA, ITALY
BORD	UNIV. DE BORDEAUX	BORDEAUX, FRANCE	MSU	MICHIGAN STATE UNIV.	EAST LANSING, MICH., USA
BOST	BOSTON UNIV.	BOSTON, MASS., USA	MUNICH	UNIV. OF MUNICH	MUNICH, GERMANY
BRAN	BRANDEIS UNIV.	WALTHAM, MASS., USA	NAGC	NAGOYA UNIV.	NAGOYA, JAPAN
BRIS	H. H. MILLS PHYS. LAB., U. OF BRISTOL	BRISTOL, ENGLAND	NAPL	UNIV. DI NAPOLI	NAPOLI, ITALY
BROW	BROWN UNIV.	PROVIDENCE, R. I., USA	NDAM	UNIV. OF NOTRE DAME	NOTRE DAME, IND., USA
BRUX	UNIV. LIBRE DE BRUXELLES	BRUXELLES, BELGIUM	NEAS	NORTHEASTERN UNIV.	BOSTON, MASS., USA
BUCH	BUCHAREST STATE UNIV.	BUCHAREST, ROMANIA	NEUC	UNIV. DE NEUCHÂTEL	NEUCHÂTEL, SWITZERLAND
BUDA	CENTRAL RESEARCH INSTITUTE OF PHYSICS	BUDAPEST, HUNGARY	NEVIS	UNIV. OF NEW YORK AT YONKERS	YONKERS, N. Y., USA
BUFF	STATE UNIV. OF NEW YORK AT BUFFALO	BUFFALO, N. Y., USA	NIMJ	R. K. UNIV. NIJMEGEN	NIJMEGEN, NETHERLANDS
CAEN	LAB. DE PHYS. CORPUSCULAIRE	CAEN, FRANCE	NORD	NORDISK INS. FOR TEOR. ATOMFYS.	COPENHAGEN, DENMARK
CAMB	AUSTRALIAN NATIONAL UNIV.	CANBERRA, AUSTRALIA	NOVO	UNIV. OF MONTPELLIER	MONTPELLIER, FRANCE
CARL	CARLTON UNIV.	OTTAWA, CANADA	NRES	NAVAL RESEARCH LABORATORY	WASHINGTON, D.C., USA
CARN	CARNEGIE-MELLON UNIV.	PITTSBURGH, PA., USA	NRL	NORTHWESTERN UNIV.	EVANSTON, ILL., USA
CASE	CASE WESTERN RESERVE UNIV.	CLEVELAND, OHIO, USA	NRS	NEW YORK UNIV.	NEW YORK, N. Y., USA
CAVE	CAMBRIDGE UNIV. OF CAMBRIDGE	CAMBRIDGE, ENGLAND	OHIO	OHIO UNIV.	OHIO, USA
CCAC	COMMUNITY COLLEGE OF ALLEGHENY COUNTY	PITTSBURGH, PENN., USA	OREG	UNIV. OF OREGON	OREGON, USA
CDEF	COLLEGE DE FRANCE	CAMBRIDGE, MASS., USA	ORSA	GIAK RIKSRIKSTANAL LAB.	STOCKHOLM, SWEDEN
CEA	EUROPEAN ORG. FOR NUC. RES.	GENÈVE, SWITZERLAND	OSLO	OSLO UNIV.	OSLO, NORWAY
CERN	UNIV. OF CHICAGO	CHICAGO, ILL., USA	OSLU	OSLO STATE UNIV.	OSLO, NORWAY
CHIC	UNIV. OF CHICAGO	CHICAGO, ILL., USA	OXF	OXFORD UNIV.	OXFORD, ENGLAND
CINC	UNIV. OF CINCINNATI	CINCINNATI, OHIO, USA	PAPO	UNIV. DI PADOVA	PADOVA, ITALY
CIT	CALIF. INSTITUTE OF TECHNOLOGY	PASADENA, CALIF., USA	PATR	UNIV. OF PATRAS	PATRAS, GREECE
CNRC	CANADIAN NATIONAL RESEARCH COUNCIL	OTTAWA, CANADA	PENN	UNIV. OF PENNSYLVANIA	PHILADELPHIA, PA., USA
COLD	COLUMBIA UNIV.	NEW YORK, N. Y., USA	PISA	UNIV. DI PISA	PISA, ITALY
COLU	CORNELL UNIV.	ITHACA, N. Y., USA	PITT	UNIV. OF PITTSBURGH	PITTSBURGH, PA., USA
CONN	CORNELL UNIV.	ITHACA, N. Y., USA	PPA	PRINCETON-PENN. PROTON ACCEL.	PRINCETON, N. J., USA
CRAC	INST. FOR NUCLEAR RESEARCH	WARSAW, POLAND	PRG	INSTITUTE OF PHYSICS, CSAV	PRAGUE, CZECHOSLOVAKIA
CUNY	CITY UNIV. OF NEW YORK	NEW YORK, N. Y., USA	PRIN	PRINCETON UNIV.	PRINCETON, N. J., USA
CURI	LABORATOIRE JOLIO-CURIE	PARIS, FRANCE	PURD	PURDUE UNIV.	LAFAYETTE, IND., USA
DARE	CARESBURY NUC. PHYS. LAB.	CARESBURY, ENGLAND	RHEL	WEIZMANN INST. OF SCI.	REHOVOTH, ISRAEL
DART	DARTMOUTH COLLEGE	HANOVER, N. H., USA	RISO	RUTHERFORD HIGH ENERGY LAB.	RUTHERFORD, BERKS., ENGLAND
DESY	DEUTSCHES ELEKTROEN-SYNCH.	HAMBURG, GERMANY	RMS	ROYAL MILITARY COLLEGE OF SCIENCE	SRIYVENHAM, ENGLAND
DUKE	DUKE UNIV.	DURHAM, N. C., USA	ROCH	UNIV. OF ROCHESTER	ROCHESTER, N. Y., USA
DURH	UNIV. OF DURHAM	DURHAM, ENGLAND	ROMA	UNIV. DI ROMA	ROMA, ITALY
DUUC	DURHAM COLLEGE	DURHAM, ENGLAND	RUTG	RUTGERS UNIV.	NEW BRUNSWICK, N. J., USA
EDIN	UNIV. OF EDINBURGH	EDINBURGH, SCOTLAND	SACLAY	CEA, D. ETUDES NUC. SAACLAY	SACLAY, FRANCE
EFT	ENRICO FERMI INST. FOR NUCL. STUDIES	CHICAGO, ILL., USA	SAGA	SAGA UNIV.	SAGA, JAPAN
EPOL	ECOLE YTECHNIQUE	ZURICH, SWITZERLAND	SEAT	SEATTLE PACIFIC COLLEGE	SEATTLE, WASH., USA
ETHZ	SWISS FEDERAL INST. OF TECHNOLOGY	ZURICH, SWITZERLAND	SEB	RESEARCH CENTER SEIBERSDORF	VIENNA, AUSTRIA
FIRZ	UNIV. DI FIRENZE	FIRENZE, ITALY	SERP	INST. OF HIGH EN. PHYS.	SERPUKHOV, USSR
FISK	FISK UNIV.	NASHVILLE, TENN., USA	SETO	SETO UNIV.	SANTA CRUZ, CALIF., USA
FLOR	UNIV. OF FLORIDA	GAINESVILLE, FLA., USA	SFLA	UNIV. OF SOUTH FLORIDA	TAMPA, FLA., USA
FNAL	FERMILAB NATIONAL ACCELERATOR LAB.	BATAVIA, ILL., USA	SHEP	UNIV. OF SHEFFIELD	SHEFFIELD, ENGLAND
FRAS	LAB. NATIONALI DEL FINESTRONE	FRASCATI, ITALY	SIMP	UNIV. OF SIMP	SOUTHAMPTON, ENGLAND
FSU	FLORIDA STATE UNIV.	TALLAHASSEE, FLA., USA	SIN	SWISS INST. FOR NUCLEAR RESEARCH	VILLIGEN, SWITZERLAND
GENA	UNIV. DI GENOVA	GENOVA, ITALY	SLAC	STANFORD LINEAR ACCEL. CENTER	STANFORD, CALIF., USA
GESC	GENERAL ELECTRIC RES. AND DEV. CENTER	SCHENECTADY, N. Y., USA	SOFI	BULGARIAN ACAD. OF SCI.	SOFIA, BULGARIA
GEVO	UNIV. DE GENÈVE	GENÈVE, SWITZERLAND	STAN	STANFORD UNIV.	STANFORD, CALIF., USA
GLAS	UNIV. OF GLASGOW	GLASGOW, SCOTLAND	STEV	STEVENS INST. OF TECH.	HOBOKEN, N. J., USA
GRAZ	UNIV. OF GRAZ	GRAZ, AUSTRIA	STLD	ST. LOUIS UNIV.	ST. LOUIS, MO., USA
GSCD	GEOLOGICAL SURVEY OF CANADA	OTTAWA, CANADA	STOH	STOCKHOLM UNIV.	STOCKHOLM, SWEDEN
HAIF	TECHNION ISRAEL INST. OF TECHNOLOGY	HAIFA, ISRAEL	STON	STATE UNIV. OF NEW YORK AT STONY BROOK	STONY BROOK, L.I., N. Y., USA
HAMB	UNIV. HAMBURG	HAMBURG, GERMANY	STRB	CENTRE DES RES. NUCLEAIRES	STRASBOURG, FRANCE
HARV	HARVARD UNIV.	CAMBRIDGE, MASS., USA	SUSX	SUSSEX UNIV.	SUSSEX, ENGLAND
HAWA	UNIV. OF HAWAII	HONOLULU, HAWAII, USA	SYRA	SYRACUSE UNIV.	SYRACUSE, N. Y., USA
HEID	UNIV. HEIDELBERG	HEIDELBERG, GERMANY	TATA	TATA INST. OF FUNDAMENTAL RESEARCH	BOMBAY, INDIA
HEL5	HELSINGIN YLIOPISTO	HELSINKI, FINLAND	TELA	UNIV. OF TEL-AVIV	TEL-AVIV, ISRAEL
IDM	INTERNATIONAL BUSINESS MACHINES	PALO ALTO, CALIF., USA	TEMP	TEMPLE UNIV.	PHILADELPHIA, PA., USA
IIT	ILLINOIS INST. OF TECH.	CHICAGO, ILL., USA	TENN	UNIV. OF TENNESSEE	KNOXVILLE, TENN., USA
ILL	UNIV. OF ILLINOIS	URBANA, ILL., USA	TORO	UNIV. OF TORONTO	TORONTO, CANADA
ILLC	UNIV. OF ILLINOIS AT CHICAGO	CHICAGO, ILL., USA	TOHO	TOHOKU UNIV.	SENDAI, JAPAN
IND	UNIV. OF INDIANA	BLOOMINGTON, IND., USA	TOYO	UNIV. OF TOKYO	TOKYO, JAPAN
INNS	PHYS. INST., UNIV. INNSBRUCK	INNSBRUCK, AUSTRIA	TORI	UNIV. DI TORINO	TORINO, ITALY
IOWA	UNIV. OF IOWA	IOWA CITY, IOWA, USA	TRST	UNIV. DI TRIESTE	TRIESTE, ITALY
IPN	INST. DE PHYS. NUCLEAIRE	ORSAY, FRANCE	TUFT	TUFTS UNIV.	MEDFORD, MASS., USA
IPNC	INSTITUT DE PHYSIQUE NUCLEAIRE	PARIS, FRANCE	UNC	UNIV. OF BRITISH COLUMBIA	VANCOUVER, CANADA
IRAD	INSTITUTE DU RADIUM	OTTAWA, CANADA	UCB	UNIV. OF CALIF. AT BERKELEY	BERKELEY, CALIF., USA
ISU	IOWA STATE UNIV.	AMES, IOWA, USA	UCI	UNIV. OF CALIF. AT IRVINE	IRVINE, CALIF., USA
ITEP	INST. FOR TECH. AND EXP. PHYS.	AMES, IOWA, USA	UCLA	UNIV. OF CALIF. AT LOS ANGELES	LOS ANGELES, CALIF., USA
IUNP	INDIANA UNIV. - PURDUE U. AT INDIANAPOLIS	INDIANAPOLIS, IND., USA	UCSC	UNIV. OF CALIF. AT SANTA BARBARA	SANTA BARBARA, CALIF., USA
JAGL	JAGELLONIAN UNIV.	CRAKOW, POLAND	UCSD	UNIV. OF CALIF. AT SAN DIEGO	SAN DIEGO, CALIF., USA
JHU	JOHNS HOPKINS UNIV.	BALTIMORE, MD., USA	UNCN	UNIV. COLLEGE	SCHENECTADY, N. Y., USA
JINR	JOINT INST. FOR NUCLEAR RESEARCH	DUBNA, USSR	UNJN	UPSALA COLLEGE	EAST ORANGE, N. J., USA
KANS	UNIV. OF KANSAS	LAWRENCE, KANSAS, USA	UTAH	UNIV. OF UTAH	SALT LAKE CITY, UTAH, USA
KERN	TECHNISCHE HOCHSCHULE KARLSRUHE	KARLSRUHE, GERMANY	VAND	VANDERBILT UNIV.	NASHVILLE, TENN., USA
KEK	NAT. LAB. FOR HIGH ENERGY PHYS., JAPAN	TSUKUBA-GUN., JAPAN	VIRG	VIRGINIA POLYTECHNIC INST.	BLACKSBURG, VA., USA
KIAC	KURCHATOV INST. OF ATOMIC ENERGY	LENINGRAD, USSR	WAR5	UNIV. OF WARSAW	WARSAW, POLAND
KNTY	UNIV. OF KENT	LENINGRAD, USSR	WASH	UNIV. OF WASHINGTON	SEATTLE, WASH., USA
LALC	LINEAR ACCELERATOR LAB. ORSAY	ORSAY, FRANCE	WIEN	UNIV. OF VIENNA	VIENNA, AUSTRIA
LANC	LANCASTER UNIV.	LANCASTER, ENGLAND	WISC	UNIV. OF WISCONSIN	MADISON, WIS., USA
LASL	U. S. LOS ALAMOS SCIENTIFIC LAB.	LOS ALAMOS, N. M., USA	WOOD	WOODSTOCK COLLEGE	WOODSTOCK, MD., USA
LAUS	UNIV. OF LAUSANNE	LAUSANNE, SWITZERLAND	WASH	WASHINGTON UNIV.	ST. LOUIS, MO., USA
LEBD	LEBERDEY PHYSICS INST.	BERKELEY, CALIF., USA	WYOM	UNIV. OF WYOMING	LARAMIE, WYOMING, USA
LEHI	LEHIGH UNIV.	BETHLEHEM, PA., USA	YALE	YALE UNIV.	NEW HAVEN, CONN., USA
LEIT	INST. LORENTZ	LEIDEN, NETHERLANDS	ZEEM	ZEEMAN LAB., UNIV. OF AMSTERDAM	AMSTERDAM, NETHERLANDS
LINZ	LINZ INSTITUT FUR PHYSIK, KEPLER HOCH.	LINZ, AUSTRIA			
LIVP	LIVERPOOL UNIV.	LIVERPOOL, ENGLAND			
LIOC	IMPERIAL COL. OF SCI. AND TECH.	LONDON, ENGLAND			
LQCM	QUEEN MARY COLLEGE	LONDON, ENGLAND			
LOUC	UNIVERSITY COLLEGE	LONDON, ENGLAND			
LOWC	WESTFIELD COLLEGE	LONDON, ENGLAND			
LPMF	LAB. DE PHYS. NUCL. ET HAUTES ENERGIES	PARIS, FRANCE			
LPTP	LAB. DE PHYS. THEOR. ET HAUTES ENERGIES	PARIS, FRANCE			
LRL	U. S. LAWRENCE BERKELEY LAB.	BERKELEY, CALIF., USA			
LSU	LOUISIANA STATE UNIV.	BATON ROUGE, LA., USA			

Stable Particles

$\gamma, \nu_e, \nu_\mu, e, \mu$

Data Card Listings

For notation, see key at front of Listings.

[gamma] 0 GAMMA(0,J=1)

0 GAMMA MASS (IN UNITS OF 10**--21 MEV)
M P (6.) OR LESS PATEL 65 SATELLITE DATA 10/69
M 6. OR LESS GINTSBURG 64 SATELLITE DATA 10/73
M 2.3 OR LESS GOLDHABER 68 SATELLITE DATA 10/69
M F (0.06) OR LESS FRANKEN 71 LOW FREQ RES CIR 3/72
M 10. OR LESS WILLIAMS 71 CNTR TESTS GAUSS LAW 3/71
M 0.73 OR LESS HOLLWEG 74 ALFVEN WAVES 7/74**
M F VALIDITY QUESTIONABLE ACCORDING TO AUTHORS AND KROLL 71. 3/72
M P SEE CRITICISM IN GOLDHABER 71 3/72

REFERENCES FOR GAMMA
GINTSBR 64 SOV. ASTR. AJ7 536 M. A. GINTSBURG (ACAD SCI, USSR)
PATEL 65 PL 14 105 V. L. PATEL (DURHAM)
GOLDHABE 68 PRL 21 567 A. GOLDHABER, M. NIETO (STONY BROOK)
FRANKEN 71 PRL 26 115 P A FRANKEN, G W AMPULSKI (MICH)
WILLIAMS 71 PRL 26 721 +FALLER, HILL (WESLEYAN)
HOLLWEG 74 PRL 32 961 J V HOLLWEG (NATL CENTER FOR ATMOS RESRCH)

PAPERS NOT REFERRED TO IN DATA CARDS
GOLDHABE 71 RMP 43 277 A S GOLDHABER, M M NIETO (STON+BOHR+UCSB)
KROLL 71 PRL 26 1395 N M KROLL (SLAC)

[nu_e] 1 E-NEUTRINO(0,J=1/2)

1 E-NEUTRINO MASS (KEV)
M (0.25) OR LESS LANGER 52 CNTR ANTI-NEUT.(TRITIUM)
M (0.50) OR LESS HAMILTON 53 CNTR ANTI-NEUT.(TRITIUM) 11/73
M (0.55) (0.28) FRIEDMAN 58 CNTR ANTI-NEUT.(TRITIUM)
M 4.1 OR LESS CL=.67 BECK 68 CNTR NEUTRINO SODIUM 22 11/73
M 0.5 OR LESS CL=.90 DARIS 69 CNTR ANTI-NEUT.(TRITIUM) 11/73
M 0.32 OR LESS CL=.90 SALGO 69 CNTR ANTI-NEUT.(TRITIUM) 11/73
M 0.06 OR LESS CL=.90 BERGKVIS 72 CNT ANTI-NEUT.(TRITIUM) 11/73
M (0.008) OR LESS COMSJK 72 THEOR. LIM. FROM COSMOLOGY 3/74
M 0.086 OR LESS CL=.90 RODE 72 CNTR ANTI-NEUT.(TRITIUM) 11/73
M D DARIS 69 VALUE .075 KEV (CL=.67) DISAGREES WITH FIG. 6. WE USE FIG. 6. 11/73
M C 450. OR LESS CL=.90 CLARK 74 ASPK KE3 DECAY 11/75*
M C LOWEST LIMIT FROM STRANGENESS CHANGING DECAY. 11/75*

1 (E-NEUTRINO) - (E-ANTINEUTRINO) MASS DIFF. (KEV)
DM 450. OR LESS CL=.90 CLARK 74 ASPK KE3 DECAY 11/75*

REFERENCES FOR E-NEUTRINO
LANGER 52 PR 88 689 L M LANGER, P J D MOFFAT (INDIANA)
HAMILTON 53 PR 92 1521 D HAMILTON, W P ALFORD, L GROSS (PRINCETON)
FRIEDMAN 58 PR 109 2214 LEWIS FRIEDMAN, LINCOLN G SMITH (SNU)
BECK 68 ZPHY 216 229 F BECK, H DANIEL (MPIH)
DARIS 69 NP A138 545 N. DARIS, C ST-PIERRE (LAVAL-QUEBEC)
SALGO 69 NP A138 417 R C SALGO, H H STAUB (ZURICH)
BERGKVIS 72 NP 839 317 KARL-ERIK BERGKVIST (UNIV STOCKHOLM)
COMSJK 72 PRL 29 669 R COMSJK, J MC CLELLAND (UCB)
RODE 72 LNC 5 139 B RODE, H DANIEL (MUNICH+MPIH)
CLARK 74 PR D9 533 +ELIOFF, FRISCH, JOHNSON, KERTH, SHEN+ (LBL)

[nu_mu] 2 MU-NEUTRINO(0,J=1/2)

2 MU-NEUTRINO MASS (MEV)
M 3.5 OR LESS BARKAS 56 EMUL 7/66
M 4.0 OR LESS DUZIAK 59 CNTR 7/66
M 3.6 OR LESS FEINBERG 63 RVUE 7/66
M 3.0 OR LESS ALLCOCK 65 RVUE 7/66
M 2.5 OR LESS BARDON 65 ASPK 7/74**
M 2.8 OR LESS CL=.90 SHAFER 65 CNTR 5/71
M 1.6 OR LESS CL=.90 BOOTH 67 CNTR 3/68
M 2.2 OR LESS CL=.90 HYMAN 67 HEBB 0. K- HE 11/67
M B M (1.2) OR LESS CL=.90 BACKENSTOSS 71 CNTR M**2=1.28+-1.24 10/71
M S 1.15 OR LESS CL=.90 SHRUM 71 CNTR M**2=1.55+-1.14 12/71
M (8 EV) OR LESS COMSJK 72 THEOR. LIM. FROM COSMOLOGY 3/74
M B M 1.15 OR LESS CL=.90 BACKENSTOSS 73 CNTR M**2=0.29+-0.90 1/73
M 0.65 OR LESS CL=.90 CLARK 74 ASPK KMUS DECAY 7/74**
M D M 1.0 OR LESS CL=.90 DAUM 76 SPEC M**2=0.23+-0.54 1/76*
M M WE CALCULATE UPPER LIMIT AT CL=.90 FROM M**2. 1/76*
M B BACKENSTOSS 73 REPLACES BACKENSTOSS 71 AND USES THEIR NEW PI- MASS. 1/73
M S SHRUM 71 USES SHAFER 67 PI- MASS VALUE AND CRANE 71 MU MASS VALUE. 1/73
M D DAUM 76 USES OUR 1974 PI- AND MU MASSES. 1/76*

2 (MU-NEUTRINO) - (MU-ANTINEUTRINO) MASS DIFF. (MEV)
DM (0.45) OR LESS CL=.90 CLARK 74 ASPK KMUS DECAY 11/75*

REFERENCES FOR MU-NEUTRINO
BARKAS 56 PR 101 778 M B BARKAS, W BIRNBAUM, F M SMITH (LFL)
DUZIAK 59 PR 114 336 W F DUZIAK, R SAGANE, J VEDDER (LRL)
FEINBERG 63 ARNS 13 431 G FEINBERG, L M LEDERMAN (COLUMBIA)
ALLCOCK 65 PPSL 85 875 G R ALLCOCK (LIVERPOOL)
BARDON 65 PRL 14 449 BARDON, NORTON, PEOPLES + (COLU+STONY BROOK)

SHAFER 65 PRL 14 923 R E SHAFER, C CROWE, JENKINS (LFL)
BOOTH 67 PL 268 39 BOOTH, JOHNSON, WILLIAMS, WORMALD (LIVERPOOL)
HYMAN 67 PL 258 376 +LOKEN, PENNITT, MCKENZIE+ (LANL+CARN+NWES)
BACKENSTOSS 71 PL 368 603 BACKENSTOSS, DANIEL+KOOCH+ (CERN+KARL+HEID)
SHRUM 71 PL 378 114 E V SHRUM, K O H ZIOCK (UNIV OF VIRGINIA)
COWSIK 72 PRL 29 669 R COWSIK, J MC CLELLAND (UCB)
BACKENSTOSS 73 PL 438 539 BACKENSTOSS, DANIEL+KOOCH+ (CERN+KARL+MUNICH)
CLARK 74 +ELIOFF, FRISCH, JOHNSON, KERTH, SHEN + (LBL)
DAUM 76 PUBL. IN PL B FEB +DUBAL, EATON, FROSCHE, MCCULLOCH+ (STN, ETHZ)

[e] 3 ELECTRON(0,5,J=1/2)

3 ELECTRON MASS (MEV)
M (.5110061(.00002) COHEN 65 RVUE
M (.5110041(.0000016) TAYLOR 69 RVUE USING NEW E/H 7/70
M .5110034 .0000014 COHEN 73 RVUE

3 ELECTRON MEAN LIFE (UNITS 10**21 YR)
T 2.0 OR MORE MDE 65 CNTR 6/66
T S 5.2 OR MORE STEINBERG 75 CNTR 2/76*
T S STEINBERG 75 SENSITIVE TO ALL DECAY MODES IN WHICH DECAY PARTICLES ESCAPE FROM DETECTOR WITHOUT DEPOSITING ENERGY. TEST OF CHARGE CONSERVATION. 2/76*
T S 2/76*

3 ELECTRON MAGNETIC MOMENT (E/2ME)
MM SEE RICH 72 FOR A REVIEW OF THEORY AND EXPERIMENTS. 3/74
MM (1.0011609) +-(24)*10**--7 SCHUPP 61 CNTR 8/66
MM (1.001159622) +-(27)*10**--9 WILKINSON 63 CNTR - 8/66
MM (1.001168) +-(22)*10**--6 RICH 66 CNTR + POSITRON 8/66
MM R (1.001159525) +-(30)*10**--9 RICH 68 CNTR - 6/68
MM (1.0011596389) +-(31)*10**--10 TAYLOR 69 RVUE 2/71
MM (1.001159644) +-(7)*10**--9 WESLEY 70 CNTR 6/70
MM (1.0011596577) +-(35)*10**--10 WESLEY 71 CNTR - 2/72
MM (1.0011603) +-(12)*10**--7 GILLELAND 72 CNTR + 2/72
MM R 1.001159667 +-(35)*10**--10 COHEN 73 RVUE 3/74
MM R RICH 68 IS REEVALUATION OF WILKINSON 63.

REFERENCES FOR ELECTRON
SCHUPP 61 PR 121 1 A A SCHUPP, R W PIDD, H R CRANE (MICH)
WILKINSON 63 PR 130 852 D T WILKINSON, H R CRANE (MICH)
COHEN 65 RMP 37 537 COHEN, JUMOND (N.A. AVIATION SCI. CENTER+CIT)
MDE 65 PR 140 B 992 M K MDE, F REINES (CASE INST TECHNOLOGY)
RICH 66 PRL 17 271 A RICH, H R CRANE (MICH)
RICH 68 PRL 20 967 A RICH (MICH)
TAYLOR 69 RMP 41 375 +PARKER, LANGENBERG (PRIN+UCI+PENN)
WESLEY 70 PRL 24 1320 J C WESLEY, A RICH (MICH)
WESLEY 71 PR 44 1341 J C WESLEY, A RICH (MICH)
GILLELAND 72 PR 45 38 J GILLELAND, A RICH (MICH)
RICH 72 RMP 44 250 A RICH, J C WESLEY (MICH)
COHEN 73 J. PHYS. CHEM. REF. DATA 2, P. 663, E. R. COHEN, B. N. TAYLOR (UMD)
STEINBERG 75 PR D12 2582 STEINBERG, KWIAKOWSKI, MAENHAUT, WALL (UMD)

[mu] 4 MUON(106,J=1/2)

4 MUON MASS (MEV)
M (105.659) (0.002) FEINBERG 63 RVUE
M (105.6599) (0.0014) TAYLOR 69 RVUE USING NEW E/H 7/70
M C (105.6597) (0.0005) CRANE 72 CNTR INCLUDED IN COHEN73 1/73
M D (105.6594) (0.0006) CRANE 72 CNTR INCLUDED IN COHEN73 2/72
M 105.65948 0.00035 COHEN 73 RVUE 3/74
M C CRANE 71 GIVES MU/ME=206.76878(85). WE USE ME=5110041(16) MEV. 1/73
M D CROWE 72 GIVES MU/ME=206.7682(5) AND USES ME=5110041(16) MEV. 1/73
M FIT 105.65948 0.00035 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*

4 MUON MEAN LIFE (UNITS 10**--6)
T 2.198 0.001 0.001 FARLEY 62 CNTR
T 2.203 0.004 LUNDY 62 CNTR CONLEV=.98 11/67
T 2.202 0.003 0.003 ECKHAUSE 63 CNTR
T 2.197 0.005 0.002 MEYER 63 CNTR +
T 2.198 0.002 0.002 MEYER 63 CNTR - 7/66
T W (2.20206) (0.00081) WILLIAMS 72 CNTR + 2/76*
T 2.1973 0.0003 DUCLOS 75 CNTR + 1/76*
T 2.19711 0.00008 BALANDIN 74 CNTR + 1/76*
T W WILLIAMS 72 MEAN LIFE MEASUREMENT WAS NOT THE PRIMARY PURPOSE OF THEIR EXPERIMENT AND DISAGREES STRONGLY WITH LATER EXPTS. NOT AVGD. 1/76*
T AVG 2.197134 0.000077 0.000077 AVERAGE ERROR INCL. SCALE FACTOR OF 1.01
T STUDENT 2.197133 0.000084 0.000084 AVG. USING STUDENT10(H/1.11) -- SEE TEXT

4 MU+MU- MEAN LIFE RATIO
DT 1.000 0.001 MEYER 63 CNTR MEAN LIFE MU+/MU- 7/66

4 MUON ANOMALOUS MAGN. MOMENT (10**--6*(2*MU MASS))
MM SEE RICH 72 FOR A REVIEW OF THEORY AND EXPERIMENTS. 3/74
MM (1162.0) (5.0) CHARPAK 62 CNTR +
MM B (1165.75) (0.71) BAILEY 68 CNTR + STOR. RINGS 5/69
MM B (1166.25) (0.24) BAILEY 68 CNTR + STOR. RINGS 5/69
MM B ERRORS STATISTICAL. VALUES COMBINED TO GIVE MU- VALUE BELOW 5/69
MM 1166.16 0.31 BAILEY 68 CNTR + STOR. RINGS 5/69
MM 1165.895 0.027 BAILEY 75 CNTR + STORAGE RING 11/75*
MM AVG 1165.897 0.027 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.01
MM STUDENT 1165.897 0.029 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Data Card Listings

Stable Particles

For notation, see key at front of Listings.

μ

4 MUON TO PROTON MAGNETIC MOMENT RATIO

Table listing muon to proton magnetic moment ratio data. Includes columns for experiment name (MMR), description, and numerical values. Key experiments include MMR (3.18651), MMR (3.18331), MMR D (3.18331), and MMR AVG (3.1833100).

4 MUON PARTIAL DECAY MODES

Table listing muon partial decay modes. Columns include mode number (P1-P5), description (e.g., MUON INTO E ANTI(E-NEU)), and branching ratios.

4 MUON BRANCHING RATIOS

Table listing muon branching ratios. Columns include mode number (R1-R4), description (e.g., MUON INTO E+2GAMMA), and numerical values. Includes references to other sections like VI A.

4 MUON DECAY PARAMETERS

RELATED TEXT SECTION VI A

Table listing muon decay parameters. Columns include parameter name (RHO, RHO P, RHO C, etc.), description, and numerical values. Includes sub-sections for RHO parameters, ETA parameters, and XSI parameters.

Table listing various decay parameters and constants. Columns include parameter name (DELTA, DEL, DEL AVG, etc.), description, and numerical values. Includes sub-sections for DELTA parameters, HELICITY OF DECAY ELECTRON, SCALAR COUPLING CONSTANT, AXIAL VECTOR COUPLING CONSTANT, PHASE BETWEEN VECTOR AND AXIAL VECTOR COUPLINGS, TENSOR COUPLING CONSTANT, and PSEUDOSCALAR COUPLING CONSTANT.

REFERENCES FOR MUON

Table listing references for muon. Columns include author names (e.g., COFFIN, BARDOON, GARWIN, etc.) and publication details (e.g., 58 PR 109 973, 58 PRL 1 38, etc.).

Stable Particles

π^\pm, π^0

Data Card Listings
For notation, see key at front of Listings.

π^\pm										
8 CHARGED PION(140,JPG=0--) I=1										
8 CHARGED PION MASS (MEV)										
M	139.37	0.20	CROME	54	CNTR	-				
M	139.68	0.15	BARKAS	56	EMUL	+				
M S	(139.577)	(0.013)	SHAFFER	67	CNTR	-	MESONIC ATOMS	6/68		
M B	(139.549)	(0.008)	BACKENSTO	71	CNTR	-	MESONIC ATOMS	10/71		
M S	139.566	0.013	SHAFFER	72	CNTR	-	MESONIC ATOMS	1/73		
M B	139.569	0.008	BACKENSTO	73	CNTR	-	MESONIC ATOMS	1/73		
M S	SHAFFER 72 UPDATES	SHAFFER 67 WITH NEW ALPHA AND NEW CALIB. LINE ENER.						1/73		
M B	BACKENSTOSS 73	CORRECTS BACKENSTOSS 71 WITH NEW VACUUM POL. CALC.						1/73		
M	AVG	139.5682	0.0068	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)					
M	STUDENT	139.5682	0.0073	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT					
M	FIT	139.5688	0.0064	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.0)				2/76*	
8 (PI+) - (MU+) MASS DIFFERENCE (MEV)										
D	34.00	0.076	BARKAS	56	EMUL					
D	33.89	0.076	BARKAS	56	EMUL					
D	145	33.881	0.035	HYMAN	67	HEBC	+ K-HE		2/71	
D		33.925	0.025	BOOTH	70	CNTR	+ MAGNETIC SPECT.		2/71	
D	AVG	33.915	0.019	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)					
D	STUDENT	33.915	0.021	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT					
D	FIT	33.9093	0.0064	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.0)				2/76*	
8 (PI+) - (PI-)/AVG., MASS DIFFERENCE (PERCENT)										
DM	0.02	0.05	AYRES	71	CNTR				3/71	
8 CHARGED PION MEAN LIFE (UNITS 10**--9)										
T	25.6	0.5	0.5	CROME	57	RVUE				
T	25.6	0.8	0.8	ANDERSON	60	CNTR				
T	8000	25.46	0.32	0.32	ASHKIN	60	CNTR	+		
T		26.02	0.04	ECKHAUSE	65	CNTR	+		9/66	
T		25.6	0.3	BARON	66	CNTR			9/66	
T		25.9	0.3	DUNAITSEV	66	CNTR			6/68	
T	N	(26.40)	(0.08)	KINSEY	66	CNTR	+		6/66	
T	N	SYSTEMATIC ERRORS IN CALIBR. IN THIS EXP. DISCUSSED BY NORDBERG 67							8/67	
T		26.67	0.24	LOBKOWICZ	66	CNTR			9/66	
T		26.04	0.05	NORDBERG	67	CNTR	+		8/67	
T		26.02	0.04	AYRES	71	CNTR	+		3/71	
T		26.09	0.08	DUNAITSEV	73	CNTR	+		3/74	
T	AVG	26.030	0.023	0.023	AVERAGE	(ERROR INCL. SCALE FACTOR OF 1.0)				
T	STUDENT	26.028	0.025	0.025	AVG. USING STUDENT10(H/1.11) -- SEE TEXT					
8 (PI+) - (PI-)/AVG., MEAN LIFE DIFF. (PERCENT)										
DT	N	THIS QUANTITY IS A MEASURE OF CPT INVARIANCE IN W.I.								
DT	L	0.23	0.40	LOBKOWICZ	66	CNTR	SEE NOTE L		9/66	
DT	L	ABOVE IS THE MOST CONSERVATIVE VALUE QUOTED BY AUTHORS							9/66	
DT		0.4	0.7	BARON	66	CNTR			7/66	
DT		-0.14	0.29	PETRUKHIN	68	CNTR			8/68	
DT		0.055	0.071	AYRES	71	CNTR			3/71	
DT	AVG	0.053	0.068	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)					
DT	STUDENT	0.053	0.073	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT					
8 CHARGED PION PARTIAL DECAY MODES										
P1	CHAR.	PION INTO MU (MU-NEU)					DECAY MASSES			
P2	CHAR.	PION INTO E (E-NEU)					105+ 0			
P3	CHAR.	PION INTO MU (MU-NEU) GAMMA					105+ 0+ 0			
P4	CHAR.	PION INTO P10 E (E-NEU)					134+ .5+ 0			
P5	CHAR.	PION INTO E NEU GAMMA					.5+ 0+ 0			
P6	CHAR.	PION INTO E NEU E+ E-					.5+ 0+ .5+ .5			
8 CHARGED PION BRANCHING RATIOS										
R1	CHAR.	PION INTO MU NEU GAMMA (UNITS 10**--4)	(P31)/(P1)							
R1	26	1.24	0.25	CASTAGNOL	58	EMUL	E(MU).LT.3.38 MV			
R2	CHAR.	PION INTO E NEU (UNITS 10**--4)	(P21)/(P1)							
R2	1.21	0.07	ANDERSON	60	CNTR					
R2 D	(1.247)	(0.028)	DI CAPUA	64	CNTR				11/75*	
R2 D	1.274	0.026	BRYMAN	75	RVUE				9/75*	
R2 D	BRYMAN 75 IS A RECALC. OF DICAPUA 64 EXPT USING LATEST PI LIFETIME.								9/75*	
R2	AVG	1.267	0.023	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)					
R2	STUDENT	1.268	0.025	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT					
R3	CHAR.	PION INTO P10 E NEU (UNITS 10**--8)	(P41)/(P1)							
R3	D	52	(1.15)	(1.22)	DEPOMMIER	63	CNTR	+	2/72	
R3	D	36	0.97	0.29	BARTLETT	64	OSPK	+		
R3	D	38	1.07	0.21	BACASTOW	65	OSPK	+		
R3	D	1.10	0.26	BERTRAM	65	OSPK	+		6/66	
R3	D	43	1.1	0.2	DUNAITSEV	65	CNTR	+	7/66	
R3	D	332	1.00	0.08	0.10	DEPOMMIER	68	CNTR	+	3/68
R3	AVG	1.023	0.069	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)					
R3	STUDENT	1.023	0.074	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT					
D	DEPOMMIER 68 STATES THAT THE RESULT OF DEPOMMIER 63 IS AT LEAST								2/72	
D	10 PERCENT TOO LARGE BECAUSE OF A SYSTEMATIC ERROR IN THE P10								2/72	
D	DETECTION EFFICIENCY. THIS MAY BE TRUE OF ALL THE PREVIOUS								2/72	
D	MEASUREMENTS ACCORDING TO DEPOMMIER 68 AND V.SOERGEL, PRIVATE								2/72	
D	COMMUNICATION, 1972.									
R4	CHAR.	PION INTO E NEU GAMMA (UNITS 10**--8)	(P51)/(P1)							
R4	143	3.0	0.5	DEPOMMIER	63	CNTR	GAM KE 50-90 MEV		6/66	
R5	CHAR.	PION INTO E NEU E+ E- (UNITS 10**--8)	(P61)/(P1)							
R5	3.4	OR LESS	CL=90	KORENCHEN	71	OSPK	+		10/71	

REFERENCES FOR CHARGED PION									
CROME	54	PR	96	470	K M CROME, R H PHILLIPS	(LFL)			
BARKAS	56	PR	101	778	W H BARKAS, W BIRNBAUM, F M SMITH	(LFL)			
CROME	57	NC	5	541	K M CROME	(STANFORD HEPL)			
CASTAGNO	58	PR	112	1779	C CASTAGNOLI, M MUCHNIK	(ROMA)			
ANDERSON	60	PR	119	2050	H L ANDERSON, T FUJII, R H MILLER +	(EFT)			
ASHKIN	60	NC	16	490	ASHKIN, FAZZINI, FIOCCARDI, IPMAN +	(CERN)			
DEPOMMIER	63	PL	5	61	DEPOMMIER, HEINTZE, RUBIA, SOERGEL	(CERN)			
DEPOMMIER	63	PL	7	285	P DEPOMMIER, HEINTZE, RUBIA, SOERGEL	(CERN)			
BARTLETT	64	PR	1368	1452	BARTLETT, DEVONS, MEYER, ROSEN	(COLUMBIA)			
DI CAPUA	64	PR	1338	1333	DI CAPUA, GARLAND, PONDROM, STRELZOFF	(COLU)			
BACASTOW	65	PR	139	8407	*GHESQUIERE, WIEGAND, LARSEN	(LRL+SLAC)			
BERTRAM	65	PR	139	B 617	BERTRAM, MEYER, CARRIGAN +	(MICH+CARNEGIE)			
DUNAITSEV	65	JETP	20	58	DUNAITSEV, PETRUKHIN, PROKOSHIN +	(DUBNA)			
ECKHAUSE	65	PL	19	348	ECKHAUSE, HARRIS, SHULER +	(WILLIAM AND MARY)			
BARON	66	PRL	16	775	BARON, DORE, DORFAN, KRIEGER +	(COLUMBIA)			
DUNAITSEV	66	PL	23	283	*KUTYIN, PROKOSHIN, RASUVAEV, SIMONOV	(DUBNA)			
KINSEY	66	PR	144	1132	KINSEY, LOBKOWICZ, NORDBERG	(ROCHESTER UNIV)			
LOBKOWICZ	66	PRL	17	548	LOBKOWICZ, MELISSINOS, NAGASHIMA +	(ROCH+BNL)			
HYMAN	67	PL	258	376	*LOKEN, PEWITT, OERICK +	(ANL+CORN+NNES)			
NORDBERG	67	PL	248	594	NORDBERG, LOBKOWICZ, BURMAN	(ROCHESTER UNIV)			
SHAFFER	67	PR	163	1451	ROBERT E. SHAFFER	(LRL)			
ALSO	65	PRL	14	923	SHAFFER, CROME, JENKINS	(LRL)			
DEPOMMIER	68	NP	84	189	DEPOMMIER, DUCLOS, HEINTZE, KLEINKNECHT +	(CERN)			
PETRUKHIN	68	JINR-P1	3862		PETRUKHIN, RYKALIN, KHAZINS, CISEK	(DUBNA)			
BOOTH	70	PL	32B	723	*JOHNSON, WILLIAMS, WORMALD	(LIVP)			
AYRES	71	PR	30	1051	*CORMACK, GREENBERG, KENNEY +	(LRL+UCSB)			
ALSO	67	PR	157	1288	AYRES, CALDWELL, GREENBERG, KENNEY, KURZ +	(LRL)			
ALSO	68	PRL	21	261	AYRES, CORMACK, GREENBERG, KENNEY +	(LRL, UCSB)			
ALSO	69	UCRL	18369		CAVID S AYRES (THIS IS)	(LRL)			
ALSO	69	PRL	23	1267	GREENBERG, AYRES, CORMACK, KENNEY +	(LRL, UCSB)			
BACKENSTOISS	71	PL	368	403	BACKENSTOISS, DANIEL, KOCH +	(CERN, KARL, HEID)			
ALSO	70	THESIS			C. VON DER MALSBURG	(HEIDELBERG)			
KORENCHEN	71	SJNP	13	189	KORENCHENKO, KOSTIN, MICEL VACHER +	(JINR)			
SHAFFER	72	PRIVATE COMM.			R. SHAFFER, 1972	(FNAL)			
BACKENSTOISS	73	PL	438	539	BACKENSTOISS, DANIEL, KOCH +	(CERN+KARL+MUNICH)			
ALSO	73	SUBMITTED TO NP			L. TAUSCHER				
DUNAITSEV	73	SJNP	16	292	DUNAITSEV, PROKOSHIN, RASUVAEV +	(SERP)			
BRYMAN	75	PR	D11	1337	*PICCOTTO	(UNIV OF VICTORIA)			
PAPERS NOT REFERRED TO IN DATA CARDS									
MERRISON	62	ADVP	11	1	A W MERRISON	(LIVERPOOL)			
SHAPIRO	62	PR	125	1022	G SHAPIRO, L M LEDEFMAN	(COLUMBIA)			
CZIRR	63	PR	130	341	JOHN B CZIRR	(LRL)			

π^0									
9 NEUTRAL PION(135,JPG=0--) I=1									
9 (PI+) - (PI-) MASS DIFFERENCE (MEV)									
D	(5.37)	(1.0)	PANOFSKY	51	CNTR	-			
D	4.50	0.31	CHINOWSKY	54	CNTR	-			
D	4.62	0.05	HADDOCK	59	CNTR	-			
D	4.60	0.04	HILLMAN	59	CNTR	-			
D	4.55	0.07	CASSELL	59	CNTR	-			
D	4.69	0.07	SAMIOS	60	HBC				2/72
D	4.6056	0.0055	CZIRR	63	CNTR	-			
D	4.59	0.03	PETRUKHIN	63	CNTR	-			9/66
D	4.6034	0.0052	VASTLEVSK	66	CNTR	-			
D	AVG	4.6043	0.0037	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.0)				
D	STUDENT	4.6043	0.0040	AVERAGE	USING STUDENT10(H/1.11) -- SEE TEXT				
9 NEUTRAL PION MEAN LIFE (UNITS 10**--16)									
T	N	76	(1.9)	(0.5)	(0.5)	GLASSER	61	EMUL	
T	N	45	(2.3)	(1.1)	(1.0)	TIETGE	62	EMUL	
T	N	88	(2.8)	(0.9)	(0.9)	KOLLER	63	EMUL	SEE STAMER 66
T	N	75	1.05	0.18	0.18	VON DARDE	63	CNTR	
T	N	75	(1.7)	(0.5)		SHME	64	EMUL	
T			0.730	0.105		BELLETTIN	65	CNTR	
T	N	67	(1.6)	(0.6)	(0.5)	EVANS	65	EMUL	6/66
T	K	232	1.0	0.5		STAMER	66	EMUL	8/67
T			0.56	0.36		BELLETTIN	70	CNTR	PRIMOFF. ON NUC 7/70
T			0.9	0.068		KRYSKIN	70	CNTR	PRIMAKOFF EFFECT 12/70
T	B		0.82	0.04		BROWMAN	74	CNTR	PRIMAKOFF EFFECT 7/75*
T	N	OLD EMULSION MEASUREMENTS NOT USED BECAUSE OF POSSIBLE SYSTEMATIC							
T	N	SHIFT TO LARGER MEAN LIFE VALUES.							
T	K	INCLUDES EVENTS OF KOLLER 63.							8/67
T	B	BROWMAN GIVES PI WIDTH=0.02+-0.42EV. MEAN LIFE IS HBAR/WIDTH.							11/75*
T	AVG	0.828	0.057	0.057	AVERAGE	(ERROR INCL. SCALE FACTOR OF 1.8)			
T	STUDENT	0.835	0.038	0.038	AVG. USING STUDENT10(H/1.11) -- SEE TEXT				
(SEE IDEOGRAM BELOW)									
9 NEUTRAL PION PARTIAL DECAY MODES									
P1	P10	INTO 2 GAMMA					DECAY MASSES		
P2	P10	INTO E+ E- GAMMA					0+ 0		
P3	P10	INTO 4 ELECTRONS					.5+ .5+ 0		.5
P4	P10	INTO 3 GAMMA					0+ 0+ 0		0
P5	P10	INTO							

Data Card Listings
For notation, see key at front of Listings.

Stable Particles

pi0, K±

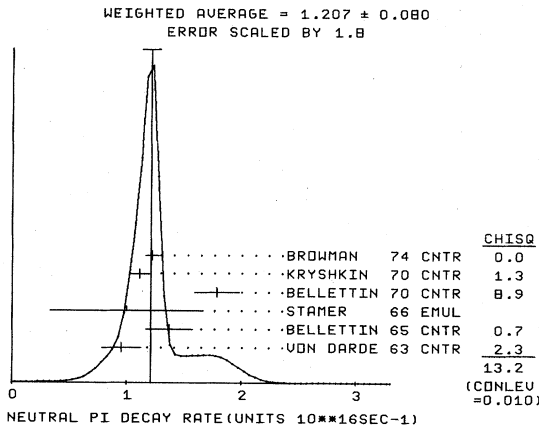


Table with 2 columns: Name, CHISQ. Rows include BROWMAN 74 CNTR (0.0), KRYSHKIN 70 CNTR (1.3), BELLETTIN 70 CNTR (8.9), STAMER 66 EMUL (0.7), BELLETTIN 65 CNTR (2.3), UDN DARDE 63 CNTR (13.2).

Table of experimental data for neutral pi0 decay rate. Columns include experiment name, parameters, units, and CHISQ values.

REFERENCES FOR NEUTRAL PION

Table of references for neutral pion, listing authors and their affiliations.



10 CHARGED K (494, JP=0-) I=1/2

10 CHARGED K MASS (MEV)

Table of charged K mass measurements. Columns include experiment name, mass value, error, and CHISQ.

10 (K+) - (K-) MASS DIFFERENCE (MEV)

Table of charged K mass difference measurements. Columns include experiment name, mass difference, error, and CHISQ.

10 CHARGED K MEAN LIFE (UNITS 10**--8)

Table of charged K mean life measurements. Columns include experiment name, mean life, error, and CHISQ.

WEIGHTED AVERAGE = 0.8084 ± 0.0021, ERROR SCALED BY 2.4

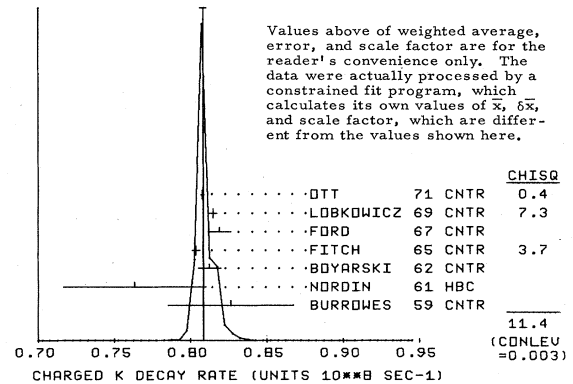


Table with 2 columns: Name, CHISQ. Rows include DTT 71 CNTR (0.4), LOBKOWICZ 69 CNTR (7.3), FORD 67 CNTR (3.7), FITCH 65 CNTR (3.7), BOYARSKI 62 CNTR (11.4), NORDIN 61 HBC (0.0033), BURROWS 59 CNTR (0.0026).

10 ((K+) - (K-))/AVG., MEAN LIFE DIFFERENCE (PERCENT)

Table of mean life difference measurements. Columns include experiment name, difference, error, and CHISQ.

10 CHARGED K PARTIAL DECAY MODES

Table of charged K partial decay modes, listing various decay channels and their branching ratios.

CHARGED K CONSTRAINED FIT

OVERALL FIT OF MEAN LIFE, WIDTHS AND BRANCHING RATIOS USES 57 DATA POINTS TO DETERMINE SIX QUANTITIES. OVERALL FIT HAS CHISQ=78-2. MAIN CONTRIBUTION (13.3) COMES FROM R19 OF HAIDT 71 (WE SEE NO REASON TO REJECT THIS EXPERIMENT AT THIS TIME)

Stable Particles

K±

Data Card Listings

For notation, see key at front of Listings.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_i δP_i), while the off-diagonal elements are the normalized correlation coefficients (δP_i δP_j) / (δP_i δP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 6 columns (P 1 to P 6) and 6 rows of branching fraction data and error correlations.

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., G_i ≡ Γ_i = Γ_total P_i, in appropriate units. In analogy to the matrix above, the diagonal elements are G_i ± δG_i, where δG_i = √(δG_i δG_i), while the off-diagonal elements are the normalized correlation coefficients (δG_i δG_j) / (δG_i δG_j). Note that, because of the error in Γ_total, the errors and correlations here are not directly derivable from those above.

Table with 6 columns (G 1 to G 6) and 6 rows of rate data and error correlations.

10 CHARGED K DECAY RATES

Table listing decay rates for various K particles, including W1, W2, W3, W4, W5, W6, W7, W8, W9, W10, W11, W12, W13, W14, W15, W16, W17, W18, W19, W20, W21, W22, W23, W24, W25, W26, W27, W28, W29, W30, W31, W32, W33, W34, W35, W36, W37, W38, W39, W40, W41, W42, W43, W44, W45, W46, W47, W48, W49, W50, W51, W52, W53, W54, W55, W56, W57, W58, W59, W60, W61, W62, W63, W64, W65, W66, W67, W68, W69, W70, W71, W72, W73, W74, W75, W76, W77, W78, W79, W80, W81, W82, W83, W84, W85, W86, W87, W88, W89, W90, W91, W92, W93, W94, W95, W96, W97, W98, W99, W100.

10 ((K+) - (K-))/AVG., DECAY RATE DIFFERENCE (PERCENT)

Table listing decay rate differences for various K particles, including D1, D2, D3, D4, D5, D6, D7, D8, D9, D10, D11, D12, D13, D14, D15, D16, D17, D18, D19, D20, D21, D22, D23, D24, D25, D26, D27, D28, D29, D30, D31, D32, D33, D34, D35, D36, D37, D38, D39, D40, D41, D42, D43, D44, D45, D46, D47, D48, D49, D50, D51, D52, D53, D54, D55, D56, D57, D58, D59, D60, D61, D62, D63, D64, D65, D66, D67, D68, D69, D70, D71, D72, D73, D74, D75, D76, D77, D78, D79, D80, D81, D82, D83, D84, D85, D86, D87, D88, D89, D90, D91, D92, D93, D94, D95, D96, D97, D98, D99, D100.

10 CHARGED K BRANCHING RATIOS

Table listing branching ratios for various K particles, including R 0, R 1, R 2, R 3, R 4, R 5, R 6, R 7, R 8, R 9, R 10, R 11, R 12, R 13, R 14, R 15, R 16, R 17, R 18, R 19, R 20, R 21, R 22, R 23, R 24, R 25, R 26, R 27, R 28, R 29, R 30, R 31, R 32, R 33, R 34, R 35, R 36, R 37, R 38, R 39, R 40, R 41, R 42, R 43, R 44, R 45, R 46, R 47, R 48, R 49, R 50, R 51, R 52, R 53, R 54, R 55, R 56, R 57, R 58, R 59, R 60, R 61, R 62, R 63, R 64, R 65, R 66, R 67, R 68, R 69, R 70, R 71, R 72, R 73, R 74, R 75, R 76, R 77, R 78, R 79, R 80, R 81, R 82, R 83, R 84, R 85, R 86, R 87, R 88, R 89, R 90, R 91, R 92, R 93, R 94, R 95, R 96, R 97, R 98, R 99, R 100.

Table listing experimental data for K particles, including R 2, R 3, R 4, R 5, R 6, R 7, R 8, R 9, R 10, R 11, R 12, R 13, R 14, R 15, R 16, R 17, R 18, R 19, R 20, R 21, R 22, R 23, R 24, R 25, R 26, R 27, R 28, R 29, R 30, R 31, R 32, R 33, R 34, R 35, R 36, R 37, R 38, R 39, R 40, R 41, R 42, R 43, R 44, R 45, R 46, R 47, R 48, R 49, R 50, R 51, R 52, R 53, R 54, R 55, R 56, R 57, R 58, R 59, R 60, R 61, R 62, R 63, R 64, R 65, R 66, R 67, R 68, R 69, R 70, R 71, R 72, R 73, R 74, R 75, R 76, R 77, R 78, R 79, R 80, R 81, R 82, R 83, R 84, R 85, R 86, R 87, R 88, R 89, R 90, R 91, R 92, R 93, R 94, R 95, R 96, R 97, R 98, R 99, R 100.

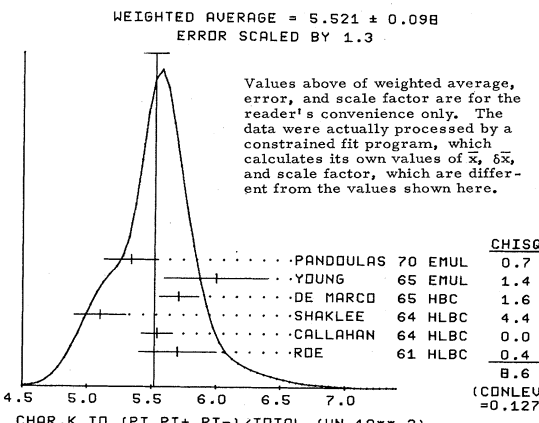
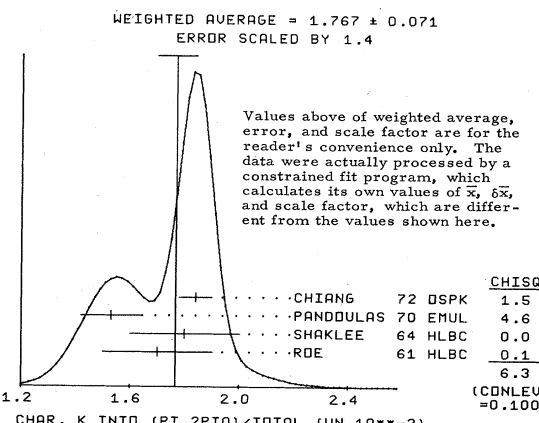


Table listing experimental data for K particles, including R 4, R 5, R 6, R 7, R 8, R 9, R 10, R 11, R 12, R 13, R 14, R 15, R 16, R 17, R 18, R 19, R 20, R 21, R 22, R 23, R 24, R 25, R 26, R 27, R 28, R 29, R 30, R 31, R 32, R 33, R 34, R 35, R 36, R 37, R 38, R 39, R 40, R 41, R 42, R 43, R 44, R 45, R 46, R 47, R 48, R 49, R 50, R 51, R 52, R 53, R 54, R 55, R 56, R 57, R 58, R 59, R 60, R 61, R 62, R 63, R 64, R 65, R 66, R 67, R 68, R 69, R 70, R 71, R 72, R 73, R 74, R 75, R 76, R 77, R 78, R 79, R 80, R 81, R 82, R 83, R 84, R 85, R 86, R 87, R 88, R 89, R 90, R 91, R 92, R 93, R 94, R 95, R 96, R 97, R 98, R 99, R 100.



Data Card Listings

Stable Particles

For notation, see key at front of Listings.

K±

R5 CHAR. K INTO (MU P10 NEU)/TOTAL (UNITS 10**--2) (P5)
R5 0 (2.8) (1.0) BIRGE 56 EMUL +
R5 0 (5.9) (1.3) ALEXANDER 57 EMUL +
R5 0 (2.8) (0.4) TAYLOR 59 EMUL +
R5 0 EARLIER EXPERIMENTS NOT AVERAGED
R5 2345 3.33 0.16 CHIANG 72 OSPK + 1.84 GEV/C K+ 9/72
R5 FIT 3.197 0.087 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

R6 CHAR. K INTO (E P10 NEU)/TOTAL (UNITS 10**--2) (P6)
R6 0 (3.2) (1.3) BIRGE 56 EMUL +
R6 0 (5.1) (1.3) ALEXANDER 57 EMUL +
R6 0 EARLIER EXPERIMENTS NOT AVERAGED
R6 429 4.7 0.3 SHAKLEE 61 HLBC + 11/67
R6 3516 4.86 0.10 CHIANG 72 OSPK + 1.84 GEV/C K+ 9/72
R6 AVG 4.849 0.093 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R6 STUDENT 4.85 0.10 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R6 FIT 4.823 0.052 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R7 CHAR. K INTO (PI2 + MU3)/TOTAL (UNITS 10**--2) (P2+P5) 11/67
R7 WE COMBINE THESE TWO MODES FOR EXPTS MEASURING THEM IN XENON BC
R7 BECAUSE OF DIFFICULTIES OF SEPARATING THEM THERE
R7 23.4 1.1 ROE 61 HLBC + 11/67
R7 886 25.4 0.9 SHAKLEE 64 HLBC + 11/67
R7 AVG 24.60 0.98 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)
R7 STUDENT 24.61 0.83 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R7 FIT 24.25 0.15 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R8 K+ INTO (PI+ PI- E- NEU)/TOTAL (UNITS 10**--7) (P8)
R8 20. OR LESS CL=.95 BIRGE 65 FBC + 8/66
R8 0 6.9 OR LESS CL=.95 ELY 69 HLBC + 10/69
R8 0 9.0 OR LESS CL=.95 SCHWEINBE 71 HLBC + 9/71

R9 K+ INTO (PI+ PI- MU+ NEU)/TOTAL (UNITS 10**--5) (P9)
R9 1 0.77 0.54 0.50 CLINE 65 FBC + 8/66

R10 K+ INTO (PI+ PI+ MU- NEU)/TOTAL (UNITS 10**--6) (P10)
R10 0 3.0 OR LESS CL=.95 BIRGE 65 FBC + 8/66

R11 CHAR. K INTO (E NEU)/TOTAL (UNITS 10**--5) (P11) 11/67
R11 160.0 OR LESS CL=.95 BORREANI 64 HBC + 11/67
R11 4 (2.8) (1.3) BOWEN 67 OSPK + 8/67
R11 BOWEN RESULT SHOULD BE CORRECTED TO 1.9(+1.7)-1.21 BECAUSE OF
R11 K+ TO E+ NEU GAMMA DECAYS BEFORE COMPARING WITH BOTTERILL 67 R28

R12 CHAR. K INTO (PI GAMMA GAMMA)/TOTAL (UNITS 10**--4) (P17)
R12 ALL VALUES GIVEN HERE ASSUME A PHASE SPACE PION ENERGY SPECTRUM 2/72
R12 -0.1 0.6 CHEN 68 OSPK + TPI1 60-90 MEV 8/71
R12 0 0.5 OR LESS CL=.90 KLEMS 71 OSPK + TPI1GT 117 MEV 9/73
R12 0 0.35 OR LESS CL=.90 LJUNG 73 HLBC + 6-102,114-127MEV 9/73

R13 CHAR. K INTO (PI P10 GAMMA)/TOTAL (UNITS 10**--4) (P13)
R13 18 2.2 0.7 CLINE 64 FBC + PI+ KE 55-80 MEV 8/66
R13 0 1.9 OR LESS CL=.90 EMERSON 69 OSPK PI+ KE 55-80 MEV 10/69
R13 M 0 1.0 OR LESS CL=.90 MALTSEV 70 HLBC + PI+ KE LT 55 MEV 12/75*
R13 A2100 2.71 0.19 ABRAMS 72 ASPK -- PI+ KE 55-90 MEV 1/73
R13 24 2.4 0.8 EDWARDS 72 OSPK PI+ KE 58-90 MEV 9/73
R13 L (1.5) (1.1) (0.6) LJUNG 73 HLBC + PI+ KE 55-80 MEV 9/73
R13 L (2.4) (1.5) (1.1) LJUNG 73 HLBC + PI+ KE 55-90 MEV 9/73
R13 L 17 6.8 3.7 2.1 LJUNG 73 HLBC + PI+ KE 55-102MEV 9/73
R13 M MALTSEV TO SELECTS LOW PI+ ENERGY TO ENHANCE DIRECT EMISSION CONTR. 1/76*
R13 L THE LJUNG 73 VALUES ARE NOT INDEPENDENT.
R13 A ABRAMS 72 OBSERVES DIRECT EMISSION BR. RATIO OF (1.56+-0.35)*10**--5 1/73
R13 A +-0.5*10**--5 ADDNL. SYST. ERROR AND INNER BREMSSTRAHLUNG BR. RATIO 1/73
R13 A OF (2.55+-0.18)*10**--4. WE QUOTE THE SUM OF THESE BR. RATIOS. 1/73
R13 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

R14 CHAR. K INTO (PI PI+ PI- GAMMA)/TOTAL (UNITS 10**--4) (P14)
R14 1.0 0.4 STAMER 65 EMUL + EGAM GT 11MEV 8/66

R15 CHAR. K INTO (PI E- E-)/TOTAL (UNITS 10**--6) (P15)
R15 1 2.45 OR LESS CL=.90 CAMERINI 64 FBC + 8/66
R15 4.4 OR LESS CL=.90 BISI 67 DBC + 11/67
R15 C (0.4) OR LESS CLINE1 67 FBC + 11/67
R15 C 0.88 OR LESS CL=.90 CLINE2 67 FBC + 2/74
R15 32.0 OR LESS CL=.90 BETIER 72 OSPK -- 9/72
R15 1.7 OR LESS CL=.90 CENCE 74 ASPK + THREE TRACK EVTS 10/74*
R15 0.27 OR LESS CL=.90 GENCE 74 ASPK + TWO TRACK EVTS 10/74*
R15 C CLINE2 REPLACES CLINE1. CLINE1 IS NOT FOR CL=.90 2/74

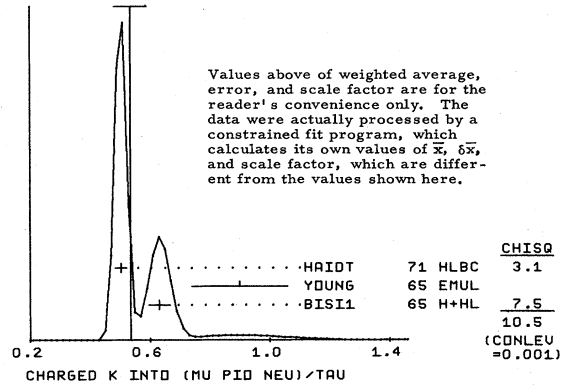
R16 CHAR. K INTO (PI MU+ MU-)/TOTAL (UNITS 10**--6) (P16)
R16 3.0 OR LESS CL=.90 CAMERINI 65 FBC + 8/66
R16 2.4 OR LESS CL=.90 BISI 67 DBC + 11/67

R17 CHAR. K INTO (PI P10)/TAU (P21)/(P3)
R17 134 3.24 0.34 YOUNG 65 EMUL + 8/66
R17 1045 3.96 0.15 CALLAHAN 66 FBC + 9/66
R17 AVG 3.84 0.27 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)
R17 STUDENT 3.86 0.16 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R17 FIT 3.768 0.033 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R18 CHAR. K INTO (PI ZP10)/TAU (P41)/(P3)
R18 2027 0.303 0.009 BISI 65 H+HL + 8/66
R18 17 0.393 0.099 YOUNG 65 EMUL + 8/66
R18 AVG 0.3037 0.0090 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R18 STUDENT 0.3037 0.0097 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R18 FIT 0.3098 0.0079 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

R19 CHAR. K INTO (MU P10 NEU)/TAU (P51)/(P3)
R19 2845 0.632 0.035 BISI 1 65 H+HL + 8/66
R19 38 0.90 0.16 YOUNG 65 EMUL + 8/66
R19 H 1505 (0.510) (0.017) EICHTEN 68 HLBC + 11/68
R19 H1505 0.503 0.019 HAIDT 71 HLBC + 12/70
R19 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
R19 AVG 0.536 0.054 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.2)
R19 STUDENT 0.527 0.025 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R19 FIT 0.572 0.016 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.536 ± 0.054
ERROR SCALED BY 3.2



CHISQ
71 HLBC 3.1
65 EMUL
65 H+HL 7.5
10.5
(CONLEV = 0.001)

R20 CHAR. K INTO (E P10 NEU)/TAU (P61)/(P3)
R20 230 0.90 0.06 BORREANI 64 HBC + 8/66
R20 37 0.90 0.16 YOUNG 65 EMUL + 8/66
R20 854 0.94 0.09 BELLOTT2 67 HLBC 11/67
R20 H 4385 (0.846) (0.021) EICHTEN 68 HLBC + 11/68
R20 H4385 0.850 0.019 HAIDT 71 HLBC + 12/70
R20 2827 0.856 0.040 BRAUN 75 HLBC + 12/75*
R20 H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.
R20 AVG 0.858 0.016 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R20 STUDENT 0.858 0.018 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R20 FIT 0.8632 0.0098 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R21 K+ INTO (PI+ PI- E+ NEU)/TAU (UNITS 10**--4) (P71)/(P3)
R21 69 6.7 1.5 BIRGE 65 FBC + 8/66
R21 269 5.83 0.63 ELY 69 HLBC + 11/68
R21 500 7.36 0.68 BURQUIN 71 ASPK 12/71
R21 106 7.0 0.9 SCHWEINBE 71 HLBC + 9/71
R21 AVG 6.64 0.40 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R21 STUDENT 6.66 0.48 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

R22 K+ INTO (PI+ PI- MU+ NEU)/TAU (UNITS 10**--4) (P91)/(P3)
R22 1 (2.5) APPROX GREINER 64 EMUL + 8/66
R22 7 2.57 1.55 BISI 67 DBC + 11/67

R23 CHAR. K INTO (E P10 NEU)/(MU2+PI2) (UNITS 10**--2) (P61)/(P1+P2)
R23 1679 5.89 0.21 CESTER 66 OSPK + 8/67
R23 5110 6.16 0.22 ESCHSTRUT 68 OSPK + 3/68
R23 AVG 6.02 0.15 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R23 STUDENT 6.02 0.17 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R23 FIT 5.977 0.067 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R24 CHAR. K INTO (PI P10)/(MU NEU) (P21)/(P1)
R24 A4917 0.3277 0.0065 AUERBACH 67 OSPK + 1/74
R24 1600 0.305 0.018 ZELLER 69 ASPK + 10/69
R24 25K 0.328 0.005 WEISSENBE 74 STRC + 7/74*
R24 A AUERBACH 67 CHANGED FROM .3253+-0.0065. SEE COMMENT WITH RATIO R26. 1/74
R24 AVG 0.3268 0.0039 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R24 STUDENT 0.3269 0.0042 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R24 FIT 0.3309 0.0029 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R25 CHAR. K INTO (E P10 NEU)/(MU NEU) (P61)/(P1)
R25 A 295 0.0791 0.0054 AUERBACH 67 OSPK + 1/74
R25 960 0.0775 0.0033 BOTTERILL 68 ASPK + 5/68
R25 561 0.069 0.006 GARLAND 68 OSPK + 4/68
R25 350 0.009 0.006 ZELLER 69 ASPK + 10/69
R25 A AUERBACH 67 CHANGED FROM .0797+-0.0054. SEE COMMENT WITH RATIO R26. 1/74
R25 A THE VALUE .0785+-0.0025 GIVEN IN AUERBACH 67 IS AN AVERAGE OF
R25 A AUERBACH 67 R25 AND CESTER 66 R23. 3/74
R25 AVG 0.0752 0.0024 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R25 STUDENT 0.0753 0.0027 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R25 FIT 0.07582 0.00091 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

R26 CHAR. K INTO (MU P10 NEU)/(MU NEU) (P51)/(P1)
R26 A 307 0.0486 0.0040 AUERBACH 67 OSPK + 1/74
R26 G 424 0.0480 0.0037 GARLAND 68 OSPK 1/74
R26 240 0.054 0.009 ZELLER 69 ASPK + 10/69
R26 A AUERBACH 67 CHANGED FROM .0602+-0.0046 BY ERRATUM WHICH BRINGS THE
R26 A MU-SPECTRUM CALCULATION INTO AGREEMENT WITH GAILLARD TO APPENDIX B. 1/74
R26 G GARLAND 68 CHANGED FROM .055+-0.004 IN AGREEMENT WITH MU-SPECTRUM 1/74
R26 G CALCULATION OF GAILLARD 70 APPENDIX B. L.G.PONDROM, PRIV.COMM.(73) 1/74
R26 AVG 0.0488 0.0026 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R26 STUDENT 0.0487 0.0028 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R26 FIT 0.0503 0.0014 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

R27 CHAR. K INTO (MU NEU)/TAU (P11)/(P3) 9/66
R27 R 427 (10.38) (0.82) YOUNG 65 EMUL +
R27 R DELETED FROM OVERALL FIT BECAUSE YOUNG 65 CONSTRAINS HIS RESULTS.
R27 R TO ADD UP TO 1. ONLY YOUNG MEASURED MU2 DIRECTLY.
R27 FIT 11.385 0.072 FROM FIT

R28 CHAR. K INTO (E NEU)/(MU NEU) (UNITS 10**--5) (P111)/(P1)
R28 10 1.9 0.7 0.5 BOTTERILL 67 ASPK + 11/67
R28 8 1.8 0.8 0.6 MACEK 69 ASPK + 4/69
R28 112 2.42 0.42 CLARK 72 OSPK + 1/73
R28 534 2.37 0.17 HEARD2 75 SPEC + 11/75*
R28 404 2.51 0.15 HEINTZE 76 SPEC + 2/76*
R28 AVG 2.42 0.11 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R28 STUDENT 2.42 0.12 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

Stable Particles

K±

Data Card Listings

For notation, see key at front of Listings.

R29	CHAR. K INTO (MU P10 NEU)/(E P10 NEU)	(P51)/(P6)	
R29	C1509 0.703 0.056	CALLAHAN 66 HLCB	6/68
R29	5601 0.667 0.017	BOTTERI 68 ASPK +	6/68
R29	H 1398 (0.634) (0.022)	EICHTEN 68 HLCB	10/68
R29	H (0.596) (0.025)	HAIDT 71 HLCB +	12/70
R29	D3480 0.698 0.025	CHIANG 72 OSPK + 1.84 GEV/C K+	9/72
R29	L 554 0.705 0.063	LUCAS 73 HBC - DALITZ PRS ONLY	11/73
R29	B 1585 (0.608) (0.014)	BRAUN 75 HLCB +	1/76*
R29	COMMENTS		
R29	C FROM CALLAHAN 66 WE USE ONLY THE MU3/E3 RATIO AND DO NOT		
R29	C INCLUDE IN THE FIT THE RATIOS MU3/TAU AND E3/TAU, SINCE THEY SHOW		
R29	C LARGE DISAGREEMENTS WITH THE REST OF THE DATA.		
R29	H HAIDT 71 IS A REANALYSIS OF EICHTEN 68.		
R29	H ONLY INDIVIDUAL RATIOS INCLUDED IN FIT (SEE R19 AND R20).		
R29	D CHIANG 72 R29 IS STATISTICALLY INDEPENDENT OF CHIANG 72 R5 AND R6.		
R29	L LUCAS 73 GIVES N(MU3)=554+-7.6PCT, N(E3)=786+-3.1PCT. WE DIVIDE.		
R29	B BRAUN 75 VALUE IS FROM FORM FACTOR FIT. ASSUMES MU-E UNIVERSALITY.		
R29	AVG 0.680 0.013 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)		
R29	STUDENT 0.680 0.015 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT		
R29	FIT 0.663 0.018 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)		
R30	CHAR. K INTO (P10 E NEU GAMMA)/(P10 E NEU) (UNITS 10**=-2)		
R30	R30 1.2 0.8 BELLOTTI 67 HLCB + EGAM GT 30MEV (P181)/(P6) 11/67		
R30	R 13 0.76 0.28 ROMANO 71 HLCB EGAM GT 10MEV 10/71		
R30	R (0.53) (0.22) ROMANO 71 HLCB + EGAM GT 30 MEV 9/73		
R30	L 16 0.48 0.20 LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R30	L (0.22) (0.19) (0.10) LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R30	L FIRST LJUNG VALUE IS FOR COS(ELCET-GAMMA)=0.9, SECOND VALUE IS		
R30	L FOR COS(ELCET-GAMMA) BETW 0.6 AND 0.9 FOR COMPARISON WITH ROMANO.		
R30	R BOTH ROMANO VALUES ARE FOR COS(ELCET-GAMMA) BETW 0.6 AND 0.9.		
R30	R SECOND VALUE IS FOR COMPARISON WITH SECOND LJUNG VALUE.		
R30	R WE USE LOWEST EGAM CUT FOR TABLE VALUE. SEE ROMANO FOR EGAM DEPEND.		
R30	R30 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)		
R31	K- INTO (PI+ E- E-)/TOTAL (UNITS 10**=-5) (P19)		
R31	TEST OF LEPTON NUMBER CONSERVATION.		
R31	1.5 OR LESS CHANG 68 HBC - 3/68		
R32	CHAR. K INTO (PI NEU NEU)/TOTAL (UNITS 10**=-6) (P20)		
R32	C (1.4) OR LESS CL=.90 KLEMS 71 OSPK + T(P1) 117-127MEV 3/74		
R32	C (0.94) OR LESS CL=.90 CABLE 73 CNTR + T(P1) 60-105 MEV 2/74		
R32	C 0.56 OR LESS CL=.90 CABLE 73 CNTR + T(P1) 60-127 MEV 2/74		
R32	L 0 57.0 OR LESS CL=.90 LJUNG 73 HLCB + 9/73		
R32	C KLEMS 71 AND CABLE 73 ASSUME PI SPECTRUM SAME AS KE3 DECAY.		
R32	C SECOND CABLE LIMIT COMBINES CABLE AND KLEMS DATA FOR VECTOR INT.		
R32	L LJUNG 73 ASSUMES VECTOR INTERACTION. 9/73		
R33	CHAR. K INTO (E NEU GAMMA)/TOTAL (UNITS 10**=-5) (P21)		
R33	M 7.1 OR LESS MACEK 70 OSPK + P(E) 234 TO 247 12/70		
R33	M ABOVE IS MEASUREMENT OF STRUCTURE-DEPENDENT DECAY ONLY.		
R34	CHAR. K INTO (PI GAMMA)/TOTAL (UNITS 10**=-6) (P22)		
R34	4.0 OR LESS CL=.90 KLEMS 71 OSPK + 8/71		
R35	CHAR. K INTO (TAU)/(TAU PRIME)		
R35	USED FOR DELTA I=1/2 TEST.		
R35	FIT 3.227 0.083 FROM FIT (P3/P4)		
R36	CHAR. K INTO (PI 3GAMMA)/TOTAL (UNITS 10**=-4) (P23)		
R36	3.0 OR LESS CL=.90 KLEMS 71 OSPK + T(P1) GT 117MEV 8/71		
R37	K+ INTO (PI+ PI+ E- NEU)/(PI+ PI- E+ NEU) (P8)/(P7)		
R37	0 0.013 OR LESS CL=.95 BOURQUIN 71 ASPK 12/71		
R38	CHAR. K INTO (P10 P10 E NEU)/KE3 (UNITS 10**=-6) (P24)/(P6)		
R38	0 37.0 OR LESS CL=.90 ROMANO 71 HLCB 12/71		
R38	2 3.8 5.0 1.2 LJUNG 73 HLCB + 9/73		
R39	K+ INTO (PI- E+ MU-)/TOTAL (UNITS 10**=-8) (P25)		
R39	K- INTO (PI+ E- MU-)/TOTAL IS ALSO INCLUDED HERE		
R39	2.8 OR LESS CL=.90 BEIER 72 OSPK +- 9/72		
R40	K+ INTO (PI+ E+ MU-)/TOTAL (UNITS 10**=-8) (P26)		
R40	K- INTO (PI- E- MU-)/TOTAL IS ALSO INCLUDED HERE		
R40	1.4 OR LESS CL=.90 BEIER 72 OSPK +- 9/72		
R41	CHAR. K INTO (MU 3NEU)/TOTAL (UNITS 10**=-6) (P27)		
R41	P 0 6.0 OR LESS CL=.90 PANG 73 CNTR + 11/73		
R41	P PANG 73 ASSUMES MU SPECTRUM FROM NEU-NEU INTERACTION OF BARDIN 70. 3/74		
R42	CHAR. K INTO (P10 MU NEU GAM)/TOTAL (UNITS 10**=-5) (P28)		
R42	0 6.1 OR LESS CL=.90 LJUNG 73 HLCB + EGAM GT 30 MEV 9/73		
R43	CHAR. K INTO (E P10 NEU)/(PI P10) (P6)/(P2)		
R43	L 786 0.221 0.012 LUCAS 73 HBC - DALITZ PRS ONLY 11/73		
R43	L LUCAS 73 GIVES N(E3)=786+-3.1PCT, N(P12)=3564+-3.1PCT. WE DIVIDE. 11/73		
R43	FIT 0.2291 0.0031 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R44	CHAR. K INTO (PI 2P10)/(PI P10) (P4)/(P2)		
R44	L 574 0.081 0.005 LUCAS 73 HBC - DALITZ PRS ONLY 11/73		
R44	L LUCAS 73 GIVES N(P1 2P10)=574+-5.9 PCT, N(P12)=3564+-3.1 PCT. 11/73		
R44	L WE QUOTE 0.5*(PI 2P10)/(PI P12) WHERE 0.5 IS BECAUSE ONLY DALITZ 11/73		
R44	L PAIR P10'S WERE USED. 11/73		
R44	FIT 0.0822 0.0022 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R45	CHAR. K INTO (MU NEU GAMMA)/TOTAL (UNITS 10**=-3) (P12)		
R45	12 5.8 3.5 WEISSENBE 74 STRC + E-GAMMA GT 9 MEV 7/74*		
R46	CHAR. K INTO (PI+ E- E-)/(PI+ PI- E+ NEU) (UNITS 10**=-3) (P15)/(P7)		
R46	B 41 7.0 1.3 BLOCH 75 SPEC + 11/75*		
R46	B BLOCH 75 QUOTES THIS RESULT MULTIPLIED BY OUR 1974 KE4 BR.FRAC. 11/75*		
R47	CHAR. K INTO (E NEU GAM)/(E NEU) (P21)/(P11)		
R47	H 56 1.05 0.25 HEARD 75 SPEC + P(E) 236 TO 247 11/75*		
R47	H ABOVE IS SENSITIVE ONLY TO STRUCTURE DECAY TERM. 11/75*		

Note on Slope Parameter for K → 3π Decays

As was discussed in Section VI B.1 of the text, for the 3π decays of the K mesons we list the slope parameter "g" which is defined, as in that section, by

$$|M|^2 \propto 1 + g \frac{(s_3 - s_0)}{m_{\pi^+}^2} + h \left(\frac{s_3 - s_0}{m_{\pi^+}^2} \right)^2 + j \frac{(s_2 - s_1)}{m_{\pi^+}^2} + k \left(\frac{s_2 - s_1}{m_{\pi^+}^2} \right)^2 + \dots \quad (1)$$

where

$$s_i = (p_K - p_i)^2 = (m_K - m_i)^2 - 2m_K T_i \quad (2)$$

$$s_0 = \frac{1}{3} \sum s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2) \quad (3)$$

p_K, p_i are the four-vectors for the K and the i^{th} pion, and the index 3 refers to the odd pion, i.e., the third pion in the decays listed below.

We refer to the three possible charged decays as $\tau, \tau',$ and τ^0 :

$$\begin{aligned} \tau^\pm & K^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp \\ \tau^\pm & K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm \\ \tau^0 & K_L^0 \rightarrow \pi^+ \pi^- \pi^0 \end{aligned}$$

The measurements of g vary considerably beyond the authors' quoted errors as can be seen in the ideograms associated with the GT+, GT-, and GTP subsections of the K± Data Card Listings and the GTO subsection of the K_L⁰ Listings. Appendix I discusses tests of the ΔI = 1/2 rule utilizing these slopes.

There is no indication of a CP-violating asymmetry in K_L⁰ decay as measured by the asymmetry parameter σ_± (equivalent to j) given in subsection A of the K_L⁰ Listings.

There is conflicting evidence regarding the presence or absence of the second order terms in Eq. (1), even when only the high-statistics experiments are considered. In K_L⁰ → π⁺π⁻π⁰, MESSNER 74 find a significant h and k, while others find them smaller or consistent with zero. The results on h for the high-statistics experiments are as follows:

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

Experiment	Number of events	Value of h
MESSNER 74	509 K	+0.079 ± 0.007
BUCHANAN 75	56 K	+0.041 ± 0.024
BUCHANAN 70	36 K	consistent with 0
ALBROW 70	29 K	-0.009 ± 0.016
BISI 74	20 K	0.000 ± 0.009

All values other than MESSNER 74 have been converted by us from the authors' parametrizations into the parametrization of Eq. (1) above.

FORD 72 (1.5M events) have studied $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ and find that the χ^2/DF goes from 1.38 to 1.20 for $DF \approx 150$ when the second order and the CP violation terms are added. However, the authors state that since their Coulomb correction is larger than the experimental errors and is not well known, it is difficult to interpret these results. LUCASI 73 (81K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\mp} \pi^{\pm}$ events) note a quadratic behavior in their data, but find that it can be explained by normal measurement error and subsequent fitting, which together result in a depletion of events in the center of the Dalitz plot and a bunching of events near the kinematic limit. Because of this they state that it is not appropriate for them to quote a value of the quadratic coefficient. MAST 69 (51K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\mp} \pi^{\pm}$ events), HOFFMASTER 72 (40K $K^{\pm} \rightarrow \pi^{\pm} \pi^{\pm} \pi^{\mp}$ events) and SMITH 75 (27K $K^{\pm} \rightarrow \pi^0 \pi^0 \pi^{\pm}$ events) also find no evidence for quadratic behavior.

HEUSSE 70 have studied the $K_L^0 \rightarrow \pi^0 \pi^0 \pi^0$ decay where only a second order term could explain deviation from uniformity of the Dalitz plot. They also get results consistent with a zero coefficient.

We list the value of g obtained when the quadratic terms were dropped from the fit unless otherwise noted in the data card footnotes.

In the literature other definitions of slope parameters have appeared. We have converted to the definition of g in Eq. (1) whatever experimental quantity has been reported. We give the conversion to the definition (1) for the most widely used parametrizations and tabulate the conversion factors for the reader's convenience.

a) For analysis of charged K's the expression often used is:

$$|M|^2 = 1 + a_y y$$

with

$$y = \frac{3T_3 - Q}{Q}, \quad Q = m_K - \sum m_i$$

The relevant formulae are:

$$y = -\frac{3}{2} \frac{s_3 - s_0}{m_K Q} + \Delta$$

with

$$\Delta = \frac{m_1 - m_3}{Q} \left(2 - \frac{m_3 + m_1}{m_K} \right)$$

and

$$g = \frac{-c_y a_y \Delta}{1 + a_y \Delta}, \quad \text{with } c_y = \frac{3}{2} \frac{m_{\pi^+}^2}{m_K Q}$$

b) For the analysis of K^0 decay the expression often used is:

$$|M|^2 = 1 + 2a_t \frac{m_K}{m_{\pi^+}^2} (2T_3 - T_{3\max})$$

with

$$T_{3\max} = \frac{m_K^2 + m_3^2 - 4m_{12}^2}{2m_K} - m_3$$

The relevant transformations are

$$T_3 = -\frac{s_3 - s_0}{2m_K} + \frac{Q}{3} (1 + \Delta)$$

and

$$g = \frac{-2a_t}{1 + a_t c_t}$$

with

$$c_t = \frac{2m_K}{m_{\pi^+}^2} \left[\frac{2}{3} Q(1 + \Delta) - T_{3\max} \right]$$

c) Other K^0 authors use the same form of matrix element as given in b) above, but define

$$T_{3\max} = \frac{2}{3} Q$$

Stable Particles

K[±]

Data Card Listings

For notation, see key at front of Listings.

The relevant transformation is then

$$g = \frac{-2a_u}{1 + a_u c_u}, \text{ with } c_u = \frac{4m_K}{3m_{\pi^+}} Q\Delta.$$

d) Older K⁰ analyses were done using

$$|M|^2 = 1 + a_v \frac{T_3}{m_K}.$$

The relevant transformation is then

$$g = \frac{-c_v a_v}{1 + d_v a_v}, \text{ with } c_v = \frac{m_{\pi^+}^2}{2m_K^2}$$

and

$$d_v = \frac{Q}{3m_K} (1 + \Delta).$$

For the reader's convenience we give a table of numerical values for Q, T_{3max}, Δ, c_y, and c_t, obtained using the masses from the current edition.

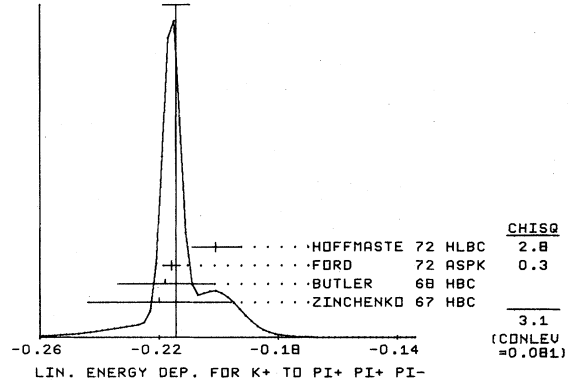
	τ [±]	τ [±]	τ ⁰
Q	75.00	84.21	83.60
T _{3max}	48.10	53.22	53.91
Δ	0.0	-0.0789	0.0798
c _y	0.7891	0.7028	0.7023
c _t	0.0963	-0.0768	0.3204
c _u	0.0	-0.2247	0.2272
c _v	0.0400	0.0400	0.0393
d _v	0.0506	0.0524	0.0605

10 CHARGED K ENERGY DEPENDENCE OF DALITZ PLOT

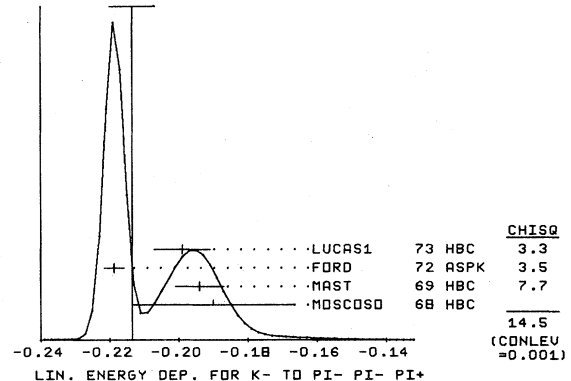
RELATED TEXT SECTION VI B.1, APPENDIX I, AND MINI-REVIEW ABOVE
MATRIX ELEMENT SQUARED = 1 + G (S3-S0)/(MPI**2)

GT+ LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS K+ INTO PI+ PI+ PI-
GT+ THESE EXPTS FIT **2=1+AY*Y. WE LIST G IN THE MAIN LISTING AND
GT+ GIVE AY AT RIGHT. G=-1.5*AY*(MPI**2)/(MK*Q). SEE NOTE ABOVE.
GT+Z 5428 -0.22 0.024 ZINCHENKO 67 HBC + AY=0.28+-03 10/69
GT+ 9994 -0.218 0.016 BUTLER 68 HBC + AY=0.277+-020 10/69
GT+ 617898 (-0.196) (0.012) GRAUMAN 70 HLBC + AY=0.228+-030 8/70
GT+Q 750K -0.2158 0.0028 FORD 72 ASPK + AY=0.2734+-0035 4/72
GT+ 39819 -0.201 0.008 HOFFMASTE 72 HLBC + INCLUDES GRAUMAN 1/71
GT+ Q THIS VALUE OF AY IS FROM A QUADRATIC FIT WITH Y**2 COEF=.030+-010. 4/72
GT+ Q A LINEAR FIT IS QUOTED ONLY FOR THEIR COMBINED K+ AND K- SAMPLE. 4/72
GT+ Q IT GIVES AY=0.2737+-0032. THE QUADRATIC FIT TO THE COMBINED 4/72
GT+ Q SAMPLE GIVES AY=0.2752+-0033 AND Y**2 COEFF=0.025+-010. 4/72
GT+ Q (CHISO/DF)=1.38 FOR LINEAR FIT AND 1.20 FOR QUADRATIC FIT. 1/73
GT+ G EMULS. DATA ADDED - ALL EVENTS INCLUDED BY HOFFMASTE 72 1/71
GT+ Z ALSO INCLUDES DBC EVENTS.
GT+
GT+ AVG -0.2144 0.0045 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)
GT+ STUDENT -0.2146 0.0029 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
(SEE IDEOGRAM BELOW)
GT- LINEAR ENERGY DEPENDENCE (G) FOR TAU DECAYS K- INTO PI- PI- PI+
GT- FOR DEFINITION OF AY SEE NOTE IN ABOVE SECTION GT+.
GT- F 1347 (-0.220) (0.035) FERRO-LUZ 61 HBC - AY=0.28+-045 10/69
GT- M 5778 -0.190 0.023 MSCOSD 68 HBC - AY=0.242+-029 10/69
GT- Q 50919 -0.194 0.007 MAST 69 HBC - AY=0.247+-009 10/69
GT-Q 750K -0.2187 0.0028 FORD 72 ASPK - AY=0.2770+-0035 4/72
GT- 81K -0.199 0.008 LUCAS1 73 HBC - AY=0.252+-011 10/72
GT- Q THIS VALUE OF AY IS FROM A QUADRATIC FIT WITH Y**2 COEF=-0.20+-010. 4/72
GT- Q SEE ALSO THE NOTE Q IN THE GT+ SECTION ABOVE.
GT- F NO RADIATIVE CORRECTIONS INCLUDED.
GT- M ALSO INCLUDES DBC EVENTS.
GT-
GT- AVG -0.2135 0.0066 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.7)
GT- STUDENT -0.2142 0.0039 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = -0.2144 ± 0.0045
ERROR SCALED BY 1.7



WEIGHTED AVERAGE = -0.2135 ± 0.0066
ERROR SCALED BY 2.7

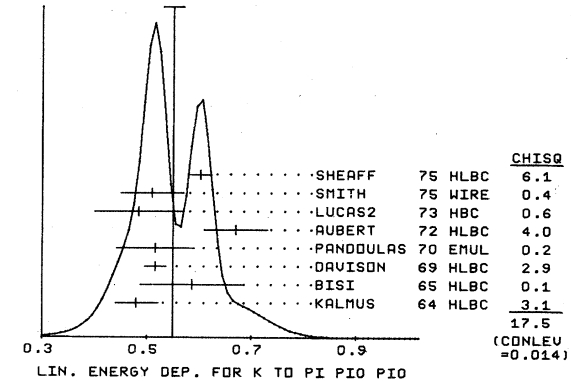


DG ((GT+)-(GT-))/((GT+)+(GT-)) IN PERCENT
DG A NON-ZERO VALUE FOR THIS QUANTITY INDICATES CP VIOLATION
DG 3.2M -0.70 0.53 FORD 70 ASPK 11/70

GTP LINEAR ENERGY DEPENDENCE (G) FOR TAU PRIME DECAY CHA.K INTO PI PIPPIO
GTP 1792 0.48 0.04 KALMUS 64 HLBC + 10/69
GTP 1874 0.586 0.098 BISI 65 HLBC + ALSO HBC 10/69
GTP 4948 0.516 0.020 DAVIDSON 69 HLBC + ALSO EMUL 10/69
GTP 198 0.516 0.074 PANDOLAS 70 EMUL + 10/70
GTP A1365 0.67 0.06 AUBERT 72 HLBC 1/73
GTP 574 0.484 0.084 LUCAS2 73 HBC - DALITZ PRS ONLY 9/73
GTP S 27K 0.510 0.060 SMITH 75 WIRE + 12/75*
GTP 5635 0.602 0.021 SHEAFF 75 HLBC + 2/76*
GTP A WE GIVE LINEAR TERM OF HIGHER ORDER FIT. EQ.1 OF APP.II,AUBERT 72. 1/73
GTP S SMITH 75 MEASURES QUADRATIC COEFFICIENT H=0.009+-0.040. 12/75*

GTP AVG 0.550 0.020 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
GTP STUDENT 0.542 0.019 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.550 ± 0.020
ERROR SCALED BY 1.6



Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

Note on $K_{\ell 3}^{\pm}$ and $K_{\ell 3}^0$ Form Factors

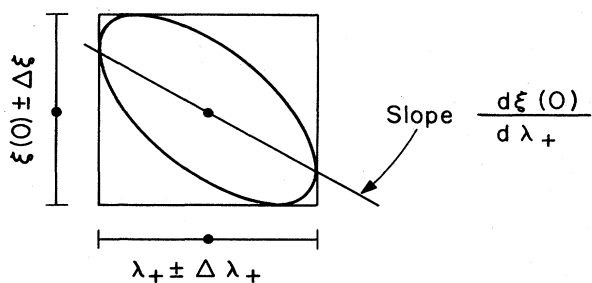
Definitions of the parameters λ_+ , $\xi(0)$, λ_0 , $|f_0/f_+|$ and $|f_T/f_+|$ and a general discussion of the methods of analysis are given in Section VI B.2 of the text.

This note describes the contents of the Data Card Listings for the two $K_{\mu 3}$ parametrizations, $(\lambda_+, \xi(0))$ and (λ_+, λ_0) , which were discussed in the text. Problems related to our data entries for individual experiments are discussed and a comparison of results is given.

$K_{\mu 3}$ Experiments

The matrix element for $K_{\mu 3}$ decay, assuming a pure vector current, is given by Eq. (2) in Section VI B.2 of the text. Most experiments appear to be compatible with the assumption that f_+ depends linearly on t and that f_- is constant. Only DALLY 72 ($K_{\mu 3}^0$) appears to require $\lambda_- \neq 0$ (by about three standard deviations). A single data bin at low q^2 seems to be responsible. The effect is not observed in the high-statistics experiment of DONALDSON 74 (also $K_{\mu 3}^0$).

λ_+ , $\xi(0)$ Parametrization: λ_+ data from $K_{\mu 3}$ decay are entered into the K^{\pm} and K_L^0 sections of the Data Card Listings in subsection L+M. The corresponding $\xi(0)$ values are entered in subsection XIA, XIB, or XIC, depending on whether Method A, B, or C, discussed below and in the text, was used. The data cards contain the values, one-standard-deviation errors $\Delta\lambda_+$ and $\Delta\xi(0)$, as well as the correlation $d\xi(0)/d\lambda_+$, all indicated on the $e^{-1/2}$ likelihood contour below. The correlations are given on the right side of the $\xi(0)$ data cards.



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λ_+, λ_0 Parametrization: This parametrization is used in recent $K_{\mu 3}$ analyses. To facilitate comparison between experiments, we convert earlier experiments from the $(\lambda_+, \xi(0))$ parametrization to (λ_+, λ_0) whenever possible (i.e., when λ_+ and $\xi(0)$ values, errors, and correlations are given). The transformation between these parametrizations is:

$$\lambda_0 = \lambda_+ + a\xi(0) ,$$

$$\Delta\lambda_0^2 = (1 + 2a \frac{d\xi(0)}{d\lambda_+}) \Delta\lambda_+^2 + a^2 \Delta\xi^2 ,$$

$$\frac{d\lambda_0}{d\lambda_+} = 1 + a \frac{d\xi(0)}{d\lambda_+} ,$$

where $a = m_\pi^2 / (m_K^2 - m_\pi^2)$. The λ_0 value, the one-standard-deviation error $\Delta\lambda_0$, and the correlation $d\lambda_0/d\lambda_+$ are given in subsection L0 of the data cards.

We also convert (λ_+, λ_0) results into the $(\lambda_+, \xi(0))$ parametrization whenever possible so that subsection L0 is essentially equivalent to the three subsections XIA, XIB, and XIC.

Individual analyses have used a variety of parametrizations. Problems arise when trying to express their results in terms of the parametrizations used here. The discussion of these problems is divided into three sections corresponding to the three methods of analyses discussed in the text.

Method A: Dalitz plot analyses and pion spectrum analyses usually determine λ_+ and $\xi(0)$ (or λ_0) values, errors, and correlation. Such measurements are entered in the L+M, XIA, and L0 subsections. They give rise to the error ellipses shown in Figs. 1 and 2.

Some analyses of this type fix λ_+ and determine $\xi(0)$, e.g., CARPENTER 66 and PEACH 73 (both $K_{\mu 3}^0$). We enter $\xi(0)$ and $d\xi(0)/d\lambda_+$ in the XIA section and give the fixed λ_+ value in the data card footnote. The $\xi(0)$ error is parenthesized because it does not include the uncertainty in the value of λ_+ . These results, transformed to λ_0 measurements, give rise to bands in Fig. 2.

In some cases, we alter an error from its published value in order to obtain an error ellipse with a width which matches the error in $\xi(0)$ for

Stable Particles

K^\pm

Data Card Listings

For notation, see key at front of Listings.

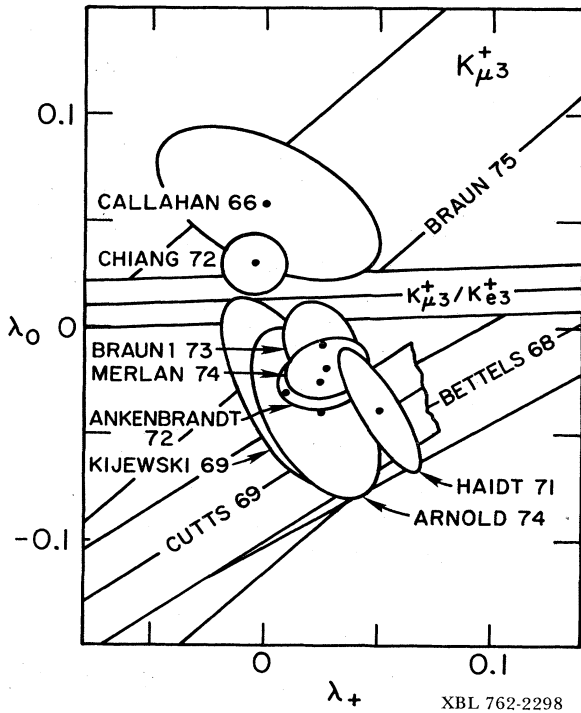


Fig. 1. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$.

fixed λ_+ . These adjustments are noted in the $\xi(0)$ data card footnotes, e.g., for CALLAHAN 66 and HAITD 71 (K^+ subsection XIA), where the published errors and correlation violate the constraint $|C_{\lambda\xi}| < 1$ on the normalized correlation coefficient $C_{\lambda\xi}$ given by

$$C_{\lambda\xi} = \frac{\Delta\lambda_+}{\Delta\xi} \frac{d\xi(0)}{d\lambda_+}$$

In some cases, e.g., BRAUN 73, the parametrization used is $\lambda_+, \xi(0), \xi(t^*)$, where t^* is the weighted average of t with weighting according to the sensitivity to ξ . In this case we do not use $\xi(0)$. It is a badly determined parameter comparable to λ_- or the slope of $\xi(t)$. Instead, we use

$$\xi(0) = \xi(t^*) (1 + \lambda_+ t^*)$$

$$\frac{d\xi(0)}{d\lambda_+} = \frac{d\xi(t^*)}{d\lambda_+} (1 + \lambda_+ t^*) + \xi(t^*) t^*$$

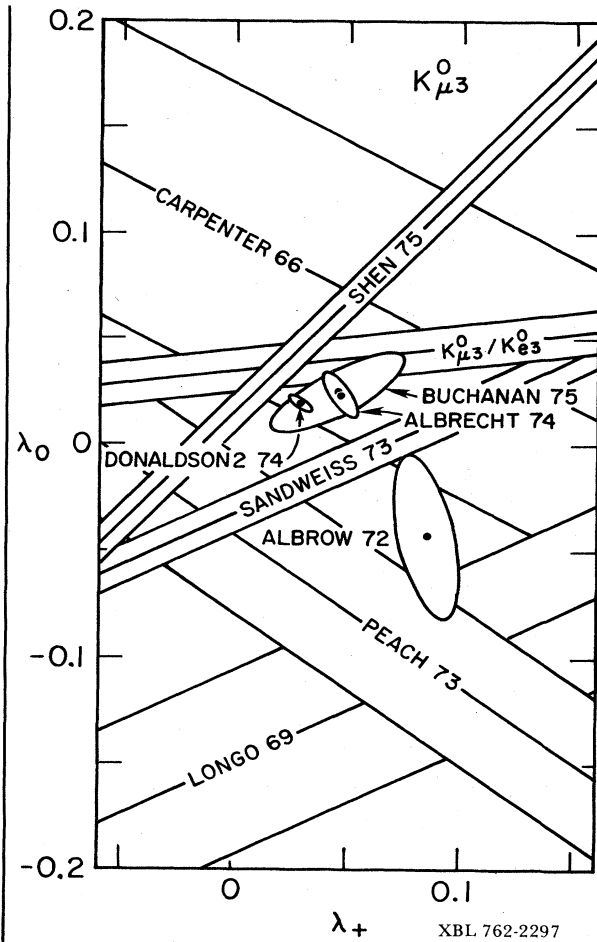


Fig. 2. One-standard-deviation ($e^{-1/2}$) likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^0$.

With the BRAUN 73 values, $\lambda_+ = 0.027$, $\xi(6.6) = -0.34 \pm 0.20$, and $d\xi(6.6)/d\lambda_+ = -14$, we obtain

$$\xi(0) = (-0.40 \pm 0.24) - 19(\lambda_+ - 0.027);$$

or for their fitted $\lambda_+ = 0.025 \pm 0.017$, we get $\xi(0) = -0.36 \pm 0.40$.

Method B: Branching ratio experiments cannot determine λ_+ and $\xi(0)$ simultaneously, but simply fix a relationship between them, given in Section VI B.2 of the text. Results are usually quoted as values of $\xi(0)$ at fixed λ_+ . We list these results in subsection XIB, but we do not average them because the λ_+ values differ. Instead, we compute a combined result by using the relations in the

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K[±]

text and our fitted values of $\Gamma(K_{\mu 3}^+)/\Gamma(K_{e 3}^+)$ and $\Gamma(K_{\mu 3}^0)/\Gamma(K_{e 3}^0)$, which include the branching ratios from these experiments. The branching ratios from our current edition and the results for $\xi(0)$ and λ_0 evaluated at $\lambda_+ = 0.030$ are

	K [±]	K _L ⁰
$\Gamma(K_{\mu 3})/\Gamma(K_{e 3})$	0.663 ± .018 (S=1.7)	0.696 ± .017
$\xi(0)$	-0.20 ± .15 (S=1.7)	+0.09 ± .13
$d\xi(0)/d\lambda_+$	-11.9	-10.3
λ_0	+0.014 ± .012 (S=1.7)	+0.038 ± .011
$d\lambda_0/d\lambda_+$	+0.04	+0.12

The scale factor S is the amount by which the error has been multiplied in order to compensate for discrepancies in the branching ratios. These λ_0 results give rise to the $K_{\mu 3}/K_{e 3}$ bands in Figs. 1 and 2.

Method C: Polarization experiments measure $\langle \xi(t) \rangle$, the weighted average of $\xi(t)$ over the t range of the experiment, where the weighting accounts for the variation with t of the sensitivity to $\xi(t)$. Such measurements are entered in subsection XIC.

To reinterpret these results in the $(\lambda_+, \xi(0))$ parametrization, we recognize that $\lambda_+ = 0$ corresponds to $\xi(t)$ constant (always assuming $\lambda_- = 0$) so that

$$\xi(0) \Big|_{\lambda_+=0} \equiv \langle \xi(t) \rangle .$$

The correlation with λ_+ is given by the following relations (valid for small λ_+):

$$\xi(0) \approx \langle \xi(t) \rangle \left(1 + \lambda_+ \left\langle \frac{t}{m_\pi^2} \right\rangle \right) ,$$

$$\frac{d\xi(0)}{d\lambda_+} \approx \langle \xi(t) \rangle \left\langle \frac{t}{m_\pi^2} \right\rangle ,$$

where $\langle t/m_\pi^2 \rangle$ is the average value of t weighted by the sensitivity to $\xi(t)$. These results, transformed to λ_0 and $d\lambda_0/d\lambda_+$ values, are entered in subsection I0 and give rise to bands in Figs. 1 and 2.

In Figs. 1 and 2, we disregard those polarization measurements for which $d\xi(0)/d\lambda_+$ is not obtainable. Also we disregard MERLAN 73 because

the signs of $\xi(0)$ and $d\xi(0)/d\lambda_+$ are opposite, whereas the above equation requires them to be the same (since $t > 0$).

Comparison of $K_{\mu 3}$ Experiments: Figures 1 and 2 show the likelihood contours in the (λ_+, λ_0) plane for $K_{\mu 3}^+$ and $K_{\mu 3}^0$ respectively.

The $K_{\mu 3}^+$ Dalitz plot results (ellipses) shown are fairly consistent and appear to cluster between the $K_{\mu 3}/K_{e 3}$ result and the polarization results of BETTELS 68 and CUTTS 69. The $K_{\mu 3}^0$ results are much less consistent with a small cluster appearing in the neighborhood of the DONALDSON2 74 result.

χ^2 fits to the results shown in Fig. 1 and Fig. 2 yield the following values for λ_+ and λ_0 . The corresponding values of $\xi(0)$ are also given.

	K _{μ3} ⁺	K _{μ3} ⁰
λ_+	+0.027 ± .008*	+0.034 ± .006*
λ_0	-0.009 ± .007*	+0.021 ± .006*
$d\lambda_0/d\lambda_+$	-0.15	-0.24
χ^2/DF	34/18	72/12
S	1.4	2.5
.....		
$\xi(0)$	-0.45 ± .14*	-0.17 ± .10*
$d\xi(0)/d\lambda_+$	-14.	-15.

*All errors have been increased by the scale factor $S = (\chi^2/DF)^{1/2}$ to take into account the discrepancies between measurements.

In view of the large χ^2/DF of these fits, especially $K_{\mu 3}^0$, the fit results should be taken with a grain of salt. The largest contributors to χ^2 in the $K_{\mu 3}^+$ case are CHIANG 72 with 10.4, and the polarization results, BETTELS 68 with 5.4 and CUTTS 69 with 4.2. In the $K_{\mu 3}^0$ case the largest contributors are the polarization results of SANDWEISS 73 with 17, LONGO 69 with 14, and SHEN 75 with 12, and the Dalitz plot results of ALBROW 72 with 11, ALBRECHT 74 with 7.3, and PEACH 73 with 5.3. All other χ^2 values were less than 4.

The DONALDSON2 74 result

$$\lambda_+ = .030 \pm .003$$

$$\lambda_0 = .019 \pm .004$$

Stable Particles

K±

clearly dominates the statistics in the $K_{\mu 3}^0$ case. The λ_+ value is consistent with the $K_{e 3}$ value of λ_+ , and with the pole approximation

$$f_+(t) = f_+(0) \frac{m_{K^*}^2}{m_{K^*}^2 - t}$$

Their $f_0(t)$ extrapolates linearly to the Callan-Treiman point. It is less than two standard deviations from the $K_{\mu 3}/K_{e 3}$ result.

$K_{e 3}$ Experiments

The f_- term of the matrix element [Eq. (2) text Section VI.B.2] can be neglected for $K_{e 3}$ because it is proportional to the lepton mass. The f_+ term is usually assumed to be linear in $t = q^2 = (P_K - P_\pi)^2$, the square of the four-momentum transfer, i.e., the effective mass of the lepton pair. We quote the linear coefficient λ_+^e (L+E on the data cards).

There has been some suggestion of departure from linearity [CHIEN 71 ($K_{e 3}^0$) and Chounet, Gaillard, and Gaillard¹ - Review] but no compelling evidence. The λ_+ results are fairly consistent and the average values are

$$K_{e 3}^+ : \lambda_+ = 0.0285 \pm 0.0043$$

$$K_{e 3}^0 : \lambda_+ = 0.0288 \pm 0.0028 (S=1.4)$$

where the $K_{e 3}^0$ error has been multiplied by the scale factor 1.4 to compensate for inconsistencies (see ideogram in K_L^0 section L+E).

Many of the changes in this section were stimulated by the comments and criticism of L. M. Chounet, and by the excellent reviews of Gaillard and Chounet¹ and Chounet, Gaillard, and Gaillard.²

References

1. M. K. Gaillard and L. M. Chounet, $K_{L 3}$ Form Factors, CERN 70-14 (May 1970), and Phys. Letters 32B, 505 (1970).
2. L. M. Chounet, J. M. Gaillard, and M. K. Gaillard, Physics Reports 4C, 199 (1972).

Data Card Listings

For notation, see key at front of Listings.

10 CHARGED K FORM FACTORS		8/67
RELATED TEXT SECTION VI B.2 AND MINI-REVIEW ABOVE.		
IN THE FORM FACTOR COMMENTS, THE FOLLOWING ABBREVIATIONS ARE USED. F+ AND F- ARE FORM FACTORS FOR THE VECTOR MATRIX ELEMENT. FS AND FT REFER TO THE SCALAR AND TENSOR TERM. FO = (D+)* F(F+)*T*(KMU3)-MFI**2 L+, L- AND LO ARE THE LINEAR EXPANSION COEFFS. OF F+, F- AND FO. L+ REFERS TO THE KMU3 VALUE EXCEPT IN THE KE3 SECTIONS. DXI/DL IS THE CORRELATION BETWEEN XI(0) AND L+ IN KMU3. DLO/DL IS THE CORRELATION BETWEEN LO AND L+ IN KMU3. T = MOMENTUM TRANSFER TO THE PI IN UNITS OF MPI**2. DP = DALITZ PLOT ANALYSIS PI = PI SPECTRUM ANALYSIS MU = MU SPECTRUM ANALYSIS POL= MU POLARIZATION ANALYSIS BR = KMU3/KE3 BRANCHING RATIO ANALYSIS E = POSITRON OR ELECTRON SPECTRUM ANALYSIS RC = RADIATIVE CORRECTIONS		
XIA	XIA = F-/F+ (DETERMINED FROM SPECTRA)	-----
XIA	76 (+1.8) (0.6) BROWN 62 XEBC + DP+BR, L+=0	1/74
XIA	87 (+0.7) (0.5) GIACOMELLI 64 EMUL + MU+BR RVUE, L+=0	1/74
XIA J	2648 (-0.08) (0.7) JENSEN 64 XEBC + DP+BR(KMU3,KE3)	1/74
XIA	2648 (0.0) (1.1) CALLAHA1 66 FRBC + MU, L+=0, T UNKN	1/74
XIA C	444 +0.72 0.93 (0.9) CALLAHA1 66 FRBC + PI, DXI/DL=-17	1/74
XIA	78 (-0.5) (0.9) EISLER 68 HLBC + PI, L+=0, NO DX/DL	1/74
XIA	K2041 -0.5 0.8 KIJEWSKI 69 OSPK + PI, DXI/DL=-26	1/74
XIA	H3240 -1.1 0.56 HADIT 71 HLBC + DP, DXI/DL=-29	1/74
XIA	A4025 -0.42 0.28 ANKENBRAN 72 ASPK + PI, DXI/DL=-12	1/74
XIA	B3480 +0.45 0.28 CHIANG 72 OSPK + DP, DXI/DL=-15	1/74
XIA	D1897 -0.36 0.40 BRAUNI 73 HLBC + DP, DXI/DL=-19	3/74
XIA	N 490 -0.8 0.8 ARNOLD 74 HLBC + DP, DXI/DL=-20	11/75*
XIA	M6527 -0.57 0.24 MERLAN 74 ASPK + DP, DXI/DL=-9	3/74
XIA J	JENSEN 64 GIVES L+M=L+E=-.020+-.027. DXI/DL UNKNOWN. INCLUDES	1/74
XIA J	SHAKLEE 64 XIB(KMU3/KE3).	1/74
XIA C	CALLAHAN 66 TABLE 1 (PI ANAL) GIVES DXI/DL=(-.72-.05)/(0-.04)=-17,	1/74
XIA	C ERROR RAISED FROM .80 TO AGREE WITH DXI/DL=.37 FOR FIXED L+.	1/74
XIA K	KIJEWSKI 69 FIG. 17 WAS USED TO OBTAIN DXI/DL AND ERRORS.	1/74
XIA H	HADIT 71 TABLE 8 (DP ANAL) GIVES DXI/DL=(-1.1+.05)/(0.050-.029)=-29,	1/74
XIA H	H ERROR RAISED FROM .50 TO AGREE WITH DXI/DL=.20 FOR FIXED L+.	1/74
XIA A	ANKENBRANDT 72 FIG. 3 WAS USED TO OBTAIN DXI/DL.	1/74
XIA B	CHIANG 72 FIG. 10 WAS USED TO OBTAIN DXI/DL.	1/74
XIA B	FIT HAD L=L+ BUT WOULD NOT CHANGE FOR L=0. L.PONDROM,PRIV.COM.74	3/74
XIA D	BRAUNI 73 GIVES XI(T)=-.34+-.20, DXI(T)/DL=-.14 FOR L+=.027, T=6.6.	1/74
XIA D	WE CALCULATE ABOVE XI(0) AND DXI(0)/DL+ FOR THEIR L+=.025+-.017.	3/74
XIA N	ARNOLD 74 FIG. 4 WAS USED TO OBTAIN XI AND DXI/DL.	11/75*
XIA M	MERLAN 74 FIG.5 WAS USED TO OBTAIN DXI/DL.	3/74
XIA	
XIA	FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS ABOVE.	
XIB	XIB = F-/F+ (DETERMINED FROM KMU3/KE3)	-----
XIB	THE KMU3/KE3 BRANCHING RATIO FIXES A RELATIONSHIP BETWEEN XI(0)	1/74
XIB	AND L+. WE QUOTE THE AUTHORS XI(0) AND ASSOCIATED L+ BUT DO NOT	1/74
XIB	AVERAGE BECAUSE THE L+ VALUES DIFFER. THE FIT RESULT AND SCALE	1/74
XIB	FACTOR GIVEN IN OUR NOTE ON KL3 FORM FACTORS ARE NOT OBTAINED	2/76*
XIB	FROM THESE XIB VALUES. INSTEAD THEY ARE OBTAINED DIRECTLY FROM THE	2/76*
XIB	FITTED KMU3/KE3 RATIO (R29).	
XIB	-0.17 0.75 0.99 SHAKLEE 64 XEBC + BR, L+=0	1/74
XIB	+0.6 0.5 BISI 1 65 HBC + BR, L+=0	1/74
XIB	500 +0.8 0.4 CUTTS 65 OSPK + BR, L+=0	1/74
XIB	636 +0.4 0.4 CALLAHA1 66 FRBC + BR, L+=0	1/74
XIB	306 +0.75 0.50 AUERBACH 67 OSPK + BR, L+=0	1/74
XIB B	5601 (-0.08) (0.15) BOTTERIL2 68 ASPK + BR, L+=.023+-.008	1/74
XIB E	1398 (-0.60) (0.20) EICHTEN 68 HLBC + BR, SEE NOTE E	1/74
XIB	986 +1.0 0.6 GARLAND 68 OSPK + BR, L+=0	1/74
XIB	+0.91 0.82 ZELLER 69 ASPK + BR, L+=.023	1/74
XIB B	-0.35 0.22 BOTTERIL 70 OSPK + BR, L+=.045+-.015	1/74
XIB E	1505 -0.81 0.27 HADIT 71 HLBC + BR, L+=.025, FIG.8	1/74
XIB	5825 0.0 0.15 CHIANG 72 OSPK + BR, L+=.03, FIG.10	1/74
XIB B	BOTTERIL 70 IS REEVALUATION OF BOTTERIL 2 68 WITH DIFFERENT L+.	1/74
XIB E	EICHTEN 68 HAS L+=.023+-.008, T=4, INDEP. OF L-. REPL. BY HADIT 71.	1/74
XIB	
XIB	FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS ABOVE.	
XIC	XIC = F-/F+ (DETERMINED FROM MU POLARIZATION IN KMU3)	-----
XIC	THE MU POLARIZATION IS A MEASURE OF XI(T). NO ASSUMPTIONS ON L+	
XIC	NECESSARY. T (WEIGHTED BY SENSITIVITY TO XI(0)) SHOULD BE SPECIFIED.	
XIC	IN L+XI(0) PARAMETERIZATION THIS IS XI(0) FOR L+=0. DXI/DL=XI* ⁺ T.	
XIC	FOR RAD. CORR. TO MUON POLARIZATION IN KMU3 SEE GINSBERG 71.	
XIC T	2100 (+1.2) (2.4) (1.8) BORREANI 65 HLBC + POLARIZATION	8/67
XIC T	500 BTWN -4.0 AND +1.7 CUTTS 65 OSPK + LONG. POL.	1/74
XIC T	397 (-1.4) (1.8) CALLAHA1 66 FRBC + TOTAL POL.	8/67
XIC T	2950 (-0.7) (0.9) (3.3) CALLAHA1 66 FRBC + LONG. POL.	8/67
XIC	86000 -1.0 0.3 BETTELS 68 HLBC + TOTAL POL. T=4.9	1/74
XIC	C3133 -0.95 0.3 CUTTS 69 OSPK + TOTAL POL. T=4.0	1/74
XIC M	40K (-0.64) (0.27) MERLAN 74 ASPK + POL,DXI/DL=-1.7	3/74
XIC	D1585 -0.25 1.20 BRAUN 75 HLBC + POL. T=4.2	1/76*
XIC T	T VALUE NOT GIVEN.	
XIC B	BETTELS 68 DXI/DL=XI* ⁺ T=-1.0*4.9=-4.9.	
XIC C	CUTTS 69 T=4.0 WAS CALCULATED FROM FIG.8. DXI/DL=XI* ⁺ T=-.95*4=-3.8.	1/74
XIC M	MERLAN 74 POLARIZATION RESULT (FIG.5) NOT POSSIBLE. SEE DISCUSSION	1/76*
XIC M	OF POLARIZATION EXPERIMENTS IN NOTE ON KL3 FORM FACTORS ABOVE.	1/76*
XIC D	BRAUN 75 DXI/DL=XI* ⁺ T=-.25*4.2=-1.0.	1/76*
XIC	
XIC	FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS ABOVE.	
IXI	IMAGINARY PART OF XI (TEST OF T REVERSAL)	-----
IXI	2648 0.0 1.0 CALLAHA1 66 FRBC + MU	1/74
IXI	397 +1.6 1.3 CALLAHA1 66 FRBC + TOTAL POL.	1/74
IXI	2950 0.5 1.4 0.5 CALLAHA1 66 FRBC + LONG. POL.	1/74
IXI	6000 -0.1 0.3 BETTELS 68 HLBC + TOTAL POL.	1/74
IXI	3133 -0.3 0.3 0.4 CUTTS 69 OSPK + TOTAL POL. FIG.7	1/74
IXI	AVG -0.09 AVERAGE ERROR INCLUDES SCALE FACTOR OF 1.0	
IXI	STUDENT -0.10 0.23 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
L+M	LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KMU3 DECAY)	-----
L+M	SEE ALSO THE CORRESPONDING ENTRIES AND FOOTNOTES IN SECTIONS XIA,	
L+M	XIC, AND LO.	
L+M	FOR RAD.CORR. OF KMU3 DP SEE GINSBURG 70 AND BECHERRAWY 70.	3/74
L+M	444 0.0 0.05 CALLAHA1 66 FRBC + PI	1/74
L+M	2041 0.009 0.026 KIJEWSKI 69 OSPK + PI	1/74
L+M	3240 0.050 0.018 HADIT 71 HLBC + DP	1/74
L+M	A4025 0.024 0.019 ANKENBRAN 72 ASPK + PI	1/74
L+M	3480 -0.006 0.015 CHIANG 72 OSPK + DP	1/74
L+M	1897 0.025 0.017 BRAUNI 73 HLBC + DP	3/74
L+M	490 0.025 0.030 ARNOLD 74 HLBC + DP	11/75*
L+M	6527 0.027 0.019 MERLAN 74 ASPK + DP	3/74
L+M	A ANKENBRANDT 72 L+ FROM FIG.3 TO MATCH DXI/DL. TEXT GIVES .024+-.022	1/74
L+M	
L+M	FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS ABOVE.	

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K±

LAMBDA 0 (LINEAR ENERGY DEPENDENCE OF FO IN KMU3 DECAY)
WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(O) INTO
VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KE3 DECAY)
FOR RAD. COR. OF KE3 DP SEE GINSBURG 67 AND BECHERRAWY 70.
L+ E 217 +0.036 +0.45 BROWN 62 XEBC + P1, NO RC
L+ E 407 -0.010 +0.29 JENSEN 64 XEBC + P1, NO RC

FS/F+ RATIO OF SCALAR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE)
.18 OR LESS CL=.90 BELLOTTE 67 HLBC 10/69
.30 OR LESS CL=.95 KALMUS 67 HLBC + 10/69

FT/F+ RATIO OF TENSOR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE)
.1 OR LESS CL=.95 KALMUS 67 HLBC + 10/69
.58 OR LESS CL=.90 BOTTERILL 68 ASPK 8/66

FTM FT/F+ RATIO OF TENSOR TO F+ COUPLINGS FOR KMU3 DECAY
1585 0.02 0.12 BRAUN 75 HLBC 1/76*

REFERENCES FOR CHARGED K

BIRGE 56 NC 4 834
LDFD 56 PR 102 927
ALEXANDER 57 NC 6 478
COHEN 57 FUND. CONS. PHYS.
EISENBERG 58 NC 8 663
BURROUGHS 59 PRL 2 117
TAYLOR 59 PR 114 359
FREDEN 60 PR 118 564
BARKAS 61 PR 124 1209
BHOWMIK 61 NC 20 857
FERRO-LU 61 NC 22 1087
NORDIN 61 PR 123 2166
ROE 61 PRL 7 346
BOYARSKI 62 PR 128 2398
BROWN 62 PRL 8 450
BARKAS 63 PRL 11 26
BORREANI 64 PL 12 123
CALLAHAN 64 PR 136 B 1463
CAMERINI 64 PRL 13 318
CLINE 64 PRL 13 101
GIACOMELLI 64 NC 34 1134
GREINER 64 PR 139 2884
JENSEN 64 PR 136 B1431
KALMUS 64 PRL 13 99
SHAKLEE 64 PR 136 B 1423
BIRGE 65 PR 139 B 1600
BIST 1 65 PR 139 B 1068
BORREANI 65 PR 139 B 440
CALLAHAN 65 PRL 15 129
CAMERINI 65 NC 37 1795
CLINE 65 PL 15 293
CUTTS 65 PR 138 8969
DE MARCO 65 PR 140 B 1430
FITCH 65 PR 140 B 1088
GREINER 65 ARNS 15 67
STAMER 65 PR 139 B 440
TRILLING 65 UCRL 16473
UPDATED FROM 1965 ARGONNE
YOUNG 65 UCRL 16362
ALSO 67 PR 156 1464
BIRGE, PERKINS, PETERSON, STORK, WHITEHEA (LRL)
ILOFF, GOLDHABER, LAMNUTTI, GILBERT (LRL)
ALEXANDER, JOHNSTON, OCEALLAIGH (DUBLIN INST)
+CROWE, DUMOND (ATOMICS INTER.+LRL+CIT)
EISENBERG, KOCH, LOHRMANN, NIKOLIC + (BERN)
BURROUGHS, CALDWELL, FRISCH, HILL + (MIT)
S TAYLOR, HARRIS, OREAR, LEE, BAUMEL (COLUMBIA)
S C FREDEN, F C GILBERT, R S WHITE (LRL)
BARKAS, DYER, HANSON, NORDIN, NICKOLS, SMIT (LRL)
B BHOWMIK, P C JAIN, P C MATHUR (DELHI UNIV)
FERRO-LUZZI, MILLER, MURRAY, ROSENFELD+ (LRL)
PAUL NORDIN JR (LRL)
ROE, SINCLAIR, BROWN, GLASER + (MICH+LRL)
BOYARSKI, LUDWIG, NEMELA, RITSON (MIT)
BROWN, KADYK, TRILLING, ROE+ (LRL, MICH)
W H BARKAS, J N DYER, H H HECKMAN (LFL)
G BORREANI, G RINAUDO, A WERBROUCK (TURIN)
A CALLAHAN, R MARCH, R STARK (WISCONSIN)
CAMERINI, CLINE, FRY, POWELL (WISCONSIN+LRL)
D CLINE, W F FRY (LRL)
GIACOMELLI, MONTI, QUARENI + (BOLOGNA, MUNICH)
D GREINER, W OSBORNE, R BARKAS (LRL)
JENSEN, SHAKLEE, ROE, SINCLAIR (MICH)
+KERNAN, J U, POWELL, DOWD (LRL, WISC)
SHAKLEE, JENSEN, ROE, SINCLAIR (MICH)
BIRGE, ELY, GIDAL, CAMERINI, CLINE + (LRL+WISC)
BISI, BORREANI, CESTER, FERRARO + (TORINO)
BORREANI, MARZARI, CHIESA, R RINAUDO+ (TORINO)
D GREINER, LUDWIG, NEMELA, CAFORIO+ (BARI, TORI)
A CALLAHAN, D CLINE (WISCONSIN)
+CLINE, GIDAL, KALMUS, KERNAN (WISC+LRL)
A CLINE, W F FRY (WISCONSIN)
CUTTS, ELIOFF, STIENING (LRL)

CALLAHAN 66 PR 150 1153
CALLAHAN 66 NC 444 90
CESTER 66 PL 21 343
ALSO 67 AUERBACH, FOOTNOTE 1.
AUERBACH 67 PR 155 1505
ERRATUM
BELLOTTI 67 HEIDELBERG CONF
BELLOTTI 67 NC 52A 1287
ALSO 66 PL 20 690
BISI 67 PL 25B 572
BOTTERILL 67 PRL 19 982
ALSO 68 BOTTERILL
BOWEN 67 PR 154 1314
CLINE1 67 HEIDELBERG CONF
CLINE2 67 HERCEG NOVI TBL.4
FLETCHER 67 PRL 19 98
FORD 67 PRL 18 1214
IMLAY 67 PRL 160 1203
KALMUS 67 PR 159 1187
ZINCHENKO 67 RUTGERS (THESES)
BETTELS 68 NC 56A 1106
ALSO 71 HAIDT
BOTTERILL 68 PR 171 1402
BOTTERILL 68 PR 174 1661
BOTTERILL 68 PRL 21 766
BUTLER 68 UCRL-18420
CHANG 68 PRL 20 510
CHEN 68 PRL 20 73
EICHEN 68 PL 27B 586
EISLER 68 PR 169 1090
ESCHSTRUP 68 PR 165 1487
GARLAND 68 PR 167 1225
MOSCOSO 68 THESIS
CUTTS 69 PR 184 1380
ALSO 68 PRL 20 955
DAVISON 69 PR 180 1333
ELY 69 PR 180 1319
EMMERSON 69 PRL 23 393
HERZO 69 PR 186 1403
KJEWESKI 69 UCRL-18433 THESIS
LDBKOWICZ 69 PR 185 1676
ALSO 66 PRL 17 548
MACER 69 PRL 22 32
MAST 69 PR 183 1200
ZELLER 69 PR 182 1420
BOTTERILL 70 PL 318 325
FORD 70 PRL 25 1370
GRAUMAN 70 PR D1 1277
ALSO 69 PRL 23 737
MACER 70 PR D1 249
MALTSEV 70 SUNJ 10 678
PANDOLA 70 PR D2 1205
BASILE 71 PL 368 619
BOURQUIN 71 PL 368 615
HAIDT 71 PR D3 10
ALSO 69 PL 29B 691
KLEMS 71 PR D4 66
ALSO 70 PRL 26 1086
ALSO 70 PRL 25 473
KUNSELMA 71 PL 368 485
OTT 71 PR D5 52
ROMANO 71 PL 368 525
SCHWEINB 71 PL 368 246
STEINER 71 PL 368 521
ABRAMS 72 PRL 29 1118
ANKENBRA 72 PRL 28 1472
AUBERT 72 NC 12A 509
BETER 72 PRL 29 678
CHIANG 72 PR D6 1254
CLARK 72 PRL 29 1274
EDWARDS 72 PR D5 2720
FORD 72 PL 388 335
HOFFMAST 72 NP 836 1
ABRAMS 73 PRL 30 500
BACKENST 73 PL 438 431
BETER 73 PRL 30 399
BRAUN1 73 PL 478 182
ALSO 75 BRAUN
BRAUN2 73 PL 478 185
ALSO 75 BRAUN
CABLE 73 PR D8 3807
LJUNG 73 PR D8 1307
ALSO 72 PRL 28 523
ALSO 72 PRL 28 1287
ALSO 69 PRL 23 326
LUCAS1 73 PR D8 719
LUCAS2 73 PR D8 727
PANG 73 PR D8 1989
ALSO 72 PL 408 699
SMITH 73 NP 860 411
SRNLD 74 PR D9 1221
BRAUN 74 PL 518 393
CENCE 74 PRL 25 2570
ALSO 73 THESIS (UNPUBL.)
KUNSELMA 74 PR C9 2469
MERLAN 74 PR D9 107
WEISENB 74 PL 488 474
BLOCH 75 PR 568 201
BRAUN 75 NP 889 210
HEARD1 75 PL 558 324
HEARD2 75 PL 558 327
SHEAFF 75 PR D12 2570
SMITH 75 NP 891 45
HEINTZE 76 PL (TO BE PUBL.)

CALLAHAN, CAMERINI + (WISC, LRL, RIVERSIDE, BARI)
A C CALLAHAN (WISCONSIN)
CESTER, ESCHSTRUP, ONEILL + (PRINCETON-PENN)
+DOBBS, MANN, MCFARLANE, WHITE+ (PENN, PRIN)
ERRATUM
BELLOTTI, PULLIA (MILAN)
BELLOTTI, FIORINI, PULLIA (MILAN)
BELLOTTI, FIORINI, PULLIA+ (MILAN)
BISI, CESTER, CHIESA, VIGONE (TORINO)
BOTTERILL, BROWN, CORBETT, CULLIGAN + (OXFORD)
BOWEN, MANN, MCFARLANE, HUGHES + (PENN-PRINCETO)
CLINE, HAGGERTY, SINGLETON, FRY+ (WISCONSIN)
D. CLINE, PROC. INT'L. SYM. ON ELEM. PART. PHYSICS
FLETCHER, BEIER, EDWARDS + (ILLINOIS)
+LEMONICK, NAUENBERG, PIROUE (PRINCETON)
IMLAY, ESCHSTRUP, FRANKLIN+ (PRINCETON)
KALMUS, KERNAN (LRL)
ZINCHENKO RUTGERS (THESES)
AACHEN-BARI-IBERGEN-CERN-EP-NIJMEGEN-ORSAY+
BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
BOTTERILL, BROWN, CLEGG, CORBETT + (OXFORD)
+BLAND, GOLDHABER, GOLDHABER, HIRATA+ (LRL)
CHANG, YODH, EHRLICH, PLAND+ (MARYLAND, RUTGERS)

CHEN, CUTTS, KJEWESKI, STIENING + (LRL, MIT)
AACHEN-BARI-CERN-EP-ORSAY-PADOVA-VALENCIA
EISLER, FUNG, MARATECK, MEYER, PLANO (RUTGERS)
ESCHSTRUP, FRANKLIN, HUGHES + (PRINCETON, PENN)
+KLEMS, MANN, MCFARLANE, ROBERTS + (PENN, TEMPLE)
+TSIPIS, DEVONS, ROSEN+ (COLUMBIA, RUTG, MICH)
M L MOSCOSO (UNIV PARIS ORSAY)
+STIENING, WIEGAND, DEUTSCH (LRL, MIT)
CUTTS, STIENING, WIEGAND, DEUTSCH (LRL, MIT)
+BACASTOW, BARKAS, EVANS, FUNG, PORTER+ (UCR)
ELY, GIDAL, HAGOPIAN, KALMUS+ (LOUCI+WISC+LRL)
EMMERSON, QUIRK (OXFORD)

+BANNER, BETER, BERTRAM, EDWARDS + (ILL)
P K KJEWESKI (LRL)
+MELISSINOS, NAGASHIMA, TEWSBURY+ (ROCH, BNL)
LDBKOWICZ, MELISSINOS, NAGASHIMA+ (ROCH+BNL)
MANN, MCFARLANE, ROBERTS + (PENN, TEMPLE)
+GERSHWIN, ALSTON-GARJOST, BANGERTER+ (LRL)
ZELLER, HADDOCK, HELLAND, PAHL+ (UCLA, LRL)
+BROWN, CLEGG, CORBETT, CULLIGAN+ (OXF)
+PIROUE, REMMEL, SMITH, SQUIDER (PRIN)
+KOLLER, TAYLOR, PANDOLA+ (STEVE, SETO, LEHI)
+KOLLER, TAYLOR, PANDOLA+ (STEVE, SETO, LEHI)
+MANN, MCFARLANE, ROBERTS (PENN)
+PESTOVA, SOLODOVNIKOVA, FADEEV + (JINR)
+TAYLOR, KOLLER, GRAUMAN + (STEVE, SETO)

+BREHIN, DIAMANT-BERGER, KUNZ+ (SACL+GEVA)
+BOYMOND, EXTERMANN, MARASCO+ (GEVA, SACL)
AACHEN+BARI-CERN+EP-NIJMEGEN+ORSAY+PADOVA+
+(AACH, BARI, CERN, EPOL, NIJMEGEN, ORSAY, PADO, TORI)
+HILDEBRAND, STIENING (CHIC, LRL)
KLEMS, HILDEBRAND, STIENING (LRL, CHIC)
KLEMS, HILDEBRAND, STIENING (LRL, CHIC)

R. KUNSELMAN (WYOMING)
OTT, PRIETZARD
+RENTON, AUBERT, BURBAN-LUTZ (BARI, CERN, ORS)
AACHEN+BELGIUM+CERN+NIJMEGEN+PADOVA+COLUMB
AACHEN+BARI-CERN+EPOL+ORSAY+NIJMEGEN+PADO+TORI
+CARROLL, KYCIA, LI, MENES, MICHAEL + (BNL)
ANKENBRANDT, LARSEN+ (BNL+LRS+FNAL+YALE)
+HEUSSE, PASCAUD, VIALLE+ (ORSAY+BRUX+EPOL)
+BUCHHOLZ, MANN, PARKER (PENNSYLVANIA)
+ROSEN, SHAPIRO, HANDLER, OLSEN+ (ROCH+WISC)

+CORK, ELIOFF, KERTH, MCREYNOLDS, NEWTON+ (LRL)
+BETER, BERTRAM, HERZO, KOESTER + (ILL)
+PIROUE, REMMEL, SMITH, SQUIDER (PRINCETON)
HOFFMASTER, KOLLER, TAYLOR+ (STEVE+SETO+LEHI)
+CARROLL, KYCIA, LI, MENES, MICHAEL + (BNL)
BACKENSTOSS, BAMBERGER+ (CERN+KARL+HEID+STIC)
+BUCHHOLZ, MANN, PARKER, ROBERTS (PENN)
AACHEN+BARI+BRUSSELS+CERN COLLABORATION
AACHEN+BARI+BRUSSELS+CERN COLLABORATION
+HILDEBRAND, PANG, STIENING (EFI+LBL)

D. LJUNG, D. CLINE (WISC)
D. LJUNG (WISC)
D. CLINE, D. LJUNG (WISC)
CAMERINI, LJUNG, SHEAFF, CLINE (WISC)
LUCAS, TAFT, WILLIS (YALE)
LUCAS, TAFT, WILLIS (YALE)
+HILDEBRAND, CABLE, STIENING (EFI+ARIZ+LBL)
CABLE, HILDEBRAND, PANG, STIENING (EFI+LBL)
+BOOTH, RENSHALL, JONES+ (GLAS+IVP+OXF+RHEL)
C L ARNOLD, S P ROE, D SINCLAIR (MICH)
+CORNELIUSSEN, MARTYN + (AACH+BARI+BRUX+CERN)
+HARLIS, JONES, MORGADD + (HAMA+LBL+WISC)
D B CLARKE (WISC)
R. KUNSELMAN (WYOM)
+KASHA, WANDERER, ADARI+ (YALE+BNL+LRS)
WEISENBERG, EGOROV, MINERVA + (ITEP+LEBD)
+BREHIN, BUNCE, DEVAUX+ (SACL+GEVA)
+CORNELIUSSEN, MARTYN+ (AACH+BARI+BRUX+CERN)
+HEINTZE, HEINZELMANN+ (CERN+HEID)
+HEINTZE, HEINZELMANN+ (CERN+HEID)
M. SHEAFF (WISC)
+BOOTH, RENSHALL, JONES+ (GLAS+IVP+OXF+RHEL)
+HEINZELMANN, IGO-KEMENES, MUNDHENKE+ (HETD)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS

BLOCK 62 CERN CONF 371 BLOCK, LENDINARA, MENARI (NMES+BOLOGNA)

PAPERS NOT REFERRED TO IN DATA CARDS

BRENE 61 NP 22 553 BRENE, EGARDT, QVIST (NORD)
BIRGE 63 PRL 11 35 BIRGE, ELY, GIDAL, CAMERINI + (LRL+WISC+BARI)
ADAIR 64 PL 12 67 ADAIR, LEIPUNER (YALE, BNL)

Stable Particles

K^\pm, K^0, K_S^0

Data Card Listings

For notation, see key at front of Listings.

CABI880 64 PL 9 352	CABI880,MAKSYMOWICZ (CERN)
ALSO 64 PL 11 360	CABI880,MAKSYMOWICZ (CERN)
ALSO 65 PL 14 72	CABI880,MAKSYMOWICZ (CERN)
CABI880 66 BERKELEY CONF 33	CABI880 (CERN)
GINSBERG 67 PR 162 1570	EDWARD S GINSBERG (U. MASS BOSTON)
WILLIS 67 HEIDELBERG 273	W J WILLIS -RAPPORTEUR TALK (YALE)
CROWIN 68 VIENNA CONF 241	RAPPORTEUR TALK (PRINCETON)
HAIDT 2 69 PL 298 696	+ (AACH,BARI,CERN,EPOL,NIJM,ORSA,PADO,TORI)
BARDIN 70 PL 328 121	BARDIN,BILENKY,PONTECORVO (JINR)
BECHERRA 70 PR D1 1452	T.BECHERRAMY (ROCH)
FEARING 70 PR D3 542	+FISCHBACK,SMITH (STON+BOHR)
GAILLARD 70 CERN 70-14	N K GAILLARD, L M CHOUNET (CERN+ORSA)
GINSBERG 70 PR D1 229	E S GINSBERG (IIT HAIFA)
GINSBERG 71 PR D4 2893	E S GINSBERG (MIT)
CHOUNET 72 PL 4C 199	(PHYS. REPTS.)CHOUNET,2*GAILLARD(ORSA+CERN)

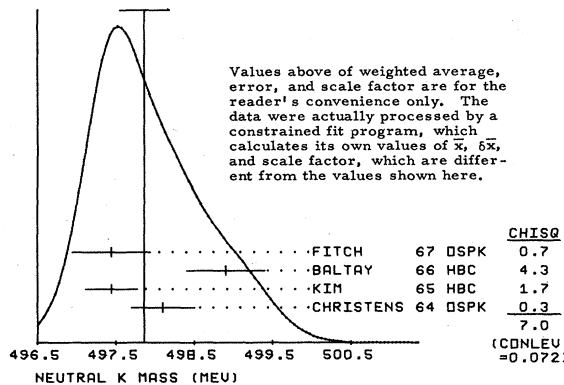
K^0

11 NEUTRAL K(498,JP=0-) I=1/2

11 NEUTRAL K MASS (MEV)

M	498.1	0.4	CHRISTENS 64 OSPK		
M	2223 497.44	0.33	KIM 65 HBC	KO FROM PBAR P	6/66
M	4500 498.9	0.5	BALTAY 66 HBC	KO FROM PBAR P	6/66
M	497.44	0.50	FITCH 67 OSPK		11/67
M	AVG	497.87	0.32	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)	
M	STUDENT	497.83	0.26	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
M	FIT	497.70	0.13	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	
				(SEE IDEOGRAM BELOW)	2/76*

WEIGHTED AVERAGE = 497.87 ± 0.32
ERROR SCALED BY 1.5



	CHISQ
FITCH 67 OSPK	0.7
BALTAY 66 HBC	4.3
KIM 65 HBC	1.7
CHRISTENS 64 OSPK	0.3
	7.0
	(CONLEV = 0.072)

11 (K0) - (K+-) MASS DIFFERENCE (MEV)

D	3.9	0.6	ROSENFELD 59 HBC -		
D	5.4	1.1	CRAWFORD 59 HBC +		
D	9	3.90	BURNSTEIN 65 HBC -	6/68	
D	7	3.71	KIM 65 HBC -	K - P TO KO N	
D	417	3.95	HILL 68 DBC +	K+D TO KOPP	
D	AVG	3.92	0.14	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
D	STUDENT	3.91	0.15	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
D	FIT	3.99	0.13	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	
					2/76*

REFERENCES FOR NEUTRAL K

CRAWFORD 59 PRL 2 112	CRAWFORD,CRESTI,GOOD,STEVENSON,TICHO (LRL)
ROSENFELD 59 PRL 2 110	A H ROSENFELD,F SOLMITZ,R D TRIPP (LRL)
CHRISTEN 64 PRL 13 138	CHRISTENSON,CROWIN,FITCH,TURLAY (PRINCETON)
BURNSTEIN 65 PR 138 B 895	R A BURNSTEIN,H A RUBIN (MARYLAND)
KIM 65 PR 140 B 1334	J K KIM,L KIRSCH,D MILLER (COLUMBIA)
BALTAY 66 PR 142 932	BALTAY,SANDWEISS,STONEHILL + (YALE+BNL)
FITCH 67 PR 164 1711	FITCH,ROTH,RUSS,VERNON (PRINCETON)
HILL 68 PR 168 1534	HILL,ROBINSON,SAKITT,CANTER (BNL,CARNEGIE)

K_S^0

12 SHORT-LIVED NEUTRAL K(498, JP=0-) I=1/2

Note on the K_S^0 Mean Life

From 1968 until 1972 our average value for the K_S^0 mean life was $(0.862 \pm 0.006) \times 10^{-10}$ second. Since then three high-precision experiments, SKJEGGESTAD 72 (Oslo, CERN, Saclay), GEWENIGER 74 (CERN, Heidelberg), and CARITHERS 75 (Columbia, NYU)

have obtained results compatible with each other which average $(0.8930 \pm 0.0023) \times 10^{-10}$ second. This is about five standard deviations above the previous average. The origin of this discrepancy is not known.

The corrections for systematic biases in SKJEGGESTAD 72 and HILL 68 (including the correction for the new value of η_{+-}) amount to +1% and 0.7%, respectively. Similar corrections, if applied to the older bubble chamber results, would probably increase the old average by only about one standard deviation and would not account for the discrepancy.

The two experiments which contribute most to the χ^2 (see ideogram below) are KIRSCH 66 and DONALD 68. The combined result of all other experiments is $(0.8914 \pm 0.0021) \times 10^{-10}$ second with a χ^2 of 11.5 for 10 degrees of freedom. This value does not differ significantly from the average of the newer (post-1971) experiments.

Since the newer experiments are in principle superior - that is, they have higher statistics, better acceptance, and easier trigger conditions - we have chosen to average them separately from the older experiments as is seen in the Data Card Listings below. In the Stable Particle Table, we quote the new value, but give the old value in a warning footnote.

12 KOS MEAN LIFE (UNITS 10^{*-10} SEC)

T	O	90	(1.07)	(0.13)	(0.13)	BOLDT	58 CC	
T	O	512	0.94	0.05	0.05	CRAWFORD	59 HBC	
T	O	63	(1.09)	(0.18)	(0.15)	BOWEN	60 CC	
T	O	OLD EXPTS WITH LOW STATISTICS NOT INCLUDED IN AVERAGE.						6/68
T	O	378	0.94	0.05	0.05	BERTANZA	62 HBC	
T	O	503	0.87	0.05	0.05	CHRÉTIEN	63 HBC	
T	O	545	0.86	0.04	0.04	KREISLER	64 OSPK	
T	O	572	0.90	0.06	0.06	ALFF-STEI	66 OSPK	9/66
T	O	4500	0.92	0.04	0.05	AUERBACH	66 OSPK	8/67
T	B	5009	(0.904)	(0.024)		BALTAY	66 HBC	6/66
T	H	19994	0.843	0.013		BOTT-BODDE	66 OSPK	9/66
T	H	20000	0.856	0.008		KIRSCH	66 HBC	6/66
T	H	20000	0.872	0.009		DONALD	68 HBC	6/68
T	H	20000	0.864	0.009		HILL	68 DBC	11/72
T	AVG	0.8641	0.0065	0.0065	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.3)			
T	STUDENT	0.8642	0.0060	0.0059	AVG. USING STUDENT10(H/1.11) -- SEE TEXT			
					(SEE IDEOGRAM BELOW)			
T	O	50K	0.8958	0.0045		SKJEGGESTAD	72 HBC	1/73
T	H	2173	(0.867)	(0.024)		FACKLER	73 OSPK	11/73
T	C	6M	0.8937	0.0048		GEWENIGER	74 ASPK	3/74
T	C	6	0.8913	0.0032		CARITHERS	75 SPEC	7/75*
T	AVG	0.8930	0.0023	0.0023	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.0)			
T	STUDENT	0.8930	0.0025	0.0025	AVG. USING STUDENT10(H/1.11) -- SEE TEXT			
T	FIT	0.8930	0.0023	0.0023	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)			
					(SEE IDEOGRAM BELOW)			

COMMENTS

T	H	HILL 68 HAS BEEN CHANGED BY THE AUTHORS FROM THE PUBLISHED VALUE	11/72
T	H	(0.865+-0.009) BECAUSE OF A CORRECTION IN THE SHIFT DUE TO ETA+-.	11/72
T	H	SKJEGGESTAD 72 AND HILL 68 GIVE DETAILED DISCUSSIONS OF SYSTEMATICS	
T	H	ENCOUNTERED IN THIS TYPE OF EXPERIMENT.	
T	B	KOS MEAN LIFE NOT THE PRIMARY QUANTITY MEASURED IN THIS EXPT.	6/68
T	F	FACKLER 73 DOES NOT INCLUDE SYSTEMATIC ERRORS.	11/73
T	C	CARITHERS 75 VALUE IS FOR KOL-KOS MASS DIFFERENCE DM=-.5348+-0.0021.	11/75*
T	C	THE DM DEPENDENCE OF THE TOTAL DECAY RATE (INVERSE MEAN LIFE) IS	11/75*
T	C	GAMMA(KOS)=[(1.122+-0.04)+.16*(DM-.5348)/DM]*10**10 /SEC.	11/75*
T	C	VALUE WOULD NOT CHANGE WITH OUR CURRENT DM=.5349+-0.022.	2/76*

Data Card Listings

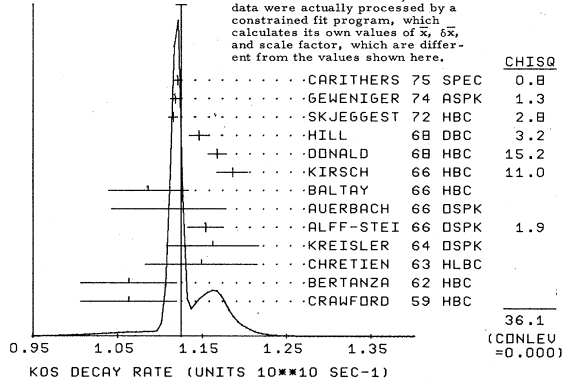
For notation, see key at front of Listings.

Stable Particles

K⁰

WEIGHTED AVERAGE = 1.1257 ± 0.0065
ERROR SCALED BY 2.5

Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of X, δX, and scale factor, which are different from the values shown here.



12 KOS PARTIAL DECAY MODES

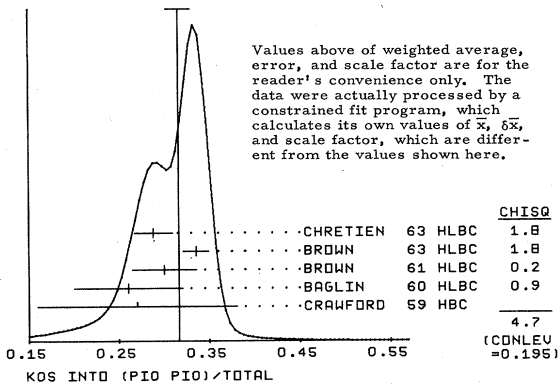
Table with columns: Mode (P1, P2, P3, P4, P5, P6, P7), Decay Masses (139+, 139, 134+, 134, 105+, 105, .5+, .5, 139+, 139+, 0, 134+, 134+, 134).

12 KOS BRANCHING RATIOS

Table with columns: Ratio (R1, R2), KOS INTO (PI+ PI-)/TOTAL, (PI1), (P2), and values. Includes sub-sections for R1, R2, R2 AVG, R2 STUDENT, R2 FIT.

WEIGHTED AVERAGE = 0.316 ± 0.014
ERROR SCALED BY 1.3

Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of X, δX, and scale factor, which are different from the values shown here.



Main data table for K⁰ particles, listing various decay modes, branching ratios, and references. Includes sub-sections for R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, R14, R15, R16, R17, R18, R19, R20, R21, R22, R23, R24, R25, R26, R27, R28, R29, R30, R31, R32, R33, R34, R35, R36, R37, R38, R39, R40, R41, R42, R43, R44, R45, R46, R47, R48, R49, R50, R51, R52, R53, R54, R55, R56, R57, R58, R59, R60, R61, R62, R63, R64, R65, R66, R67, R68, R69, R70, R71, R72, R73, R74, R75, R76, R77, R78, R79, R80, R81, R82, R83, R84, R85, R86, R87, R88, R89, R90, R91, R92, R93, R94, R95, R96, R97, R98, R99, R100, R101, R102, R103, R104, R105, R106, R107, R108, R109, R110, R111, R112, R113, R114, R115, R116, R117, R118, R119, R120, R121, R122, R123, R124, R125, R126, R127, R128, R129, R130, R131, R132, R133, R134, R135, R136, R137, R138, R139, R140, R141, R142, R143, R144, R145, R146, R147, R148, R149, R150, R151, R152, R153, R154, R155, R156, R157, R158, R159, R160, R161, R162, R163, R164, R165, R166, R167, R168, R169, R170, R171, R172, R173, R174, R175, R176, R177, R178, R179, R180, R181, 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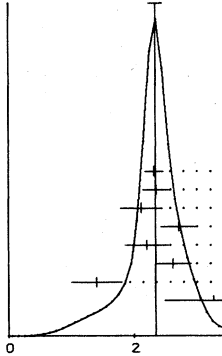
Data Card Listings

For notation, see key at front of Listings.

Stable Particles

K⁰

WEIGHTED AVERAGE = 2.34 ± 0.11
ERROR SCALED BY 1.2



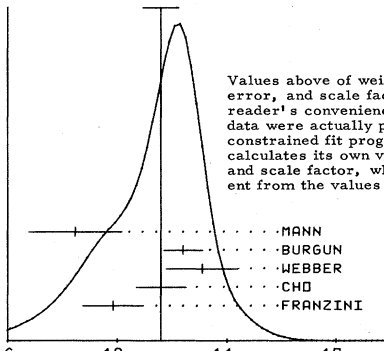
Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

Table listing particle names and their corresponding values: BALDOCE DL 75 HLBC 0.0, JAMES 72 HBC 0.0, MEISNER 71 HBC 0.4, CHD 71 DBC 1.7, WEBBER 70 HBC 0.2, BEHR 66 HLBC 1.0, FRANZINI 65 HBC 5.5, ANDERSON 65 HBC.

CHI SQ 0.0, B.9 (CONLEU = 0.177)

Table with columns W3, W4, W5 and rows for KOL INTO PI E NEUTRINO, KOL INTO CHARGED (3-BODY), and KOL INTO LEPTONIC (KMU3+KE3) with various particle names and values.

WEIGHTED AVERAGE = 11.60 ± 0.65
ERROR SCALED BY 1.5



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

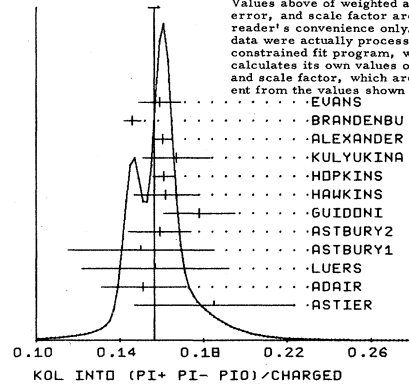
Table listing particle names and their corresponding values: MANN 72 HBC 3.4, BURGUN 72 HBC 1.3, WEBBER 71 HBC 1.3, CHD 70 DBC 0.0, FRANZINI 65 HBC 2.5.

CHI SQ 3.4, B.6 (CONLEU = 0.072)

Table with columns W6, W7, W8 and rows for KOL INTO PI MU NEUTRINO UNITS 10**6 SEC-1 and 13 KOL BRANCHING RATIOS.

Table with columns R2, R3 and rows for KOL INTO (PI+ PI- P10)/CHARGED, listing particle names like ASTIER, ADAIR, LUERS, ASTBURY1, ASTBURY2, GUIDONI, HOPKINS, HAWKINS, HOPKINS, KULYUKINA, ALEXANDER, BRANDENBU, EVANS.

WEIGHTED AVERAGE = 0.1564 ± 0.0026
ERROR SCALED BY 1.2



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of x, sigma_x, and scale factor, which are different from the values shown here.

Table listing particle names and their corresponding values: EVANS 73 HLBC 0.1, BRANDENBU 73 HBC 6.7, ALEXANDER 73 HBC 1.2, KULYUKINA 68 CC 0.4, HOPKINS 67 HBC 0.8, HAWKINS 66 HBC 0.1, GUIDONI 65 HBC 1.6, ASTBURY2 65 CC 0.0, ASTBURY1 65 CC, LUERS 64 HBC, ADAIR 64 HBC 0.1, ASTIER 61 CC.

CHI SQ 0.1, 11.1 (CONLEU = 0.194)

Large table with columns R3, R4, R5, R6, R7, R8, R9 and rows for KOL INTO (PI MU NEUTRINO)/CHARGED, KOL INTO (PI E NEUTRINO)/CHARGED, KOL INTO (PI E NEU)/(PI E NEU+PI MU NEU), KOL INTO (PI+ PI- P10)/TOTAL, KOL INTO (LEPTON PI NEUTRINO)/TOTAL, KOL INTO (2 GAMMA)/TOTAL (UN. 10**4), and AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0).

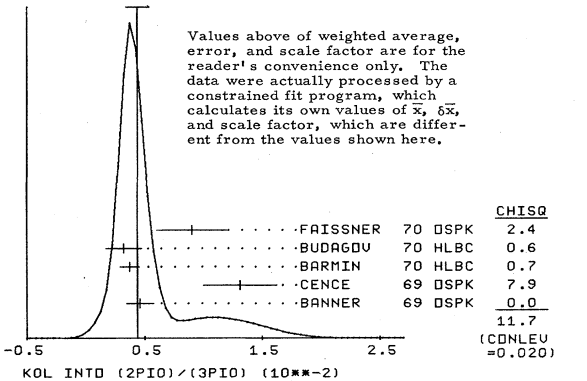
Stable Particles

K_L^0

Data Card Listings

For notation, see key at front of Listings.

R10	KOL INTO (PI MU NEU)/(PI E NEU)	(P3)/(P4)		
R10	0.81	0.19	ADAIR	64 HBC
R10	0.82	0.10	DEBOUARD	67 OSPK
R10	273	0.7	HANKINS	67 HBC
R10	0.81	0.08	HOPKINS	67 HBC
R10	770	0.71	BUDAGOV	68 HLBC
R10 K	(0.67)	(0.13)	KULYUKINA	68 CC
R10 B	569	(0.71)	(0.04)	BEILLIERE 69 HLBC
R10	1309	(0.648)	(0.030)	EVANS 69 HLBC REPL. BY EVANS 73
R10	3548	0.68	0.08	BASILE 70 OSPK
R10	5700	0.741	0.024	BRANDENBU 73 HBC
R10	1309	0.662	0.030	EVANS 73 HLBC
R10	10K	0.662	0.037	WILLIAMS 74 ASPK
R10 K	KULYUKINA 68 R10 IS NOT MEASURED INDEPENDENTLY FROM R2 AND R4.			
R10 B	BEILLIERE 69 IS A SCANNING EXPT USING SAME EXPOSURE AS BUDAGOV 68			
R10	AVG	0.695	0.019	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
R10	STUDENT	0.695	0.021	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R10	FIT	0.696	0.017	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
R11	KOL INTO (MU+MU-)/CHARGED (UNITS 10**=-6)	(P61)/(P2+P3+P4)		
R11	100.0	OR LESS	ANIKINA	65 CC
R11	250.0	OR LESS	CL=.90	ALFF-STEI 66 OSPK
R11	5	OR LESS	CL=.90	BOTT-BODE 67 OSPK
R11	35.0	OR LESS	CL=.90	FITCH 67 OSPK
R12	KOL INTO (PI+ PI- GAMMA)/TOTAL (UNITS 10**=-3)	(P10)		
R12	15.0	OR LESS	ANIKINA	65 CC
R12	0	OR LESS	BELLOTTI	66 HLBC
R12	1	3.0	OR LESS	NEFKENS 66 OSPK
R12	0.4	OR LESS	CL=.90	THATCHER 68 OSPK
R12	3.2	OR LESS	CL=.90	BOBISUT 74 HLBC
R12 D	24	0.062	0.021	DONALD S1 74 SPEC
R12	0.46	OR LESS	CL=.90	WOD 74 SPEC
R12 D	USES KOL TO PI+PI-PI0/ALL KOL DECAYS = 0.126			
R13	KOL INTO (E+ E-)/CHARGED (UNITS 10**=-6)	(P71)/(P2+P3+P4)		
R13	1000.0	OR LESS	ANIKINA	65 CC
R13	200.0	OR LESS	CL=.90	ALFF-STEI 66 OSPK
R13	23.0	OR LESS	CL=.90	BOTT-BODE 67 OSPK
R14	KOL INTO (E MU)/CHARGED (UNITS 10**=-4)	(P81)/(P2+P3+P4)		
R14	10.0	OR LESS	ANIKINA	65 CC
R14	1.0	OR LESS	CL=.90	CARPENTER 66 OSPK
R14	0.1	OR LESS	CL=.90	BOTT-BODE 67 OSPK
R14	0.08	OR LESS	CL=.90	FITCH 67 OSPK
R15	KOL INTO (E+ PI- NEU)/(E- PI+ NEU)			
R15 O	97	(0.90)	(0.18)	NEAGU 61 CC
R15 O	(1.01)	(0.16)	(0.18)	LUERS 64 HBC
R15 O	894	(0.92)	(0.023)	KULYUKINA 66 CC
R15 O	1539	(1.06)	(0.05)	VERHEY 66 OSPK
R15 O	LOW PRECISION EXPTS NOT AVERAGED. FOR MORE PRECISE VALUE, SEE S13A2 (BENNETT 70, MARX 70)			
R16	KOL INTO (MU+ PI- NEU)/(MU- PI+ NEU)			
R16	1M	1.0081	0.0027	DORFAN 67 OSPK
R16	SEE ALSO S13A2 AND S13AL IN THE CP VIOLATION SECTION			
R17	KOL INTO (PI0 PI0)/TOTAL (UNITS 10**=-3)	(P11)		
R17 C	7	(1.2)	(1.5)	(1.2) CRIEGEE 66 OSPK
R17 C	CRIEGEE EXPT NOT DESIGNED TO MEASURE 2 PI0 DECAY MODE			
R17 G	189	(2.5)	(0.8)	E00=3.6+-0.6
R17 G	LATEST RESULT OF THIS EXPERIMENT GIVEN BY FAISSNER TO R19			
R17	0.94	0.19	FROM FIT	
R18	KOL INTO (3PI0)/(PI+PI-PI0)	(P11)/(P2)		
R18	188	2.0	0.6	ALEKSANYA 64 FBC
R18	1010	1.80	0.13	BUDAGOV 68 HLBC
R18	883	(1.65)	(0.07)	BARMIN 72 HLBC
R18	ERROR STAT. ONLY			
R18	AVG	1.81	0.13	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R18	STUDENT	1.81	0.14	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R18	FIT	1.747	0.070	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
R19	KOL INTO (2PI0)/(3PI0)	(P11)/(P1)		
R19 C	109	(1.89)	(0.31)	CRONIN 1 67 OSPK
R19 C	(1.36)	(0.18)	(0.18)	CRONIN 2 67 OSPK
R19 C	CRONIN 2 IS FURTHER ANALYSIS OF CRONIN 1, NOW BOTH WITHDRAWN			
R19	NO EVENTS SEEN			
R19	57	0.46	0.11	BARTLETT 68 OSPK
R19	133	1.31	0.31	BANNER 69 OSPK
R19	29	0.37	0.08	CENCE 69 OSPK
R19	30	0.32	0.15	BARMIN 70 HLBC
R19	172	0.90	0.30	BUDAGOV 70 HLBC
R19 F	FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69 R17			
R19	AVG	0.439	0.098	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)
R19	STUDENT	0.425	0.065	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT
R19	FIT	0.439	0.088	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.5)
				(SEE IDEOGRAM BELOW)
				WEIGHTED AVERAGE = 0.439 ± 0.098
				ERROR SCALED BY 1.7



R20	KOL INTO (PI+ PI-1)/(KE3 + KMU3) (UNITS 10**=-3)	(P51)/(P3+P4)		
R20 O	309	(2.51)	(0.23)	DEBOUARD 67 OSPK
R20 O	525	(2.35)	(0.19)	FITCH 67 OSPK
R20 O	OLD EXPERIMENTS EXCLUDED FROM FIT. SEE SUBSECTION E+ BELOW FOR			
R20 O	AVERAGE ETA+- OF THESE EXPERIMENTS AND FOR NOTE ON DISCREPANCY.			
R20	3.041	0.090	FROM FIT	
R21	KOL INTO (2GAMMA)/(3 PI0) (UNITS 10**=-3)	(P91)/(P1)		
R21	16	2.5	0.7	ARNOLD 68 HLBC
R21	BANNER 69 IS NEW EXPT. NOT TO BE CONF WITH R8 OF CRONIN 67			
R21	115	2.24	0.28	BANNER 69 OSPK
R21	28	2.13	0.43	BARMIN 71 HLBC
R21	AVG	2.24	0.22	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R21	STUDENT	2.24	0.24	AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

Note on the $K_L^0 \rightarrow \mu^+ \mu^-$ Controversy

The $K_L^0 \rightarrow \mu^+ \mu^-$ branching ratio upper limit of CLARK 71 does not agree with the results of subsequent experiments.

A total of twelve examples of this decay have been observed in three experiments: CARITHERS1 73 (Columbia-CERN-NYU), CARITHERS2 73 (Columbia-BNL-CERN), and FUKUSHIMA 76 (Princeton-Univ. of Mass.-Amherst). Background from all sources (primarily from $K_L^0 \rightarrow \pi \mu \nu$ with the pion either penetrating the muon counter or decaying to $\mu \nu$ in flight) has been estimated to be negligible for these experiments. The rates from these three experiments and the upper limits of FOETH 69 and DARRIULAT 70 are all quite consistent (see subsection R22 below). The combined result of these five experiments is

$$R = \frac{\Gamma(K_L^0 \rightarrow \mu^+ \mu^-)}{\Gamma(K_L^0 \rightarrow \pi^+ \pi^-)} = (5.0^{+1.6}_{-1.4}) \times 10^{-6}$$

where we have divided the observed number ($12^{+3.8}_{-3.3}$) by the sum of the effective denominators (which include the $\pi \pi / \mu \mu$ relative acceptances) given in the data card comments field.

CLARK 71 found no events and obtained the limit $R < 1.53 \times 10^{-6}$ at 90% confidence level. This result is from the FIELD 74 reanalysis of the experiment. The reanalysis uncovered no major flaws, but did find a few problems which increased the upper limit from 1.2×10^{-6} to the above. The combined result of the other five experiments leads us to expect 7.5 events for the sensitivity of the CLARK 71 experiment (9.5 events before reanalysis).

The theoretical lower limit for $\Gamma(K_L^0 \rightarrow \mu^+ \mu^-) / \Gamma(K_L^0 \rightarrow \pi^+ \pi^-)$ based on unitarity considerations is 3×10^{-6} , in agreement with the combined result, but above the CLARK 71 upper limit.

Stable Particles

K_L⁰

Data Card Listings

For notation, see key at front of Listings.

XIC XIC = F/F+ (DETERMINED FROM MU POLARIZATION IN KMU3)
XIC THE MU POLARIZATION IS A MEASURE OF XI(T). NO ASSUMPTIONS ON L+
XIC NECESSARY, T (WEIGHTED BY SENSITIVITY TO XI(O)) SHOULD BE SPECIFIED.

IXI IMAGINARY PART OF XI (TEST OF T REVERSAL)
IXI -0.02 0.08 ABRAMS 68 OSPK POLARIZATION 10/69
IXI 2.2M -0.060 0.045 SANDWEISS 73 CNTR POL, T=3.3 1/74

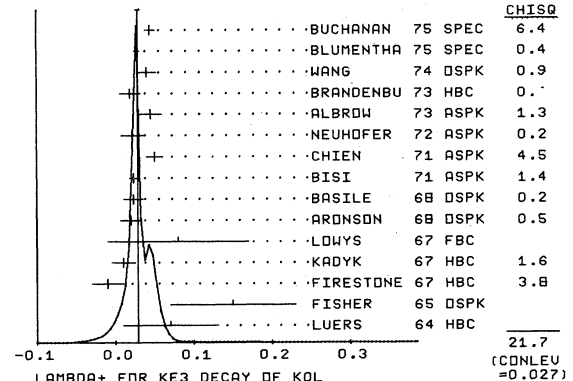
L+M LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KMU3 DECAY)
L+M SEE ALSO THE CORRESPONDING ENTRIES AND NOTES IN SECTION XIA AND LO.
L+M FOR RAD. CORR. OF KMU3 DP SEE GINSBURG 70 AND BECHERRAWY 70.

LO LAMBDA 0 (LINEAR ENERGY DEPENDENCE OF F0 IN KMU3 DECAY)
LO WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(O) INTO
LO VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

L+E LAMBDA + (LINEAR ENERGY DEPENDENCE OF F+ IN KO E3 DECAY)
L+E FOR RAD. CORR. OF KE3 DP SEE GINSBURG 67 AND BECHERRAWY 70.
L+E 153 +0.07 .06 LUERS 64 HBC DP, NO RC 8/67

FS FS/F+ RATIO OF SCALAR TO F+ COUPLINGS FOR KE3 DECAY (ABS. VALUE)
FS 5600 0.19 OR LESS CL=.68 KULYUKINA 67 CC 10/69
FS 25K 0.04 OR LESS CL=.68 BLUMENTHA 75 SPEC 7/75*

WEIGHTED AVERAGE = 0.0288 ± 0.0028
ERROR SCALED BY 1.4

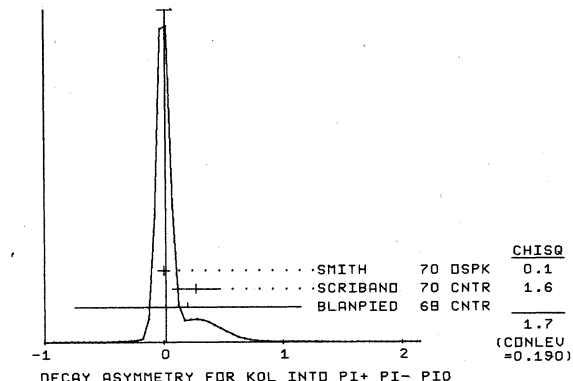


13 CP VIOLATION PARAMETERS IN KOL DECAYS
RELATED TEXT SECTION VI B.3 AND MINI-REVIEW BELOW

13 CHARGE ASYMMETRY IN TAU DECAYS
TEXT SECTION VI B.3 B

SEE SCRIBANO 70 FOR DEFINITION (HIS SIGMA+1, A=1 FOR MAX ASYMMETRY
(M)*2 = 1+ SIG+ (2/SQRT(3)) * ((T+)-(T-))/TMAX) AS SCRIBANO 70
A DECA ASYMMETRY PARAMETER FOR PI+ PI- PIO (UNITS 10**=-2)
A .3M 0.2 0.95 BLANPIED 68 CNTR .4/70
A .3M 0.27 0.2 SCRIBANO 70 CNTR .12/70
A 4400 0.000 0.050 SMITH 70 OSPK .10/70

WEIGHTED AVERAGE = 0.016 ± 0.063
ERROR SCALED BY 1.3



13 CHARGE ASYMMETRY IN LEPTONIC DECAYS (PERCENT)
TEXT SECTION VI B.3 C

SUCH ASYMMETRY VIOLATES CP. IT IS RELATED TO REAL(EPSILON).
A1 D KOL INTO ((MU+PI-NU)-((MU-PI+NU)))/((MU+PI-NU)+((MU-PI+NU))) (PERCENT)
A1 D 1M (0.403) (0.134) DORFAN 67 OSPK DERIVED FROM R16 11/67
A1 D 1M 0.57 0.17 PACIOTTI 69 OSPK 1/73
A1 7.7M 0.278 0.051 PICCIONI 72 ASPK 1/73
A1 4.1M 0.60 0.19 MCCARTHY 73 CNTR 6/73
A1 15M 0.313 0.029 GEMENIGI 74 ASPK 7/74*

Data Card Listings

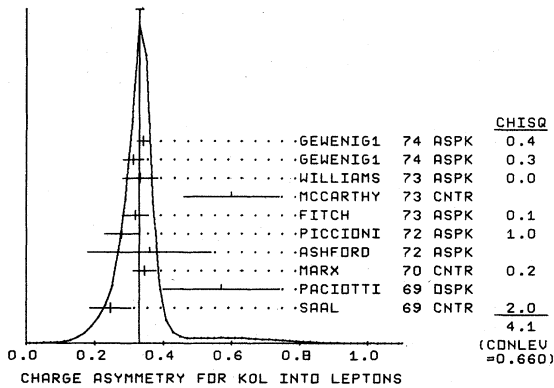
Stable Particles

For notation, see key at front of Listings.

K_L^0

AL	KOL INTO ((L+)-(L-))/((L+)+(L-)) (COMBINED A1 AND A2) (PERCENT)			
AL B	10M	0.246	0.059	SAAL 69 CNTR KE3 2/71
AL D	1M	0.57	0.17	PACIOTTI 69 DSPK KMU3 1/73
AL	10M	0.346	0.033	MARX 70 CNTR KE3 2/71
AL	600K	0.36	0.18	ASHFORD 72 ASPK KE3 2/72
AL	7.7M	0.278	0.051	PICCIONI 72 ASPK KMU3 1/73
AL	40M	0.318	0.038	FITCH 73 ASPK KE3 12/73
AL	4.1M	0.60	0.14	MCCARTHY 73 CNTR KMU3 6/73
AL	33M	0.333	0.050	WILLIAMS 73 ASPK KMU3*KE3 12/73
AL	15M	0.313	0.029	GEWENIGER 74 ASPK KMU3 7/74*
AL	34M	0.341	0.018	GEWENIGER 74 ASPK KE3 7/74*
AL	SEE FOOTNOTES IN SECTIONS A1 AND A2 ABOVE.			
AL	AVG	0.330	0.012	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AL	STUDENT	0.330	0.013	AVERAGE USING STUDENT(10/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.330 ± 0.012
ERRDR SCALED BY 1.0



-----13 PARAMETERS FOR KOL INTO 2PI DECAY-----
TEXT SECTION VI B.3 D
ETA+- = A(KL TO PI+PI-)/A(KS TO PI+PI-)
ETA0 = A(KL TO PIP0I0)/A(KS TO PIP0I0)

Note on $|\eta_{+-}|$

There is a very large discrepancy between old and new results for $|\eta_{+-}|$, the ratio of amplitudes $A(K_L^0 \rightarrow \pi^+\pi^-)/A(K_S^0 \rightarrow \pi^+\pi^-)$. The average of the older $|\eta_{+-}|$ results of CHRISTENSON 64, GALBRAITH 65, BASILE 66, BOTT-BODENHAUSEN 66, DE BOUARD 67, and FITCH 67 (see subsections R9, R20, and E+- of the K_L^0 Data Card Listings) is

$$|\eta_{+-}| = (1.95 \pm 0.03) \times 10^{-3}$$

with very good consistency.

In 1973 a CERN-Heidelberg K_L^0, K_S^0 interference experiment and a Colorado-SLAC-Santa Cruz branching ratio measurement gave the following results:

$$|\eta_{+-}| = (2.30 \pm 0.035) \times 10^{-3} \quad \text{GEWENIGER2 74,}$$

$$|\eta_{+-}| = (2.23 \pm 0.05) \times 10^{-3} \quad \text{MESSNER 73,}$$

which are around 11 standard deviations above the previous average but are in good agreement with each

other. GEWENIGER2 74 also made a check measurement in order to confirm their interference result by measuring the branching ratio $K_L^0 \rightarrow \pi^+\pi^-/K_L^0 \rightarrow \pi e \nu$. They obtained the result $|\eta_{+-}| = (2.30 \pm 0.06) \times 10^{-3}$, which, although systematically less reliable than the interference result, is in good agreement with it. A Columbia experiment¹ and two Chicago experiments (Cronin group² and Telegdi group³) have also yielded unpublished measurements which corroborate the higher values of $|\eta_{+-}|$.

The origin of the discrepancy between old and new results is not known. The effect of the change in the K_S^0 mean life value (see note in K_S^0 section of the Stable Particle Data Card Listings) is insufficient to explain it, raising $|\eta_{+-}|$ by only about 2%. The MESSNER 73 result above was evaluated for the old K_S^0 mean life and branching fractions for $K_S^0 \rightarrow \pi^+\pi^-$ and $K_L^0 \rightarrow \pi^+\pi^0$. The new K_S^0 mean life and the current branching fractions increase the MESSNER 73 $|\eta_{+-}|$ value slightly to $(2.246 \pm 0.032) \times 10^{-3}$.

We are troubled by this large unexplained discrepancy. We feel that our normal procedure of averaging and increasing the error by a scale factor S to account for the discrepancy is not adequate for this case. The two new results, when combined with the average of the earlier results by that procedure, give 2.15 ± 0.11 ($S = 6.0$). While this value-and-error makes some sense in that it nearly spans both incompatible sets of data, we choose not to quote it. Instead, since the newer experiments are in principle superior (higher statistics, better acceptance, easier trigger conditions), we have chosen to average them separately from the earlier experiments as is seen in the E+- subsection of the Data Card Listings below.

The entry referenced as GKL/GKS 71 is the average of all seven experiments before 1971. This average is determined from the branching ratios quoted for these experiments in subsections R9 and R20 above. It was quoted by us as the $|\eta_{+-}|$ value in our 1971 edition. The average and fit values at the end of the E+- subsection below do not include

Stable Particles

K^0

the pre-1971 $|\eta_{+-}|$ results. The reader may thus utilize the pre-1971 and post-1971 results as desired.

The fitted values in the EOS, E+-, and ER subsections below are the result of a fit to the unparenthesized values of $|\eta_{00}|$, $|\eta_{+-}|$, and $|\eta_{00}/\eta_{+-}|$ in these subsections. We quote the fitted values of $|\eta_{+-}|$ and $|\eta_{00}|$ in the Addendum to the Stable Particle Table but warn in a footnote that they exclude pre-1971 $|\eta_{+-}|$ results.

See also D. Nygren's "Review of K^0 Decays"⁴ for additional discussion of this and other K^0 controversies.

References

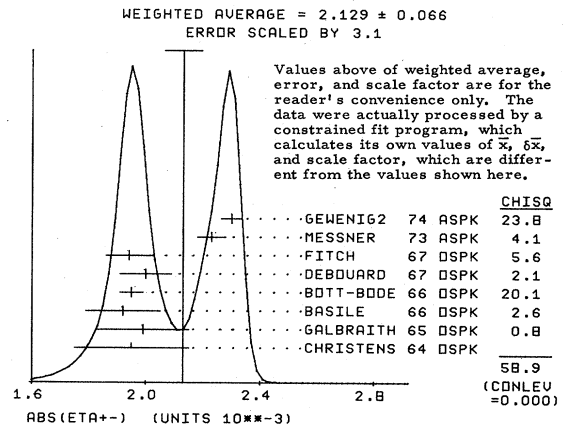
1. D. R. Nygren, private communication.
2. R. DeVoe, private communication.
3. S. Aronson, private communication to D. Nygren reported in reference 4.
4. D.R. Nygren, Review of K^0 Decays, presented at the Berkeley A.P.S. Meeting of the Division of Particles and Fields, Berkeley, California, August 13-17, 1973; LBL-2407.

THE FITTED VALUES OF η_{+-} AND η_{00} GIVEN BELOW ARE THE RESULTS OF A FIT TO η_{+-} , η_{00} AND η_{00}/η_{+-} RESULTS. THE VALUES LISTED BELOW WHICH ARE NOT PARENTHESIZED ENTER THE FIT AS SHOWN. THE VALUES WHICH ARE PARENTHESIZED AND BEAR THE FOOTNOTE X DO NOT ENTER THE FIT AS SHOWN. THESE EXPERIMENTS GIVE BRANCHING RATIOS AND ENTER THE FIT VIA THE QUANTITY ACTUALLY MEASURED -- BRANCHING RATIOS R_9 , R_{20} AND R_{27} (η_{+-}) AND R_{17} AND R_{19} (η_{00}). THESE BRANCHING RATIOS ARE COMBINED WITH CURRENT NORMALIZATIONS AND CURRENT K_L AND K_S MEAN LIVES TO OBTAIN $\pi^+\pi^-$ RATES. THE η_{+-} AND η_{00} VALUES OBTAINED FROM THESE RATES ARE ENTERED BELOW WITH THE NAME 'GKL/GKS'.

EOS	(ETA00)**2 = (A(KL TO 2PI0)/A(KS TO 2PI0))**2 (UNITS 10***-6) ---	
EOS X	0 (-2.) (7.0) BARTLETT 68 DSPK	10/69
EOS X	57 (4.9) (1.2) BANNER 69 DSPK	2/72
EOS X	133 (14.1) (3.4) CENCE 69 DSPK	10/69
EOS XF	180 (13.1) (4.1) GAILLARD 69 DSPK	10/69
EOS X	29 (4.08) (0.9) BARMIN 70 HLBC	12/70
EOS X	30 (3.61) (1.9) BUDAGOV 70 HLBC	10/70
EOS C	8.7 3.7 CHOLLET 70 DSPK CU REG.,4 GAMMAS	2/72
EOS XF	172 (9.9) (3.4) FAISSNER 70 DSPK	12/70
EOS C	56 (4.0) 2.0 WOLFF 71 DSPK CU REG.,4 GAMMAS	12/71
EOS X	5.2 1.1 GKL/GKS 76 RVUE BR SCALE FACTOR=1.5	2/76*
EOS X	SEE NOTE ABOVE REGARDING FITTED VALUES OF η_{+-} AND η_{00} .	
EOS C	CHOLLET 70 GIVES $\eta_{00}=(1.23+0.24i)*(\text{REGEN AMPL}, 2\text{GEV}/C \text{ CUI}/10000\text{MB})$	2/72
EOS C	WOLFF 71 GIVES $\eta_{00}=(1.13+0.12i)*(\text{REGEN AMPL}, 2\text{GEV}/C \text{ CUI}/10000\text{MB})$	2/72
EOS C	WE COMPUTE BOTH η_{00} VALUES FOR (REGEN AMPL, 2GEV/C CUI)=24+-2MB.	2/72
EOS C	THIS REGEN AMPL RESULTS FROM AVERAGING OVER FAISSNER 69.	2/72
EOS C	EXTRAPOLATED USING OPTICAL MODEL CALCULATIONS OF BOHM ET AL.	2/72
EOS C	PL 278 594 (1968) AND THE DATA OF BALATS 71. (FROM H. FAISSNER,	2/72
EOS C	PRIVATE COMMUNICATION)	2/72
EOS F	FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69	
EOS	
EOS AVG	5.90 0.93 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
EOS STUDENT	5.9 1.1 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
EOS FIT	5.41 0.22 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	2/76*
EOO	THIS FIT VALUE CORRESPONDS TO $\eta_{00}=2.325+0.093i$	2/76*
E+-	η_{+-} = A(KL TO $\pi^+\pi^-$)/A(KS TO $\pi^+\pi^-$) UNITS 10***-3 -----	
E+- X	45 (1.95) (0.20) CHRISTENS 64 DSPK	2/76*
E+- X	54 (1.99) (0.16) GALBRAITH 65 DSPK	2/76*
E+- X	(1.92) (0.13) BASILE 66 DSPK	2/76*
E+- X	(1.95) (0.04) BOTT-BODE 66 DSPK	2/76*
E+- X	(2.00) (0.09) DEBOUARD 67 DSPK	2/76*
E+- X	(1.94) (0.08) FITCH 67 DSPK	2/76*
E+- AX	(1.95) (0.03) GKL/GKS 71 RVUE EXPTS. BEFORE 71	2/76*
E+- A	AVERAGE OF ABOVE EXPTS. EXCLUDED FROM FIT. SEE TYPED NOTE ABOVE.	2/76*
E+- X	4200 (2.23) (0.05) MESSNER 73 ASPK	11/75*
E+-	2.30 0.035 GENIEG2 74 ASPK	3/74
E+- B X	2.246 0.032 GKL/GKS 76 RVUE BR EXP. AFTER 71	2/76*
E+- B	CURRENTLY INCLUDES ONLY MESSNER 73 + NEW KL TAU AND KS $\pi^+\pi^-$ RATES	2/76*
E+- X	SEE NOTE ABOVE REGARDING FITTED VALUES OF η_{+-} AND η_{00} .	
E+-	
E+- AVG	2.271 0.027 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)	
E+- STUDENT	2.270 0.027 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
E+- FIT	2.272 0.023 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	2/76*
E+-	(SEE IDEOGRAM BELOW)	
ER	RATIO OF η_{00} OVER η_{+-}	
ER	124 1.03 0.07 BANNER1 72 DSPK	8/72
ER	167 1.00 0.06 HOLDER 72 ASPK	8/72
ER	
ER AVG	1.013 0.046 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
ER STUDENT	1.013 0.049 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
ER FIT	1.023 0.040 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	2/76*

Data Card Listings

For notation, see key at front of Listings.



Note on $K_L^0 \rightarrow 2\pi$ and K_S^0 Regeneration

Some experiments obtain ϕ_{+-} (the phase of η_{+-}) using K_S^0 , $K_L^0 \rightarrow \pi^+\pi^-$ interference behind a regenerator. In these interference experiments the measured quantity is the difference of ϕ_{+-} and the regeneration phase ϕ_R , as shown in the expression below. After the regenerator, the intensity of the $\pi^+\pi^-$ decays in the forward direction is

$$I(t, p) = S(p) \left[|R(p)|^2 e^{-\Gamma_S t} + |\eta_{+-}|^2 e^{-\Gamma_L t} + 2|R(p)||\eta_{+-}| \times e^{-(\Gamma_S + \Gamma_L)t/2} \times \cos(\Delta m t + \phi_R(p) - \phi_{+-}) \right], \quad (1)$$

where:

t is the decay time in the K^0 rest frame,
 $\Delta m = m_L - m_S$, and $m_L, \Gamma_L, m_S, \Gamma_S$ are the masses and decay rates of the long- and short-lived K^0 ,

$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}}$ is the ratio of decay amplitudes $A(K_L^0 \rightarrow \pi^+\pi^-)/A(K_S^0 \rightarrow \pi^+\pi^-)$,

$S(p)$ is proportional to the K_L momentum spectrum, and

$R(p) = |R(p)| e^{i\phi_R(p)}$ is the transmission-regenerated K_S amplitude (relative to the K_L):

$$R(p) = \pi N i \frac{[f_0(p) - \bar{f}_0(p)]}{p} \left\{ \frac{-\frac{1}{2}\Gamma_S \ell(p) [1 - 2i\Delta m/\Gamma_S]}{1 - e^{-\frac{1}{2}\Gamma_S \ell(p) [1 - 2i\Delta m/\Gamma_S]}} \right\}, \quad (2)$$

where

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Stable Particles

K_L^0

For notation, see key at front of Listings.

$\lambda(p)$ is the thickness of regenerator measured in units of the mean decay length of K_S ,

N is the number of nuclei per cubic centimeter,

Λ is the K_S mean decay length, and

$f_0(p)$, $\bar{f}_0(p)$ are the forward scattering amplitude of K^0 and \bar{K}^0 .

From (1) above it is clear that the value of ϕ_{+-} is highly correlated with the value of Δm and ϕ_R . In addition, a less obvious but significant correlation exists between ϕ_{+-} and Γ_S . Usually Δm is a parameter of the fit and ϕ_R is determined by some other means (optical model calculations, time dependence of the charge asymmetry in K_{e3} decay, etc.).

We list ϕ_{+-} and give in comment cards both the value of ϕ_R used by the authors and the Δm dependence of ϕ_{+-} .

F+-	PHASE OF ETA +- (DEGREES)	-----	
F+-	THE DEPENDENCE OF THE PHASE ON THE KOL-KOS MASS DIFFERENCE IS GIVEN FOR EACH EXPERIMENT IN THE COMMENTS BELOW, WHERE DM IS (MASS DIFF./HBAR) IN UNITS 10**10 SEC-1. WE HAVE EVALUATED THESE MASS DEPENDENCES USING OUR APRIL 1976 VALUE, DM=0.5349+-0.0022 TO OBTAIN THE VALUES AND AVERAGE QUOTED BELOW.		
F+- O	(45.0) (50.0)	FITCH 65 OSPK BE REGEN	11/67
F+- O	(30.0) (45.0)	FIRESTONE 66 HBC	11/67
F+- O	(70.0) (21.0)	BOTT-BODE 67 OSPK C REGEN	11/67
F+- O	(25.0) (35.0)	MISCHKE 67 OSPK CU REGEN	7/68
F+- O	OLD EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.		
F+- N	(51.0) (11.0)	BENNETT2 68 CNTR CU REG. USES	8/68
F+- C	34.2 10.0	BENNETT 69 CNTR CU REGEN	2/71
F+- B	45.3 12.0	BOHM 69 OSPK VACUUM REGEN	2/71
F+- F	45.2 7.4	FAISSNER 69 ASPK CU REGEN	2/71
F+- J	40.6 4.2	JENSEN 70 ASPK VACUUM REGEN	2/71
F+- D	37.2 12.0	BALATS 71 OSPK CU REGEN	9/71
F+- P	36.2 6.1	CARNEGIE 72 ASPK CU REGEN	1/73
F+- G	46.5 1.6	GEWENIG2 74 ASPK VACUUM REGEN	3/74
F+- H	45.5 2.8	CARITHERS 75 SPEC C REGEN	7/75*
F+-
F+- AVG	45.0 1.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
F+- STUDENT	45.1 1.4	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
F+- FIT	45.0 1.2	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	2/76*
F+-	COMMENTS		
F+- N	BENNETT 69 IS A REEVALUATION OF BENNETT2 68.		
F+- C	BENNETT 69 USES MEASUREMENT OF (F+-)-(PHIF) OF ALFF-STEINBERGER 66.		
F+- B	BOHM 69 F+-=(34.9+-10.0)+69*(DM-.545) DEG. FR=-49.9+-5.4 DEG.		
F+- F	FAISSNER 69 ERROR ENLARGED TO INCLUDE ERROR IN REGENERATOR PHASE.		
F+- J	JENSEN 70 F+-=(42.4+-4.0)+576*(DM-.538) DEG. FR=-42.7+-5.0 DEG.		
F+- D	BALATS 71 F+-=(39.0+-12.0)+198*(DM-.544) DEG. FR=-43.0+-4.0 DEG.		
F+- P	CARNegie 72 F+- IS INSENSITIVE TO DM. FR=-56.2+-5.2 DEG.		
F+- G	GEWENIG2 74 F+-=(49.4+-1.0)+565*(DM-.540) DEG. FR=-40.9+-2.6 DEG.		
F+- H	CARITHERS 75 F+-=(45.5+-2.8)+224*(DM-.5348) DEG.		
F00	PHASE OF ETA 00 (DEGREES)	-----	
F00	FIRST QUADRANT PREFERRED		
F00 C	51.0 30.0	CHOLLET 70 OSPK CU REG.,4 GAMMAS	10/70
F00 W	56 38.0 25.0	WOLFF 71 OSPK CU REG.,4 GAMMAS	12/71
F00 C	CHOLLET 70 USES REGENERATOR PHASE FR=-46.5+-4.4 DEG.		
F00 W	WOLFF 71 USES REGENERATOR PHASE FR=-48.2+-3.5 DEG.		
F00
F00 AVG	43.3 19.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
F00 STUDENT	43.3 20.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
F00 FIT	48.3 13.2	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	2/76*
DF	PHASE DIFFERENCE F00 - F+- (DEGREES)	-----	
DF B	7.6 18.0	BARBIELLI 73 ASPK	7/73
DF B	INDEPENDENT OF REGENERATOR MECHANISM,DM,AND LIFETIMES.		
DF
DF FIT	3.3 13.1	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	2/76*

Superweak Model Predictions for ϕ_{+-} and $Re\epsilon$

The superweak model of Wolfenstein, Phys. Letters 13, 562 (1964) predicts that

$$\phi_{+-} = \phi_{00} = \tan^{-1}\left(\frac{2\Delta m\tau_S}{h}\right)$$

and

$$Re\epsilon = |\eta_{+-}| \left[1 + \left(\frac{2\Delta m\tau_S}{h}\right)^2 \right]^{-1/2}$$

These expressions and the values of the $K_L^0 - K_S^0$ mass difference $\Delta M = (0.5349 \pm 0.0022) \times 10^{10} h \text{ sec}^{-1}$, the K_S^0 mean life $\tau_S = (0.8930 \pm 0.0023) \times 10^{-10} \text{ sec}$, and the magnitude of the $K_L^0 \rightarrow \pi^+\pi^-/K_S^0 \rightarrow \pi^+\pi^-$ amplitude ratio $|\eta_{+-}| = (2.272 \pm 0.023) \times 10^{-3}$, all from the current edition, result in the predictions that

$$\phi_{+-} = \phi_{00} = (43.69 \pm 0.13)^\circ$$

and

$$Re\epsilon = (1.643 \pm 0.017) \times 10^{-3}$$

These can be compared with the experimental values

$$\phi_{+-} = (45.0 \pm 1.2)^\circ$$

$$\phi_{00} = (48.3 \pm 13.2)^\circ$$

$$Re\epsilon = (1.624 \pm 0.088) \times 10^{-3}$$

where $Re\epsilon$ has been computed using the relation

$$Re\epsilon = \frac{\delta}{2} \left(\frac{|1-x|^2}{|1-|x||^2} \right)$$

and our current values of the charge asymmetry parameter for leptonic K_L^0 decay

$\delta = (0.330 \pm 0.012)\%$ and the $\Delta S = -\Delta Q$ amplitude $(Re\epsilon, Im\epsilon) = (0.008 \pm 0.020, -0.003 \pm 0.027)$.

The superweak predictions are in good agreement with the data.

The values of τ_S and $|\eta_{+-}|$ have undergone major revision in the past few years (see notes on these parameters in the K_S^0 and K_L^0 sections of the Data Card Listings). We have utilized only the post-1971 experimental values of these parameters and have made necessary changes in the related parameters ΔM and ϕ_{+-} (see footnotes in the related data card subsections D and F+- above).

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

 K^0 , η

FRESTON 67 PRL 18 176
 FITCH 67 PR 164 1711
 HAWKINS 67 PR 156 1444
 HILL 67 PRL 19 668

 HOPKINS 67 PRL 19 185
 KADYK 67 PRL 19 597
 KULYUKIN 67 PREPRINT
 LOWYS 67 PL 248 75
 MICHKE 67 PRL 18 138
 NEFKENS 67 PR 157 1233
 TODOROFF 67 THESIS

 ABRAMS 68 PR 176 1603
 ARNOLD 68 PL 288 56
 ARONSON 68 PRL 20 287
 ALSO 69 PR 175 1708
 BALATZ 68 PL 268 320
 BARTLETT 68 PRL 21 558

 BASILE 68 PL 268 542
 BASILE2 68 PL 288 58
 BENNETT1 68 PL 278 244
 BENNETT2 68 PL 278 248
 BLANNETD 68 PRL 21 1650
 BUDAGOV 68 NC 57A 182
 ALSO 68 PL 288 215

 CARNEGIE 68 PRINC 1744 THESIS1
 JAMES 68 NP 88 365
 ALSO 68 PRL 21 257
 KULYUKIN 68 PRL 26 20
 KUNZ 68 THESIS (PU 46)
 MELHOP 68 PR 172 1613
 THATCHER 68 PR 174 1674

 BANNER 69 PR 188 2033
 ALSO 68 PRL 21 1103
 ALSO 68 PL 21 1107
 BEILLIER 69 PL 308 202
 BENNETT 69 PL 298 517
 BOHM 69 PR 89 605
 ALSO 68 PL 278 321

 BOTT-BOD 69 CERN 69-7 329
 CENCE 69 PRL 22 1210
 EVANS 69 PRL 23 427
 FAISSNER 69 PL 308 204
 FOETH 69 PL 308 282

 GAILLARD 69 NC 59A 453
 ALSO 67 PRL 18 20
 GOBBI 69 PRL 22 685
 LITTENBERG 69 PRL 22 654
 LONGO 69 PR 181 1808
 PACIOTTI 69 THESIS, UCRL 19446
 SAAL 69 THESIS

 ALBROW 70 PL 338 516
 ARONSON 70 PRL 25 1057
 BARNIN 70 PL 338 377
 BASILE 70 PR D2 78
 BUCHANAN 70 PL 338 523
 ALSO PRIVATE COMMUNICATION
 BUDAGOV 70 PR D2 815
 ALSO 68 PL 288 215

 CHIEN 70 PL 338 627
 ALSO PRIVATE COMMUNICATION
 CHO 70 PR D1 3031
 ALSO 67 PRL 19 668
 CHOLLET 70 PL 318 458
 CULLEN 70 PL 328 523

 DARRIULA 70 PL 338 249
 FAISSNER 70 NC 70A 57
 JENSEN 70 THESIS
 ALSO 69 PRL 23 615
 MARX 70 PL 328 219
 ALSO 70 THESIS, NEVIS 179

 SCRIBANO 70 PL 328 224
 SMITH 70 PL 328 133
 WEBBER 70 PR D1 1967
 ALSO 69 UCRL 19226 THESIS 18

 BALATS 71 S JNP 13 53
 BAPMIN 71 PL 358 604
 BIS 71 PL 368 533
 BURGUN 71 LNC 2 1169
 CARNEGIE 71 PR D4 1
 CHAN 71 LBL-350 THESIS
 CHIEN 71 PL 358 261
 ALSO 72 DALLY

 CHO 71 PR D3 1557
 CLARK 71 PRL 26 1667
 ALSO 70 UCRL 19709-THESIS
 ALSO 71 UCRL 20264-THESIS
 ALSO 74 SLAC-PUB-1498
 ENSTROM 71 PR D4 2629
 ALSO 70 THESIS (SLAC 125)

 HILL 71 PR D4 7
 JAMES 71 PL 358 265
 MEISNER 71 PR D3 59
 PEACH 71 PL 358 351

 REPELLIN 71 PL 368 603
 WEBBER 71 PR D3 64
 ALSO 68 PRL 21 498
 ALSO 69 UCRL 19266-THESIS
 WOLFF 71 PL 368 517

 ALBROW 72 NP 844 1
 AIFORD 72 PL 348 47
 BANNER1 72 PRL 28 1597
 BANNER2 72 PRL 29 237
 BARNIN1 72 S JNP 15 636
 BARNIN2 72 S JNP 15 638
 BURGUN 72 NP 850 194
 CARNEGIE 72 PR D6 2335

 FRESTONE, KTH, LACH, SANDWEISS,+ (VALE,BNL)
 FITCH, ROTH, RUSS, VERNON (PRINCETON)
 C J B HAWKINS (VALE)
 HILL, LUERS, ROBINSON, CANTER+ (BNL,CARNEGIE)

 HOPKINS, BACON, EISLER (BNL)
 KADYK, CHAN, DRIJARD, OREN, SHELDOON (LRL)
 KULYUKINA+MESTVIRISHVILI+NEAGU+ (JINR)
 LOWYS, AUBERT, CHOUNET, PASCAUD+ (EPCL, ORSA)
 MICHKE, ABASHIAN, ABRAMS+ (ILLINOIS)
 *ABASHIAN, ABRAMS, CARPENTER, FISHER+ (ILL)
 JOHN A TODOROFF (ILLINOIS)

 *ABASHIAN, MICHKE, NEFKENS, SMITH+ (ILLINOIS)
 ARNOLD, BUDAGOV, CUNDY AUBERT+ (CERN-ORSA)
 S H ARONSON, K W CHEN (PRINCETON)
 S H ARONSON, K W CHEN (PRINCETON)
 BALATZ, BERZIN, VISHNEVSKY, GALANINA+ (ITEP)
 BARTLETT, CARNEGIE, FITCH+ (PRINCETON)

 BASILE, CRONIN, THEVENET, TURLAY+ (SLAC)
 *CRONIN, THEVENET, TURLAY, ZYBERAJCH+ (SLAC)
 BENNETT, NYGREN, STEINBERGER+ (COLUMBIA+CERN)
 BENNETT, NYGREN, STEINBERGER+ (COLUMBIA+CERN)
 BLANNIED, LEVIT, ENGELS+ (CAS+HARV+MGI)
 BUDAGOV, BURMESTER, CUNDY+ (CERN-ORSA, IPNP)
 *CUNDY, MYATT, NEZRICK+ (CERN, ORSA, EPOL)

 R, K, CARNEGIE (PRINCETON)
 F JAMES, H BRIAND (IPNP, CERN)
 HELLAND, LONGO, YOUNG (UCLA, MICH)
 *KULYUKINA, MESTVIRISHVILI, NEAGU+ (PRINCETON)
 P F KUNZ (PRINCETON)
 MELHOP, MURTY BOWLES, BURNETT+ (LA JOLLA)
 THATCHER, ARABASHIAN, ABRAMS, CARPENTER+ (ILL)

 *CRONIN, LIU, PILCHER (PRINCETON)
 BANNER, CRONIN, LIU, PILCHER (PRINCETON)
 BANNER, CRONIN, LIU, PILCHER (PRINCETON)
 BEILLIERE, BOUTANG, LIMON (EPOL)
 *NIGREN, SAAL, STEINBERGER+ (COLU+BNL+CERN)
 *DARRIULAT, GROSSO, KAFTANOV+ (CERN)
 BOHM, DARRIULAT, GROSSO, KAFTANOV (CERN)

 BOTT-RODENHAUSEN, DE BOUJARD, CASSEL+ (CERN)
 CENCE, JONES, PETERSON, STENGER+ (HAWAII, LPL)
 EVANS, GOLDEN, MUIR, PEACH+ (EDINBURGH, CERN)
 *FOETH, STAUDE, TITTEL+ (AACH, CERN, TORI)
 *HOLDER, RADERMACHER+ (AACH, CERN, TORI)

 *GALBRAITH, HUSSRI, JANE+ (CERN, RHEL, AACHEN)
 *KRIESEN, GALBRAITH, HUSSRI+ (CERN+RHEL+AACH)
 *GREEN, MAKEL, MOFFETT, ROSEN, GOZ+ (ROCH+IRUTG)
 LITTENBERG, FELD, PICCIONI, MEHLHOP+ (UCS)
 H J LONGO, K K YOUNG, J A HELLAND (MICH, UCLA)
 (LRL)
 H J SAAL (COLUMBIA)

 *ASTON, BARBER, BIRD, ELLISON+ (MCHS+DARE)
 *EHRICH, HOFER, JENSEN+ (EFI, ILL, SLAC)
 *BARYLON, BORISOV, BYSHEVA+ (ITEP, JINR)
 *CRONIN, THEVENET, TURLAY, ZYBERAJCH+ (SLAC)
 *DRICKEY, RUDDICK, SHEPARD+ (SLAC, JHU, UCLA)
 B + COX, FEB, 71
 *CUNDY, MYATT, NEZRICK+ (CERN, ORSA, EPOL)
 *CUNDY, MYATT, NEZRICK+ (CERN, ORSA, EPOL)

 C-Y, CHEN, COX, ETTLINGER+ (JHU+SLAC+UCLA)
 B. COX, FEB, 71.

 *GAILLARD, CANTER, ENGLER, FISK+ (CERN, BNL, CASE)
 HILL, LUERS, ROBINSON, SAKITT+ (BNL, CERN)
 *GAILLARD, JANE, RATCLIFFE, REPELLIN+ (CERN)
 *DARRIULAT, DEUTSCH, FOETH+ (AACH, CERN, TORI)

 *FERREDO, GROSSO, HOLDER+ (AACH, CERN, TORI)
 *REITHLER, THOME, GAILLARD+ (AACH, CERN, RHEL)
 D A JENSEN (EFI)
 JENSEN, ARONSON, EHRICH, FRYBERGER+ (EFI, ILL)
 *NYGREN, PEOPLES, STEINBERGER+ (COLU, HARV, CERN)
 JAY MARX (COLUMBIA)

 *MANNELLI, PIERAZZINI, MARX+ (PISA, COLU, HARV)
 *WANG, WHATLEY, ZORN, HORNBOSTEL (UMD, BNL)
 *SOLMITZ, CRAWFORD, ALSTON-GARNJUST (LRL)
 B WEBBER (LRL)

 *BEREZIN, VISHNEVSKII, GALANINA+ (ITEP)
 *BARYLOV, VESELOVSKY, DAVIDENKO+ (ITEP)
 *DARRIULAT, FERREDO, RUBBIA+ (AACH, CERN, TORI)
 *LESQUY, MULLER, PAULI+ (SLAC+CERN+OSLO)
 *CESTER, FITCH, STROVINK, SULAK (PRIN)
 J-HIONG-SING CHAN (LBL)
 *COX, ETTLINGER, RESVANIS+ (JHU, SLAC, UCLA)

 *DRALLE, CANTER, ENGLER, FISK+ (CERN, BNL, CASE)
 *ELIOPF, FELD, FRISCH, JOHNSON, KERTH+ (LRL)
 ROLLAND JOHNSON (LRL)
 HENRY FRISCH (LRL)
 R-C, FELD (SLAC)
 *AKAVIA, COOMBS, DORFAN+ (SLAC, STAN)
 J E ENSTROM (STANFORD)

 *SAKITT, SKJEGGSTAD, CANTER+ (BNL, CERN, CASE)
 *MONTANET, PAUL, PAULI+ (CERN+BNL+OSLO)
 *MANN, HERTZBACH, KOFLER+ (MASA+BNL+VALE)
 *EVANS, MUIR, BUDAGOV, HOPKINS+ (EDIN, CERN)

 *NDLFF, CHOLLET, GAILLARD, JANE+ (ORSA, CERN)
 *SOLMITZ, CRAWFORD, ALSTON-GARNJUST (LRL)
 WEBBER, SOLMITZ, CRAWFORD, ALSTON-GARNJUST (LRL)
 B WEBBER (LRL)
 *CHOLLET, REPELLIN, GAILLARD+ (ORSA, CERN)

 *ASTON, BARBER, BIRD, ELLISON+ (MCHS+DARE)
 *BROWN, MASER, MAUNG, MILLER, RUDERMAN+ (UCSD)
 *CRONIN, HOFFMAN, KNAPP, SHOCHET (PRINCETON)
 *CRONIN, HOFFMAN, KNAPP, SHOCHET (PRINCETON)
 *DAVIDENKO, DEMIDOV, DOLGO ENKO+ (ITEP)
 *BARYLOV, DAVIDENKO, DEMIDOV+ (ITEP)
 *LESQUY, MULLER, PAULI, + (SLAC+CERN+OSLO)
 *CESTER, FITCH, STROVINK, SULAK (PRINCETON)

DALLY 72 PL 418 647
 ALSO 70 CHIEN
 ALSO 71 CHIEN
 GRAHAM 72 NC 9A 166
 HOLDER 72 PL 408 141
 JAMES 72 NP 849 1
 KRENZ 72 LNC 4 213

 MANN 72 PR D6 137
 MANTSCH 72 NC 9A 160
 METCALF 72 PL 408 703
 NEUHOFER 72 PL 418 642
 PICOZZI 72 PRL 29 1412
 ALSO 74 PR D9 2939
 VOSBURGH 72 PR D6 1834
 ALSO 71 PRL 26 866

 ALBROW 73 NP 858 22
 ALEXANDER 73 NP 865 301
 BARBIELL 73 PL 438 529
 BRANDBERG 73 PR D8 1978
 CARITHEI 73 PRL 30 1336
 CARITHEI 73 PRL 31 1025
 EVANS 73 PR D7 36
 ALSO 69 PRL 23 427

 FACKLER 73 PRL 31 847
 FITCH 73 PRL 31 1524
 ALSO 72 COO-3072-13
 GJESDAL 73 2ND AIX CONF
 HART 73 NP 866 317
 MALLARY 73 PR D7 1953
 ALSO 70 PRL 25 1214

 MCCARTHY 73 PR D7 687
 MCCARTHY, BREWER, BUDNITZ, ENTIS, GRAVEN+ (LBL)
 ALSO 71 THESIS LBL-550
 MESSNER 73 PRL 30 876
 PEACH 73 PL 438 441
 SCANDWEIS 73 PRL 30 1002
 WILLIAMS 73 PRL 31 1521

 ALBRECHT 74 PL 488 393
 BIS 74 PL 508 504
 ROBIUSIT 74 PL 428 291
 DONALD51 74 PRL 33 554
 EVANS 74 DONALDSON 3
 DONALD52 74 PR D9 2960
 ALSO 73 PRL 31 337
 DONALD53 74 SLAC 184-THESIS

 GEWENIG1 74 PL 488 483
 ALSO 74 CERN INT. REPT.
 GEWENIG2 74 PL 488 487
 ALSO 74 PL 528 119
 GEWENIG3 74 PL 528 108
 GJESDAL1 74 PL 528 113

 MESSNER 74 PRL 33 1458
 NIEBERGA 74 PL 498 103
 SLONE 74 SLAC 181-THESIS
 WANG 74 PR D9 540
 WILLIAMS 74 PRL 33 240
 WOO 74 LNC 10 38

 BALDOCEO 75 NC 25A 688
 BLUMENTHAL 75 PRL 34 164
 BUCHANAN 75 PR D11 457
 CARITHERS 75 PRL 34 1244
 KHEN 75 LBL-4275 THESIS
 FUSHIM 76 PRL 36 348

 ALEXANDE 62 PRL 9 69
 JOVANOVIC 63 BNL CONF 42
 STERN 64 PRL 12 459
 BEHR 65 ARGONNE CONF 59
 MESTVIRI 65 JINR P 2449
 TRILLING 65 UCRL 16473
 UPDATED FROM 1965 ARGONNE
 GINSBERG 67 PR 162 1570

 RUBBIA 67 PL 248 531
 ALSO 1 66 PL 20 207
 ALSO 2 66 PL 21 595
 ALSO 3 66 PL 23 167
 SCHMIDT 67 NEVIS 160(THESIS)
 CRONIN 68 VIENNA CONF P.281
 BECHERRA 70 PR D1 1452
 GINSBERG 70 PR D1 229
 HEUSSE 70 LNC 3 449
 GINSBERG 73 PR D8 3887

 *C. RUBBIA, J. STEINBERGER (CERN+COLU)
 ALFF-STERNBERGER, HEUER, KLEINNECHT+ (CERN)
 ALFF-STERNBERGER, HEUER, KLEINNECHT+ (CERN)
 C. RUBBIA, J. STEINBERGER (CERN+COLU)
 P. SCHMIDT (COLUMBIA)
 CRONIN, RAPPORTEURS TALK (PRINCETON)
 T BECHERRA Y (ROCH)
 E S GINSBERG (IIT HALFA)
 *AUBERT, PASCAUD, VIALLE (ORSA)
 E S GINSBERG, J SMITH (MIT+STON)

14 ETA (549, JPG0-0) 1=0

FOR C. BALATZ'S REVIEW OF THE ETA MESON, SEE PROC. UNIV. OF PENN. CONF. ON NEUTRON SPECTROSCOPY (W.A. BENJAMIN, N.Y., 1968)

14 ETA MASS (MEV)

M	53 549.0	1.2	BASTIEN	62 HBC
M	35 546.0	4.0	ALFF	62 HBC
M	91 548.0	1.0	PICKUP	62 HBC
M	549.3	2.9	DELICOURT	63 CNTR
M	148 549.0	0.7	FJELSCHE	64 HRC
M	325 552.0	3.0	KRAEMER	64 DBC
M	548.2	0.65	FOSTER	65 HBC
M	250 555.0	2.0	JAMES	66 HBC
M	548.82	0.56	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
M	STUDENT 548.72	0.44	AVERAGE USING STUDENTS(H1, I1)	-- SEE TEXT (SEE IDEOGRAM BELOW)

7/66

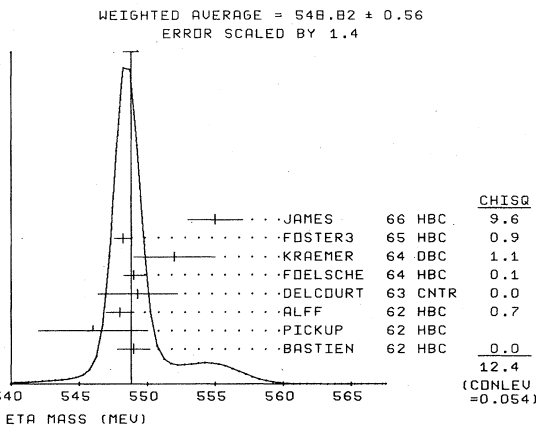
6/66

Stable Particles

η

Data Card Listings

For notation, see key at front of Listings.



14 ETA DECAY RATES

Note on Conflicting Results for $\eta \rightarrow \gamma\gamma$ Decay Rate

The BROWMAN 74 result for $\Gamma_{\gamma\gamma}$ is three standard deviations below the earlier BEMPORAD 67 result. This is because the BEMPORAD 67 analysis finds a negligible strong-production amplitude and attributes all of their observed η photoproduction cross section to Coulomb production. BROWMAN 74, on the other hand, find that the strong-production contribution to the cross section is significant, especially at lower energies, and that the Coulomb contribution is smaller, resulting in a smaller value of $\Gamma_{\gamma\gamma}$.

The η photoproduction data of BROWMAN 74 were taken at incident γ -ray energies of 5.8, 9.0, and 11.45 GeV on beryllium, aluminum, copper, silver, and uranium targets, while the BEMPORAD 67 data were taken at 4.0 and 5.5 GeV on lead, silver, and zinc. The higher energies, higher statistics, better angular resolution, and larger range of atomic weights of the BROWMAN 74 experiment result in a more reliable separation of the Coulomb amplitude from the strong-production amplitude. In addition, the bad approximation mentioned in BROWMAN 74 footnote 4 was present in the theoretical work utilized by BEMPORAD 67 and this may account for some of the discrepancy.

Browman et al. state that the BROWMAN 74 result is compatible with the BEMPORAD 67 data, although the agreement is not as good as that given by the BEMPORAD 67 fit to these data. On the other hand, the BEMPORAD 67 result appears to be incompatible with the 9 GeV data (especially uranium) shown in Fig. 2 of BROWMAN 74.

We quote the BROWMAN 74 result.

14 ETA WIDTH

Mode	Branching Fraction	Decay Rate (Units KeV)	Source
91	(10.0) OR LESS	ALFF	62 HBC
148	(10.0) OR LESS	FDELSCHKE	64 HBC
31	(12.0) OR LESS	JAMES	66 HBC
(4.0) OR LESS	BALTAY	66 DBC	
(.9) OR LESS	CL=.95	JONES	66 CNTR

ETA WIDTH DETERMINED FROM MASS SPECTRUM (UNITS KEV)

THIS IS THE PARTIAL DECAY RATE (W1) FOR THE MODE (ETA INTO 2GAMMA) DIVIDED BY THE FITTED BRANCHING FRACTION (P1) FOR THAT MODE.

FIT 0.85 0.12 FROM FIT

14 ETA PARTIAL DECAY MODES

Mode	Decay Masses
P1	ETA INTO 2GAMMA
P2	ETA INTO 5PI0
P3	ETA INTO P1+ P1- P10
P4	ETA INTO P1+ P1- GAMMA
P5	ETA INTO E+ E- P10 (VIOLATES C IN E.M.I.)
P6	ETA INTO E+ E- P1+ P1-
P7	ETA INTO P10 2GAMMA
P8	ETA INTO E+ E- GAMMA
P9	ETA INTO 2P10 GAMMA (VIOLATES C)
P10	ETA INTO P1+ P1- P10 GAMMA
P11	ETA INTO P1+ P1- 2GAMMA
P12	ETA INTO MU+ MU- GAMMA
P13	ETA INTO MU+ MU- GAMMA
P14	ETA INTO MU+ MU- P10
P15	ETA INTO P1+ P1-

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (\delta P_i \cdot \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 7	P 8
P 1	.3799+-0098					
P 2	-.2691	.2990+-0106				
P 3	-.3224	-.2353	.2358+-0056			
P 4	-.2866	-.2095	.8201	.0489+-0013		
P 7	-.4271	-.5781	-.0939	-.0801	.0314+-0109	
P 8	-.0434	-.0326	-.0494	-.0501	-.0036	.0050+-0012

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e., $G_i = \Gamma_i = \Gamma_{total} P_i$, in appropriate units. In analogy to the matrix above, the diagonal elements are $G_i \pm \delta G_i$, where $\delta G_i = \sqrt{(\delta G_i)^2}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta G_i \delta G_j) / (\delta G_i \cdot \delta G_j)$. Note that, because of the error in Γ_{total} , the errors and correlations here are not directly derivable from those above.

	G 1	G 2	G 3	G 4	G 7	G 8
G 1	.3240+-0460					
G 2	-.9386	.2550+-0385				
G 3	.9689	.9513	.2011+-0294			
G 4	.9646	.9471	1.0097	.0417+-0061		
G 7	-.3661	-.3028	.4240	.4226	.0268+-0104	
G 8	-.5178	-.5082	.5233	.5199	.2218	.0043+-0012

Mode	Rate (Units KeV)	Source	Notes
W1	ETA INTO 2GAMMA (UNITS KEV)	(G1)	
W1 B	(1.00) (0.22)	BEMPORAD 67 CNTR	PRIMAKOFF EFFECT 11/75*
W1	0.324	0.046	BROWMAN 74 CNTR PRIMAKOFF EFFECT 7/74*
W1 B	BEMPORAD 67 GIVES W1=1.21+-26 KEV ASSUMING THAT W1/TOTAL=0.314.		11/75*
W1 B	BEMPORAD PRIVATE COMMUNICATION GIVES MORE GENERAL RESULT AS		11/75*
W1 B	W1*W1/TOTAL=.380+-083. WE EVALUATE THIS USING W1/TOTAL=.38+-01.		11/75*
W1 B	NOT INCLUDED IN AVERAGE BECAUSE THE UNCERTAINTY RESULTING FROM THE		2/76*
W1 B	SEPARATION OF THE COULOMB AND NUCLEAR AMPLITUDES HAS APPARENTLY		2/76*
W1 B	BEEN UNDERESTIMATED. SEE NOTE ON DISCREPANCY ABOVE.		2/76*
W1	.324 0.046 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)		

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

n, p, n

BAGLINI 67 PL 248 637
BAGLINI 67 BAPS 12 567
BALTAY 67 PRL 19 1495
BALTAY 67 PRL 19 1498
BEMPORAD 67 PL 258 380
ALSO PRIVATE COMMUNICATION

ARNOLD 68 PL 278 466
BAZIN 68 PRL 20 895
BULLOCK 68 PL 278 462
GORMLEY 68 PRL 21 402
WEHMANN 68 PRL 20 748

BAGLIN 69 PL 298 445
ALSO 70 NP 822 66
HYAMS 69 PL 298 128
JACQUET 69 NC 58 743
MULLER 69 THESIS

BAGLIN 70 NP 822 66
BUTTRAM 70 PRL 25 1358
CARPENTER 70 PR D1 1303
COX 70 PRL 24 534
DANBURG 70 PR D2 2564
DEVONS 70 PR D1 1936
GORMLEY 70 PR D2 501
ALSO 70 NEVIS 181 (THESIS)

BASTIEN 62 PRL 8 114
CARMONY 62 PRL 8 117
ROSENFELD 62 PRL 8 293

16 PROTON(938, J=1/2) I=1/2

16 PROTON MASS (MEV)
M (938.256) (0.005) COHEN 65 RVUE 7/66
M (938.2592) (0.0052) TAYLOR 69 RVUE USING NEW E/H 7/70
M 938.2796 0.0027 COHEN 73 RVUE 3/74

16 PROTON MEAN LIFE (UNITS 10**26 YR)
T (0.00001) OR MORE GOLDHABE 54 TH 232 FISS. MODE INDEPEN 12/75*
T (0.002) OR MORE FLEROV 57 TH 232 FISS. MODE INDEPEN 12/75*
T B (1.5) OR MORE BACKENSTOSS 60 CNTR 12/75*
T B (60.0) OR MORE KRUPF 65 CNTR 12/75*
T (200.0) OR MORE GURR 67 CNTR DEP. ON DECAY MODE 12/75*
T (1300.0) OR MORE BERGAMASCO 74 CNTR 12/75*
T R 20000.0 OR MORE REINES 74 CNTR 12/75*
T B KRUPF 73 VALUE SENSITIVE TO PARTICULAR DECAY MODES OF PROTON 12/75*
T B REINES 74 LOOKS AT DECAY CHAINS ENDING IN MUON DECAY. 12/75*

16 PROTON MAGNET. MOMENT(E/2MP)
MM (2.79276) (.00002) COHEN 65 RVUE 7/70
MM (2.792782) (.000017) TAYLOR 69 RVUE USING NEW E/H 7/70
MM 2.7928456 .0030011 COHEN 73 RVUE 3/74

16 ANTI-PROTON MAGNETIC MOMENT (E/2MP)
MM1 O (-1.8) (1.2) BUTTON 62 CNTR 11/75*
MM1 R (-2.83) (0.10) FOX 72 CNTR 7/75*
MM1 O -2.818 0.356 ROBERTS 74 CNTR 7/75*
MM1 O OLD EXPERIMENT WITH LARGE ERROR. NOT AVERAGED.
MM1 R ROBERTS 74 IS REANALYSIS OF FOX 72 DATA. REPLACES OLD FOX VALUE. 7/75*

16 PROTON ELECTRIC DIPOLE MOMENT (UNITS 10**23 E CM)
NONZERO VALUE IMPLIES VIOLATION OF T AND P IN EM INTERACTION
EDM 16 700. 900. HARRISON 69 MBR 10/69

REFERENCES FOR PROTON
GOLDHABE 54 PR 96 1157 FN022 GOLDHABER, F. REINES (LOS ALAMOS, BNL)
FLEROV 57 SOV PHYS DOK 3 79 FLEROV, KLIOCHKOV, SKOBEIN, TERENTEV (USSR)
BACKENSTOSS 60 NC 16 749 BACKENSTOSS, FRAUENFELDER, HYAMS (CERN)
BUTTON 62 PR 127 1297 J BUTTON, B MAGLIC (LBL)
COHEN 65 RMP 37 537 *DUMOND (IN AMER. AVIATION SCIENCE CENT., CIT)
KRUPF 65 PR 137 8 740 W R KRUPF, F REINES (CASE INST TECHNOLOGY)
GURR 67 PR 158 1321 GURR, KRUPF, REINES, MEYER (CASE, JOHANNESBURG)

HARRISON 69 PRL 22 1263 HARRISON, SANDARS, WRIGHT (CLARENDON OXFORD)
TAYLOR 69 RMP 41 375 *PARKER, LANGENBERG (PRIN+UCI+PENN)
FOX 72 PRL 29 193 *BARNES, EISENSTEIN (BNL+CARN+VPI+WILL+HYGM)
COHEN 73 J. PHYS. CHEM. REF. DATA 2, P. 663, E. R. COHEN, B. N. TAYLOR
BERGAMASCO 74 LNC 11 636 BERGAMASCO, PICCHI (TORI+FRAS)
REINES 74 PRL 32 493 *KROUCH (UCI+CASE)
ROBERTS 74 PRL 33 1181 *COX, FCKHAUSE (WILL+VPI+CARN+HYDM+CIT+BNL)
ALSO 75 PR D12 1232 ROBERTS, COX (WILL+VPI+CARN+HYDM+CIT+BNL)

17 NEUTRON(939, J=1/2) I=1/2

17 NEUTRON MASS (MEV)
M T (939.5527) (0.0052) TAYLOR 69 RVUE USING NEW E/H 7/70
M T 939.5731 0.0027 COHEN 73 RVUE 3/74
M T THESE DETERMINATIONS OF NEUTRON MASS NOT INDEPENDENT OF 7/70
M T NEUTRON-PROTON MASS DIFFERENCE MEASUREMENTS BELOW. 7/70

17 (NEUTRON) - (PROTON) MASS DIFFERENCE (MEV)
D M (1.29344) (0.00007) MATTAUCH 65 RVUE 3/71
D M 1.29343 0.00004 COHEN 74 RVUE 3/74
D M WE HAVE CONVERTED MATTAUCH NEUTRON-HYDROGEN MASS DIFFERENCE TO 3/71
D M NEUTRON-PROTON MASS DIFFERENCE USING CURRENT VALUE OF ELECTRON MASS 3/71
D M AND A HYDROGEN BINDING ENERGY OF 13.6 EV. 3/71

17 NEUTRON MAGNETIC MOMENT (MAGNETONS, 938.2 MEV)
MM -1.913148 0.000066 COHEN 56 RVUE 7/66

17 NEUTRON ELECTRIC DIPOLE MOMENT (UNITS 10**23 E CM)
TEST OF CP OR T VIOLATION IN THE EM INTERACTION
EDM (5.) OR LESS BAIRD 69 MBR INCLUDED IN DRESS 73 10/69
EDM 0.32 0.75 DRESS 73 MBR .LT. 10**23 CL=.80 6/73

17 NEUTRON MEAN LIFE (UNITS 10**3 SEC)
THE MEASUREMENT OF THE NEUTRON MEAN LIFE BY SOGNSOVSKII 59 HAS BEEN DISCARDED SINCE 1. IT DISAGREES WITH THE BETTER AND MORE RECENT RESULT OF CHRISTENSEN 67. 2. THE VALUE OF GA/GV DERIVED FROM THE NEW VALUE OF THE MEAN LIFE AGREES WELL WITH THE GA/GV VALUE OBTAINED FROM THE FREE NEUTRON DATA.
T (1.012) (0.021) SOGNSOVSKI 59 PILE SEE NOTE E 7/68
T (0.935) (0.014) CHRISTENSEN 67 PILE REPL BY CHRISTENSEN 73 3/68
T 0.918 0.014 CHRISTENSEN 73 PILE 6/72
E ERROR CHANGED BECAUSE ERROR IN CROSS SECTION FOR NEUTRON ABSORPTION IN GOLD HAS BEEN REDUCED.

17 NEUTRON BETA DECAY COUPLING CONSTANTS

RELATED TEXT SECTION VI D.1
AV GA/GV (SEE TEXT FOR SIGN CONVENTION)
AV (-1.250) (0.044) CONFORTO 67 RVUE SEE NOTE C BELOW
AV EP (-1.23) (0.01) CHRISTENSEN 67 CNTR N DECAY FT VALUE 11/68
AV P (-1.22) (0.08) GRIGOREV 68 CNTR E-NEU ANG CORREL 10/71
AV P (-1.26) (0.02) CHRISTENS 70 CNTR PE/NEUT SPIN CORREL 10/71
AV EP (-1.27) (0.025) EROZOLIMS 71 CNTR PE/NEUT SPIN CORREL 10/71
AV EP (-1.239) (0.011) CHRISTENSEN 72 CNTR N DEC. + FT VALUE 1/73
AV P (-1.265) (0.016) KRUPF 73 RVUE N DECAY ALONE 1/73
AV P -1.250 0.009 KRUPF 73 RVUE N DEC. + FT VALUE 1/73
AV E -1.250 0.036 DOBROZEMSK 75 CNTR E-NEU ANG CORREL. 12/75*
AV E CONFORTO 67 COMBINES FREE NEUTRON DATA TO 1967. REPL. BY KROPP 73. 1/73
AV E THESE EXPERIMENTS MEASURE THE ABSOLUTE VALUE OF GA/GV ONLY 10/71
AV P KROPP 73 VALUE OBTAINED BY FITTING ALL DATA THROUGH 1972. 1/73
AV
AV AVG -1.2500 0.0087 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AV STUDENT -1.2500 0.0094 AVERAGE USING STUDENT(S/H/1.11) -- SEE TEXT
F PHASE ANGLE OF GA RELATIVE TO GV (DEGREES)
F C (176.1) (6.4) CONFORTO 67 RVUE 11/68
F P (181.3) (1.3) EROZOLIMS 70 CNTR POLAR. NEUTRON 10/69
F P 181.1 1.3 KRUPF 73 RVUE N DECAY 1/73
F C CONFORTO 67 COMBINES FREE NEUTRON DATA TO 1967. REPL. BY KROPP 73. 1/73
F P KROPP 73 VALUE OBTAINED BY FITTING ALL DATA THROUGH 1972. 1/73

REFERENCES FOR NEUTRON
COHEN 56 PR 104 283 V W COHEN, CORNGOLD, RAMSEY (BNL+HARVARD)
SOGNSOVSKI 59 JETP 9 717 SOGNSOVSKII, SPIVAK, PROKOFEV (IAE MOSCOW)
MATTAUCH 65 NP 67 1 *THIELE, WAPSTRA (MAX PLANCK INST. CHEM.)
CHRISTENSEN 67 PL 268 11 CHRISTENSEN, NIELSON, BAHNSEN, BROWN (RISO)
CONFORTO 67 AP4H 22 15 G. CONFORTO (CERN)
GRIGOREV 68 SJNP 6 239 *GRISHIN, VLADIMIRSKII, NIKOLAEVSKII + (ITEP)
BAIRD 69 PR 179 1285 *MILLER, DRESS, RAMSEY (ORNL, HARV)
TAYLOR 69 RMP 41 375 *PARKER, LANGENBERG (PRIN+UCI+PENN)
CHRISTENSEN 70 PR D5 1628 CHRISTENSEN, KROPP, RINGO (ANL)
ERZOLIM 70 SJNP 11 583 EROZOLIMSKI, BONDARENKO, + (KURC MOSCOW)
ALSO PL 278 557 EROZOLIMSKI, BONDARENKO + (KURC IN MOSCOW)
ERZOLIM 71 JETPL 13 252 EROZOLIMSKI, BONDARENKO + (KURC MOSCOW)
CHRISTENSEN 72 PR D5 1628 CHRISTENSEN, NIELSON, BAHNSEN, BROWN (RISO)
COHEN 73 J. PHYS. CHEM. REF. DATA 2, P. 663, E. R. COHEN, B. N. TAYLOR
DRESS 73 PR D7 3147 DRESS, MILLER, RAMSEY (ORNL+HARV)
KROPP 73 ZPHY TO BE PUBL. A KROPP, H PAUL (LINZ)
ALSO 70 NP A154 160 KRUPF, H PAUL (WIEN)
COHEN 74 PRIVATE COMM. E. R. COHEN (ROCKWELL INTERNATL. SCI. CENTER)
DOBROZEMSK 75 PR D11 510 DOBROZEMSKY, KERSCHBAUM, MORAW, PAUL + (SEIR)

PAPERS NOT REFERRED TO IN DATA CARDS
JACKSON 57 PR 106 517 JACKSON, TREIMAN, WYLD (PRINCETON)
COHEN 65 RMP 37 537 *DUMOND (IN AMER. AVIATION SCIENCE CENT., CIT)
BHALLA 66 PL 19 691 C P BHALLA (ALABAMA)

Stable Particles

Data Card Listings

A

For notation, see key at front of Listings.

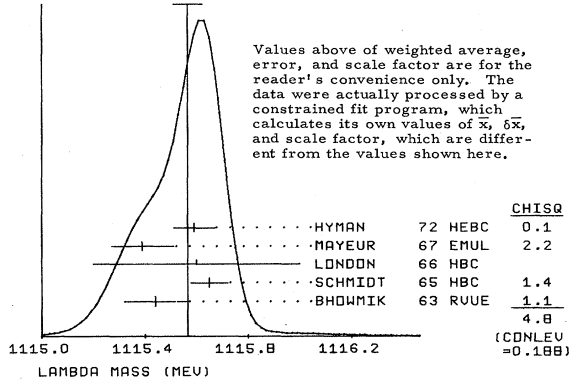


18 LAMBDA(1115,JP=1/2+) I=0

18 LAMBDA MASS (MEV)

M N SINCE OUR FINAL VALUES FOR THE SIGMA AND LAMBDA MASSES COME FROM... M L ABOVE LAMBDA MASS HAS BEEN RAISED 35 KEV TO ACCOUNT FOR 46 KEV... M S 635(1115.86) (0.09) BALTAY 65 HBC ERROR IS STATIS. 6/66... M S 488 1115.65 (0.07) SCHMIDT 65 HBC SEE NOTE N 3/74... M S 1147(1115.74) (0.04) CHIEN 66 HBC 6.9 PBAR P 9/67... M S 972(1115.69) (0.05) CHIEN 66 HBC 6.9 PBAR PANTIL 9/67... M S 1115.6 (0.4) LONDON 66 HBC 6/66... M S (1116.0) (0.2) BADIER 67 HBC 2.4 PBAR P,LLBAR 8/67... M S 195 1115.39 0.12 MAYEUR 67 EMUL 11/67... M B 1524(1115.52) (0.03) BOHM 70 EMUL 3/72... M S 935 1115.59 0.08 HYMAN 72 HBC 11/71... M B AVERAGE OF VERY INCONSISTENT DATA. ERROR STATISTICAL ONLY. AUTHORS 3/72... M B DETECT SYSTEMATIC EFFECT OF ABOUT .15 MEV, WHICH THEY ATTRIBUTE 3/72... M B TO ERROR IN RANGE-ENERGY RELATIONS. IN REGION BETA=0.6-0.7, 6X, 3/72... M B THIS EFFECT, IF CONFIRMED, WOULD AFFECT VERY LITTLE THE VALUES OF 3/72... M B BHOWMIK 63 AND MAYEUR 67... M S ERROR PURELY STATISTICAL... M AVG 1115.566 0.056 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) 4/72... M STUDENT 1115.568 0.053 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT 4/72... M FI 1115.600 0.048 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2) 2/76* (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 1115.566 ± 0.056
ERROR SCALED BY 1.3

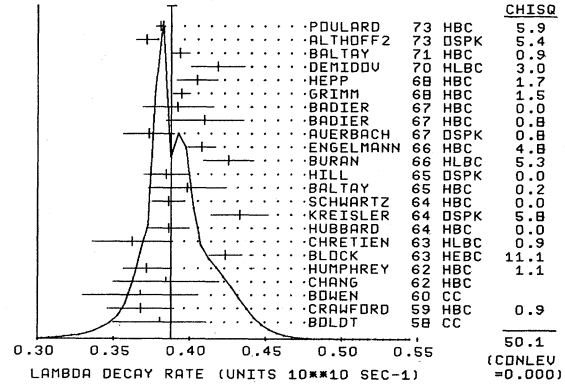


18 LAMDA - ANTI LAMBDA MASS DIFFERENCE (MEV)
DM 0.05 0.06 CHIEN 66 HBC 6.9 PBAR P 9/67
DM 0.29 0.15 BADIER 67 HBC 2.4 PBAR P 8/67
DM AVG 0.083 0.063 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
DM STUDENT 0.080 0.063 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

18 LAMBDA MEAN LIFE (UNITS 10**-10)

T 189 2.63 0.21 0.21 BOLDT 58 CC
T 825 2.72 0.16 0.16 CRAWFORD 59 HBC
T 140 2.72 0.29 0.27 BOWEN 60 CC
T 186 2.60 0.28 0.20 CHANG 62 HBC
T 799 2.69 0.11 0.11 HUMPHREY 62 HBC
T 2239 2.36 0.06 0.06 BLOCK 63 HBC
T 706 2.76 0.20 0.20 CHRETIEN 63 HLBC
T 794 2.59 0.09 0.09 HUBBARD 64 HBC
T 2260 2.51 0.10 0.10 KREISLER 64 OSPK
T 1378 2.59 0.07 0.07 SCHWARTZ 64 HBC
T 635 2.51 0.16 0.16 BALTAY 65 HBC 6/66
T 2534 2.6 0.1 65 OSPK
T 916 2.35 0.09 0.09 BURAN 66 HLBC 6/66
T S 1147 (2.50) (0.14) CHIEN 66 HBC 6.9 PBAR P 9/67
T S 972 (2.70) (0.20) CHIEN 66 HBC 6.9 PBAR P,ANTI 9/67
T 2213 2.452 0.056 0.054 ENGELMANN 66 HBC 9/66
T 585 2.68 0.13 0.11 AUERBACH 67 OSPK 8/67
T 2.44 0.15 67 HBC 6/68
T 2.45 0.15 67 HBC 2.4 PBAR P,ANTI 6/68
T 8342 2.535 0.35 0.35 GRIMM 68 HBC 6/68
T 2690 2.47 0.08 0.08 HEPP 68 HBC 8/68
T 1059 2.35 0.10 0.10 DEMODOV 70 HLBC PI-P, 3.86 GEV/C 12/70
T 4872 2.54 0.04 0.04 BALTAY 71 HBC K-P AT REST 6/71
T 6582 2.69 0.05 0.05 ALTHOFF2 73 OSPK PI+N TO K+LAMBDA 2/74
T 36K 2.626 0.020 0.020 POULARD 73 HBC K-P,KMOM 4T02.3 9/73
T S ERROR PURELY STATISTICAL.
T AVG 2.578 0.021 0.021 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.6)
T STUDENT 2.577 0.020 0.020 AVG. USING STUDENT10(H/1.11) -- SEE TEXT (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.3879 ± 0.0031
ERROR SCALED BY 1.6



18 (LAMBDA - ANTI LAMBDA)/AVG., MEAN LIFE DIFFERENCE

DT 0.044 0.085 BADIER 67 HBC 2.4 PBAR P 8/67

18 LAMBDA MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

MM -1.5 0.5 COOL 62 DSPK
MM 0.0 0.6 KERNAN 65 CC
MM 8553 -1.39 0.72 ANDERSON 64 HBC
MM 151 -0.5 0.28 CHARRIERE 65 EMUL
MM 49 -0.67 0.31 0.37 BARKOV 71 EMUL PRELIM. RESULT 2/72
MM 1300 -0.66 0.07 DAHJENSE 71 EMUL MAG FIELD=200G 6/71
MM 3868 -0.73 0.18 HILL 71 OSPK 10/71
MM AVG -0.672 0.061 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
MM STUDENT -0.670 0.067 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

18 LAMBDA ELECTRIC DIPOLE MOMENT (UNITS 10**-14 E CM)

NONZERO VALUE IMPLIES VIOLATION OF T AND P
EDM B 5.0 OR LESS CL=.95 GIBSON 66 EMUL 2/72
EDM B 1.0 OR LESS CL=.95 BARONI 71 EMUL 2/72
EDM B BARONI MEASURES (-5.9+2.9)*10**-15 E CM 2/72

18 LAMBDA PARTIAL DECAY MODES

DECAY MASSES
P1 LAMBDA INTO PROTON PI- 938+ 139
P2 LAMBDA INTO NEUTRON P0 939+ 134
P3 LAMBDA INTO PROTON MU- NEUTRINO 938+ 105+ 0
P4 LAMBDA INTO PROTON E- NEUTRINO 938+ .5+ 0
P5 LAMBDA INTO PROTON PI- GAMMA 938+ 139+ 0

18 LAMBDA BRANCHING RATIOS

R1 LAMBDA INTO (P PI-)/((P PI-)+(N P0)) (P1)/(P1+P2)
R1 0.627 0.031 CRAWFORD 59 HBC
R1 0.65 0.05 COLUMBIA 60 HBC
R1 U (0.685) (0.017) ANDERSON 62 HBC
R1 903 0.645 0.016 HUMPHREY 62 HBC
R1 U 6736 0.635 0.007 DOYLE 69 HBC PI-P TO LAM. KO 2/71
R1 4572 0.646 0.008 BALTAY 71 HBC K-P AT REST 6/71
R1 U ANDERSON RESULT NOT PUBLISHED, EVENTS ADDED TO DOYLE SAMPLE. 2/71
R1 AVG 0.6399 0.0049 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R1 STUDENT 0.6398 0.0055 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R1 FIT 0.6419 0.0049 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
R2 LAMBDA INTO (N P0)/((P PI-)+(N P0)) (P2)/(P1+P2)
R2 0.23 0.09 EISLER 57 HLBC
R2 0.43 0.14 CRAWFORD 59 HBC
R2 0.28 0.08 BAGLIN 60 HLBC
R2 0.35 0.05 BROWN 63 HLBC
R2 75 0.291 0.034 CHRETIEN 63 HLBC
R2 AVG 0.304 0.025 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
R2 STUDENT 0.304 0.028 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R2 FIT 0.3581 0.0049 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)
R3 LAMBDA INTO (P E- NEU)/TOTAL (UNITS 10**-3) (P4)/(P1+P2)
R3 O 15 (2.0) (0.5) HUMPHREY 61 RUUE
R3 O 8 (2.9) (1.5) (1.2) AUBERT 62 FBC
R3 N 150 (0.82) (0.12) ELY 63 FBC K- AT REST
R3 N 102 (0.78) (0.12) (0.13) BAGLIN 64 FBC K- AT 1.45 GEV/C
R3 O 20 (1.51) (0.34) LIND 64 HBC
R3 N 143 (0.80) (0.08) MALONEY 69 HBC
R3 N 86 (0.78) (0.09) CANTNER 71 HBC K-P AT REST 10/69
R3 N 218 (0.88) (0.10) LINDQUIST 71 OSPK PI- P TO KO LAM 4/71
R3 N THESE VALUES HAVE BEEN CHANGED BY US INTO RATIOS TO PROTON PI-, 3/72
R3 N BECAUSE THAT IS THE DIRECTLY MEASURED QUANTITY. SEE R5 BELOW 3/72
R3 O LOW STATISTICS EXPERIMENTS. NOT AVERAGED 7/70

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Λ , Σ^+

Table with columns R4, R5, R6 and data for LAMBDA INTO (P MU- NEU)/TOTAL (UNITS 10**4) (P31)/(P1+P2). Includes sub-headers R4, R5, R6 and student averages.

Table with columns R4, R5, R6 and data for RELATED TEXT SECTION VI D AND APPENDIX III. Includes sub-headers A-, A+, A0 and student averages.

Table with columns R4, R5, R6 and data for PH I ANGLE (SIN(PHI)/COS(PHI)=BETA/GAMMA) (DEGREES). Includes sub-headers F-, F+, F0 and student averages.

WEIGHTED AVERAGE = -0.65B ± 0.054
ERRDR SCALED BY 1.2

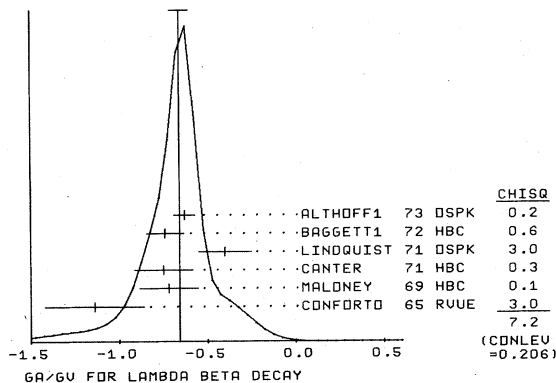


Table with columns R4, R5, R6 and data for REFERENCES FOR LAMBDA. Lists various researchers and their associated particle data, including EISLER, BAGLIN, BOWEN, CORK, COLUMBIA, HUMPHREY, ANDERSON, AUBERT, CHANG, COOL, GOOD, HUMPHREY, ALSTON, BHOWMIK, BLOCK, BROWN, CHRETIEN, CRONIN, ELY, KERNAN, ANDERSON, BAGLIN, HUBBARD, KERNAN, KREISLER, LIND, RONNE, SCHWARTZ, BAGLIN, BALTAJ, BARLOW, CHARRIERE, CONFORTO, ELY, HILL, SCHMIDT, BERGE, BURAN, CATTEN, ENGELMAN, GIBSON, LONDON, AUERBACH, BADIER, CLELAND, MAYEUR, OVERSETH, GRIMM, HEPP, MERRILL, DAUBER, DOYLE, MALONEY, BOHM, DEMIDOV, OLSEN, ALTHOFF1, ALTHOFF2, BALTAY, BARKOV, BARONI, CANTER, DAHLJENS, HILL, LINDQUIST, BAGGETT1, BAGGETT2, BAGGETT3, CLELAND, HYMAN, ALTHOFF1, ALTHOFF2, POULARD, ARMENTEROS, BALTAY, BERGE.

Table with columns R4, R5, R6 and data for CHISQ and CDNLEU. Includes a box with the symbol Σ^+ and the text '19 SIGMA+(1189,JP=1/2+ I=1)'. Below this, it says '19 SIGMA+ MASS (MEV)'. A note reads 'M N SEE NOTE PRECEDING LAMBDA MASS LISTINGS'. A table lists sources and their contributions to the CHISQ: M 144 1189.38 0.15 BARKAS 63 EMUL + SEE NOTE 5 BELOW, M 58 1189.48 0.22 BHOWMIK 64 EMUL + SEE NOTE 5 BELOW, M S ABOVE SIGMA+ MASSES HAVE BEEN RAISED 30 KEV TO ACCOUNT FOR 46 KEV, M S INCREASE IN PROTON MASS AND 21 KEV DECREASE IN PION MASS, M 4205 1189.61 0.08 SCHMIDT 65 HBC SEE NOTE N 3/74, M 1189.16 0.12 HYMAN 67 HBC 6/68, M R 607 1189.33 0.04 89HM 72 EMUL 12/73, M B BOHM 72 UPDATED WITH PDG APR. 73 X-, PI- AND P10 MASSES. 12/73, M M AVG 1189.371 0.060 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8), M STUDENT 1189.354 0.041 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT, M FIT 1189.366 0.057 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8) (SEE IDEOGRAM BELOW) 2/76*

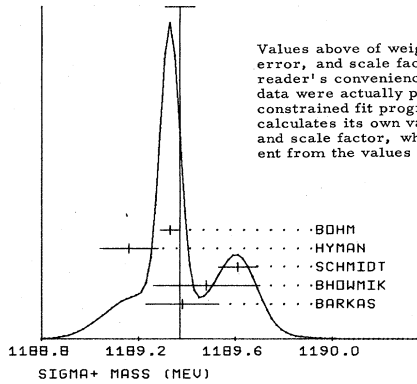
Stable Particles

Σ^+

Data Card Listings

For notation, see key at front of Listings.

WEIGHTED AVERAGE = 1189.371 ± 0.060
ERROR SCALED BY 1.8



Values above of weighted average, error, and scale factor are for the reader's convenience only. The data were actually processed by a constrained fit program, which calculates its own values of \bar{x} , $\bar{\sigma}$, and scale factor, which are different from the values shown here.

Table with 3 columns: Name, Particle Type, CHISQ. Includes entries for BDHM (72 EMUL, 1.0), HYMAN (67 HEBC, 3.1), SCHMIDT (65 HBC, 8.9), BHDWMIK (64 EMUL, 0.2), BARKAS (63 EMUL, 0.0), and a total of 13.3 (CDNLEV = 0.010).

19 SIGMA+ MEAN LIFE (UNITS 10** -10)

Table listing mean life measurements for Sigma+ particles. Columns include name, value, error, and source. Includes entries for GLASER, PUSCHEL, EVANS, FREDEN, KAPLON, CHIESA, BERTHELOT, BARKAS, CARAYAN, HUMPHREY, BHDWMIK, BARLOUTAU, EISELE, and BAKKER.

19 SIGMA+ MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

Table listing magnetic moment measurements for Sigma+ particles. Columns include name, value, error, and source. Includes entries for COOK, KOTELCHUC, SULLIVAN, COMBE, MAST, ALLEY, SAHA, and student averages.

19 SIGMA+ PARTIAL DECAY MODES

Table listing partial decay modes for Sigma+ particles. Columns include mode name and branching ratio. Includes entries for PROTON P10, NEUTRON P1+, LAMBDA E+ NEU, and NEUTRON MU+ NEUTRINO.

19 SIGMA+ BRANCHING RATIOS

Table listing branching ratios for Sigma+ particles. Columns include mode name, value, error, and source. Includes entries for NEUTRON P1+, LAMBDA E+ NEU, and NEUTRON MU+ NEUTRINO.

Table listing Sigma+ data for various experiments. Columns include name, value, error, and source. Includes entries for CARRARA, BAZIN, QUARENI, ANG, and GERSHWIN.

WEIGHTED AVERAGE = 0.240 ± 0.035
ERROR SCALED BY 1.4

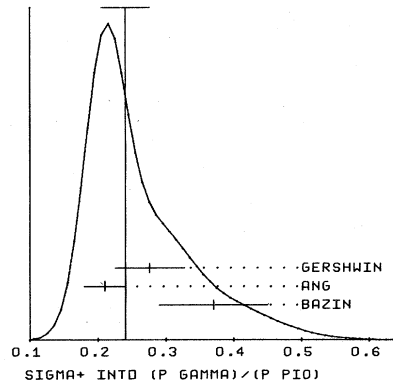


Table with 3 columns: Name, Particle Type, CHISQ. Includes entries for GERSHWIN (69 HBC, 0.5), ANG (69 HBC, 1.0), and BAZIN (65 HBC, 2.6), with a total of 4.1 (CDNLEV = 0.126).

Table listing Sigma+ data for various experiments. Columns include name, value, error, and source. Includes entries for COURANT, MURPHY, NAUENBERG, BIERMAN, EISELEZ, and NORTON.

Table listing Sigma+ data for various experiments. Columns include name, value, error, and source. Includes entries for GALTIERI, COURANT, NAUENBERG, BAGGETT, and NORTON.

Table listing Sigma+ data for various experiments. Columns include name, value, error, and source. Includes entries for BAGGETT, NORTON, ANG, and EBENHOH.

19 SIGMA+ DECAY PARAMETERS

Table listing decay parameters for Sigma+ particles. Columns include parameter name, value, error, and source. Includes entries for ALPHA+ALPHA0, CORK, TRIPP, BANGERTER, BERLEY, REUCROFT, and BERLEY.

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Σ^+ , Σ^-

ALPHA SIGMA0 (SIG+ INTO P10 PROTON)
AD 0.80 0.16 BEALL 62 CNTR
AD (-0.90) (0.25) TRIPP 62 HBC REPLAC. BY BANGE
AD O 5200 (-0.986) (0.072) BANGERTER 66 HBC K-P TO SIG+ PI- 7/66
AD 32000 -0.999 0.022 BANGERTER 69 HBC 10/69
AD H 1335 -0.98 0.05 0.02 HARRIS 70 OSPK 5/70
AD 16K -0.940 0.045 BELLAMY 72 ASPK PI+P TO SIG+ K+ 11/72
AD L 1259 -0.945 0.055 0.042 LIPMAN 73 OSPK PI+P TO SIG+ 7/73
AD L DECAY PROTONS SCATTERED OFF ALUMINUM.
AD H DECAY PROTONS SCATTERED OFF CARBON.
AD AVG -0.979 0.016 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
AD STUDENT -0.979 0.018 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

PHI+ ANGLE (SIG+ INTO N PI) SIN(PHI)/COS(PHI)=BETA/GAMMA (DEGREE)
F+ O 370 (180.) (30.) BERLEY 66 HBC + NEUTRON RESCATT. 9/66
F+ 560 143. 29. BANGERTI 69 HBC 10/69
F+ C105% 184. 24. BERLEY 70 HBC K-P AT 400 MEV/C 11/69
F+ C CHANGED FROM 176 TO 184 TO AGREE WITH SIGN CONVENTION.
F+ AVG 167.3 20.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)
F+ STUDENT 167.5 21.2 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

ALPHA SIGMA0 (SIG+ INTO PROTON GAMMA)
AG 61 -1.03 0.52 0.42 GERSHWIN 69 HBC K-P TO SIG PI 11/69
PHI0 ANGLE (SIG+ INTO P10 PROTON) SIN(PHI)/COS(PHI)=BETA/GAMMA (DEG)
FO H 32.0 90.0 HARRIS 70 OSPK PI+P TO SIG+ K+ 5/70
FO L 1259 38.1 35.7 37.1 LIPMAN 73 OSPK PI+P TO SIG+ K+ 7/73
FO L DECAY PROTON SCATTERED OFF ALUMINUM.
FO H DECAY PROTONS SCATTERED OFF CARBON.
FO AVG 35.8 33.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
FO STUDENT 35.8 36.3 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

REFERENCES FOR SIGMA+

CORK 60 PR 120 1000 CORK, KERTH, WENZEL, CRONIN, COOL (LRL+PRIN+BNL)
EVANS 60 NC 15 873 BRIST+BRUSS+IAS-UJ, COL-DUBL IN+LON+MILAN+PAD
FREDEN 60 NC 16 611 S FREDEN+H KORNBUM, R WHITE (LRL)
KAPLON 60 ANP 9 139 M KAPLON, A MELISSINOS, YAMANOUCHI (ROCH)
PUSCHEL 60 NP 20 254 W PUSCHEL (MAX PLANCK INST)
BARKAS 61 PR 124 1209 BARKAS, DYER, HASON, NICHOLS, SMITH (LRL)
BERTHELO 61 NC 21 693 BERTHELOT, DAUDIN, GOUSSU (SACLAY+ORSAY)
CHIESA 61 NC 19 1171 CHIESA, QUASSIATI, RINAUDO (INFN-TURIN)
BEALL 62 PRL 8 75 BEALL, CORK, KEEFE, MURPHY, WENZEL (LRL)
GRARD 62 PR 127 407 F GRARD, G A SMITH (LRL)
GALTIERI 62 PRL 9 26 GALTIERI, BARKAS, HECKMAN, PATRICK, SMITH (LRL)
HUMPHREY 62 PR 127 1305 W E HUMPHREY, R R ROSS (LRL)
TRIPP 62 PRL 9 66 R D TRIPP, M B WATSON, M FERRO-LUZZI (LRL)
BARKAS 63 PRL 11 26 W H BARKAS, J N DYER, H H HECKMANN (LRL)
ALSO 61 UCRL 9450 JOHN DYER (THEISIS, BERKELEY)
BHOWMIK 64 NP 53 22 B BHOWMIK, P JAIN, P MATHUR, LAKSHMI (DELHI)
CARRARA 64 PL 12 72 CARRARA, CRESTI, GRIGOLETTO, PERUZZO (PADOVA)
COURANT 64 PR 136 8 1791 COURANT, FILTHUTH+ (CERN+HEID+UMD+NRL+BNL)
MURPHY 64 PR 134 8 188 C THORNTON MURPHY (WISCONSIN)
NAUENBERG 64 PRL 12 679 NAUENBERG, MARATECK+ (COLU+RUTG+PRIN)
WILLIS 64 PRL 13 291 WILLIS+; COURANT, ENGLISHAN+ (BNL, CERN, HEID, UMD)
BALTAI 65 PR 140 8 1027 BALTAI, SANDHEISS, CULWICK, KOPP + (YALE+BNL)
BAZIN 65 PRL 14 154 BAZIN, BLUMENFELD, NAUENBERG + (PRIN+COLU)
BAZIN 65 PRL 14 154 BAZIN, BLAND, SCHMIDT+ (PRIN, RUTG, COLU)
CARAYAN 65 PR 138 8 433 CARAYAN+POPOULOS, TAUFEST, WILLMAN (PURDUE)
QUARENI 65 NC 40 4 928 QUARENI, CARTACCI + (BGN, FIRZ, GENG, PARMA)
SCHMIDT 65 PR 140 8 1328 P SCHMIDT (COLUMBIA)
BANGERTER 66 PRL 17 495 BANGERTER, GALTIERI, BERGE, MURRAY+ (LRL)
BERLEY 66 PRL 17 1071 +HERZBACH, KOFLER, YAMAMOTO + (BNL+MAS+YALE)
CHANG 66 PR 151 1081 CHUNG YUN CHANG (COLUMBIA)
ALSO 65 NEVIS 145 THESIS CHUNG YUN CHANG (COLUMBIA)
CHIEN 66 PR 152 1171 +LAI, SANDHEISS, TAIT, YEH, OREN + (YALE+BNL)
COOK 66 PRL 17 223 V COOK, EWART, MASEK, ORR, PLATNER (WASHINGTON)
BAGGETT 67 PRL 19 1458 BAGGETT, DA V, GLASSER, KEHOE, KNOP+ (MARYLAND)
ALSO 68 VIENNA ABS. 374 BAGGETT, KEHOE (MARYLAND)
ALSO 68 PRIVATE COMM. N. BAGGETT
BARASH 67 PRL 19 181 BARASH, DAY, GLASSER, KEHOE, KNOP + (MARYLAND)
EISELE 67 ZPHYS 205 409 +ENGELMANN, FILTHUTH, FOHLISCH, HEPP+ (HEID)
HYMAN 67 PL 25 8 376 +LOKEN, PEWITT, MCKENZIE, + (ANL+CARM+MWS)
KOTELCHU 67 PRL 18 1166 KOTELCHUCK, GOZA, SULLIVAN, ROSS (VANDERBILT)
SULLIVAN 67 PRL 18 1163 SULLIVAN, MCINTURFF, KOTELCHUCH (VANDERBILT)
ALSO 64 PRL 13 246 A D MCINTURFF, C E ROOS (VANDERBILT)
BIERMAN 68 PRL 20 1459 BIERMAN, KOUNOSU, NAUENBERG + (PRINCETON)
COMBE 68 NC 57A 54 CERN-BRISTOL-LAUSANNE-MUNICH-ROME-COLLABOR
MAST 68 PRL 20 1312 MAST, GERSHWIN, ALSTON-GARNJUST + (LRL)
+EBENHOH, EISELE, ENGELMANN, FILTHUTH+ (HEID)
N V BAGGETT (THEISIS) (UMD)
BALTAI, FRANZINI, NEWMAN, NORTON+ (COLU, STON)
ROGER ODELL BANGERTER (THEISIS) (LRL)
BANGERTI 69 PR 187 821 BANGERTER, GARNJUST, GALTIERI, GERSHWIN+ (LRL)
BARLOUTA 69 NP 814 153 BARLOUTA, BELLEFON, GRANET+ (SACL+GERN+HEID)
EISELE 69 ZPHYS 221 1 +ENGELMANN, FILTHUTH, FOHLISCH, HEPP+ (HEID)
EISELE 69 ZPHYS 221 401 +ENGELMANN, FILTHUTH, FOHLISCH, HEPP+ (HEID)
GERSHWIN 69 PR 188 2077 +ALSTON-GARNJUST, BANGERTER + (LRL)
ALSO UCRL 19246 THESIS LAWRENCE K GERSHWIN
NORTON 69 NEVIS 175 (THEISIS) HERBERT NORTON (COLUMBIA)
BERLEY 70 PR D1 2015 +YAMIN, HERTZBACH, KOFLER + (BNL, MAS, YALE)
EISELE 70 ZPHY 238 372 +FILTHUTH, HERTZBACH, PRESER, ZIEH (HEIDELBERG)
HARRIS 70 PRL 24 165 +OVERSETH, PONDROM, DETTMANN (MICH, WISC)
ALLEY 71 PR 03 75 +BENBROOK, COOK, GLASS, GREEN, HAGUE + (WASH)
BAKKE 71 LNC 1 37 +SABRE COLLAB. (ZEEM+SACL+BGNA+REHO+EPDL)
COLE 71 PR D4 431 +LEE-FRANZINI, LOVELESS, BALTAI+ (STON, COLU)
TOVEE 71 NP 833 493 LOUC, BELGRADE, BERL, BRUX, DOUBLIN, WARS COLLAB
BELLAMY 72 PL 398 299 +ANDERSON, CRAWFORD, OSMON+ (UMD)
BOHM 72 NP 848 1 BERLIN+BELGRADE+BRUX+DUBL IN+LUC+WARSAW
ALSO 73 TIME-73.2 NOV BRUSSELS BULLETIN, SAME COLLABORATION
EBENHOH 73 ZPHY 264 413 +EISELE, FILTHUTH, HEPP, LEITNER, THOU+ (HEID)
LIPMAN 73 PL 43B 89 +UTO, WALKER, MONTGOMERY+ (RHEL+SUSS+LOWC)
SAHA 73 PR D7 3295 +FEYKOWICH, HEINTZELMAN, MELTZER + (GARM)
SCHIZOR 73 PR D8 12 B. SECHI-ZORN, G. SNOW (UMD)
EBENHOH 74 ZPHY 266 367 +EISELE, ENGELMANN, FILTHUTH, HEPP + (HEID)
REUCROFT 76 PREPRINT +ROOS, WATERS, WEBSTER, HANSL+ (VAND, MPIH)

PAPERS NOT REFERRED TO IN DATA CARDS
GLASER 58 CERN CONF 270 GLASER, GOOD, MORRISON (MICH+LFL)
QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS
TRIPP 62 PRL 8 175 R TRIPP, M WATSON, M FERRO-LUZZI (LRL) P
ALFF 63 SIENA CONF 1 205 ALFF, NAUENBERG, KIRSCH, + (COLU+RUTG+BNL)
ALSO 65 PR 137 8 1105 ALFF, GELFAND, BRUGGER, BERLEY+ (COLU+RUTG+BNL)
COURANT 63 SIENA CONF 1 73 COURANT, FILTHUTH, BURNSTEIN, DAY+ (CERN+UMD)

20 SIGMA- (1198, JP=1/2+) I=1

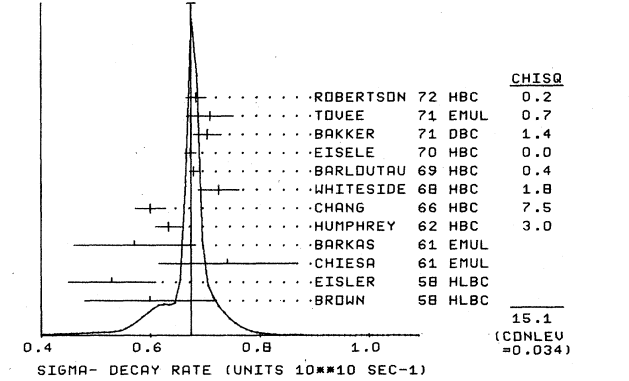
20 SIGMA- MASS (MEV)
M N SEE NOTE PRECEDING LAMBDA MASS LISTINGS
M 3000 1197.43 0.08 SCHMIDT 65 HBC SEE NOTE N 3/74
M FIT 1197.35 0.06 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*

20 (SIGMA-) - (SIGMA+) MASS DIFFERENCE (MEV)
D 87 8.25 0.40 BARKAS 63 EMUL -
D 2500 8.25 0.25 DOSCH 65 HBC
D 86 7.91 0.23 BOHM 72 EMUL 1/73
D AVG 8.09 0.16 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
D STUDENT 8.10 0.18 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
D FIT 7.98 0.38 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2) 2/76*

20 (SIGMA-) - (LAMBDA) MASS DIFFERENCE (MEV)
DL N SEE NOTE PRECEDING LAMBDA MASS LISTINGS.
DL 85 81.70 0.19 BURNSTEIN 64 HBC 9/66
DL 2279 81.80 0.13 SCHMIDT 65 HBC SEE NOTE N 3/74
DL 81.64 0.09 HEPP 68 HBC 8/68
DL AVG 81.693 0.069 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
DL STUDENT 81.692 0.077 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
DL FIT 81.750 0.054 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0) 2/76*

20 SIGMA- MEAN LIFE (UNITS 10**--10)
T 1.67 0.40 0.28 BROWN 58 HLBC
T 1.89 0.33 0.25 EISLER 58 HLBC
T 45 1.35 0.32 0.17 CHIESA 61 EMUL
T 41 1.75 0.39 0.30 BARKAS 61 EMUL
T 1208 1.58 0.06 0.06 HUMPHREY 62 HBC STOP. K- 6/66
T C 3267 1.666 0.075 CHANG 66 HBC STOP. K- 9/67
T S 61 (2.08) (0.22) CHIEN 66 HBC + 6.9 PBAR P, ANTI 9/67
T S 64 (1.46) (0.31) CHIEN 66 HBC + 6.9 PBAR P, ANTI 9/67
T 506 1.38 0.07 WHITESIDE 68 HBC STOP. K- 6/68
T 10253 1.472 0.016 BARLOUTA 69 HBC K-P 1.4-1.2 GEV/C 11/69
T 1M 1.485 0.022 EISELE 70 HBC K-P AT REST 2/71
T 1383 1.42 0.05 BAKKER 71 DBC -K-N TO SIG- 2PI 10/71
T 1.41 0.09 TOVEE 71 EMUL 12/71
T 2400 1.463 0.039 ROBERTSON 72 HBC K-P +25 GEV/C 3/74
T C CHANG ERROR 0.018 RAISED BY US. SEE 1970 EDITION, RMP 42, 1231 (1970) 1/73
T S ERROR PURELY STATISTICAL.
T AVG 1.482 0.017 0.017 AVERAGE (ERROR INCL. SCALE FACTOR OF 1.5)
T STUDENT 1.479 0.013 0.013 AVG. USING STUDENT10(H/1.11) -- SEE TEXT
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 0.6746 ± 0.0076
ERRDR SCALED BY 1.5



20 SIGMA- MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)
MM R, BTWN -1.6 AND +0.8 FIX 73 CNTR SIG-ATOM FINE ST 3/74
MM R -1.48 0.37 ROBERTS 74 CNTR SIG-ATOM FINE ST 12/75*
MM R ROBERTS 74 INCLUDES DATA FROM FOX 73. 12/75*

Stable Particles

Σ^-, Σ^0

Data Card Listings

For notation, see key at front of Listings.

20 SIGMA- PARTIAL DECAY MODES
P1 SIGMA- INTO NEUTRON PI-
P2 SIGMA- INTO NEUTRON PI- GAMMA
P3 SIGMA- INTO NEUTRON MU- NEUTRINO
P4 SIGMA- INTO NEUTRON E- NEUTRINO
P5 SIGMA- INTO LAMBDA E- NEUTRINO

20 SIGMA- BRANCHING RATIOS
R1 SIGMA- INTO (N MU- NEU)/(N PI-) (UNITS 10**=-3) (P3)/(P1)
R2 SIGMA- INTO (N E- NEU)/(N PI-) (UNITS 10**=-3) (P4)/(P1)
R3 SIGMA- INTO (LAMBDA E- NEU)/(N PI-) (UNITS 10**=-4) (P5)/(P1)

20 SIGMA- BRANCHING RATIOS (continued)
R2 SIGMA- INTO (N E- NEU)/(N PI-) (UNITS 10**=-3) (P4)/(P1)
R3 SIGMA- INTO (LAMBDA E- NEU)/(N PI-) (UNITS 10**=-4) (P5)/(P1)

20 SIGMA- BRANCHING RATIOS (continued)
R3 SIGMA- INTO (LAMBDA E- NEU)/(N PI-) (UNITS 10**=-4) (P5)/(P1)

20 SIGMA- BRANCHING RATIOS (continued)
R4 SIGMA- INTO (N PI- GAMMA)/(N PI-) (UNITS 10**=-3) (P2)/(P1)
R5 SIGMA- INTO (N PI- GAMMA)/(N PI-) (UNITS 10**=-3) (P2)/(P1)

20 SIGMA- DECAY PARAMETERS
RELATED TEXT SECTION VI D AND APPENDIX III
A- ALPHA SIGMA-
A- (-0.16) (0.21) TRIPP 62 HBC REPL BY BANGERTE

20 SIGMA- DECAY PARAMETERS (continued)
F- PHI ANGLE (SIN(PHI)/COS(PHI))=BETA/GAMMA (DEGREES)
F- 0 1006 (+22.1) (30.1) BERLEY 67 HBC K-P TO SIG- PI+ 11/67

20 SIGMA- DECAY PARAMETERS (continued)
AV GV/GV FOR SIGMA TO LAMBDA BETA DECAY (TEXT SEC VI D.1 FOR SIGN CONV)
AV PREDICTED TO BE ZERO BY CONSERVED VECTOR CURRENT THEORY 11/67

20 SIGMA- DECAY PARAMETERS (continued)
AVI GA/GV FOR SIGMA TO NEUTRON BETA DECAY (TEXT SEC VI D.1 FOR SIGN CONV)
AVI 57 (0.051) (0.23) (0.32) GERSHWIN 68 HBC REPLACED BY GER.69 6/68

REFERENCES FOR SIGMA-
BROWN 58 CERN CONF 270 (MICH)
EISLER 58 NC SER10 10 150 EISLER, BASSI, CONVERS I* (COLU, BNL, BGN, PISA)

BARKAS 61 PR 124 1209
CHIESA 61 NC 19 1171
HUMPHREY 62 PR 127 1305
TRIPP 62 PRL 9 66

BARKAS 63 PRL 11 26
BURNSTEIN 64 PRL 13 66
COURANT 64 PR 136 B 1791
MILLER 64 PR 11 262

BARASH 67 PRL 19 181
BIERMAN 68 PR 20 1459
GERSHWIN 68 PR 20 1270
HEPP 68 ZPHY 214 71

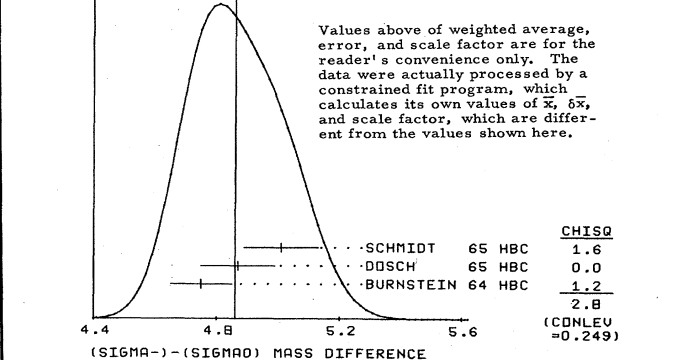
BERLEY 70 PR D1 2015
BOGERT 70 PR D2 6
EBENHOH 70 KIEV CONF
EISELE 70 ZPHY 238 372

BAKKER 71 LNC 1 37
COLE 71 PR D4 631
TOVEE 71 NP 839 493
BALTHAY 72 PR D5 1569

BROWN 57 PR 108 1036
NIETO 68 RMP 40 140
J BROWN, D GLASER, M PERL (MICH+BNL)
M NIETO (STON)

21 (SIGMA-) - (SIGMA0) MASS DIFFERENCE (MEV)
D1 N SEE NOTE PRECEDING LAMBDA MASS LISTINGS.
D1 18 4.75 0.1 BURNSTEIN 64 HBC
D1 37 4.87 0.12 DOSCH 65 HBC
D1 12 5.01 0.12 SCHMIDT 65 HBC SEE NOTE N 3/74

WEIGHTED AVERAGE = 4.860 ± 0.076
ERROR SCALED BY 1.2



Data Card Listings

Stable Particles

For notation, see key at front of Listings.

Σ^0, Ξ^-

21 (SIGMA) - (LAMBDA) MASS DIFFERENCE (MEV)

DL N SEE NOTE PRECEDING LAMBDA MASS LISTINGS.

Table with columns DL, N, and values for SCHMIDT, COLAS, DL AVG, DL STUDENT, DL FIT.

21 SIGMAO MEAN LIFE (UNITS 10**(-14))

T (1.0) OR LESS DAVIS 62 EMUL

21 SIGMAO PARTIAL DECAY MODES

Table with columns P1, P2, P3, SIGMAO INTO LAMBDA GAMMA, DECAY MASSES.

21 SIGMAO BRANCHING RATIOS

Table with columns R1, R2, SIGMAO INTO (LAMBDA E-)/TOTAL, QUANTUM ELECT., SIGMAO INTO (LAMBDA GAMMA GAMMA)/(LAMBDA GAMMA), CL=90.

REFERENCES FOR SIGMAO

Table listing references for SIGMAO, including FEINBERG, DAVIS, BURNSTEIN, DOSCH, SCHMIDT, COLAS.

COURANT 63 PRL 10 409 COURANT, FILTHUTH, FRANZINI+ (CERN+UMD+NRL)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA CARDS

ALFF 65 PR 137 B1105 ALFF, GELFAND, NAUENBERG+ (COLUMBIA+RUTG+MILP)

22 XI-(1321, JP=1/2) I=1/2

22 XI- MASS (MEV)

Table with columns M, H, WANG, FOWLER, SCHNEIDER, BAUFER, CHIEN, LONDON, GOLDHASSE, WILQUET, DL AVG, DL STUDENT, DL FIT.

22 ANTI-XI+ MASS (MEV)

Table with columns M1, M1 S, M1, M1 S, BROWN, CHIEN, SHEN, STONE, VOTRUBA, DL AVG, DL STUDENT, DL FIT.

22 (XI-) - (ANTI-XI+) MASS DIFFERENCE (MEV)

DM 1.0 1.1 CHIEN 66 HBC 6.9 PBAR P 9/67

22 XI- MAGNETIC MOMENT (MAGNETONS, 938.26 MEV)

Table with columns MM, MM STUDENT, BINGHAM, COOL, DL AVG, DL STUDENT.

22 XI- MEAN LIFE (UNITS 10**(-10))

Table with columns T, H, WANG, FOWLER, SCHNEIDER, CARMONY, HUBBARD, PUJERROU, CHIEN, LONDON, DAUBER, MAYEUR, BALTAY, COOL, DL AVG, DL STUDENT.

22 ANTI-XI+ MEAN LIFE (UNITS 10**(-10))

Table with columns T1, S, CHIEN, SHEN, STONE, VOTRUBA.

22 XI- PARTIAL DECAY MODES

Table with columns P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13, XI- INTO LAMBDA PI-, XI- INTO LAMBDA E- NEUTRINO, XI- INTO NEUTRON PI-, XI- INTO SIGMAO MU- NEUTRINO, XI- INTO SIGMAO MU+ NEUTRINO, XI- INTO SIGMAO E- NEUTRINO, XI- INTO SIGMAO MU- NEUTRINO, XI- INTO SIGMAO MU+ NEUTRINO, XI- INTO PROTON PI- PI-, XI- INTO PROTON PI- E- NEUTRINO, XI- INTO PROTON PI- MU- NEUTRINO, XI- INTO XI0 E- NEUTRINO.

22 XI- BRANCHING RATIOS

Table with columns R1, R2, R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, XI- INTO (LAMBDA E- NEU)/(LAMBDA PI-) (UNITS 10**(-3)), XI- INTO (NEUTRON PI-)/(LAMBDA PI-) (UNITS 10**(-3)), XI- INTO (LAMBDA MU- NEUTRINO)/TOTAL (UNITS 10**(-3)), XI- INTO (SIGMAO E- NEUTRINO)/TOTAL (UNITS 10**(-3)), XI- INTO (SIGMAO MU- NEU)/(LAMBDA PI-) (UNITS 10**(-3)), XI- INTO (SIGMAO MU+ NEU)/(LAMBDA PI-) (UNITS 10**(-3)), XI- INTO (SIGMAO E- NEU + LAMBDA E- NEU)/TOTAL (UNITS 10**(-3)), XI- INTO (N E- NEU)/(LAMBDA PI-) (UNITS 10**(-3)), XI- INTO (SIGMAO E- NEU + LAMBDA E- NEU)/(LAMBDA PI-) (UNITS 10**(-4)), XI- INTO (P PI- PI-)/(LAMBDA PI-) (UNITS 10**(-4)), XI- INTO (P PI- E- NEU)/(LAMBDA PI-) (UNITS 10**(-4)), XI- INTO (P PI- MU- NEU)/(LAMBDA PI-) (UNITS 10**(-4)), XI- INTO (XI0 E- NEU)/(LAMBDA PI-) (UNITS 10**(-3)).

Stable Particles

Ξ^-, Ξ^0

Data Card Listings

For notation, see key at front of Listings.

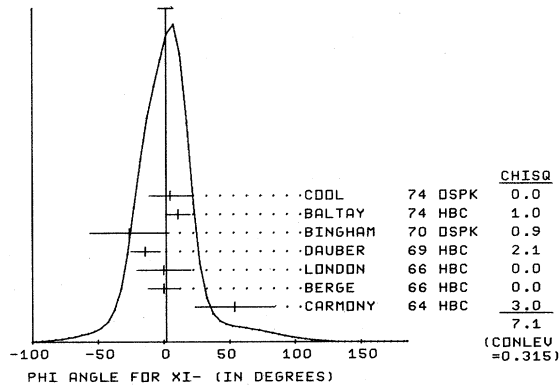
22 XI- DECAY PARAMETERS

RELATED TEXT SECTION VI D AND APPENDIX III

Table with columns for particle type (A, L, M, D), parameters (ALPHA XI-, etc.), values, and references (JAUNEAU, SCHNEIDER, etc.).

Table with columns for particle type (F, D, L, M, D), parameters (PHI ANGLE, etc.), values, and references (JAUNEAU, SCHNEIDER, etc.).

WEIGHTED AVERAGE = 2.0 ± 5.7
ERROR SCALED BY 1.1



REFERENCES FOR XI- table listing various researchers and their publications, such as FOWLER 61 PRL 6 134, WANG 61 JETP 13 512, etc.

Table with columns for particle type (U, V, W, Y), parameters (NP, etc.), values, and references (FREYTAG, HEINTZE, HEINZELMAN, etc.).

Table 23 XIO MASS (MEV) with columns for particle type (M), mass values, and references (PALMER, HILQUET, etc.).

Table 23 (XI-) - (XIO) MASS DIFFERENCE (MEV) with columns for particle type (D), mass differences, and references (JAUNEAU, CARMONY, etc.).

Table 23 XIO MEAN LIFE (UNITS 10**=-10) with columns for particle type (T), mean life values, and references (JAUNEAU, CARMONY, etc.).

Table 23 XIO PARTIAL DECAY MODES with columns for particle type (P), decay modes, and decay masses.

Table 23 XIO BRANCHING RATIOS with columns for particle type (R), branching ratios, and references (TICHO, HUBBARD, etc.).

Table 23 XIO INTO (SIGMA+ E- NEU)/(LAMBDA P) (UNITS 10**=-3) with columns for particle type (R), ratios, and references (TICHO, HUBBARD, etc.).

Table 23 XIO INTO (SIGMA+ E- NEU)/(LAMBDA P) (UNITS 10**=-3) with columns for particle type (R), ratios, and references (TICHO, HUBBARD, etc.).

Data Card Listings

For notation, see key at front of Listings.

Stable Particles

Ξ^0, Ω^-

23 XIO DECAY PARAMETER

RELATED TEXT SECTION VI D AND APPENDIX III

Table with columns for particle name, alpha value, and source. Includes entries for ALPHA XI 0, PIERROU, BERGE, LONDON, MERRILL, DAUBER, BRIDGEWATER, MAYEUR, BALTAY.

WEIGHTED AVERAGE = -0.441 ± 0.078
ERROR SCALED BY 1.3

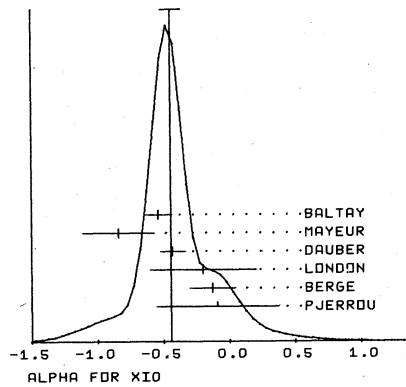


Table with columns for source name, number of events, and chi-squared value. Includes BALTAY, MAYEUR, DAUBER, LONDON, BERGE, PIERROU.

Table with columns for phi angle, sin(phi)/cos(phi) = beta/gamma, and degrees. Includes entries for BERGE, MERRILL, DAUBER, BRIDGEWATER, BALTAY.

REFERENCES FOR XIO

Table listing references for XIO, including authors like ALVAREZ, JAUNEAU, TICHOU, CARMONY, HUBBARD, PJERROU, BERGE, PALMER, DAUBER, BRIDGEWATER, MAYEUR, WILQUET, BALTAY, YEH, GEMENIGE.



24 OMEGA-(1675, JP=3/2+) I=0
QUANTUM NUMBERS ASSIGNED FROM SUB

Table with columns for particle name, mass (MEV), and source. Includes entries for EISENBERG, FRY1, FRY2, ABRAMS, PALMER, SCHULTZ, SCOTTER, SPETH, ABCLV.

24 ANTI-OMEGA+ MASS (MEV)

Table with columns for particle name, mass (MEV), and source. Includes FIRESTONE 71 HBC.

24 OMEGA- MEAN LIFE (UNITS 10** -10 SEC)

Table with columns for particle name, mean life, and source. Includes ABRAMS, BARNES, COLLEY, RICHARDSON, ABCLV, SCHULTZ, SCOTTER.

24 OMEGA- PARTIAL DECAY MODES

Table with columns for decay mode and source. Includes OMEGA- INTO LAMBDA K-, OMEGA- INTO XIO PI-, OMEGA- INTO XI- P10.

24 OMEGA- BRANCHING RATIOS

Table with columns for branching ratio, particle name, and source. Includes OMEGA- INTO LAMBDA K-, OMEGA- INTO XIO PI-, OMEGA- INTO XI- P10.

24 OMEGA- DECAY PARAMETERS

Table with columns for decay parameter, particle name, and source. Includes ALPHA FOR OMEGA- TO K- LAMBDA.

REFERENCES FOR OMEGA-

Table listing references for Omega minus, including authors like EISENBERG, FRY1, FRY2, ABRAMS, BARNES, COLLEY, RICHARDSON, SAMIOS.

Stable Particles

 Ω^- , PARTICLE SEARCHES

ABCLV CO 68 NUC PHYS 84 326	AACHEN+BERLIN+CERN+LONDON IMP. COLL.+VIENNA
ALLISON 68 PRIV. COMM.	JOHN ALLISON (LANCASTER)
PALMER 68 PL 26 B 323	PALMER, RADDJICIC, RAU, RICHARDSON+ (BNL, SYRA)
SCHULTZ 68 PR 168 1509	SCHULTZ+ (ILL, ARGONNE, NORTHWESTERN, WISC)
SCOTTER 68 PL 26 B 474	SCOTTER+ (BIRM, GLASGOW, LOIC, MUNICH, OXF)
SPETH 69 PL 29 B 252	SPETH+ (AACHEN, BERLIN, CERN, LOIC, VIEN)
FRESTON 71 PRL 26 410	*GOLDHABER, LISSAUER, SHELTON, TRILLING (LBL)
ABCLV 73 NP B61 102	AACHEN+BERLIN+CERN+LONDON+VIENNA COLLABOR.
ALVAREZ 73 PR D8 702	LUIS W. ALVAREZ (LBL)
KOCHER 74 PL 51 B 193	KOCHER, WERNHARD (INNS+VIEN)

PARTICLE SEARCHES

Heavy Leptons

Two types of searches for heavy leptons have been made. The first type uses the fact that we can calculate the transition rate for the decay of a heavy lepton to states involving other leptons, with or without pions. Production mechanisms are also considered to be well known, proceeding via the electromagnetic interaction. BERNARDINI 73 and BARISH 73 are of this type. The second method (BUSHNIN 73) is to look for charged, weakly interacting, long-lived particles.

Several experiments have observed events which could arise from the decay of a heavy lepton. PERL 75 in a SLAC-LBL experiment have observed 64 events of the form

$$e^+ + e^- \rightarrow e^\pm + \mu^\mp + \geq 2 \text{ undetected particles}$$

at or above 4 GeV c.m. energy. Such events have no conventional explanation, but could arise from the production and decay of heavy lepton pairs or charmed boson pairs.

Dimuon events such as those of the Harvard, Pennsylvania, Wisconsin, Fermilab experiment (see BENVENUTI1-5 75 entries in the charmed hadron searches section) could involve heavy lepton production and decay. However, the characteristics of their events do not favor the heavy lepton interpretation.

The Kolar Gold Mines cosmic ray experiment (KRISHNASWAMY 75) reported 5 events which typically have three observed tracks emerging from a common origin about 70 cm from the rock walls of the mine. The authors suggest that these may be massive, long-lived particles, possibly heavy leptons produced by neutrino interactions in the rock walls. Another hypothesis (see DE RUJULA 75) which was put forth to account for their apparent copious production but slow decay is that the observed particle is a

Data Card Listings

For notation, see key at front of Listings.

very long-lived neutral heavy lepton which is a decay product of a charged heavy lepton, itself pair-produced by cosmic rays in the upper atmosphere.

An FNAL neutrino experiment (BENVENUTI 75) roughly simulated the Kolar Gold Mines experiment for the most likely case that the new particle is neutral. Their result appears to contradict (at least fails to confirm) the KRISHNASWAMY 75 result.

Intermediate Bosons (W^{\pm} and ϕ^0)

Experimentally, W bosons could be produced by hadrons, muons, or neutrinos. They can decay into leptons or into hadrons. At present, there is no positive evidence for the existence of W bosons, although there have been several searches. We list upper limits on the mass and on production cross section.

All gauge theories postulate, in addition to the vector boson W, at least one scalar boson (Higgs' scalar), ϕ^0 . The couplings of the Higgs' scalar are model-dependent, and limits on its mass depend on M_W . We list these limits and add in comment cards the assumptions made.

Quarks

Three main techniques have been used in the experimental search for quarks: accelerators, cosmic rays, and searches in stable matter. The most recent limits using each of these techniques are given in the Data Card Listings. Accelerator experiments generally measure quark production cross sections (we quote these in Section C) and differential cross sections (D). Cosmic ray experiments measure quark flux (F), and searches in stable matter measure quark concentration (RHO). Some of the accelerator and cosmic ray experiments have looked for fractionally charged particles, and some have looked for high-mass, low-velocity particles. We give the fractional charge, mass ranges, velocity ranges, and other information in the notes on the right-hand side of the card and below the data cards.

Magnetic Monopoles

Searches have been made for magnetic monopoles using accelerators, cosmic rays, and searches in matter.

In accelerator searches strong magnetic fields

Stable Particles
PARTICLE SEARCHES

Data Card Listings

For notation, see key at front of Listings.

HEAVY LEPTON PRODUCTION CROSS SECTION (E+ E-) (UNITS 10**--35 CM**2)
PERL 75 SPEC M=1.6-2.0 GEV 2/76*
MORE MISSING PARTICLES. CS IS FOR EVENTS ARE E-GEV AND THETA=50-130DEG. 2/76*
CROSS SECTION RISES FROM 5X10**36 AT E=4 TO ABOVE MAX. THEN DROPS 2/76*
TO 6E-36 AT E=7.5. AUTHORS SAY THESE EVENTS HAVE NO CONVENTIONAL 2/76*
EXPLANATION. SUGGEST HEAVY LEPTON OR CHARMED HADRON, M=1.6-2.0GEV. 2/76*

QUARK SEARCHES

QUARK PRODUCTION CROSS SECT. FROM ACCELERATOR EXPTS (CM**2)
Y 0 3.2E-39 OR LESS CL=.90 ALLBAY 69 CNTR Q=-1/3 M=2GEV 1/76*
Y 0 1.0E-35 OR LESS CL=.90 ALLBAY 69 CNTR Q=+1/3 M=2GEV 1/76*
Y 0 1.0E-35 OR LESS CL=.90 ALLBAY 69 CNTR Q=+2/3 M=2GEV 1/76*
Y 0 4. E-37 OR LESS CL=.90 ANTIPOV1 69 CNTR Q=-2/3 M=0-5GEV 2/74
A 0 3. E-37 OR LESS CL=.90 ANTIPOV2 69 CNTR Q=-1/3 M=4.5-5GEV 1/76*
A 0 1. E-39 OR LESS CL=.90 ANTIPOV2 69 CNTR Q=+1/3 M=0-5GEV 1/76*
V 0 1. E-36 OR LESS CL=.90 ANTIPOV 71 CNTR Q=-4/3 M=4+6GEV 1/76*
B 0 3. E-34 OR LESS BOTB-BODE 72 CNTR Q=+1/3 M=0-22GEV 2/74
B 0 6. E-34 OR LESS BOTB-BODE 72 CNTR Q=+2/3 M=0-13GEV 2/74
P 0 1. E-32 OR LESS CL=.90 ALPER 73 SPEC Q= 2/3 M=4-24 GEV 1/76*
P 0 1. E-35 OR LESS LEIPUNER 73 CNTR Q= 1/3 M=0-12GEV 2/74
L 0 1. E-35 OR LESS LEIPUNER 73 CNTR Q= 2/3 M=0-12GEV 2/74
L 0 5. E-31 OR LESS LEIPUNER 73 CNTR Q= 4/3 M=0-12GEV 2/74
Y ALLBAY 69 IS A CERN 27 GEV P+BE EXPT. STUDYS MASSES 0-2.7GEV 1/76*
Y ASSUMING NN=NNQQ. CROSS SECTIONS ASSUME ISOTROPIC PROD. IN CM. 1/76*
Y CROSS SECTIONS AT 2GEV ARE GIVEN HERE. SEE FIG-9 FOR MASS DEPEN. 1/76*
ANTIPOV1 69 IS A SERPUKHOV 70 GEV P EXPT. MASS LIMIT FROM NN=NNQQ. 2/74
ANTIPOV1 69 AND ANTIPOV2 69 ARE SERPUKHOV 70GEV P EXPTS. ANTIPOV2 1/76*
GIVES RESULTS FOR M=2-5GEV ASSUMING NN--NNQQ, HADRONIC OR LEPTONIC 1/76*
QUARKS. WE QUOTE TYPICAL VALUES. 1/76*
ANTIPOV 71 IS A SERPUKHOV 70 GEV P+AL EXPT. STUDIES DIQUARK MASSES 1/76*
1.9-4.4GEV. WE SHOW 4GEV VALUES. SEE THEIR FIG-2 FOR MASS DEPEN. 1/76*
BOTT-BODENHAUSEN 72 IS A CERN ISR 26+26 GEV P+P EXPERIMENT. 2/74
ALPER 73 IS CERN ISR 26+26 GEV P+P EXPT. ASSUMES ISOTROPIC C.M. 1/76*
P PRODUCTION. SENSITIVE TO ANY Q2/3. 1/76*
LEIPUNER 73 IS AN AL 300 GEV P EXPERIMENT. 2/74

REFERENCES FOR HEAVY LEPTON SEARCHES
+GRACHEV,KHODYREV,KUBAROVSKY+ (SERP)
+PARTI,PEVED,SALVINI,STELLA+ (ROMANFRAS)
+BARTLETT,BUCHHOLZ,HUMPHREY+ (CIT+FNAL)
+BOLLINI,BRUNINI+ (CERN+BGNA+FRAS)
ALLES-BORELLI,BERNARDINI,BOLLINI+ (CERN)
DUNAYTZEY,GOLOVKIN,KUBAROVSKY+ (SERP)
GOLOVKIN,GRACHEV,SHODYREV+ (SERP)
+DEDEN+IAACH+BELG+CERN+EPOL+M1LA+LALO+LOUCI
+GERSHTEIN,KAPTANOV,KUBANTZEV,LAPINA (SERP)
+BARTLETT,BUCHHOLZ,HERRITT+ (CIT+FNAL)
+FRISCH,SIOCHET,BOYNDON,MERMOD+ (EP1+PRIN)
+VISENTIN,CERARDINI,CONVERSI+ (FRAS+ROMA)
BENVENUTI,CLINE,FORD+ (HARV+PENN+WISC)
BINT INGER, CURRY+ (EP1+HARV+PENN+WISC)
KRISHNASWAMY,HENDON+ (BOMBAY+OSAKA)
DE RUJULA,GEORGI,GLASHOW
+ABRAMS,BOYARSKI,BREIDENBACH+ (LBL+SLAC)

QUARK PROD. DIFFERENTIAL CROSS SEC, ACCELERATOR EXPTS(CM**2/SR-GEV)
D 0 1.5E-36 OR LESS DORFAN 65 CNTR BE TARG M=3-7GEV 2/74
D 0 3. E-36 OR LESS DORFAN 65 CNTR FE TARG M=3-7GEV 2/74
D 0 7.2E-39 OR LESS CL=.90 ALLBAY 69 CNTR Q=-1/3 THETA= 0 MR 1/76*
D 0 5.2E-38 OR LESS CL=.90 ALLBAY 69 CNTR Q=-2/3 THETA=6.5MR 1/76*
D 0 2.6E-35 OR LESS CL=.90 ALLBAY 69 CNTR Q=+1/3 THETA=44 MR 1/76*
D 0 1.3E-35 OR LESS CL=.90 ALLBAY 69 CNTR Q=+2/3 THETA=44 MR 1/76*
D 0 1.7E-39 OR LESS CL=.90 ANTIPOV2 69 CNTR Q=-1/3 M=0-5GEV 1/76*
D 0 4. E-38 OR LESS CL=.90 ANTIPOV2 69 CNTR Q=-2/3 M=0-5GEV 1/76*
D 0 1.6E-36 OR LESS CL=.90 ANTIPOV 71 CNTR Q=-4/3 THETA=47 MR 1/76*
D 0 3.8E-36 OR LESS CL=.90 ANTIPOV 71 CNTR Q=-4/3 THETA=47 MR 1/76*
D 0 5. E-34 OR LESS CL=.90 JOVANOVI 75 CNTR Q=1/3, M=7-15 GEV 2/76*
D 0 2. E-34 OR LESS CL=.90 JOVANOVI 75 CNTR Q=1/3, M=15-26 GEV 11/75*
D 0 1.3E-34 OR LESS CL=.90 JOVANOVI 75 CNTR Q=2/3, M=10-26 GEV 11/75*
D 0 8. E-35 OR LESS CL=.90 JOVANOVI 75 CNTR Q=4/3, M=10-26 GEV 11/75*
D DORFAN 65 IS A 30 GEV/C P EXPERIMENT AT BNL. 2/74
Y SEE FOOTNOTE Y IN SUBSECTION C ABOVE. 2/76*
A SEE FOOTNOTE A IN SUBSECTION C ABOVE. 2/76*
V FIRST ANTIPOV 71 VALUE IS FOR M=1.9-2.3, 2.7-4.4GEV, SECOND IS FOR M=2.3-2.7GEV. SEE ALSO NOTE V IN SECTION C ABOVE. 1/76*
J JOVANOVI 75 FIG.4 COVERS RANGES Q=1/3 TO 2 AND M=3 TO 26 GEV. 11/75*
J THIS IS A CERN ISR 26+26, 22+22 GEV P+P EXPERIMENT. 2/76*

INTERMEDIATE BOSON SEARCHES

W BOSON MASS LIMITS (GEV)
M B 1.7 OR MORE CL=.99 BERNARDINI 65 HYBR + NEU N, CERN 2/74
M B 0 2.0 OR MORE CL=.90 BURNS 65 OSPK + NEU N, BNL 2/74
M C 0 3.8 OR MORE CL=.90 BARISH 73 ASPK + W TO LEP+NEU=.2 2/74
M C 0 4.5 OR MORE CL=.90 BARISH 73 ASPK + W TO LEP+NEU=.5 2/74
M C 0 4.7 OR MORE CL=.90 BARISH 73 ASPK + W TO LEP+NEU=.8 2/74
M E 0 5.0 OR MORE CL=.95 BERGESON 73 ELEC 1/76*
M B LOOKED FOR (NEU N) TO (W+ MU- N), W+ TO (MU+ NEU, E+ NEU, OR HOPNS) 2/74
M C BARISH 73 LOOKED FOR (NEU N) TO (W+ MU- N), W+ TO (MU+ NEU) AT NAL. 2/74
M C RESULT GIVEN FOR THREE ASSUMED BR.FRACS. W+ TO (LEPTON NEU)/ALL. 2/74
M E BERGESON 73 LOOKED AT ENERGY DISTR OF NEU-INDUCED MUON FLUX UNDER- 1/76*
GROUND. SCALE INVARIANCE OF THE INELASTIC STRUCT FN ASSUMED.

QUARK FLUX FROM COSMIC RAY EXPERIMENTS (NUMBER/CM**2-SR-SEC)
F C 0 3. E-10 OR LESS KASHA 68 CNTR M=5GEV OR MORE 2/74
F 0 2.4E-8 OR LESS FRANZINI 68 CNTR V=.5-.9E M=2-5GEV UP 2/74
F 0 5. E-11 OR LESS CL=.90 FUKUSHIMA 69 CNTR Q=1/3 SEA LEVEL 2/74
F 0 7.5E-10 OR LESS CL=.90 FUKUSHIMA 69 CNTR Q=2/3 SEA LEVEL 2/74
F M 1 EVENT CLMED MCCUSKER 69 CC Q=2/3 2/74
F B 0 1. E-10 OR LESS CL=.90 BOHM 72 CNTR Q=1/3 2/74
F B 0 1. E-10 OR LESS CL=.90 BOHM 72 CNTR Q=2/3 2/74
F H 0 1.7E-8 OR LESS CL=.90 HICKS 73 CNTR Q= 1/3 1/76*
F H 0 1.7E-8 OR LESS CL=.90 HICKS 73 CNTR Q= 2/3 1/76*
F C BJORNBOE 68 LOOKED FOR DELAYED PARTICLES AFTER AIR SHOWERS. 2/74
F M MCCUSKER 69 CLAIMS 1 CANDIDATE. LATER SIMILAR EXPTS. SEE NONE. 2/74
F B BOHM 72 IS FLUX IN 10**14 TO 10**15 EV AIR SHOWERS 2/74
F H HICKS 73 LOCKED AT LARGE ZENITH ANGLES, THUS USING THE ATMOSPHERE 1/76*
F H AS AN EXTENDED FILTER FOR HADRONIC QUARKS. THEIR SEARCH PUTS AN 1/76*
F H UPPER LIMIT ON LEPTONIC QUARK FLUX IN COSMIC RAYS.

REFERENCES FOR INTERMEDIATE BOSON SEARCHES
+BIENLEIN,BOHM,DARDEL,FAISSNER+ (CERN)
+GOUL LANGS,HYMAN,LEDERMAN,LEE+ (COLU+BNL)
ANKENBRANDT,LARSEN,LEIPUNER+ (BNL+YALE)
+BARTLETT,BUCHHOLZ,HUMPHREY+ (CIT+FNAL)
+CASSIDAY,HENDRICKS (UTAH)
+D'ANGELO,GATTO,PAULUZI (RCMA)

REFERENCES FOR QUARK SEARCHES
DORFAN 65 PRL 14 999 +EADES,LEDERMAN,LEE,TING (COLU)
CHUPKA 66 PRL 17 60 +SCHIFFER,STEVENS (ANL)
GALLINAR 66 PRL 23 609 GALLINARO,MORPURGO (CERN)
STOVER 67 PR 1746 1599 +MORAN,TRISCHKA (SVRA)
BJORNBOE 68 NC 853 241 +DAMGARD,HANSEN,CHATTERJEE+ (BOHR+BEN)
FRANZINI 68 PRL 21 1013 FRANZINI,SULMAN (COLU)
KASHA 68 PR 172 1297 +STEFANSKI (BNL+YALE)
RANK 68 PR 176 1635 D.W.RANK (MICH)
ALLBAY 69 NC 64A 75 +BIANCHINI,DI DENDS,DORINSON,HARTUNG+ (CERN)
ANTIPOV1 69 PL 298 245 +KARPOV,KHROMOV,LANDSBERG,LAPSHIN+ (SERP)
ANTIPOV2 69 PL 308 576 +BOLOTOV,DEVISHEV,DEVISHEVA,ISAKOV+ (SERP)
COOK 69 PR 188 2092 +DEPASQUALI,FRAUENFELDER,PEACOCK+ (ILL)
FUKUSHIMA 69 PR 178 2058 FUKUSHIMA,KIFUNE,KONDO,KOSHIBA+ (TOKY)
MCCUSKER 69 PRL 23 658 MCCUSKER,CAIRNS (SYDNEY)
ANTIPOV 71 NP 827 374 +KACHANDU,KUTJIN,LANDSBERG,LEBDEEV+ (SERP)
BOHM 72 PRL 28 326 +DIEMONT,FAISSNER,FASOLD,KRISOR+ (AACH)
BOTT-BOD 72 PL 403 693 +CALDWELL,ANDRUS,SMITH,PEAK+ (CERN+PMHM)
ALPER 73 PL 468 265 BRITISH-SCANDINAVIAN COLLABORATION
HICKS 73 NC 144 65 +PLINT,STANDIL (MANI)
LEIPUNER 73 PRL 31 1226 +LARSEN,SESSOMS+SMITH,WILLIAMS+ (BNL+YALE)
JOVANOVI 75 PL 56B 105 JOVANOVI+ (MANI+YACH+CERN+GENO+HARV+TORTI)

Data Card Listings
For notation, see key at front of Listings.

Stable Particles
PARTICLE SEARCHES

MAGNETIC MONOPOLE SEARCHES

CHARMED HADRON SEARCHES

C MONOPOLE PROD. CROSS SECTION - ACCELERATOR EXP. (CM**2)/NUCLEON
A 0 1 E=40 OR LESS CL=.95 AMALDI 63 EMUL M=0 TO 5 GEV 12/75*

CE CHARMED HADRON PRODUCTION CROSS SECTION(E+ E-) (UNITS 10**-35 CM**2)
CE B 0 18. OR LESS CL=.90 BOYARSKI 75 SPEC K+ PI+, K+ PI- 2/76*

REFERENCES FOR MAGNETIC MONOPOLE SEARCHES
+BARONI,MANFREDINI,BRADNER+(ROMA+USC+D+ERN)

REFERENCES FOR CHARMED HADRON SEARCHES
AUBERT 75 PRL 35 416 +BECKER,BIGGS,BURGER,CHEN+ (MIT+BNL)

Stable Particles
PARTICLE SEARCHES

Data Card Listings

For notation, see key at front of Listings.

OTHER STABLE PARTICLE SEARCHES

SEARCHES FOR STABLE PARTICLES (I.E. PARTICLES IMMUNE TO DECAY VIA THE STRONG INTERACTION) NOT AMONG ABOVE SEARCHES ARE LISTED HERE.

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D HEAVY PARTICLE PRODUCTION DIFFERENTIAL CROSS SECTION (CM**2/SR-GEV)
D A 0 1. E-31 OR LESS CL=.90 APPEL 74 CNTR +- M=3.2-7.2 GEV 2/76*
D J 0 8. E-35 OR LESS CL=.90 JOVANDVIC 75 CNTR +- M=15-26 GEV 2/76*
D J 0 1.5E-34 OR LESS CL=.90 JOVANDVIC 75 CNTR Q=+-2, M=3-10 GEV 2/76*
D J 0 6. E-35 OR LESS CL=.90 JOVANDVIC 75 CNTR Q=+-2, M=10-26 GEV 2/76*
D A APPEL 74 IS NAL 300 GEV P+W EXPERIMENT. STUDIES FORWARD PRODUCTION 2/76*
D A OF HEAVY (UP TO 24 GEV) CHARGED PARTICLES WITH MOMENTA 24-200GEV(-) 2/76*
D A AND 40-150GEV (+CHG). ABOVE TYPICAL VALUE IS FOR 75 GEV. 2/76*
D J JOVANDVIC 75 IS A CERN ISR 26+26 AND 15+15 GEV P+P EXPERIMENT. 2/76*
D J FIG.4 COVERS RANGES Q=1/3 TO 2 AND M=3 TO 26 GEV. 2/76*

F HEAVY PARTICLE FLUX IN COSMIC RAYS (NUMBER/CM**2-SEC-SR) 1/76*
F Y 5 6 E-9 OR MORE YOCK 74 CNTR M GT 6 GEV 1/76*
F Y YOCK 74 EVENTS COULD BE TRITONS. 1/76*

C LIGHT (BETWEEN MU AND E MASSES) PARTICLE MASS (UNITS-ELECTRON MASSES)
C 0 NONE BETWEEN 2 AND 13 BLAGOV 75 CNTR SPINOR,TAU=2E-10SEC 2/76*
C 0 NONE BETWEEN 2 AND 10.6 BLAGOV 75 CNTR SCALAR,TAU=2E-10SEC 2/76*

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REFERENCES FOR OTHER STABLE PARTICLE SEARCHES

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APPEL 74 PRL 32 428 +BOURQUIN,GAINES,LEDERMAN,PAAR+ (COLU+NAL)
YOCK 74 NP 876 175 P.C.M.YOCK (UNIV OF AUCKLAND)
BLAGOV 75 YAD.FIZ. 21,300 +KOMAR,MURASHOVA,SYREISHCHIKOVA+ (LEBD)
JOVANDVI 75 PL 56B 105 JOVANDVIC+ (MANI+AACH+CERN+GEND+HARV+TORI)

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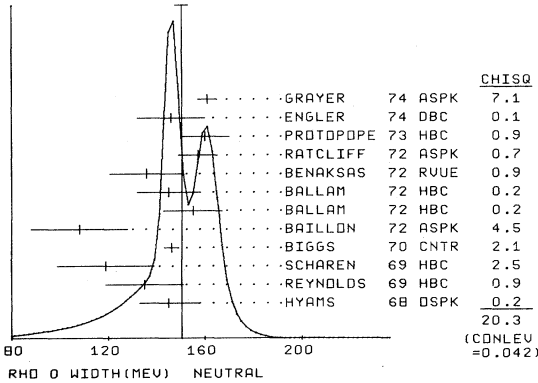
Mesons

$\rho(770)$

Data Card Listings

For notation, see key at front of Listings.

WEIGHTED AVERAGE = 150.3 ± 2.7
ERROR SCALED BY 1.4



9 RHO PARTIAL DECAY MODES

Mode	Decay Masses
P1 RHO INTO 2PI	139+ 139
P2 RHO INTO 4PI	139+ 139+ 139+ 139
P3 RHO INTO PI GAMMA	139+ 0
P4 RHO INTO E+ E-	0+ 0
P5 RHO INTO PI ETA (VIOLATES G)	139+ 548
P6 RHO INTO MU+ MU-	105+ 105
P7 RHO INTO PI+ PI- P (VIOLATES G)	139+ 139+ 134
P8 RHO INTO ETA GAMMA	548+ 0

9 RHO BRANCHING RATIOS

Mode	Ratio	Reference
R1 RHO INTO 4PI/2PI	(P2)/(P1)	
R1 RHO+ INTO (PI+ PI+ PI- P) / (PI+ P)IO	3-5 PI- P	6/66
R1 (0.01) OR LESS	DEUTSCHMA 66 HBC	6/66
R1 (0.002) OR LESS	FERBEL 66 HBC	11/66
R1 (0.0035 ± 0.004)	JAMES 66 HBC	11/66
R1 RHO INTO (PI+ PI- PI+ PI-) / (PI+ PI-)	0 2.1 PI+P	6/66
R1 (0.008) OR LESS	JAMES 66 HBC	6/66
R1 (0.002) OR LESS	CHUNG 68 HBC	7/67
R1 (0.002) OR LESS	HUSON 68 HLBC	10/66
R1 (0.002) OR LESS	GERMAN 69 HBC	10/66
R2 RHO INTO PI GAMMA/2PI (UNITS 10**3)	(P3)/(P1)	
R2 M (2.0) OR LESS	LANZEROTT 65 CNTR	11/66
R2 (5.) OR LESS	FIDECCARO 66 OSP	10/66
R2 (7.) OR LESS	HUSON 66 HLBC	6/66
R2 M (2.) OR LESS	GERMAN 69 HBC	12/75*
R2 C 0.24 ± 0.07	GOBBI 74 OSP	12/75*
R2 C GOBBI 74 QUOTES THE PARTIAL WIDTH FOR RHO- INTO PI- GAMMA		
R2 C AS (35 ± 10) KEV		
R2 M ONE PION EXCHANGE MODEL USED IN THIS ESTIMATION		

Note on $\rho^0 \rightarrow e^+e^-$

Extraction of a ratio for $\rho^0 \rightarrow e^+e^-$ is complicated by interference with ω decay. In photoproduction, $\gamma A \rightarrow e^+e^-A$, there is substantial interference between the allowed $(\rho^0, \omega) \rightarrow e^+e^-$ decays. The interference in the colliding-beam reaction $e^+e^- \rightarrow \pi^+\pi^-$ is due to G-parity-violating mixing of the overlapping ρ^0 and ω resonances; it alters the results for the rate $\Gamma(\rho^0 \rightarrow e^+e^-)$ only by a small amount. Therefore at present we average only the values from the $e^+e^- \rightarrow \pi^+\pi^-$ experiments.

R3 RHO INTO E+ E-/(PI+PI-) (UNITS 10**4)	(P4)/(P1)	
R3 P 94 (0.65) (0.14)	ASBURY 1 67 CNTR	PHOTOPRODUCTION 9/67
R3 P POSSIBLY LARGE RHO-OMEGA INTERFERENCE		
R3 H (0.65) (1.1) (0.5)	HERTZBACH 67 OSPK	ASSUME SU(3)+MIXING. 10/66
R3 H NOT SEPARATED FROM CMEGA DECAY.		
R3 A 33 (0.59) (0.11)	ASTVACAT 66 OSPK	ASSUME SU(3)+MIXING 6/68
R3 A NOT SEPARATED FROM CMEGA DECAY. ERROR STATISTICAL ONLY.		
R3 0.50 ± 0.10	AUSLENDER 69 OSPK	E+E- COLLID.BEAM 9/68
R3 F (0.49) (0.12) (0.15)	BIGGS 70 CNTR	PHOTOPRODUCTION 6/70
R3 F ASSUMING RHO WIDTH 140 MEV. ERROR STATISTICAL ONLY.		
R3 0.41 ± 0.05	BENAKSAS 72 OSPK	E+E- COLL.BEAMS 12/72
R3		
R3 AVG 0.428 ± 0.049	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R3 STUDENT 0.428 ± 0.049	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	

R4 RHO INTO (PI ETA)/(2PI)	(P5)/(P1)	
R4 (0.03) OR LESS	DEUTSCHMA 66 HBC	+ 8.0 PI+ P 6/66
R4 (0.008) OR LESS	FERBEL 66 HBC	+ PI+ P ABOVE 2.5 11/66
R5 RHO INTO (MU+ MU-)/(PI+ PI-) (UNITS 10**4)	(P6)/(P1)	
R5 SEE NOTE UNDER RHC INTO E+E- ABCVE		
R5 0.97 ± 0.31	0.33 HYAMS	47 OSPK 11 PI- LI H 6/67
R5 H HYAMS MASS RESOL. IS 20 MEV. THE OMEGA REGION WAS EXCLUDED.		
R5 R 0.82 ± 0.16	0.36 ROTHWELL 69 CNTR	PHOTOPRODUCTION 4/70
R5 R POSSIBLY LARGE RHO-CMEGA INTERFERENCE LEADS US TO INCREASE		
R5 R THE MINUS ERROR		
R5 WEHMANN 69 OSPK	12 PI- CN C.FE	7/69
R5 W RESULT CONTAINS (11 ± 1) PER CENT CORRECTION USING SU(3)		
R5 W FOR CENTRAL VALUE. THE ERROR ON THE CORRECTION TAKES ACCOUNT		
R5 W OF POSSIBLE RHO-OMEGA INTERFERENCE AND THE UPPER LIMIT AGREES		
R5 W WITH THE UPPER LIMIT OF (CMEGA INTO MU+ MU-) FROM THIS EXP.		
R5 AVG 0.67 ± 0.12	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R5 STUDENT 0.67 ± 0.14	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R6 RHO INTO (PI+ PI- P)IO/(PI+ PI-)	(P7)/(P1)	
R6 G (0.01) OR LESS	CL=.84 ABRAMS	71 HBC 0 3.7 PI+ P 11/71
R6 G MODEL DEPENDENT, ASSUMES I = 1,2,OR 3 FOR THE 3PI SYSTEM		11/71
R7 RHO INTO (ETA GAMMA)/TCTAL	(P8)	
R7 (0.01) OR LESS	CL=.90 NORDBERG 74 ASPK	7-9 GAMMA+CU 12/75*

REFERENCES FOR RHO

ANDERSON 61 PRL 6 365
 ANDERSON, BANG, BURKE, CARMGAY, SCHMITZ (LRL)
 ERWIN 61 PRL 6 628
 A.R., R. MARCH, W. C. WALKER, E. WEST (TWISCO)
 KENNEDY 62 PR 126 736
 V P KENNEY, M D SHEPHARD, C D GALL (KENTUCKY)
 SANDIOS 62 PR 125 139
 SAMIOS, BACHMANN, LEAS, BROWN, COLUCCI, KENNY
 XUONG 62 PR 128 1849
 NGUYEN HUU XUONG, GERALD R LYNCH (LRL)

ABOLINS 63 PRL 11 381
 ABOLINS, LANDER, MEHLHOP, NGUYEN, YAGER (UCSD)
 ALLITTI 63 PRL 10 62
 ALLITTI, BATON, ARNEMISE (SACL+ORS+COLU+RUTG)
 CHADWICK 63 PRL 10 62
 CHADWICK, DAVIES, DERRICK, CRESTI + (OXF+PADO)
 GUIRAGOS 63 PRL 11 85
 ZAVEN GUIRAGOSIAN (LRL)
 SACLAY 63 SIENA CONF 1 239
 SACLAY+CRSAY+BARI + BOLOGNA- COLLABORATION

BONDAR 64 NC 31 729
 BONDAR+ (AACCHEN+BRN+BDNN+DES+Y+OIC+MPI+)
 CARMONY 64 DUBNA CONF 1 486
 CARMONY, HOA, LANDER, NG, H. XUCNG, YAGER (UCSD)
 GOLDHABER 64 PRL 12 336
 GOLDHABER, BROWN, KADYK, SHEK+ (LRL+UCB)

ALYEA 65 PL 15 82
 ALYEA, CRITTENDEN, MARTIN, RHODE + (INDIANA)
 ARMENISE 65 NC 37 361
 SACLAY+ORSAY+BARI+BOLOGNA COLLABORATION
 BLIEDEN 65 PL 19 444
 CERN MISSING MASS SPECTROMETER GROUP (CERN)
 CLARK 65 PR 139 8 1556
 A CLARK, CHRISTENSON, CROMIN, TURLAY (PRINCETON)
 GUTAY 65 NC 39 381
 GUTAY, LANUTT, TULLI (FSU)
 LANZEROTT 65 PRL 15 210
 LANZEROTT, BLUMENTHAL, EHN, FAISSLER + (HARV)
 ZDANIS 65 PRL 14 721
 ZDANIS, MADANSKY, KRAEMER + (JHU+BNL)

ACCENSI 66 PL 20 557
 ACCENSI, ALLES-BORELLI, FRENCH, FRISK+ (CERN)
 ALF-STEIFENBERGER, BELLEY, BRUGER+ (COLU+RUTG)
 BALTAY 66 PR 145 1072
 +FRANZINI, LUJENIS, SEVERINS, TYCKO+ (COLUMBIA)
 BLIEDEN 66 NC 43 71
 CERN MISSING MASS SPECTROMETER GROUP (CERN)
 CAMBRIDGE 66 PR 146 994
 CAMBRIDGE BUBBLE CHAMBER GROUP (MIT+HARV+)
 CASON 66 PR 148 1282
 N H CASON (MISCONSIN)
 DEUTSCHMA 66 PL 20 82
 DEUTSCHMANN, STEINBERG + (AACCHEN+BERLIN+ CERN)
 FERBEL 66 PL 21 111
 FERBEL (ROCHESTER)
 FIDECCARO 66 PL 23 163
 G+M FIDECCARO, J. POIRIER, P. SCHIAVON (CERN)
 HAGOPIAN 66 PR 145 1128
 HAGOPIAN, SELOVE, ALITTI, BATON+ (PENNSYLVANIA)
 HAGOPIAN, PAN (PENNSYLVANIA, LRL-BERKELEY)
 HUSON 66 PL 20 91
 HUSON, ALLARD, DRIJARD, HENNESSY+ (ORSAY+POL)
 JACOBS 66 UCLR-16877
 L.O. JACOBS
 JAMES 66 PR 142 896
 F E JAMES, KRABYBIL (VALE+ BROOKHAVEN)
 WEST 66 PR 149 1089
 WEST, BOYD, ERWIN, WALKER (MISCONSIN)

ALLES-BORELLI 67 NC 50 A 776
 ALLES-BORELLI, FRENCH, FRISK, + (CERN+BNL)
 ASBURY 1 67 PRL 19 869
 ASBURY+BECKER+BERRAM+JODS+JORDAN+ (DESY+CGU)
 ASBURY 2 67 PRL 19 865
 ASBURY+BECKER+BERRAM+JODS+JORDAN+ (DESY+CGU)
 BACON 67 PR 157 1263
 +FICKINGER, HILL, HOPKINS, ROBINSON+ (BNL)
 BANNER 67 PL 25 B 300
 +FAYOUX, HAMEL, ZEMBERY, CHEZE+ (SACLAY+CAEN)
 BARLOW 67 NC 50A 701
 +LILLETOL+MONTANET+ (CERN+COEF+TRAD+LIVP)
 BATON 67 PL 25 B 419
 J. BATON, G. LAURENS, J. REIGNIER (SACLAY)
 ALSO 67 NC 3 349
 J. BATON, G. LAURENS, J. REIGNIER (SACLAY)
 CLEAR 67 NC 9 399
 +JOHNSTON+COOPER+MANNER+ (TNTO+ANL+TSC)
 DANYSZ 67 NC 51 A 801
 DANYSZ+FRENCH+SIMAK (CERN)
 EISNER 67 PR 164 1699
 +JOHNSON+KLEIN+PETERS+SAHNI+YEN+ (PURDUE)
 FRENCH 67 NC 52A 442
 +KINSON+MCDONALD+RIDDFORD+ (CERN+BRN)
 HERTZBACH 67 PR 155 1461
 HERTZBACH, KRAEMER, MADANSKI, ZDANIS+ (JHU+PENN)
 ALSO 65 ZDANIS

HUWE 67 PL 248 252
 HUWE+OPPENHEIMER+SCHULTZ+WILSON (COLU)
 HYAMS 67 PL 248 634
 +KOCH+PELLETT+POTTER+VON LINDERN+ (CERN+MPI)
 MILLER 67 PR 153 1423
 MILLER, GUTAY, JOHNSON, LOEFLER + (PURDUE)
 POIRIER 67 PR 163 1462
 +BISWAS, CASON, DERADO, KENNEY+ (NDAM+PENN)

ABC COLL 68 NP 84 501
 AACHEN+BERLIN+CERN COLLABORATION+
 ARNEMISE 68 NC 54A 999
 +GHIDINI, FORINO+ (BAR+BGNA+FRZ+ORSAY)
 ASTVACAT 68 PL 27 B 45
 ASTVACATUROV, AZIMOV, BALDIN+ (JINR-MOSCOW)
 BATON 68 PR 176 1574
 J.P. BATON, G. LAURENS (SACLAY)
 BLECHSCH 68 NC 53 A 1045
 BLECHSCHMIDT, DOND, ELSNER, + (DESY+MCHS)
 ALSO 67 NC 52 A 1348

CHUNG 68 PR 165 1491
 S-U, CHUNG, O. I. DAHL, J. KIRZ, D. H. MILLER (LRL)
 DONALD 68 NP 8 B 174
 +EDWARDS, FRODESSEN, BETTINI+ (LIVP+OSLO+PADO)
 FOSTER 68 NP 8 B 107
 +GAVILLET+LABROSSE+MCNATANET+ (CERN+COEF)
 HUSON 68 PL 28B 208
 +LUBATTI, SIX, VEILLET, + (ORSAY+MILA+UCLA)
 HYAMS 68 NP 8 7 1
 +KOCH, POTTER, WILSON, VON LINDERN+ (CERN+MPI)
 JONES 68 PR 166 1405
 +BLEULIG, CALDWELL, ELSNER, HARTING+ (CERN)
 JOHNSON 68 PR 176 1651
 +POIRIER, BISWAS, GUTAY+ (NDAM+PURDUE+SLAC)
 KEY 68 PR 166 1430
 +PRENTICE+COOPER+MANNER+ (TNTO+ANL+TSC)
 LAMSA 68 PR 166 1395
 +CASON+BISWAS+DERACC+GROVES+ (NOTREDAME)
 LANZEROTT 68 PR 166 1424
 LANZEROTT, BLUMENTHAL, EHN, FAISSLER + (HARV)
 MARATECK 68 PRL 21 1613
 +MAGCPIAN, + (PENNSYLVANIA+COLG+PURDUE+TNTO+MPI)
 PISUT 68 NP 8 B 325
 J. PISUT, M. ROOS (CERN)

AUGUST 69 PL 28 B 508
 +EIZOT+BUONHAISSINSKI+ALANNE+ (ORSAY)
 AUGUST 69 LNC 2 214
 +LEFRANCOIS, LEHMANN, MARTIN, + (ORSAY)
 AUSLENDER 69 SJNP 9 69
 AUSLENDER, BUDKER, PANTUSOVA, PESTOV+ (NOVO)
 GERMAN 69 PR 188 2060
 GERMAN BUBBLE CHAMBER COLL. (DESY)
 HASSINS 69 ARGONNE CONF. 373
 J. HASSINSKI, BLUMENFELD, KOPPELMAN, LIBBY, + (ISSU-COLJ)
 JUMALA 69 PR 184 1461
 +LEAGOCK, RHODE, KOPPELMAN, LIBBY, + (ISSU-COLJ)
 MALAMUD 69 ARGONNE CONF. P. 93
 E. MALAMUD, P. SCHLEIN (UCLA)
 MILLER 69 PR 178 2061
 R. MILLER, LICHTMAN, WILLMANN (PURDUE)
 MOTT 69 PR 177 1966
 +AMMAR, DAVIS, KRUPAC, SLATE, CAGAN+ (MNS+ANL)
 REYNOLDS 69 PR 184 1424
 +ALBRIGHT, BRADLEY, BRUCKER+HARMS+ (FSU)
 ROOS 69 NP 10 563
 M. ROOS, J. PISUT (CERN+BRATISLAVA)
 ROTHWELL 69 PRL 23 1521
 +CHASE, EARLES, GELTNER, GLASS, WEINSTEIN+ (NEAS)
 SCHAREN 69 ARGONNE CONF. 306
 SCHARENGUEVEL (PURDUE)
 WEHMANN 69 PR 178 2095
 +ENGELS, WILSON, + (HARV+CASE+SAC+CDRN+MGCT)

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho(770)$, $\omega(783)$

Table of particle data cards for various mesons, including fields like PRAL, NP, PL, PR, and names of researchers.

$\omega(783)$

Table for $\omega(783)$ listing OMEGA MASS (MEV), OMEGA FULL WIDTH (MEV), and OMEGA PARTIAL DECAY MODES with associated branching fractions.

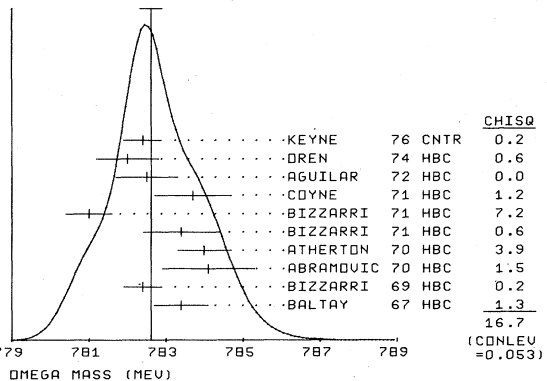


Table for OMEGA FULL WIDTH (MEV) showing various decay channels and their widths.

Table for OMEGA PARTIAL DECAY MODES showing branching fractions for various decay channels.

Table for FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS, including a matrix and a list of fitted parameters.

Table for OMEGA BRANCHING RATIOS showing ratios for various decay channels.

Table for OMEGA INTO (PI+ PI-)/(PI+ PI- P) and other ratios, showing branching fractions and correlations.

Mesons

$\omega(783)$, $M(940)$

Data Card Listings

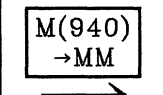
For notation, see key at front of Listings.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters. Includes entries for Omega and Cmega into various gamma and pi transitions.

***** REFERENCES FOR CMEGA

Table of references for CMEGA, listing authors, journal names, volumes, and page numbers.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters. Includes entries for various meson decays and production cross-sections.



66 M(940) EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters.

Table listing experimental data for mesons, including columns for experiment name, particle type, mass, width, and other parameters.

***** REFERENCES FOR M(940)

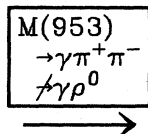
Table of references for M(940), listing authors, journal names, volumes, and page numbers.

Data Card Listings

For notation, see key at front of Listings.

Mesons

M(953), η'(958)



59 M(953, JPC= 0-+)
WHILE MASS AND WIDTH ARE CONSISTENT WITH ETA PRIME(958), THE PI+ PI- GAMMA DECAY SHOWS NO RHO SIGNAL. POSSIBLY SEEN IN MMS. OMITTED FROM TABLE.

Table with 2 columns: M, M(953.0) 2.0 AGUILAR 70 HBC 3.9-4.6K-P, P K-M 1/71

Table with 2 columns: W, M(10.0) OR LESS CL=.95 AGUILAR 70 HBC 3.9-4.6K-P, P K-M 1/71

Table with 2 columns: P1, M INTO PI+ PI- GAMMA 139+ 139+ 0

Table with 2 columns: R1, M INTO (RHOO GAMMA)/(ALL PI+ PI- GAMMA) (P2)/(P1) 1/71

Table with 2 columns: AGUILAR 70 PRL 25 1635 REFERENCES FOR M

η'(958)

2 ETA PRIME(958, JPC=0-+) I=0

Note on the J^P Assignment of η'(958)

From the Dalitz plot analyses of the η' → πππ and η' → π+π-γ decays, and from the observation of a η' → γγ decay mode, all assignments except J^PC = 0-+ and 2-+ are excluded.

Table with 2 columns: M, O ONLY EXPERIMENTS GIVING ERROR LESS THAN 2 MEV KEPT FOR AVERAGING 12/75*

Table with 2 columns: M, O 957. 1. RITTENBER 69 HBC 1.7-2.7 K- P 9/69

Table with 2 columns: W, M(8.0) OR LESS CL=.90 DAUBER 64 HBC 1.95 K-P 11/71

Table with 2 columns: P1, ETA PRIME INTO PI+ PI- ETA DECAY MASSES 139+ 139+ 548

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i^2/δP_i^2, where δP_i = sqrt(δP_i δP_i), while the off-diagonal elements are the normalized correlation coefficients (δP_i δP_j)/(δP_i δP_j).

Table with 2 columns: P 1 .4352+- .0174 P 2 -.5783 .2407+- .0198

2 ETA PRIME BRANCHING RATIOS

Note on η'(958) Branching Fractions

In our calculation of the branching fractions of the η'(958), we use the decay modes ηππ (including ηπ^0π^0), ρ^0γ, and γγ. It is assumed that the rate η → neutrals is 71.1%.

In the fit we do not use the constraint

R = Γ(η' → ηπ+π-) / Γ(η' → ηπ^0π^0) = 2

from I-spin conservation. The result of the fit is in agreement with it: R = 1.8 ± 0.2.

Table with 2 columns: R1, ETA PRIME INTO (PI+ PI-ETA (NEUTRAL DEC.))/TOTAL (PIN) 2.7 K-P 10/66

Mesons
eta'(958), delta(970)

Data Card Listings

For notation, see key at front of Listings.

Table of particle data for eta prime and delta mesons, including decay rates, branching ratios, and fit parameters. Rows include R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, R14, R15, R16, R18, R19, R20, R21, R22, R23, R24, R25, R26.

Table of particle data for eta prime meson, including decay rates and fit parameters. Rows include R27, R28, R29.

Table of decay asymmetry parameters for eta prime meson. Rows include A, A, A, A, A, A, A, A.

Table of references for eta prime meson, listing authors and publication details. Rows include DAUBER, GOLDBERG, KALBFLEI, etc.

delta(970) -> eta pi pi

Under this entry, we list three types of I=1 peaks near KK threshold:
1) Missing-mass peaks, mostly controversial.
2) eta pi decays, peaking slightly below KK threshold. This defines I^G = 1^- and J^P = Normal.
3) Threshold enhancements in the (KK)^+ system with I=1. The Q value is low and J^P therefore probably 0^+.

Data Card Listings

For notation, see key at front of Listings.

Mesons

delta(970), H(990), S*(993)

In listing these types of peaks together under a common entry we do not imply that they are necessarily all related. However, the K-K bar threshold enhancement may be due to a virtual bound state that could also be responsible for the eta pi peaks (ASTIER 67). In a coupled-channel analysis of the most significant eta pi and K-K bar data, MORGAN 75 shows that the delta(970) probably is strongly coupled to both channels. It then constitutes the isovector member of the 0+ nonet. MORGAN 75 gets a width much larger than what is observed in the eta pi decays because of the strong K-K bar coupling.

36 DELTA(970) MASS (MEV)

Table with columns for mass (MEV), missing mass experiments, and references. Includes entries for K BAR ONLY, SCAT. LENGTH, and AVERAGE ERROR.

36 DELTA(970) WIDTH (MEV)

Table with columns for width (MEV), missing mass experiments, and references. Includes entries for K BAR ONLY, SCAT. LENGTH, and AVERAGE ERROR.

36 DELTA(970) PARTIAL DECAY MODES

Table with columns for decay modes and decay masses.

36 DELTA(970) BRANCHING RATIOS

Table with columns for branching ratios and references.

REFERENCES FOR DELTA(970)

Table of references for Delta(970) from various experiments and authors.

H(990)

35 H(990, JPG=A -) I=0

THE EVIDENCE OF BENSON 66 HAS DISAPPEARED AFTER RE-ANALYSIS (CHAUDHARY 70). NO SIGNIFICANT OTHER EVIDENCE HAS BEEN PUBLISHED. OMITTED FROM TABLE.

REFERENCES FOR H

Table of references for H(990) from various experiments and authors.

S*(993)

3 S*(993, JPG=0++) I=0

UNDER THIS ENTRY WE LIST PARAMETERS OF THE POLE IN THE ISOSCALAR S-WAVE. FOR A MINI-REVIEW SEE UNDER EPSILON. POSSIBLE EVIDENCE OF D-WAVE PI-PI INTERACTIONS IN THIS REGION IS LISTED SEPARATELY UNDER ETA N(1080).

FOR EARLY WORK USING BREIT-WIGNER OR SCATTERING LENGTH PARAMETRIZATION IN FITS TO THE (K BAR) MASS SPECTRUM SEE REFERENCE SECTION AND OUR 1972 EDITION.

3 REAL PART OF THE S* POLE POSITION (MEV)

Table with columns for real part of pole position (MEV) and references.

Mesons
S*(993), phi(1020)

Data Card Listings

For notation, see key at front of Listings.

3 NEGATIVE IMAG. PART OF THE S* PCLE POSITION (MEV)
CORRESPONDS TO HALF-WIDTH, NOT FULL WIDTH.
W H 40. 40. 60. HOANG 69 DSPK 4. PI-P,K,S KS N 1/73
W H 30. 30. 70. HOANG 69 DSPK 5. PI-P,K,S KS N 1/73
W B (13.) BEUSCH 70 DSPK 4.6 PI-P 1/73
W P 27. 8. PROTOPOPE 73 HBC 7. PI+ P 1/74
W P ANOTHER SOLUTION HAS 52 MEV AND NO EPSILON PCLE.
W S 24. 7. BINNIE 73 CNTR PI- P,S* N 1/74
W S SEE NOTE S ABOVE
W (5.) ESTABROOK 73 ASPK 17 PI-P,PI+PI-N 12/75*
W (16.) GRAYER 73 ASPK 17 PI-P,PI+PI-N 12/75*
W 15. (5.) HYAMS 73 ASPK 0 17 PI-P,N,PI+PI- 1/74
W A (19.) (3.) FUJII 75 RVUE 17 PI-P,PI+PI-N 12/75*
W A S* AMPLITUDES PARAMETRIZED IN TERMS OF POLE POSITIONS USING
W A HYAMS 73 PHASE SHIFTS
W AVG 20.0 3.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W STUDENT 20.1 4.2 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

3 S* WIDTH (MEV)
(180.) APPROX. MORGAN 75 RVUE 12/75*

***** REFERENCES FOR S* *****

WANG 61 JETP 13 323 WANG TSU-TSENG, VEKSLER, VRANA, + (JINR)
BIGI 62 CERN CONF 247 A BIGI, S BRANDT, R CARRARA + (CERN)
BINGHAM 62 CERN CONF 240 H BINGHAM, H BLOCH + (EPOL+CERN)
ERMIN 62 PR 163 1377 *HARDY, HESSE, KIRZMILLER (LRL)
ERMIN, MOYER, MARCH, WALKER, WAMPLER (WISCS+BNL)
BALTAY 64 DUBNA CONF 1 409 *BARNES, JENNELL, FLAMINIO, GOLDBERG, + (BNL)
BARMIN 64 DUBNA CONF 1 433 BARMIN, DOLGICENKO, YEROFEEV, KRESTINI + (ITEP)
CRENNELL 66 PRL 16 1025 CRENNELL, KALBFLEISCH, LAI, SCARR, SCHU + (BNL)
*CAHL+HARDY+KIRZ+MILLER (LRL)
BARLOW 67 NC 50A 701 *LILLESTOL+MONTANET+ (CERN+CDEF+IRAD+LIVP)
BEUSCH 67 PL 25 B 357 *FISCHER, GOBBI, ASTBURY + (ETHZ+CERN)
DAHL 67 PR 163 1377 *HARDY, HESSE, KIRZMILLER (LRL)
ALITTI 68 PRL 21 1705 *BARNES, JENNELL, FLAMINIO, GOLDBERG, + (BNL)
LAI 68 PHILAD. CONF. P. 303 KWAN WU LAI (BNL)
PHELAN 68 THESIS JAMES J. PHELAN (ANL+ST. LOUIS UNIV)
ALSO 68 PRL 21 316 HOANG, EARTLY, PHELAN, ROBERTS + (ANL+CHIC+NDAM)
AGUILAR 69 PL 29 B 241 M. AGUILAR-BENITEZ, J. BARLOW, + (CERN+CDEF)
ALSO 67 BARLOW
ALSO 69 NP 8 14 195 M. AGUILAR-BENITEZ, J. BARLOW, + (CERN+CDEF)
HOANG 69 NC 61 A 325 T.-F. HOANG (ANL)
HOANG 69 PR 184 1363 *EARTLY, PHELAN, ROBERTS, + (ANL+ILLC)

BADIER 70 NP 8 22 512 *BCNNET, DREVILLON, BAUBILLIER, + (EPOL+IPNP)
BATON 70 PL 33 B 528 *LAURENS, REIGNIER (SACLAY)
BEUSCH 70 PHILA. CONF. P. 185 M. BEUSCH
HYAMS 70 PHILA. CONF. P. 41 *KOCH, BEUSCH, + (CERN+MPIM+ETHZ+LUIC+HAWAII)
ALSO 70 NP 8 22 189 HYAMS, KOCH, POTTER, VON LINDERN, + (CERN+MPIM)
OH 70 PR D 1 2494 *GARRINKEL, MORSE, WALKER, PRENTICE(WISCS+TNTD)
ALSTON-G 71 PL 36 B 152 ALSTON-GARNJUST, BARBARO-GALTIERI, + (LBL)
BASDEVAN 72 PL 41 B 178 BASDEVANT, FROGGATT, PETERSEN (CERN)
DAMERI 72 NC 9 A 1 *BORZATTA, GOUSSU, + (GENO+MILA+SACL)
DUBOC 72 NP 8 46 429 *GOLDBERG, MAKOWSKI, DCNALD, + (LNP+LIVP)
FLATTE 72 PL 38 B 232 *ALSTON-GARNJUST, BARBARO-GALTIERI, + (LBL)
GRAYER 72 PHIL. CONF. PROC. 5 *HYAMS, JONES, SCHLEIN, BLUM, DIETL, + (CERN+MPIM)
P. K. WILLIAMS
WILLIAMS 72 PR D 6 3178 P. K. WILLIAMS (FISU)
BINNIE 73 PRL 31 1534 *CARR, DEBENHAM, DUANE, GARBUTT, + (LOIC+SHMP)
DIAMOND 73 PR D 7 1977 *BINKLEY, + (WISCS+DUKE+GOLD+TNTD+OHIO)
ESTABROOK 73 TALLAHASSEE ESTABROOKS, MARTIN, GRAYER, HYAMS + (CERN+MPIM)
FUJII 73 NC 13 A 311 Y. FUJII, M. KATO (TOKYO)
GRAYER 73 TALLAHASSEE *HYAMS, JONES, BLUM, DIETL, KOCH + (CERN+MPIM)
HYAMS 73 NP 8 64 134 *JONES, WEILHAMMER, BLUM, DIETL, + (CERN+MPIM)
OGHS 73 THESIS W. OGHS (MPIM)
PROTOPOD 73 PR D 7 1280 PROTOPODESCU, GARNJUST, GALTIERI, FLATTE+(LBL)

phi(1020)
4 PHI(1020, JPG=1--) I=0

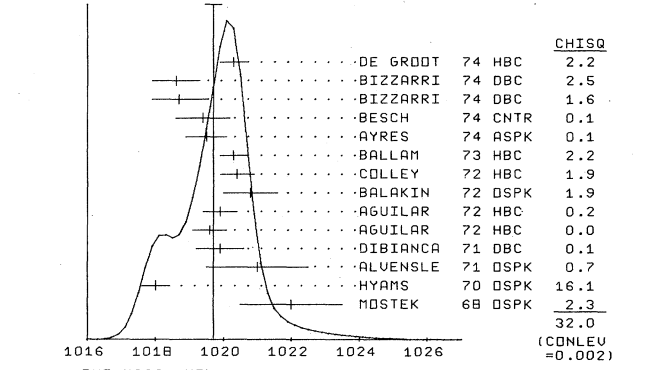
4 PHI MASS (MEV)

Table with columns for mass (M), width (S), and various decay modes and parameters for phi(1020).

W AVG 20.0 3.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
W STUDENT1019.84 0.20 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
(O SEE IDEOGRAM BELOW)

M S INSIGNIFICANT DATA WITH SMALL STATISTICS NO LONGER AVERAGED
M D MASS ERRORS ENLARGED BY US TO WIDTH/SQRT(N)

WEIGHTED AVERAGE = 1019.70 +/- 0.24
ERROR SCALED BY 1.6



4 PHI WIDTH (MEV)

Table with columns for width (W), standard deviation (S), and various decay modes and parameters for phi width.

4 PHI PARTIAL DECAY MODES

Table listing partial decay modes (P1-P14) and their corresponding decay masses.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i +/- delta P_i, where delta P_i = sqrt(delta P_i^2 + delta P_j^2), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of branching fractions for P1, P2, P3, and P4.

4 PHI BRANCHING RATIOS

Table with columns for branching ratios (R1-R4) and various decay modes and parameters.

Data Card Listings

For notation, see key at front of Listings.

Mesons

phi(1020), M(1033), B1(1040)

Table of particle data listings for phi(1020), M(1033), and B1(1040). Includes columns for particle name, mass, width, and various fit parameters.

Table of particle data listings for phi(1020), M(1033), and B1(1040). Includes columns for particle name, mass, width, and various fit parameters.

M(1033) -> MM

B1(1040) -> omega pi

REFERENCES FOR PHI

List of references for phi meson data, including authors and publication details.

REFERENCES FOR M(1033)

List of references for M(1033) meson data, including authors and publication details.

Mesons

$B_1(1040)$, $\eta_N(1080)$, $A_1(1100)$

Data Card Listings

For notation, see key at front of Listings.

DEFIOX 73 PL 43 B 141 +GDBRZYNSKI,ESPIGAT,NASCIMENTO,+ (CDEF)
 DIAZ 74 PRL 32 260 +DIBIANCA,FICKINGER,ANDERSCN,+ (CASE+GARN)

REFERENCES FOR $B_1(1040)$

$\eta_N(1080)$
 $\rightarrow \pi\pi$

30 ETA N(1080, JPC=N +) I=0 J GREATER THAN 1
 SOME EXPERIMENTS SUGGEST J=2.
 OMITTED FROM TABLE

Note on $\pi^+\pi^-$ Peaks Called $\eta_N(1080)$

The $\eta_N(1080)$ is seen in $\pi^-p \rightarrow \pi^+\pi^-n$ predominantly at backward decay angles, $\cos\theta < -0.75$. OH 70 state that this "bump is almost certainly the result of P-D interference."

Note that the selection made in some bubble chamber experiments to reduce the background under the $\eta_N(1080)$ in the reaction $\pi^-p \rightarrow \pi^+\pi^-n$ may lead to a sample of events ambiguous with $\pi^-p \rightarrow p\pi^-\pi^0\pi^0$ (BATON 70 and private communications from G. Laurens).

30 ETA N MASS (MEV)						
M	1060.0	15.0	MILLER	68 HBC	4.0 PI- P	9/68
M	70 1085.0	10.0	WHITEHEAD	68 ASPK	3.1-3.6 PI-P	10/67
M	1120.0	100.0	OH	69 HBC	7.PI- P,PI+ D	9/69
M	1112.0	16.0	CLAYTON	70 HBC	2.5 PBAR P, 4 PI	1/71
M	(1080.0)		DIAZ	70 HBC	0. PBAR P, 4 PI	5/70
M	1070.0	20.0	REYNOLDS	70 HBC	2.26-2.36 PI- P	1/71
M				
M	AVG 1083.3	9.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)			
M	STUDENT 1083.0	8.3	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

30 ETA N WIDTH (MEV)						
W	(70.0)	DR LESS	MILLER	68 HBC	4.0 PI- P	9/68
W	(25.0)	DR LESS	WHITEHEAD	68 ASPK	3.1-3.6 PI-P	10/67
W	150.0	100.0	OH	69 HBC	7.PI- P,PI+ D	9/69
W	(80.0)		CLAYTON	70 HBC	2.5 PBAR P, 4 PI	1/71
W	(80.0)		DIAZ	70 HBC	0. PBAR P, 4 PI	5/70
W	85.0	35.0	REYNOLDS	70 HBC	2.26-2.36 PI- P	1/71
W				
W	AVG 98.0	31.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W	STUDENT 97.7	34.5	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

REFERENCES FOR ETA N

MILLER 68 PRL 21 1489 +GUTAY,JOHNSON,KENNEY+ (PURDUE+NDAM+SLAC)
 WHITEHEA 68 NC 53 A 817 C.WHITEHEAD+ (AERE+SHMP+LOIC)
 OH 69 PRL 23 331 +WALKER,CARROLL,FIREBAUGH,+ (WISC+TNT0)
 BATON 70 PL 33 B 528 +LAURENS,ZEINIER (SACLAY)
 CLAYTON 70 NP 8 22 85 +NASON,MUIRHEAD,RIGOPOULOS,+ (LIVP+ATEN)
 DIAZ 70 NP 8 16 239 +GAVILLET,LABROSSE,MCNTANET,+ (CERN+CDEF)
 REYNOLDS 70 NP 8 21 77 +ALBRIGHT,BRADLEY,+ (OHIO+FSU+MNN+COLO)
 WHITEHEA 72 NP 8 48 365 WHITEHEAD,AULD,+ (AERE+RHEL+SHMP+LOUC)
 DIAMOND 73 PR D 7 1977 +BINKLEY,+ (WISC+DUKE+COLD+TNT0+OHIO)

$A_1(1100)$

10 A1(1100, JPC=1+-) I=1

The $A_1 \rightarrow \rho\pi$ bump has been mainly observed in the diffraction-like process $\pi N \rightarrow (\pi\pi\pi)N$ without quantum number exchange and at small momentum transfer. There are also observations of structure in the A_1 mass region in reactions where additional mesons are produced, in backward production (see however ABASHIAN 75), and in $\bar{p}p$ annihilations (see the Data Card Listings).

There is a special category of "diffractive mesons" consisting of statistically significant peaks like A_1 , A_3 , Q , and L (see the corresponding mini-reviews), which are not far above the respective $\rho\pi$, $f\pi$, $K^*(892)\pi$, and $K^*(1420)\pi$ thresholds. Because the behavior near threshold in these channels may be described by the Deck effect, a resonance interpretation is questionable. These peaks are included in the Meson Table, but we do not mean to imply that they are necessarily genuine resonances.

Other meson systems such as $K\omega$, $K\phi$, $K^*\bar{K}$ also exhibit low mass peaks when produced in reactions which may proceed without quantum number exchange (see for example DAVIS 72, CHUNG 74, THEOCHAROPOULOS 74, CARNEY 75, OTTER 75).

Partial-wave analyses of multi-meson systems in reactions like $\pi N \rightarrow (\pi\pi\pi)N$ are becoming available (ASCOLI 70). Several important assumptions are made in such analyses (see HANSEN 74, HERNDON 75 for detailed discussions), amongst which:

- (a) for a given t , the 3π vertex is independent of the NN vertex;
- (b) the 3π decay proceeds through quasi-two-body states ($\rho\pi, \epsilon\pi, \dots$) in the spirit of the isobar model (a partial-wave analysis with unitarity corrections to the isobar model is applied to the 3π system in ASCOLI 75; see also AARON 75, AITCHISON 75, and references therein).

The dominant effect in the A_1 mass region is a broad $J^P=1^+$ $\rho\pi$ S-wave enhancement, with a maximum intensity at ~ 1100 MeV and a width ~ 300 MeV (ANTIPOV1 73, KRUSE 74, OTTER 74, TABAK 74, THOMPSON 74, EMMS1 75). The phase of the $J^P=1^+$ wave shows little variation relative to various other "background" waves (see however the A_3 and Q mini-reviews). In contrast the A_2 peak has been confirmed as a resonance with a Breit-Wigner-like phase change of the $J^P=2^+$ partial wave (ASCOLI 70, ANTIPOV1 73, OTTER 74, TABAK 74, THOMPSON 74).

These results suggest that at most a small part of the A_1 enhancement corresponds to a $J^P=1^+$ resonance. Indeed ASCOLI 73 and ASCOLI 74 show that the partial-wave structure of the 3π

Data Card Listings

For notation, see key at front of Listings.

Mesons

$A_1(1100)$

system is qualitatively reproduced by a Reggeized pion exchange model.

BOWLER 75 attempt to explain the lack of variation of the 1^+ phase by allowing for a phase difference between the Deck amplitude and a direct " A_1 resonance" production amplitude. Good fits to the data are obtained, but the A_1 is predicted to have $M \sim 1300$ MeV, $\Gamma \sim 250$ MeV.

Recent analyses of the $(3\pi)^0$ system produced by charge exchange find no evidence for A_1 production. WAGNER 75 study the reaction $\pi^+ p \rightarrow \pi^+ \pi^- \pi^0 \Delta^{++}$ at 7 GeV/c and are able to set an upper limit of 2 μ b on the A_1 production cross section if $\Gamma < 150$ MeV. Analyzing $\pi^+ n \rightarrow \pi^+ \pi^- \pi^0 p$ at 4 GeV/c, EMMS2 75 observe no A_1 with $M \sim 1100$ MeV, but cannot rule out resonance production above 1300 MeV.

References Not Included in the Data Card Listings

- R. Aaron et al., Phys. Rev. D12, 1984 (1975).
- I. J. R. Aitchison and R. J. A. Golding, Phys. Lett. 59B, 288 (1975).
- G. Ascoli and H. W. Wyld, Phys. Rev. D12, 43 (1975).
- J. N. Carney et al., Phys. Lett. 55B, 117 (1975).
- S. U. Chung et al., Phys. Lett. 51B, 412 (1974).
- P. J. Davis et al., Nucl. Phys. B44, 344 (1972).
- J. D. Hansen et al., Nucl. Phys. B81, 403 (1974).
- D. J. Herndon et al., Phys. Rev. D11, 3165 (1975).
- G. Otter et al., Nucl. Phys. B87, 189 (1975).
- G. Otter et al., Nucl. Phys. B89, 201 (1975).
- G. Otter et al., Nucl. Phys. B96, 365 (1975).
- P. Theocharopoulos et al., Nucl. Phys. B83, 1 (1974).

10 A1 MASS (MEV)	
M	PRODUCED BY PI +
M	(1080.0)
M	(1080.) APPROX.
M	(1040.0)
M	(1128.) (8.)
M	PRODUCED BY PI -
M	(1060.)
M	(1089.0) (12.0)
M	(1090.) APPROX.
M	(1055.0) (6.0)
M	(1119.) (30.)
M	S SHOULDER ON A2 ONLY
M	(1069.0) (7.0)
M	(1120.0)
M	F T (1150.)
M	T MASS AND WIDTH SEEN TO DEPEND ON T, UNIQUE DET. IMPOSSIBLE

M	PRODUCED BY PIONS, BACKWARDS SCATT.	ANDERSON	69 MMS	- 16 PI- P, BACKW9	8/69
M	(1115.0) (20.0)				
M	PRODUCED BY PBARS, SEE TYPED NOTE.	DANYSZ	67 HBC	+ 3.3-6 PBAR P	7/67
M	(1054.) (7.)	FRIDMAN	68 HBC	+ 5.7 PBAR P	6/68
M	(1042.) (21.)	ATHERTON	73 HBC	+ 5.7 PBAR P	1/74
M	A (1076.) (5.)				
M	A JP ANALYSIS GIVES SOME EVIDENCE FOR RHO PI D-WAVE				1/73
M	PRODUCED BY K-, SEE TYPED NOTE.	ALLISON	67 HBC	+ 6 K-P, LAM +5 PI	1/68
M	(1111.) (10.)	ALLISON	67 HBC	+ 6 K-P, LAM +4 PI	1/68
M	(1117.) (35.)	JUHALA	67 HBC	0 4.6-5 K-P, 5BCDY	1/68
M	(1060.) (15.)				
M	PRODUCED BY K+, SEE TYPED NOTE.	ALEXANDER	69 HBC	+ 9 K+P	9/69
M	(1060.0) (20.0)	BERLINGHI	69 HBC	+ 0 12.7 K+ P	9/69
M	K+ (1030.0) (20.0)				
M	K+ FOR CONTRADICTIONARY EVIDENCE SEE RABIN 70.				
M	F FROM A FIT TO JP=1+ RHO PI PARTIAL WAVE				
M	AVERAGING NOT MEANINGFUL				

10 A1 WIDTH (MEV)

10 A1 PARTIAL DECAY MODES	
W	PRODUCED BY PIONS, RESONANCE INTERP. CONFUSED BY DECK EFFECT
W	PRODUCED BY PI +
W	(80.0)
W	(130.) APPROX.
W	(50.0) OR LESS
W	F (300.) APPROX.
W	F (367.) (30.)
W	PRODUCED BY PI -
W	(140.0) (31.0)
W	(125.) APPROX.
W	(77.0) (17.0)
W	K (76.) (46.)
W	K SHOULDER ON A2 ONLY
W	(99.0) (15.0)
W	F T (300.)
W	T MASS AND WIDTH SEEN TO DEPEND ON T, UNIQUE DET. IMPOSSIBLE
W	PRODUCED BY PIONS, BACKWARDS SCATT.
W	(98.0) (45.0) (20.0)
W	PRODUCED BY PBARS, SEE TYPED NOTE.
W	(33.) (19.)
W	(130.) APPROX.
W	A (36.) (20.) (15.)
W	A JP ANALYSIS GIVES SOME EVIDENCE FOR RHO PI D-WAVE
W	PRODUCED BY K-, SEE TYPED NOTE.
W	(50.) (50.)
W	(50.) (25.)
W	(120.) (15.)
W	PRODUCED BY K+, SEE TYPED NOTE.
W	(160.0) (20.0)
W	B (120.0) (30.0)
W	K+ FOR CONTRADICTIONARY EVIDENCE SEE RABIN 70.
W	(130.0) (20.0)
W	F FROM A FIT TO JP=1+ RHO PI PARTIAL WAVE
W	AVERAGING NOT MEANINGFUL

10 A1 PARTIAL DECAY MODES

10 A1 BRANCHING RATIOS		
P1	A1 INTO RHO PI	DECAY MASSES
P2	A1 INTO KBAR K	772+ 139
		493+ 497

10 A1 BRANCHING RATIOS

R1	A1 INTO (KBAR K)/(RHO PI)	DAHL	67 HBC	(P2)/(P1)	
R1	(0.0025) OR LESS			- 4.0 PI- P	.10/66

REFERENCES FOR A1

BELLINI	63 NC 29 896	BELLINI, FIORINI, HERZ, NEGRI, RATTI (MILAN)
ADERHOLZ	64 PL 10 226	AACH+BERL+BIHM+BONN+DESY+HAMBURG+LOIC+MPIN
GOLDBABE	64 PRL 12 336	GOLDBABER, BROWN, KADYK, SHEN+ (LRL+SLAC)
LANDER	64 PRL 13 346 A	LANDER, ABOLINS, CARMONY, HENDRICKS + (UCSD) JP
ABOLINS	65 ATHENS(OHIO)CONF.	+CARMONY, LANDER, XUONG, YAGER (LA JOLLA)I=1
ALITTI	65 PL 15 69	ALITTI, BATON, DELER, CRUSSARD+ (SACL+BGNA)
ALLARD	66 NC 46A 737	+DRIJARD+HENNESSY+ (ORSAY+MILAN+SACL+JUG)
DEUTSCHM	66 PL 20 82	DEUTSCHMANN, STEINBERG + (AACH+BERLIN+CERN)
HESS	66 UCRL-16832	R I FESS (THEISIS, BERKELEY)
ALLISON	67 PL 25B 619	+CRUZ+ (OXF+MPIH+BIHM+RHEL+GLAS+LOIC)
DAHL	67 PR 163 1377	+HARDY+HESS+KIRZ+MILLER (LFL)
DANYSZ	67 NC 51 A 801	DANYSZ+FRENCH+SIMAK (CERN)
JUHALA	67 PRL 19 1355	+LEACOCK+RHODE+KOPPELMAN+ (TOWA+GOLO)
SLATTERY	67 NC 50A 377	+KRAYBILL+FORMAN+FERBEL (YALE+ROCK) JP
ARMENISE	68 PL 26 B 336	+FORING+CARTACCI+ (BARI+BGNA+FRIZ+ORSAY)
ASCOLI	68 PRL 21 113	+CRAWLEY, KRUSE, MORTARA, SCHAFFER+ (ILLINOIS)
BALLAM	68 PRL 21 934	+BRODY, CHADWICK, FRIES, GUIRAGOSSIAN+ (SACL)JP
BOESEBECK	68 NP 8 501	BOESEBECK, DEUTSCHMANN+ (AACHEN+BERLIN+CERN)
CASO	68 NC 54 A 983	+CCNTE+CORDS+DIJAZ+ (GENOVA+HAMB+MILA+SACL)
CHUNG	68 PR 165 1491	S.U.CHUNG, O.DAHL, J.KIRZ, D.H.MILLER (LRL)
CNDPS	68 PRL 21 1609	+HOUGH, COHN, BUGG+ (BNL+ORNL+UCND+TENN+PENN)
FRIDMAN	68 PR 167 1268	+MAURER, HICHALON, OUDET+ (HEID+STRASBOURG)
JUNKMANN	68 NP 88 471	+COCCONI+ (AACH+BERL+BGNN+CERN+WARS)
KEY	68 PR 166 1430	+PRENTICE+COOPER+MANNER+ (TNTD+ANL+WISCI)
ALEXANDE	69 PR 183 1168	G.A.ALEXANDER, A.FRESTOCNE, G.GOLDBABER (LRL)
ALLAY	69 PL 29B 198	+BINON+DIDDENS+DUTELL+KLOVNING+... (CERN)
ANDERSON	69 PRL 22 1390	+COLLINS+ (BNL+CARN)
BERLINGH	69 PRL 23 42	BERLINGHIERT, FARBET, + (ROCK)
DONALD	69 NP 8 11 551	+EDWARDS, BURAN, BETTINI, + (LIVP+OSLO+PADO)

Mesons

A₁(1100), M(1150), A_{1,5}(1170), ε(1200)

Data Card Listings

For notation, see key at front of Listings.

Table listing meson properties and references. Includes entries for FAYOLLE, JUHALA, KENYON, ARMENISE, ASCOLI, BRANDENB, CASO, CRENELL, CHIEN1, CHIEN2, GARELTICK, RABIN, SHIH, ASCOLI, BEMPRAD, BERGER, RINAUDO, BRENVI, BLOODWDR, DIEBOLD, LAMSA, MORSE, ANTIPV1, ANTIPV2, ARNOLD, ASCOLI, ASCOLI, ATHERTON, READ, BOWLER, ASCOLI, KRUSE, LICHTMAN, OTTER, TABAK, THOMPSON, THOMPSON, ABASHIAN, BEUSCH, BOWLER, DIAZ, EMMS, EMMS, HORNE, KANE, WAGNER.

M(1150) -> MM

68 M(1150) EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

Table with 7 columns: M, N, 65, 1148.3, 3.3, JACOBEL, 72, MMS, 0, 2.4, PI-, P, N, MM, 12/72, 12/75*

Table with 7 columns: M, 65, 15.0, 9.0, 11.7, JACOBEL, 72, MMS, 0, 2.4, PI-, P, N, MM, 12/72

Table with 2 columns: JACOBEL, 72, PRL, 29, 671; BUTTRAM, 75, PRL, 35, 970

A_{1,5}(1170) -> 3π

44 A 1.5(1170, JPG = -) I=1 THIS ENTRY LISTS REFERENCES TO PEAKS OF LOW STATISTICAL SIGNIFICANCE IN THE 3 PI SYSTEM BETWEEN THE A1 AND THE A2. OMITTED FROM TABLE.

***** REFERENCES FOR A 1.5

Table listing references for A 1.5: BUTTERW, CASON, ASCOLI, DONALD, VCN, JUNKMANN, ARMENISE, GALLOWAY, MORSE.

ε(1200) 14 PI PI S WAVE, CALLED EPSILON

S-Wave ππ Interactions in the Region 280-1800 MeV

In this note, we discuss information on the isoscalar ππ S-wave in terms of its phase shift δ₀⁰, from threshold to 1800 MeV.

The threshold behavior of elastic ππ scattering involves S and P waves which can be sufficiently well described by the scattering lengths a₀⁰, a₁¹, and a₀². In spite of many attempts (see PILKUH 73, BONNIER 74, PASCUAL 74, RIESTER 75, SRINIVASAN 75), the determination of these parameters still meets with great difficulties (BASDEVANT 73, 75). The parameters a₀⁰ and a₀² are strongly correlated and must lie in a narrow band in the (a₀⁰, a₀²)-plane (MORGAN2 70). Thus if a₀⁰ is fixed, a₀² and a₁¹ are determined within small uncertainties (BASDEVANT 72, 75). However, all one knows is a region of finite extent within which a₀⁰ must lie. It is therefore not established whether, e.g., the Weinberg predictions are supported by experiment.

Near threshold the S wave shows no resonant behavior. The so-called ABC and DEF effects (BOOTH 63, HALL 69, BRODY 70, 72, BANAIGS 71, 73) occur only on nuclear targets (d, H³, He³) and move when kinematical conditions change. Thus they must be kinematical effects (DUBAL 71, BRODY2 72, RISSER 73, BAR-NIR 75, BARRY 75).

The region of elastic ππ scattering is known to extend from threshold to about 990 MeV, near the K-K̄ threshold (BATON 70, CARROLL 72, PROTOPODESCU 73, HYAMS 73, OCHS 73).

Up to the ρ meson mass region, δ₀⁰ is (qualitatively) uniquely determined; it rises monotonically and reaches a value of 60° to 70° near 700 MeV (SONDEREGGER 69, BATON 70, BAILLON 72, CARROLL 72, FRENKIEL 72, GAIDOS 72, PROTOPODESCU 73, HYAMS 73, OCHS 73, ENGLER 74, ESTABROOKS 74, 75, GRAYER 74); see Fig. 1.

In the mass region of 700 to 900 MeV, all energy-independent phase-shift analyses using the constraint η₀⁰ = 1 find two solutions ("up-down ambiguity"). This ambiguity was resolved in favor of the "down" solution by the observation of a very rapid decrease in the S-wave amplitude

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\epsilon(1200)$

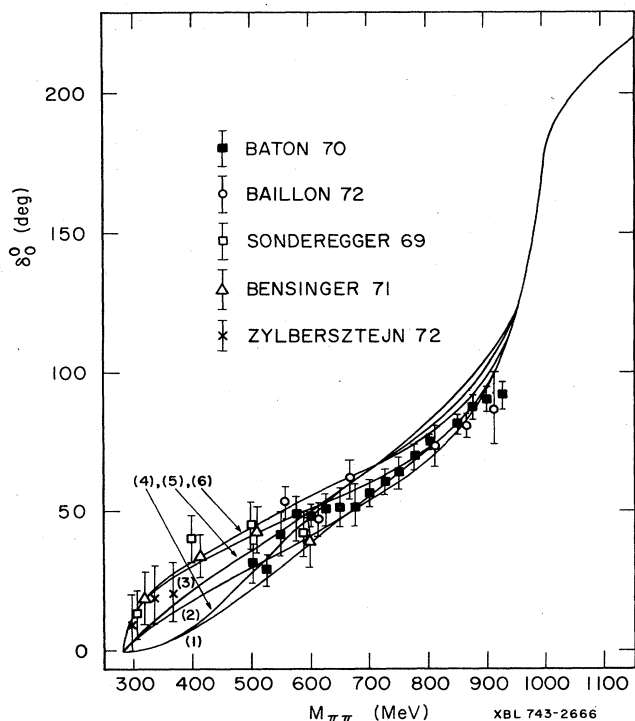


Fig. 1. The S-wave, $I=0$ $\pi\pi$ phase shift, as determined by several experiments. The curves are from BASDEVANT 73; solutions (1-3) were obtained by fitting the data of BATON 70, while (4-6) are from fits to the data of PROTOPOPESCU 73.

between 950 and 980 MeV (FLATTE 72, GAIDOS 72, HYAMS 73, BINNIE 73, ENGLER 74). The size of the observed drop corresponds to a change from nearly the unitarity limit to zero; see Fig. 2.

Independent evidence for the correctness of this "down" solution comes from studies of the $\pi^0\pi^0$ system (APEL 72, BRAUN 73, SKUJA 73, RIESTER 75). They observe a wide $\pi^0\pi^0$ enhancement of ~ 800 MeV which is much better described by the "down" solution than the "up" solution. Furthermore, indirect information from elastic $\pi\pi$ scattering in the crossed channel (ELVEKJAER 72, NIELSEN 70, 72) is compatible with the "down", but not with the "up", solution.

The ambiguities of the phase-shift solutions stem from the fact that there are more helicity amplitudes than observables. Thus in the absence of polarization measurements one is obliged to make some supplementary assumptions (see, e.g., DONOHUE 75). Analyzing the same data (GRAYER 74) by different methods, HYAMS 75 find four solutions in the region 1.0-1.8 GeV, ESTABROOKS 74 find eight solutions above 1.2 GeV and four solutions below, whereas FROGGATT 75 find arguments to favor one of these solutions [solution B of ESTABROOKS 74 (see Fig. 3), essentially the same as HYAMS 75 solution +--].

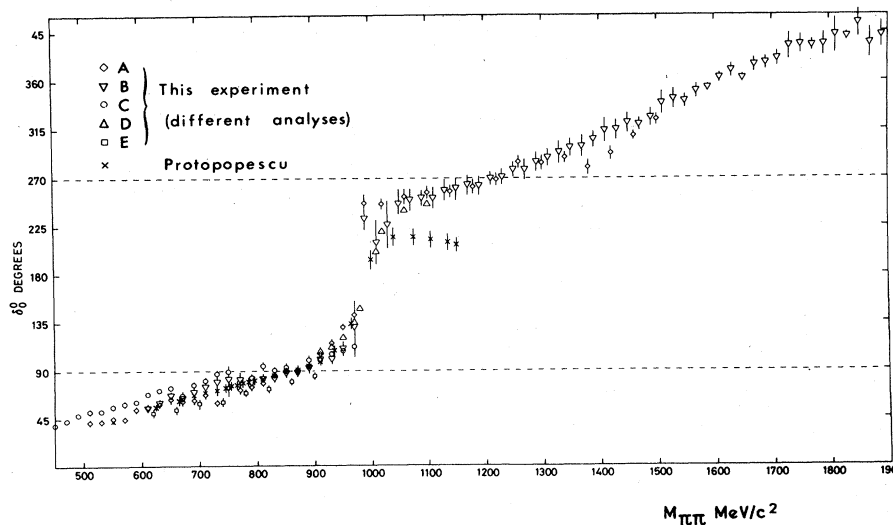


Fig. 2. The S-wave, $I=0$ $\pi\pi$ phase shift, as determined by various analyses of the data of GRAYER 74, compared with the previous results of PROTOPOPESCU 73. (Figure from GRAYER 74.)

Mesons

$\epsilon(1200)$

Data Card Listings

For notation, see key at front of Listings.

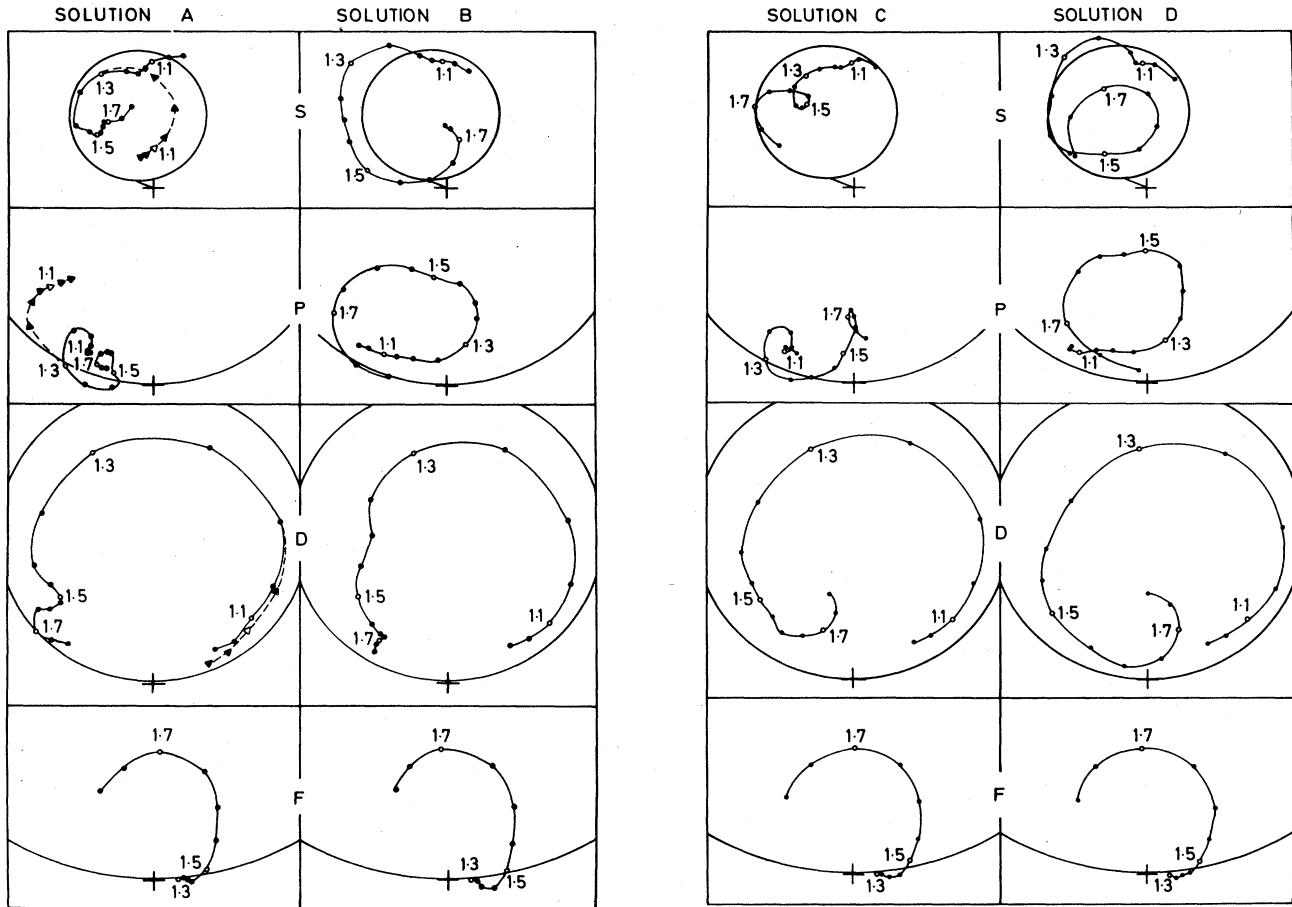


Fig. 3. Argand plots for $\pi\pi$ S, P, D, and F partial-wave amplitudes. Shown are solutions A-D of ESTABROOKS 75, obtained with the data of GRAYER 74. The points on the curves are at 50 MeV intervals. The diameters of the S-, P-, D-, and F-wave unitarity circles shown are given by $2/3$, $\sqrt{3}$, $(2/3)\sqrt{5}$, and $\sqrt{7}$, respectively. An elastic $I = 2$ S-wave phase shift of 25° independent of $M_{\pi\pi}$ has been assumed. (Figure from ESTABROOKS 75.)

Near 1.2 GeV all δ_0^0 solutions based on the GRAYER 74 data, as well as independent analyses of other experiments (CARROLL 72, 74, ENGLER 75), exhibit $\delta_0^0 = 270^\circ$.

It has now been made plausible (MORGAN 75) that all available data are compatible with the existence of just two poles: the $S^*(993)$ connected with the rapid variation of δ_0^0 near the $K\bar{K}$ threshold, and the $\epsilon(1200)$. The $S^*(993)$ is also responsible for the large $K\bar{K}$ $I=0$ S-wave scattering length.

The $\epsilon(1200)$ "replaces" the pole previously listed below 600 MeV. Note that although many analyses in the past found solutions with an

$\epsilon(600)$ pole, they also found solutions without such a pole (PROTOPODESCU 73, HYAMS 73).

Thus we have just the right number of isoscalar S-wave resonances to make up an SU(3) nonet, presumably together with the $\delta(970)$ and the $\kappa(1250)$ mesons (MORGAN 75).

Note that, although there is general consensus about the position of the S^* pole on the second sheet, it is not clear whether it is accompanied by a companion pole on the third sheet or not (FUJII 75, MORGAN 75). The contrast between the width of 180 MeV of MORGAN 75, and the value of 40 MeV obtained as twice the imaginary part, reflects just this uncertainty.

Data Card Listings

For notation, see key at front of Listings.

Mesons

ε(1200), B(1235)

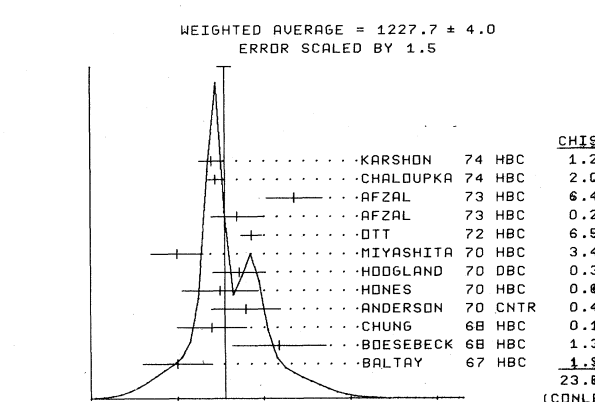
Table listing meson properties including mass (MEV), width (MEV), and various references for particles like ε(1200) and B(1235).

Table listing meson properties for particles like ANJOS, BANAGS1, BASDEVAN, etc., including mass and width.

***** REFERENCES FOR EPSILON *****

B(1235) 11 B(1235, JPC=1+-) I=1

Table listing meson properties for B(1235) including mass (MEV), width (MEV), and various references.



Mesons

B(1235), rho'(1250), f(1270)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various fit parameters. Includes sub-sections for B WIDTH (MEV), PARTIAL DECAY MODES, and BRANCHING RATIOS.

Table listing experimental data for rho'(1250) and f(1270), including sources like POLS, WERBRUGC, and AFZAL, and their respective parameters.

Table listing experimental data for B(1235), including sources like ABOLINS, GOLDHABER, and BALTAY, and their respective parameters.

Table listing experimental data for rho(1250) and f(1270), including sources like FRENKIEL, CHUNG, and BRANON, and their respective parameters.

Table listing references for B(1235) from various sources like ABOLINS, BONDAR, and ANDERSON.

Table listing references for rho(1250) and f(1270) from various sources like ANDERSON, PODOBSKY, and ALLES.

Data Card Listings

For notation, see key at front of Listings.

Mesons

f(1270)

Table with columns for particle name, mass, width, and other properties. Includes entries for SELOVE, JACOBS, RABIN, etc.

WEIGHTED AVERAGE = 181.1 ± 5.6
ERROR SCALED BY 1.7

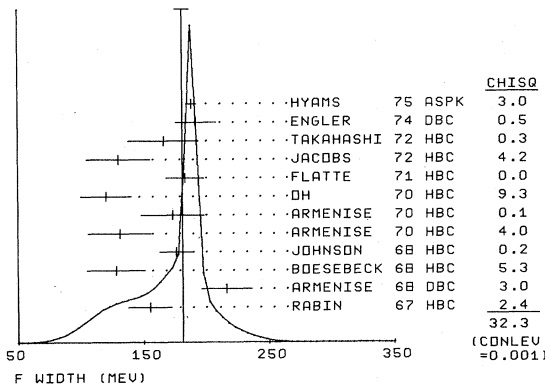


Table with columns for decay modes (P1, P2, P3, P4, P5, P6, P7) and decay masses.

Table with columns for branching ratios (R10, R11, R12, R13, R14, R15, R16, R17, R18, R19, R20) and partial wave amplitudes.

Table with columns for F INTO (2P1+2P2)/(PI PI) and F INTO (PI+PI-2P10)/(PI PI) with associated decay constants and errors.

Table with columns for F INTO (K KBAR)/(PI PI) and F INTO (ETA PI)/(PI PI) with associated decay constants and errors.

Table with columns for REFERENCES FOR F, listing various researchers and their contributions to the field.

Mesons
D(1285), A2(1310)

Data Card Listings

For notation, see key at front of Listings.

D(1285)

8 D(1285, JP=C= +) I=0
(JP=0-+1+2- WITH 1+ FAVORED.)

Table with columns: M, (1290.), APPROX., BARLOW, 67 HBC, 1.2 PBAR P, 4 PFS, 5/67. Includes student average and error.

Table with columns: W, R, (35.0), (10.0), DAHL, 67 HBC, 1.6-4.2 PI-P, 11/71. Includes student average and error.

Table with columns: P1, D INTO K KBAR PI, 497+ 497+ 134. Includes decay masses.

Table with columns: R1, D INTO (PI PI RHO) / (K KBAR PI), (P2)/(P1), 10/66. Includes branching ratios.

Table with columns: R4, D INTO (2PI+ 2PI- (INCL. RHO PI PI))/(ETA PI+ PI-), (P5)/(2/3P), 11/71. Includes partial decay modes.

REFERENCES FOR D
+BARLOW, ADAMSON, + (CDEF+ CERN+ IRAD+ LVP)
+CHUNG, DAHL, HESS, HARDY, KIRZ, + (LRL+UCB)

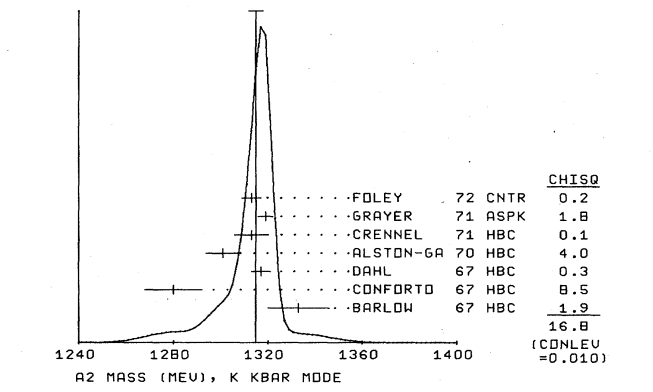
A2(1310)

12 A2(1310, JP=G=2+-) I=1
WE LIST THE A2 AS AN ORDINARY BREIT-WIGNER RESONANCE. FOR DISCUSSION OF THE REPORTED SPLITTING SEE CUR APRIL 72 AND APRIL 73 EDITION.

Table with columns: M, (1320.0), (13.0), ADERHOLZ, 64 HBC, 4.0 PI+P, 12/75*. Includes student average and error.

Table with columns: MK, 80(1317.0), (3.0), BARLOW, 67 HBC, 1.2 PBAR P, KK, 2/72. Includes branching ratios.

WEIGHTED AVERAGE = 1315.0 +/- 3.1
ERRA CALCD BY 1.7



Data Card Listings

For notation, see key at front of Listings.

Mesons

A₂(1310)

Table with columns for mass (MEV), error (ETA PI MODE), and particle identification. Includes entries for ALSTON-GA, CASO, DZIERBA, JOHNSTON, ESPIGAT, KEY, CONFORTO, and MMS.

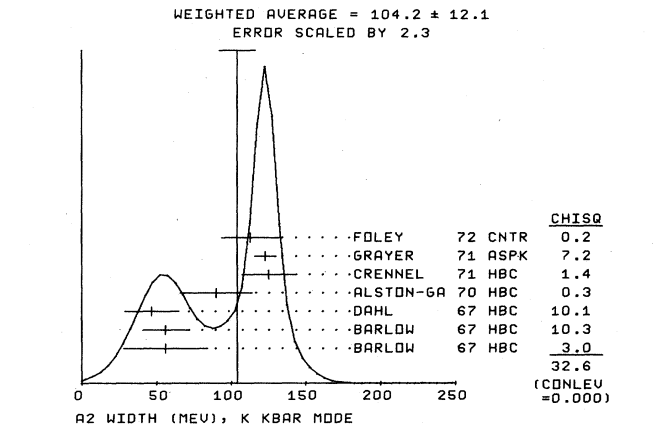
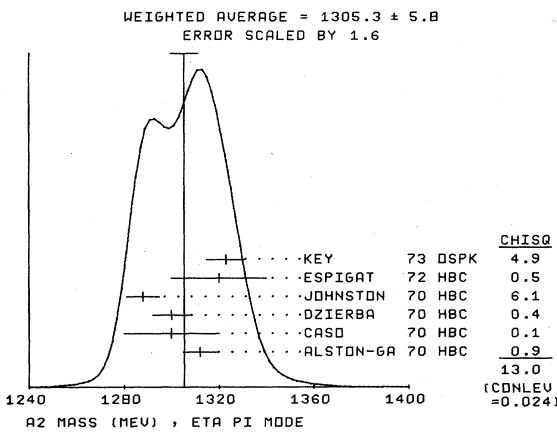


Table with columns for A2 WIDTH (MEV), 3PI MODE, and particle identification. Includes entries for ADERHOLZ, GOLDHABER, LEBEVRIS, BARNES, BENSON, LEVRAT, CHIKOVANI, ARNEISE, BOESEBECK, CHUNG, VON KROGH, GORDON, ANDERSON, ARMENISE, EISENBERG, ALSTON-GA, BOCKMANN, DIAZ, GARFINKEL, GORDON, BARNHAM, BINNIE, BINNIE, BOWEN, BOWEN, BOWEN, BLOODWORTH, ANTIPOVI, CHALUPKA, EMMS, WAGNER.

Table with columns for A2 WIDTH (MEV), ETA PI MODE, and particle identification. Includes entries for ALSTON-GA, CASO, DZIERBA, JOHNSTON, ESPIGAT, KEY, CONFORTO, and MMS.

Table with columns for A2 PARTIAL DECAY MODES and DECAY MASSES. Includes entries for A2 INTO RHO PI, A2 INTO K KBAR, A2 INTO ETA PI, A2 INTO OMEGA PI PI, A2 INTO PI+ PI- PI0 EXCL. RHO PI, A2 INTO PI+ PI- PI- EXCL. RHO PI, A2 INTO PI GAMMA, A2 INTO ETA PRIME PI, and SMALL, NOT USED IN THE FIT.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS. The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± 5P_i, where 5P_i = √(5P_i5P_i), while the off-diagonal elements are the normalized correlation coefficients (5P_i5P_j)/(5P_i5P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns for A2 WIDTH (MEV), K KBAR MODE, and particle identification. Includes entries for BARLOW, BEUSCH, CONFORTO, DAHL, DAHL, CRENNEL, ALSTON-GA, CRENNEL, GRAYER, FOLEY, and THE NEUTRAL MODE CAN INTERFERE WITH THE F MESON.

Table with columns for A2 BRANCHING RATIOS and particle identification. Includes entries for A2 (CHARGED ONLY) INTO (K KBAR)/(RHO PI), THE NEUTRAL MODE CAN INTERFERE WITH F, A2 INTO (ETA PI)/(RHO PI + K KBAR + ETA PI), and FROM FIT.

Mesons

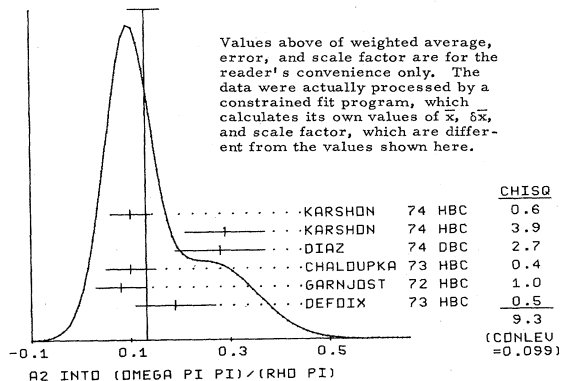
A₂(1310)

Data Card Listings

For notation, see key at front of Listings.

R3	A2 INTO (ETA PI) / (RHO PI)	(P3)/(P1)		
R3	0.3	0.2	ADERHOLZ 64 HBC	4.0 PI+P
R3	0.22	0.09	CONTE 67 HBC	11.0 PI-P
R3	0.25	0.08	ASGOLI 68 HBC	5 PI-P
R3	0.12	0.08	CHUNG 68 HBC	3.2 PI-P
R3	0.16	0.10	KEY 68 HBC	3 PI-P
R3	0.3	0.13	ABRAMVIC 70 HBC	3.93 PI-P
R3	0.18	0.06	VETLITSKY 69 HBC	3.3 PI-P
R3	0.25	0.09	BOCKMANN 70 HBC	5.0 PI+P
R3	0.34	0.17	BOCKMANN 70 HBC	5.0 PI+P
R3	0.39	0.07	DZIERBA 70 HBC	8. PI-P
R3	0.246	0.042	ALSTON-GA 71 HBC	7.0 PI+P
R3	0.211	0.044	CHALDUPKA 73 HBC	3.9 PI-P, P A2
R3	0.22	0.05	ANTIPOV 73 CNTR	40. PI-P, P A2-
R3	0.221	0.021	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R3	0.222	0.023	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R3	0.212	0.015	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R4	A2 INTO (ETA PRIME PI) / TOTAL	(P8)		
R4	0.1	OR LESS	CHUNG 65 HBC	3.2 PI-P
R4	0.02	OR LESS	BARNHAM 71 HBC	3.7 PI+P
R5	A2 INTO (ETA PRIME PI)/(RHO PI)	(P8)/(P1)		
R5	0.04	0.03	BOCKMANN 70 HBC	5.0 PI+P
R5	0.15	0.09	DZIERBA 70 HBC	8. PI-P
R5	0.04	OR LESS	ALSTON-GA 71 HBC	7.0 PI+P
R5	0.011	OR LESS	EISENSTEIN 73 HBC	5.1 PI-P, P 6P1
R6	A2 INTO (PI+ PI- P1) / (RHO PI)	(P5)/(P1)		
R6	0.17	OR LESS	BENSON 66 HBC	3.7 PI+D
R7	A2 INTO (ETA PI)/(K KEAR)	(P3)/(P2)		
R7	0.3	OR LESS	FOSTER 68 HBC	PBAR P,PBA REST.
R7	E SUPERSEDED BY ESPIGAT 72 (SEE UNDER R2 AND R8)			
R7	3.19	0.46	FROM FIT	
R8	A2 INTO (K KBAR)/(RHO PI + K KBAR + ETA PI)	(P2)/(P1+P2+P3)		
R8	0.06	0.03	BARNHAM 71 HBC	3.7 PI+P, KSK+P
R8	0.020	0.004	ESPIGAT 72 HBC	0. PBAR, P
R8	NOT AVERAGED BECAUSE OF DISCREPANCY BETWEEN MASSES			
R8	FROM (K KBAR) AND (RHO PI) MODES			
R8	0.03	0.02	DAMERI 72 HBC	11. PI-P
R8	0.05	0.02	TOET 73 HBC	5. PI+P, P K+ KO
R8	0.09	0.04	TOET 73 HBC	5.1 PI+P, P K K B
R8	0.048	0.012	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R8	0.048	0.014	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R8	0.0520	0.0058	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	
R9	A2 INTO (PI+ PI- PI-)/(RHO PI-)	(P6)/(P1C)		
R9	0.23	OR LESS	CL=90	ABRAMVIC 70 HBC
R11	A2 INTO (PI GAMMA)/TOTAL	(P7)		
R11	0.005	0.005	[0.003]EISENBERG 72 HBC	
R11	PION EXCHANGE MODEL USED IN THIS ESTIMATION			
R12	A2 INTO (OMEGA PI PI)/(RHO PI)	(P4)/(P1)		
R12	0.19	0.08	DEFOIX 73 HBC	0.7 PBAR P, 7 P1
R12	DECAYS TO B1(1040) PI, B1 INTO OMEGA PI			
R12	D ERROR INCREASED TO ACCOUNT FOR POSSIBLE SYST. ERRORS			
R12	D OF COMPLICATED ANALYSIS.			
R12	0.08	0.05	GARNJOST 72 HBC	0.7 PI+P
R12	0.10	0.05	CHALDUPKA 73 HBC	3.9 PI-P, P A2
R12	0.28	0.09	DIAZ 74 HBC	6. PI+P, P(SPI)0
R12	0.29	0.08	KARSHON 74 HBC	4.9 PI+P, DEL+ A2
R12	0.10	0.04	KARSHON 74 HBC	4.9 PI+P, P A2
R12	K KARSHON 74 SUGGEST AN ADDITIONAL I=0 STATE, STRONGLY COUPLED			
R12	K TO OMEGA PI PI COULD EXPLAIN DISCREPANCIES IN BRANCHING RATIOS			
R12	K AND MASSES.			
R12	0.131	0.032	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
R12	0.127	0.027	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R12	0.131	0.029	FRM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)	
R12	(SEE IDEOGRAM BELOW)			

WEIGHTED AVERAGE = 0.131 ± 0.032
ERROR SCALED BY 1.4



ADERHCLZ 64 PL 10 248	(AACHEN+BERLIN+ERIM+BCNN+AMBURG+LOIC+MPJM)
CHUNG 64 PRL 12 621	+DAHL+HARDY+HESS+KALBFLEISCH+KIRZ (LRL)
GOLDBABE 64 DUBNA CONF 1 480	+GOLDBABE'S GOLDBABER+D'ALLORAN+SHEN(LRL)
LANDER 64 PRL 13 346	+ABOLINS, CARMONY, HENDRICKS, XUONG* (LA JOLLA)
ABOLINS 65 ATHENS(CHIO)CONF.	+CARMONY, LANDER, XUONG, YAGER (LA JOLLA)=1
ADERHOLZ 65 PR 138 B 897	(AACHEN+BERL+BIRM+BOHN+HAMB+LOIC+MPJM)
ALITTI 65 PL 15 69	ALITTI, BATON, DELER, CRUSSARD* (SACLAY+BGNA) JP
CHUNG 65 PRL 15 325	+CAHL+HARDY+HESS, JACOBS+KIRZ, MILLER (LRL)
FORINO 65 PL 19 68	+GESSAROLI+ (BGNA+BARI+FRIZ+ORSAY+SACL)
LEFEVRE 65 PL 19 434	CERN MISSING MASS SPECTROMETER GROUP (CERN)
SEIDLITZ 65 PRL 15 217	L SEIDLITZ, D O CAHL, D H MILLER (LRL)
BARNES 66 PRL 16 41	BARNES, FOWLER, LAI, ORENSTEIN + (BNL+CUNY)
BENSON 66 MICH COO-1112-4	G.C. BENSON, THESIS (MICH)
ALSO 66 PRL 16 1177	G BENSON, LOVELL, MARQUIT, ROE + (MICH)
EHRLTCH 66 PR 152 1194	R. EHRLTCH, W. SELOVE, H. YUTA (PENN)
FERBEL 66 PL 21 111	FERBEL (ROCHESTER)
LEVRAT 66 PL 22 714	CERN MISSING MASS SPECTROMETER GROUP (CERN)
ARMENISE 67 PL 258 53	ARMENISE, FORINO, + (BARI+BGNA+FRIZ+ORSAY)
BALTAY 67 PL 258 160	+KIRSCH+KUNG+YER+RABIN (COLU+BNL+RUTGERS)
BARLOW 67 NC 50A 701	+ILLIESTOL+MCNTANET* (CERN+COEF+IRAD+LVP)
BARTSCH 67 PL 258 48	+DEUTSCHMANN+GROTE+COCCONI* (AACH+BERL+CERN)
BEUSCH 67 PL 25 8 357	+FISCHER, GOBBI, ASTBURY* (ETHZ+CERN)
CASON 67 PRL 18 880	+LAMS, BISWAS, DERADO, GROVES, + (NOTREDAME)
CHIKOVAN 67 PL 258 44	CERN MISSING MASS SPECTROMETER GROUP (CERN)
CHUNG 67 PRL 18 100	+DAHL, HARDY, HESS, KIRZ, MILLER (LRL)
ALSO 66 UCRL-16832	RICHARD I HESS--THESIS, BERKELEY (LRL)
COHN 67 NP B1 57	+MCCULLOCH+BUGG+CONDO (ORNL+UNIV. TENN.)
CONFORTO 67 NP 83 469	+MARECHAL, MCNTANET* (CERN+ANL+LVP)
CONTE 67 NC 51 A 175	+TOMASINI, CORDA, GENOVA+HAMB+MILANO+SACLAY)
DAHL 67 PR 163 1377	+HARDY+HESS+KIRZ+MILLER (LRL)
DANYSZ 67 NC 51 A 801	DANYSZ+FRENCH+SIMAK (CERN)
SLATTERY 67 NC 50A 377	+KRAYBILL+FORMAN+FERBEL (YALE+ROCH) JP
ARMENISE 68 PL 268 336	ARMENISE, FORINO, + (BARI+BGNA+FRIZ+ORSAY)
ASCOLI 68 PRL 20 1321	+CRAWLEY, MORTARA, SHAPIRO, BRIDGES+ILLINOIS) JP
BALLAM 68 PRL 21 934	+BRODY, CHADWICK, FRIES, GUIRGOSSIAN* (SLAC)
BENI 68 PL 28 B 233	CERN MISSING MASS SPECTROMETER GROUP (CERN)
BOESEBECK 68 NP B 501	BOESEBECK, DEUTSCHMANN, + (AACHEN+BERLIN+CERN)
CASO 68 NC 54 A 983	+CONTE+CORDS+DIAZ* (GENOVA+HAMB+MILAN+SACL)
CHUNG 68 PR 165 1491	S.U. CHUNG, O. DAHL, J. KIRZ, D. H. MILLER (LRL)
CRENNELL 68 PRL 20 1318	+KARSHON+KWAN LAI, SCARR, SKILLICORN (BNL)
DONALD 68 PL 26 B 327	+FROESE+BETTINI* (LIVERPOOL+OSLO+PAU)
FOSTER 68 NP B 8 174	+GAVILLET, LABROSSE, MCNTANET* (CERN+COEF)
FRIDMAN 68 PR 167 1268	+MAURER, MICHALON, OUDET+HEID +STRASBOURG)
JUNKMANN 68 NP 88 471	+COCCONI, + (AACH+BERL+BOHN+CERN+WAR)
KEY 68 PR 166 1430	+PRENTICE+COOPER+MANNERY* (INTO+ANL+LVP)
LAMSA 68 PR 166 1395	+CASON+BISWAS+DERADO+GROVES* (NOTREDAME)
VGN KRGG 68 PL 27 B 253	+MIYASHITA, KOPPELMAN, MARSHALL LIBBY (COLO)
ADERHOLZ 69 NP B 11 259	+BARTSCH, + (AACH+BERL+CERN+FRIZ+ORSAY)
AGUILAR 69 PL 29 B 62	+BARLOW, JACOBS, DELLA NEGRA* (CERN+COEF+LVP)
AGUILAR 269 PL 29 B 241	M. AGUILAR-BENITEZ, J. BARLOW, + (CERN+COEF)
ANDERSON 69 PRL 22 1390	+COLLINS, + (BNL+CERN)
ARMENISE 69 LNC 2 501	+G. IDINI, FORINO, CARTACCI* (BARI+BGNA+FRIZ)
CHIKOVAN 69 PL 28 B 526	CERN MISSING MASS SPECTROMETER GROUP (CERN) JP
CRENNELL 69 PRL 22 1327	+KARSHON+KWAN WU LAI, + (BNL) JP
DONALD 69 NP B 12 325	+EDWARDS, FOSTER, MOORE (LIVERPOOL)
EISENBERG 69 PRL 23 1322	EISENBERG, HABER, BALLAM, CHADWICK+ (REHO+SLAC)
VETLITSKY 69 SJNP 9 596	VETLITSKY, GRI GOREYEV, GRISHIN, + (ITEP)
ABRAMOVICI 70 NP B 29 466	ABRAMOVICH, BLUMENFELD, BRUYANT, + (CERN) JP
ALSTON-GATO 70 PL 33 B 607	+BARBARO, BÜHL, DERENZO, EPPERSON, FLATTE* (LRL)
ASCOLI 70 PRL 25 962	+BRÖCKMAY, CRAWLEY, EISENSTEIN, HANFT, + (ILL) JP
BASILE 70 LNC 4 838	+CALPIAZ, FRABETTI, MASSAM, + (CERN+BGNA+STR)
BAUDI 70 PL 31B 397	CERN BOSON SPECTROMETER GROUP (CERN)
BAUD2 70 PHILAD. CONF. P. 311	CERN BOSON SPECTROMETER GROUP (CERN)
BAUD3 70 PL 31 B 401	CERN BOSON SPECTROMETER GROUP (CERN)
BOCKMANN 70 NP B 16 221	+MAJOR, POLS, + (BOHN+DURH+NIJ+EPOL+TORI)
BUTLER 70 UCRL 19845	THESIS (LRL)
CAROLL 70 PRL 25 1393	+FIREBAUGH, GARFINKEL, MORSE, OH, + (WISC+INTO)
CASO 70 LNC 3 707	+CONTE, TOMASINI, CORDS+ (GENOVA+HAMB+MILAN+SACL)
DIAZ 70 NP B 16 239	+GAVILLET, LABROSSE, MCNTANET* (CERN+COEF) JP
DZIERBA 70 PR 2 2544	+SHEPARD, BISWAS, CASCN, JOHNSON, KENNEY (NDAM)
ALSO 68 LAMSA	
GARFINK 70 PL 33 B 536	GARFINKEL, AMMANN, CARMCNY, YEN (PURO) JP
GORDON 70 COO 1195 179	THESIS, ILLINOIS (ILL)
JOHNSTON 70 NP B 24 253	+KEY, PRENTICE, YOON, GARFINKEL, + (INTO+WISC)
KRUSE 70 PHILAD. CONF. P. 359	U. KRUSE, PARTIAL WAVE ANALYSIS (ILL) JP
NEF 70 THESIS+PRIV. COMM.	CERN BOSON SPECTROMETER GROUP (CERN)
SUTHERLA 70 PHILAD. CCF. P. 369	G. SUTHERLAND, INTERFERING RESONANCE (GLASGOW)
AGUILAR 71 PR D 4 2583	AGUILAR-BENITEZ, EISNER, KINSON (BNL)
ALSTON-GA71 PL 34 B 156	+BARBARO, BÜHL, DERENZO, EPPERSON, FLATTE* (LPL) JP
BARNHAM 71 PRL 26 1494	+ABRAMS, BUTLER, COYNE, GOLDFABER, HALL, + (LBL)
BEKETOV 71 SJNP 4 765	+SOMBRONSKY, KONVALOV, KRUTSCHIN, + (ITEP)
BINNIE 71 PL 36 B 257	+CAMILLETTI, DUANE, FARUQI, BURTON, + (LOIC+SHMP)
BINNIE2 71 PL 36 B 537	+CAMILLETTI, DUANE, FARUQI, BURTON, + (LOIC+SHMP)
BOWEN 71 PRL 26 1663	+FARLES, FAISSLER, BLIEDEN, + (NEAS+STON)
CRENNEL 71 PL 35 B 185	+GORDON+KWAN WU LAI, SCARR (BNL)
FARBER 71 NP B 25 237	+DE PINTO, BISWAS, CASON, DEERY, KENNEY, + (NCAN)
FOLY 71 PRL 26 413	+LOVE, OZAKI, PLATNER, LINDENBAUM, + (BNL+CUNY)
GRAYER 71 PL 34 B 333	+HYAMS, JONES, SCHLEIN, BLUM, DIETL+ (CERN+MPIM)
LYNCH 71 UCRL 20022 AND 71	G. LYNCH (LBL)
RINAUDO 71 NC 5 A 239	+BOECKMANN, MAJOR+ (TORI+BOHN+DURH+NIJ+EPOL) JP
ANKENBRA 72 PRL 29 1688	ANKENBRANDT, BRABSON, CRITTENDEN, HEINZ, + (IND)
BERENYI 72 NP B 37 621	+PRENTICE, STEENBERG, YOON, WALKER (INTO+WISC)
BLOODWORTH 72 NP B 37 203	BLOODWORTH, JACKSON, PRENTICE, YOON (INTO)
DAMERI 72 NC 9 A 1	+BORIATTA, GOUSSE, + (GENOVA+MILAN+SACL)
DAMGAARD 72 UNPUBLISHED MEMO	+LECHANOINE, MARTIN (BOHR+GEVA)
DEIBOLD 72 BATAV. CONF.	R. DIEBOLD, RAPPORTEUR TALK
EISENBERG 72 PR D 5 15	EISENBERG, BALLAM, DAGAN, + (REHO+SLAC+TELA)
ESPICAT 72 NP B 36 29	+CHESNOLIER, LILLESTOL, MCNTANET (CERN+COEF)
FOLY 72 PR D 6 747	+LOVE, OZAKI, PLATNER, LINDENBAUM, + (BNL+CUNY)
GARNJOST 72 PRIV. COMM.	M. ALSTON-GARNJOST (LBL)
GETTNER 72 PREPRINT NUB2145	M. GETTNER (NEAS)
KIENZLE 72 UNPUBLISHED MEMO	W. KIENZLE (LASSILA+YOUNG)
LASSILA 72 PRL 28 1491	LASSILA+YOUNG (INDA)
MORSE 72 NP B 43 77	+OH, WALKER, JOHNSTON, YOON (WISC+INTO)

REFERENCES FOR A2

Mesons

f'(1514), F₁(1540), ρ'(1600)

Data Card Listings

For notation, see key at front of Listings.

Table with 5 columns: W, B, S, F, E. Rows include background estimation, superseeding by Aguilar 72, and error increased by US. See typed note on K* mass.

13 F PRIME PARTIAL DECAY MODES

Table with 2 columns: P1-P7 (Prime into Pi+ Pi-, K Kbar, etc.) and Decay Masses.

13 F PRIME BRANCHING RATIOS

Table with 4 columns: R1-R6 (Prime into various channels), R1-M (Model dependent), and branching ratios.

REFERENCES FOR F PRIME

BARNES 65 PRL 15 322, CULWICK, GUIDONI, KALBFLEISCH, GOZ+ (BNL+SYRA), etc.

F₁(1540) → KKπ

47 F1(1540, JPC= 1 1-1 JP = 2-, 1+ FAVORED.

47 F1 MASS (MEV)

Table with 4 columns: M, B, S, B. Rows include mass values and background subtraction.

47 F1 WIDTH (MEV)

Table with 4 columns: W, B, S, B. Rows include width values and background subtraction.

47 F1 PARTIAL DECAY MODES

Table with 2 columns: P1-P3 (F1 into K Kbar Pi, etc.) and Decay Masses.

REFERENCES FOR F1

ADERHOLZ 69 NP 8 11 259, BARTSCH, (AAACH+BERL+CERN+CRAC+WARS), etc.

ρ'(1600) → 4π

65 RHC PRIME(1600, JPC=1-+) I=1

The ρ' was first seen in

γ (real or virtual) → ρ'0 → ρ0π+π-

with the π+π- pair apparently in an S wave (BINGHAM 72, DAVIER 73, SCHACHT 74, ALEXANDER 75).

A peak seen in the cross section of the reaction e+e- → π+π-π+π- has also been claimed as evidence for the ρ' (BARBARINO 72, CONVERSI 74); however, exchange mechanisms can be invoked that reproduce these results without any ρ' (HIRSCHFELD 74).

Evidence for a 2π decay has been looked for in several phase-shift analyses of one π-p → π-π+n experiment (GRAY 74). The first phase-shift solution found yielded clear evidence for ρ' → π+π- (HYAMS 73, OCHS 73). Since then, however, it has been recognized that at least four distinct solutions fit the data, and at most two of these exhibit any ρ' effect (ESTABROOKS 74, HYAMS 75). In an attempt to constrain the solutions further (by requiring them to be compatible with fixed-t and fixed-u analyticity for negative values of t and u), FROGGATT 75 conclude that only one solution is favored: essentially solution B of ESTABROOKS 74 or solution +- of HYAMS 75. This solution shows considerable evidence for the ρ' → ππ signal.

Note that none of the solutions has any evidence for a ρ'(1250). Finally, there is indirect evidence for the ρ' from a model-independent analysis of the pion form factor data in the space-like and time-like regions (LANG 75).

The mass and width values listed below are only indicative, because for such a broad peak they are extremely dependent on the parametrization chosen (SMADJA 72). Note also that the mass dependence of the width will be strongly affected by the inelastic channels with their rather high thresholds. For the photoproduction data, possibly diffractive background could explain the observed energy dependence of the width (SCHACHT 74).

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\rho'(1600)$, $A_3(1640)$

65 RHO PRIME MASS (MEV)

M C	(1600.)	APPROX.	BARBARINO	72	OSP	0 E+ E- TO 4 PI	1/73
M C	400	1430.	50.	BINGHAM	72	HBC	0 9.3 GAM P,P 4PI
M	1590.	20.	HYAMS	73	ASP	0 17 PI-P,N PI+PI-	12/74
M	1550.	60.	CONVERSI	74	OSP	0 E+ E- TO 4PI	12/75*
M	160	1550.	50.	SCHACHT	74	STR	05.5-9 G P,P 4PI
M	340	1450.	130.	SCHACHT	74	STR	09-18 G P,P 4PI
M	65	1570.	60.	ALEXANDER	75	HBC	07.5 GAM P,P 4PI
M P	(1600.)	(50.)	FROGGATT	75	RVUE	17 PI-P,PI+PI-N	12/75*
M P	FROM PHASE SHIFT ANALYSIS OF HYAMS 73 DATA						
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)						

65 RHC PRIME WIDTH (MEV)

W	400	650.	100.	BINGHAM	72	HBC	0 9.3 GAM P,P 4PI
W	180.	50.	50.	HYAMS	73	ASP	0 17 PI-P,N PI+PI-
W	360.	100.	100.	CONVERSI	74	OSP	0 E+ E- TO 4PI
W E	160	400.	120.	SCHACHT	74	STR	05.5-9 G P,P 4PI
W E	340	850.	200.	SCHACHT	74	STR	09-18 G P,P 4PI
W E	65	340.	160.	ALEXANDER	75	HBC	07.5 GAM P,P 4PI
W P	(220.)	(70.)	(70.)	FROGGATT	75	RVUE	17 PI-P,PI+PI-N
W P	FROM PHAS SHIFT ANALYSIS OF HYAMS 73 DATA						
W E	WIDTH ERRORS ENLARGED BY US TO 4*WIDTH/SORT(N)						
W	AVERAGE MEANINGLESS (SCALE FACTOR = 2.3)						

65 RHO PRIME PARTIAL DECAY MODES

P1	RHO PRIME INTO RHO 0 RHO 0 PI	139*	139*	773
P2	NEUTRAL RHO PRIME INTO ALL 4 CHARGED PI MODES	139*	139*	139*
P3	RHO PRIME INTO RHO RHO	773*	773*	
P4	RHO PRIME INTO PI PI	139*	139*	
P5	RHO PRIME INTO KBAR K	492*	493*	
P6	RHO PRIME INTO PI OMEGA	139*	783*	
P7	RHO PRIME INTO PI+ PI- PI0 PI0	139*	139*	134*

65 RHC PRIME BRANCHING RATIOS

R1 S	RHO PRIME INTO (RHO 0 RHO 0 PI+ PI-)/(4 PI, ALL CHARGED)	(P1)/(P2)	0 E+ E- TO 4 PI	1/73
R1 S	DOMINANT (1.80)	BINGHAM	72	HBC
R1 S	500	0.7	0.1	SCHACHT
R1 S	THE PI PI SYSTEM IS IN S WAVE			
R2	RHO PRIME INTO (RHO 0 RHO 0)/(RHO 0 PI+ PI-)	(P3)/(P1)	0 9.3 GAM P,P 4PI	1/73
R2	NCNE (FORBIDDEN BY I=1)	BINGHAM	72	HBC
R3	RHO PRIME INTO (PI+ PI-)/(4 PI, ALL CHARGED)	(P4)/(P2)	0 9.3 GAM P,P 2PI	1/73
R3	(0.14) OR LESS ESTIMATE DAVIER	73	STR	0 6-18 G P,P 4PI
R4	RHO PRIME INTO (KBAR K)/(4 PI, ALL CHARGED)	(P5)/(P2)	0 9.3 GAM P	1/73
R4	(0.04) OR LESS 2 SIGMA	BINGHAM	72	HBC
R5	RHO PRIME INTO (PI+PI-)/TOTAL	(P4)	0 5 PI+ P,DEL++2PI	1/74
R5 E	(0.15) OR LESS	EISENBERG	73	HBC
R5 E	ESTIMATED USING OPE MODEL.			
R5 P	0.25	0.05	HYAMS	73
R5 P	(1.35)	(.10)	FROGGATT	75
R5 P	FROM PHASE SHIFT ANALYSIS OF HYAMS 73 DATA			
R6	RHO PRIME INTO (PI+ PI- PI0 PI0)/(4PI, ALL CHARGED)	(P7)/(P2)	0 6-18 G P,P 4PI	1/74
R6	(1.) OR LESS ESTIMATE DAVIER	73	STR	0 6-18 G P,P 4PI
R7	RHO PRIME INTO (PI+ PI- + NEUTRALS)/(4PI, ALL CHARGED)	(P7+...)/(P2)	0 9.3 GAMMA P	12/75*
R7 U	(2.6)	(0.4)	BALLAM	74
R7 U	UPPER LIMIT. BACKGROUND NOT SUBTRACTED			

REFERENCES FOR RHO PRIME

ALVENSLE 71 PRL 26 273	ALVENSLEBEN, BECKER, BERTRAM, CHEN, + (DESY+MIT) G
BAUON 71 NP 830 213	*FRIDMAN, GERBER, GIVERNAUD, + (STRASBOURG) G
BULOS 71 PRL 26 149	*BUSZA, KEHOE, BENISTON, + (SLAC+UMD+IBM+LBL) G
BACCII 72 PL 388 551	*PENSO, SALVINI, STELLA, BALDINI-(CEROMA+FRAS) JPC
BARBARINO 72 LNC 3 689	BARBARINO, CERADINI, + (FRAS+ROMA+PADO+UMD)IGJP
BARTOLI 72 PR D 6 2374	*FELICETTI, OGGREN, + (FRAS+ROMA+NAPL)IGJP
BINGHAM 72 PL 418 635	*RABIN, ROSENFIELD, SMOJJA, YCST+(LBL,UCB,SLAC)IGJP
BRAMON 72 LNC 3 693	*GRECO (THEORETICAL PAPER) (FRASCATI) J
DIEBOLD 72 BATAV.CONF.	R. DIEBOLD RAPPORTEUR TALK (ANL)
EISENBERG 72 PR D 5 15	EISENBERG, BALLAM, DAGAN, + (REHO+SLAC+TELA)
LYSSAC 72 NC 10A 407	J. LAYSSAC, F. M. RENARD (MONT)
SMAJJA 72 PHIL. CONF. PROC349	*BINGHAM, FRETTER, BALLAM, CHADWICK+(LBL+SLAC)
CERADINI 73 PL 43 8 341	*CONVERSI, EKSTRAND, GRILLI, + (ROMA+FRAS+PADO)IGJP
DAVIER 73 NP 8 58 31	*DERADO, FRIES, LIU, MOZLEY, ODIAN, PARK, + (SLAC)
EISENBERG 73 PL 43 8 149	EISENBERG, KARSHCH, MIKENBERG, PITLUCK, + (REHO)
HYAMS 73 NP 8 64 134	*JONES, WEILHAMMER, BLUM, DIETL, + (CERN+MPIIM)
KREUZER 73 PR D 8 1431	H. J. KREUZER, A. N. KAMAL (UNIV. OF ALBERTA)
OCHS 73 THESIS	J. C. H. PARK (MPIIM) JP
PARK 73 NP 8 58 45	
BALLAM 74 NP 876 375	*CHADWICK, BINGHAM, FRETTER + (SLAC+LBL+MPIIM)
BERNABEI 74 LNC 11 261	*D. ANGULO, SPILLANTINI, VALENTE (ROMA+FRAS)
CONVERSI 74 PL 528 493	*PAGLUZZI, CERADINI, GRILLI + (ROMA+FRAS)
ESTABROO 74 NP 879 301	P. ESTABROOKS, A. C. MARTIN (DURH)
FERBEL 74 PR D9 824	T. FERBEL AND P. SLATTERY (ROCH)
GRAYNER 74 NP 8 75 189	G. GRAYNER, HYAMS, BLUM, DIETL, + (CERN+MPIIM)
HIRSHFEL 74 NP 874 211	A. C. HIRSHFELD, G. KRAMER (HAMB)
SCHACHT 74 NP 881 205	*DERADO, FRIES, PARK, YOUNG (MPIIM)
ALEXANDE 75 PL 578 487	ALEXANDER, BENARY, GANDSMAN, LISSAUER, + (TELA)
ALLES 75 NC 30A 136	ALLES-BORRELLI, BERNARDINI, + (CERN+BGNA+FRAS)
ESTABROO 75 NP 895 322	P. ESTABROOKS, A. C. MARTIN (DURH)
FROGGATT 75 NP 891 454	C. D. FROGGATT, J. L. PETERSEN (GLAS+NORD)
HYAMS 75 NP 8100 205	*JONES, WEILHAMMER, BLUM, DIETL, + (CERN+MPIIM)
LANG 75 PL 588 450	C. B. LANG, I. S. STEFANESCU (KARL)
LANGACKER 75 PREPRINT UPR00497	P. LANGACKER, G. SEGRE (PENN)
ROOS 75 NP 8 97 165	M. ROOS (HELS)

A₃(1640)

The $A_3(1640)$ is seen as a bump in the diffractive-like process $\pi N \rightarrow (\pi\pi\pi)N$. The dominant effect is a ~ 300 MeV wide enhancement in the $J^P=2^-$ $f\pi$ S-wave system, starting from $f\pi$ threshold. Neither additional (narrower) structure in the 3π mass distribution, nor other decay modes, have been clearly established. The situation would appear to be similar to that of the A_1 , but there are certain additional complications:

- Experiments with incident π^- observe little variation of the $J^P=2^-$ $f\pi$ phase in the A_3 mass region (ANTIPOV1 73, ASCOLI1 73), whereas experiments with incident π^+ show evidence for variations consistent with a resonance interpretation (OTTER 74, THOMPSON 74).
- The A_3 region is not well described by the Deck-like model of Ascoli et al., although the agreement could probably be improved by a series of more or less well motivated adjustments to the model (ASCOLI2 73, ASCOLI 74). Nevertheless, relative phase variations of $\sim 40^\circ$ through the A_3 mass region are predicted for the 2^- $f\pi$ wave.
- Other partial waves contribute strongly in the A_3 region. The 2^- $\rho\pi$ P wave may exhibit a broad enhancement; no phase variation relative to 2^- $f\pi$ is observed (ANTIPOV1 73, ASCOLI1 73, OTTER 74, THOMPSON 74). ANTIPOV1 73 show some evidence for an enhancement in the 2^+ $f\pi$ P wave with $M \sim 1750$ MeV, $\Gamma \sim 200$ MeV. The relative phases are not inconsistent with a resonance interpretation.

34 A3 MASS (MEV)

M	30(1600.0)	FORING	65	DBC	04.5 PI+ D	10/66
M	20(1630.0)	VELTITSKY	66	HBC	- 4.7 PI- P	12/75*
M	(1630.)	BALTAY	68	HBC	+ 7. 8.5 PI+ P	12/75*
M	(1600.0)	BARTSCH	68	HBC	+ 8. PI+ P, 3PI P	12/75*
M	(1610.)	LAMSA	68	HBC	- 8.0 PI-P, PI-F	12/75*
M	297(1673.0)	ARMENISE	69	DBC	+ 5.1 PI+D, 3PI++	12/75*
M	(1680.)	CASO	69	HBC	- 11 PI- P	5/70
M	(1660.0)	CASO	69	HBC	- 11 PI-P, PI-F	12/75*
M	(1645.0)	CRENNELL	70	HBC	- 6. PI- P, PI-F	12/75*
M	(1633.0)	MIYASHITA	70	HBC	- 6.7 PI-P, PI-F	1/71
M	BACKGROUND SUBTRACTION DIFFICULT.					
M	(1672.0)	BEKETOV	71	HBC	- 4.45 PI- P	11/71
M	(1600.)	PAER	71	DBC	+ 13. PI+ D, D13PI++	12/75*
M	263(1660.)	CASO	72	HBC	+ 11.7 PI+ P	12/75*
M	(1658.)	HARRISON	72	HBC	- 13..20. PI- P	12/72
M	F FIT ASSUMES AN ADDITIONAL PEAK AT 1830 MEV.					
M	EVIDENCE FOR A SUBSTANTIAL DECAY INTO S PI CLAIMED					
M	(1650.)	ANTIPICV1	73	CNTR	- 25..40. PI- P	12/75*
M	(1600.)	ASCOLI	73	HBC	- 5..25. PI- P, P A3	12/75*
M	(1600.)	THOMPSON	74	HBC	+ 13. PI+ P, P A3*	12/75*
M	EVIDENCE FOR A ROTATION OF THE PHASE CLAIMED.					
M	575(1640.)	KALEKAR	75	HBC	+ 15 PI+P, PI+ F	12/75*
M	FROM A FIT TO JP=2- F PI PARTIAL WAVE					
M	AVERAGING NOT MEANINGFUL					

Mesons

A₃(1640), ω(1675)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various decay modes and branching ratios. Includes entries for A3(1640) and ω(1675).

Table titled '34 A3 PARTIAL DECAY MODES' showing decay channels like A3 INTO 3 PI and A3 INTO RHO PI with associated decay masses.

Table titled '34 A3 BRANCHING RATIOS' providing detailed branching ratios for various decay channels of the A3 meson.

Table titled 'REFERENCES FOR A3' listing various scientific publications and sources related to the A3 meson.

Table listing various meson decays and production methods, including entries for A3(1640) and ω(1675) decaying into different pi and rho meson combinations.

ω(1675) → ρ0π0
45 OMEGA(1675, JPC=3--) I=0.
THIS RESONANCE OVERLAPS IN ITS 3PI MODE WITH THE A3, BUT IN SOME EXPERIMENTS ONE CAN ESTABLISH THE DECAY MODE RHO0 PI0, THUS I=0, WAGNER 74 FINES JP=3- UNIQUELY. THE DECAYS INTO 5PI AND OMEGA PI+PI- NEED FURTHER CONFIRMATION (SEE ALSO XI16901).

Table titled '45 OMEGA(1675) MASS (MEV)' showing mass measurements from various experiments like ARMENISE, KENYON, MATTHEWS, DIAZ, and WAGNER.

Table titled '45 OMEGA(1675) WIDTH (MEV)' showing width measurements from various experiments like ARMENISE, KENYON, MATTHEWS, DIAZ, and WAGNER.

Table titled '45 OMEGA(1675) PARTIAL DECAY MODES' showing decay channels and masses for the Omega(1675) meson.

Table titled '45 OMEGA(1675) BRANCHING RATIOS' providing branching ratios for various decay channels of the Omega(1675) meson.

Table titled 'REFERENCES FOR OMEGA(1675)' listing scientific publications and sources related to the Omega(1675) meson.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$g(1680)$

$g(1680)$

15 G(1680.JPG = 3+-) 1=1

This entry contains the 2π , 4π , $\omega\pi$, $K\bar{K}$, and $K\bar{K}\pi$ peaks in the region of 1700 MeV. The spin-parity determination and the mass and width in the Meson Table come from the 2π mode. An elasticity of 24% is found at resonance in the $\pi\pi$ elastic partial-wave analysis (HYAMS 75); this is consistent with the assumption that at least some of the effects listed are due to g decay into various channels. On the other hand, some discrepancies in masses, widths, and branching ratios reported indicate that there may be more than one $I_G=1^+$ meson in this region (BARNHAM 70, HOLMES 72, THOMPSON 74). Although we have collected all the data here under a common entry, we do not imply that they are necessarily all related.

15 G MASS (MEV)

M	PI+ PI- MODE					
M	(1700.0)	(100.0)	BELLINI	65 HLBC	0 6.1 PI-P	12/75*
M	(1640.0)		FORINO	65 CBC	0 4.5 PI+D	6/66
M	1670.0	30.0	GOLDBERG	65 HBC	0 6 PI+D, 8 PI-P	
M	(1683.1)	(13.1)	ARMENISE	68 DBC	0 5.1 PI+D	6/68
M	1729.0	20.0	CRENNELL	68 HBC	0 6.0 PI-P	12/68
M	(1655.0)	(10.0)	JOHNSTON	68 HBC	0 7.0 PI-P	6/68
M	1737.0	23.0	ARMENISE	70 DBC	0 9 PI-N	1/71
M	1687.	21.	STUNTEBEC	70 HBC	0 8. PI-P, 5.4 PI+D	2/72
M	1679.	12.	MATTHEWS	71 DBC	0 7. PI+N	2/72
M	1652.	13.	MATTHEWS	71 HBC	0 7. PI-P	2/72
M	E 600 1690.	7.	ENGLER	74 CBC	0 6. PI+N, PI+PI-N	12/75*
M	G 1693.	8.	GRAY	74 ASPK	0 17 PI-P, PI+PI-N	2/74
M	GI (1692.)	(12.)	ESTABROOK	75 RVUE	17 PI-P, PI+PI-N	12/75*
M	I (1722.)	(13.)	HYAMS	75 ASPK	0 17 PI-P, PI+PI-N	12/75*

M E MASS ERRORS ENLARGED BY US TO WIDTH/SQRT(N)
M I FROM PHASE-SHIFT ANALYSIS
M ERROR TAKES ACCOUNT OF SPREAD OF DIFFERENT PHASE-SHIFT SOLUTIONS
M G USES SAME DATA AS HYAMS 75
M
M AVG 1687.8 6.4 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)
M STUDENT1688.7 4.9 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
M (SEE IDEOGRAM BELOW)

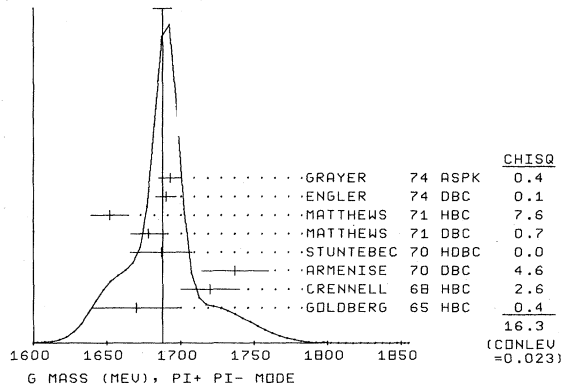
M	(2PI)++ MODE					
M	1640.0	25.0	CRENNELL	68 HBC	- 6.0 PI-P	12/68
M	122 1650.0	35.0	BARTSCH	70 HBC	+ 8 PI+ P, 2 PI	5/70
M	(1632.)	(5.)	THOMPSON	74 HBC	+ 13 PI+ P	12/75*
M	AVG 1643.4	20.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
M	STUDENT1643.4	21.9	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

M	K KBAR + K KBAR PI MODE						
M	(1675.)		EHRlich	66 HBC	+- 7.9 PI-P, K KBAR	2/72	
M	(1700.)		FRENCH	67 HBC	0 3, 3.6 PBAR P	7/67	
M	K OBSERVED IN NEUTRAL(K* KBAR) MODE (G-PARITY UNKNOWN)		FRENCH	67 HBC	(KO K+-) 3-4 PBAR P	7/67	
M	F (1740.)		FRENCH	67 HBC	(KO K+-) 3-4 PBAR P	7/67	
M	SEE FIG. 9 CF FRENCH 67						
M	1640.0	20.0	25.0	CRENNELL	68 HBC	+- 6.0 PI-P, KBAR K	12/68
M	1690.0	16.0		ADERHOLZ	69 HBC	+ 8 PI+ P, KBAR K	8/69
M	1692.	6.		BLUM	75 ASPK	0 18.4 PI-P, N K+-	11/75*
M	AVG 1688.7	8.6	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)				
M	STUDENT1689.5	6.1	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT				

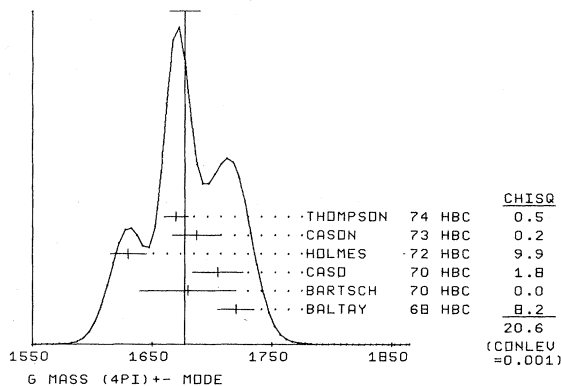
M	(4PI)++ MODE					
M	1720.	15.	BALTAY	68 HBC	+ 7, 8.5 PI+ P	6/68
M	(1675.0)	(10.0)	JOHNSTON	68 HBC	- 7.0 PI-P	6/68
M	144 1680.0	40.0	BARTSCH	70 HBC	+ 8 PI+ P, 4 PI	4/71
M	90(1640.0)	(20.0)	BARTSCH	70 HBC	+ 8 PI+ P, 2 PI	4/71
M	F 102(1689.0)	(20.0)	BARTSCH	70 HBC	+ 8 PI+ P, 2 RHO	4/71
M	1705.0	21.0	CASO	70 HBC	- 11.2PI-P, RHO 2PI	5/70
M	300(1710.)		ARMENISE	72 HBC	- 9.1 PI- P, P 4PI	12/72
M	1630.	15.	HOLMES	72 HBC	+ 10, -12, K+ P	1/73
M	167.	20.	CASO	73 HBC	+ 8, 18.5 PI-P	1/74
M	F (1685.)	(14.)	CASO	73 HBC	- 8, 18.5 PI-P	1/74
M	F 66(1733.)	(9.)	KLIGER	74 HBC	- 4.5 PI-P, P 4PI	12/75*
M	1670.	10.	THOMPSON	74 HBC	+ 13 PI+ P	12/75*

M F FROM (RHO+- RHO0) MODE
M
M AVG 1677.1 13.0 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)
M STUDENT1677.9 9.0 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
M (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE = 1687.8 ± 6.4
ERROR SCALED BY 1.5



WEIGHTED AVERAGE = 1677.1 ± 13.0
ERROR SCALED BY 2.0



M	RHO RHO MODE					
M	(1700.0)		MAURER	70 HBC	05.7 PBAR P, 7 PI	2/71
M	(1700.0)		BRAUN	71 HBC	05.7 PBAR P, 7 PI	11/71

M	OMEGA PI MODE					
M	1654.	24.	BARNHAM	70 HBC	+ 10 K+ P, OMEGA PI	6/70
M	1630.0	11.0	CASO	70 HBC	- 11.2PI-P, PI OMEG	5/70
M	(1666.)	(50.)	CASO	73 HBC	- 8, 18.5 PI-P	1/74
M	1686.	9.	THOMPSON	74 HBC	+ 13 PI+ P	12/75*
M	AVG 1662.8	18.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.8)			
M	STUDENT1665.4	14.0	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT			

M R PEAKS FROM MMS. (FOR DIFFICULTIES WITH MMS EXPTS. SEE THE A2 MINI-REVIEW IN THE 1973 EDITION)
M NR1 (1632.) (15.) FOCACCI 66 MMS - 7-12 PI-P, P MMS 12/72
M NR2 (1700.) (15.) FOCACCI 66 MMS - 7-12 PI-P, P MMS 12/72
M NR3 (1748.) (15.) FOCACCI 66 MMS - 7-12 PI-P, P MMS 12/72
M N NCT SEEN BY BOWEN 72
M R (1700.0) (47.0) ANDERSON 69 MMS - 16 PI- P, BACKW 8/69

15 G WIDTH (MEV)

M	PI+ PI- MODE					
M	(40.0)		FORINO	65 DBC	0 4.5 PI+D	6/66
M	180.0	40.0	GOLDBERG	65 HBC	0 6 PI+D, 8 PI-P	
M	188.	45.	ARMENISE	68 DBC	0 5.1 PI+D	6/68
M	200.0	100.0	CRENNELL	68 HBC	0 6.0 PI-P	12/68
M	(80.0)	(20.0)	JOHNSTON	68 HBC	0 7.0 PI-P	6/68
M	171.0	65.0	ARMENISE	70 DBC	0 9 PI+D	1/71
M	267.	72.	46.	STUNTEBEC	70 HBC	0 8. PI-P, 5.4 PI+D
M	156.	36.	MATTHEWS	71 DBC	0 7. PI+N	2/72
M	167.	40.	MATTHEWS	71 HBC	0 7. PI-P	2/72
M	E 600 167.	40.	ENGLER	74 CBC	0 6. PI+N, PI+PI-P	12/75*
M	G 200.	18.	GRAY	74 ASPK	0 17 PI-P, PI+PI-N	2/74
M	GI (240.)	(30.)	ESTABROOK	75 RVUE	17 PI-P, PI+PI-N	12/75*
M	I (267.)	(30.)	HYAMS	75 ASPK	0 17 PI-P, PI+PI-N	12/75*

M I FROM PHASE-SHIFT ANALYSIS
M ERROR TAKES ACCOUNT OF SPREAD OF DIFFERENT PHASE-SHIFT SOLUTIONS
M G USES SAME DATA AS HYAMS 75
M
M AVG 177.2 15.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
M STUDENT 181.2 14.6 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Data Card Listings

Mesons

For notation, see key at front of Listings.

X(1690), X⁻(1795), A₄(1900), S(1930)

X(1690)
-> omega TTT

64 X(1690)
THIS ENTRY CONTAINS (GMEGA PI PI) PEAKS AROUND 1690 MEV. EVIDENCE NOT COMPELLING. OMITTED FROM TABLE.

Table with columns for mass (MEV), width (MEV), and various experimental data points for X(1690).

Table with columns for width (MEV) and various experimental data points for X(1690).

REFERENCES FOR X(1690)

Table listing references for X(1690) including authors like DANYSZ, YOST, BARNES, and OREN.

X⁻(1795)

63 X-(1795, JPG=) I=1
SEEN AS A (PBAR N) BOUND STATE IN PBAR D ANNIHILATIONS AT REST. NEEDS FURTHER CONFIRMATION. OMITTED FROM TABLE.

Table with columns for mass (MEV) and various experimental data points for X-(1795).

Table with columns for width (MEV) and various experimental data points for X-(1795).

REFERENCES FOR X-(1795)

Table listing references for X-(1795) including authors like GRAY, BOGDANOV, and GAGY.

A4(1900)

43 A4(1900, JPG= -) I=1
OMITTED FROM TABLE.

THIS ENTRY CONTAINS THE DIFFRACTIVE-LIKE 3PI AND 5PI BUMPS IN THE REGION OF 1900 MEV, AS WELL AS VARIOUS PEAKS NEARBY. NOTE THAT THE EXISTENCE OF AN S-WAVE GPI THRESHOLD BUMP (AS A CONTINUATION TO A1 AND A3) IS NOT UNEXPECTED. OMITTED FROM TABLE.

Table with columns for mass (MEV) and various experimental data points for A4(1900).

Table with columns for width (MEV) and various experimental data points for A4(1900).

43 A4 PARTIAL DECAY MODES

Table listing partial decay modes for A4(1900) such as A4 INTO 3PI, A4 INTO RHO PI, etc.

43 A4 BRANCHING RATIOS
R1 A4 INTO (G PI)/(ALL 3PI)
R1 DOMINANT KALELKAR 75 HBC + (P4)/(P1) + 15 PI+P, P 3PI 12/75*

REFERENCES FOR A4

Table listing references for A4 including authors like DANYSZ, FRENCH, HUSON, BEMPORAD, CLAYTON, HARRISON, SALZBERG, BASTIEN, OREN, DEUTSCHMANN, and KALELKAR.

S(1930) REGION

31 S(1930, JPG=)

This entry contains the structure observed in s-channel NN annihilations, as well as various peaks claimed in this region by production experiments. The resonant interpretation of the broad structure observed in pp backward elastic scattering (CLINE 70, D'ANDLAU 75) has been criticized (PINSKY 71, BIZZARRI 72). Although accurate measurements of both total and elastic scattering have produced evidence for a narrow bump in this region (CARROLL 74, MARZANO 76), the charge-exchange reaction shows no such structure (GARNJOST 75, MARZANO 76). Measurements of exclusive channel cross sections and of the spectator momentum distribution in pd annihilations have suggested more complicated structure (DEFOIX 75, KALOGEROPOULOS 75). We prefer to wait for further clarification before entering the S meson into the table of established resonances.

31 S MASS (MEV)

Table listing mass (MEV) and various experimental data points for S(1930) from different experiments and authors.

Mesons

S(1930), h(2040), rho(~2100), T(2200)

Data Card Listings

For notation, see key at front of Listings.

31 S WIDTH (MEV)
W S CHANNEL NBAR N
W C (49.2) (9.) CLINE 70 HBC 0 -25--74 PBAR P 2/72
W B (35.) (9.) BENVENUTI 71 HBC 0 -1--.8 PBAR P 2/72
W S (9.) (4.) (3.) CARROLL 74 CNTR S CHAN.PBAR P,D 12/75*

REFERENCES FOR S
CHIKOVAN 66 PL 22 233 CERN MISSING MASS SPECTROMETER GROUP (CERN)
FDCACCI 66 PRL 17 890 CERN MISSING MASS SPECTROMETER GROUP (CERN)
BOESEBECK 68 NP B 4 501 BOESEBECK,DEUTSCHMANN,+ (AACHEN+BERLIN+CERN)
CLINE 68 PRL 21 1268 *ENGLISH,REEDER,TERRELL,TWITTY (WISCONSIN)

h(2040) 16 H(2040, JPG=++) I=0
APEL 75 AND BLUM 75 ESTABLISH JP AS 4+.
SINCE BOTH 2P10 AND K BAR DECAY MODES ARE SEEN, I=0.
ADDITIONAL EVIDENCE FOR THE H MESON IS REPORTED
IN WAGNER 74.

16 H MASS (MEV)
M 700 2030. 30. APEL 75 CNTR 40. PI-P,N 2P10 11/75*
M 2050. 25. BLUM 75 ASPK 18.4 PI-P,N K+K- 11/75*
M AVG 2041.8 19.2 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
M STUDENT2041.8 20.9 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

16 H WIDTH (MEV)
W 700 180. 60. APEL 75 CNTR 40. PI-P,N 2F10 11/75*
W 225. 120. 70. BLUM 75 ASPK 18.4 PI-P,N K+K- 11/75*
W AVG 192.8 50.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W STUDENT 192.8 54.8 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

16 H PARTIAL DECAY MODES
P1 H INTO PI P1 139+ 139
P2 H INTO K KBAR 497+ 497
REFERENCES FOR H
WAGNER 74 LONDON CONF. F. WAGNER, RAPORTEURS TALK (MPIM)
APEL 75 PL 578 398 *AUGENSTEIN+ (KARL+PISA+SERP+WIEN+CERN) JP
BLUM 75 PL 578 403 *CHABAUD,DIETL,GARELICK,GRAYEY+ (CERN+MPIM) JP

rho(~2100) REGION
51 RHO(2100, JPG= +)
THIS ENTRY CONTAINS REFERENCES TO PARTIAL WAVE ANALYSES OF THE REACTION PBAR P TO PI+ PI- OVER THE MASS RANGE 2-2.5 GEV. WITHOUT POLARISATION DATA IT SEEMS UNLIKELY THAT SUCH ANALYSES CAN PROVIDE UNIQUE SOLUTIONS.WE PREFER NCT TO LIST RESONANCE PARAMETERS YET. OMITTED FROM TABLE.

T(2200) REGION
32 T(2200, JPG= 1)
THIS ENTRY CONTAINS THE BUMP OBSERVED IN S CHANNEL NBAR N, AND VARIOUS PEAKS NEAR 2200 MEV. OMITTED FROM TABLE. FOR A REVIEW SEE ASTBURY 74.

32 T MASS (MEV)
M S CHANNEL NBAR N ABRAMS 70 CNTR S CHANNEL PBAR N 1/73
M B 2190. 10. SEEN AS BUMP IN I=1 STATE. SEE ALSO COOPER 68.
M B BRICMAN (69) SEES NO BUMP, SPIN LESS THAN 5 IS SO EXCLUDED
M B PEASLEE 75 CONFIRM PBAR P RESULTS OF ABRAMS 70 NC NARROW STRUCTURE
M K (2190.0) KALBFLEIS 69 HBC 0 S-CHANNEL PBAR P 7/69
M K SEEN IN PBAR P TO RHO RHO P10. I=1--
M K NOT SEEN BY BACCN 73.
M S (2141.) 2. ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
M S SEEN IN FINAL STATE (OMEGA PI+ PI-) DONALD 73 HBC 0 S CHANNEL PBAR P 1/74
M AVG 2192.9 2.0 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
M STUDENT2192.9 2.1 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

32 T WIDTH (MEV)
W S CHANNEL NBAR N ABRAMS 67 CNTR S CHANNEL PBAR N 7/67
W B (85.) SEE NOTE B UNDER T(2200) MASS ABOVE.
W K BETWEEN 20 AND 80 MEV KALBFLEIS 69 HBC 0 S-CHANNEL PBAR P 7/69
W S (14.) 8. ALSPECTOR 73 CNTR S CHANNEL PBAR P 1/74
W S SEEN IN FINAL STATE (OMEGA PI+ PI-) DONALD 73 HBC 0 S CHANNEL PBAR P 1/74
PEAKS FROM PRODUCTION EXPERIMENTS
W N (13.0) OR LESS CHIKOVANI 66 MMSP - 12.0 PI-P 8/66
W 62. 52. ALLES-BOR 67 HBC 0 5.7 PBAR P 12/66
W T (150.) ANDERSON 69 MMS - 16 PI-P,BACKW 12/75*
W T (25.) OR LESS ANDERSON 69 MMS - 16 PI-P,BACKW 12/75*
W T TWO SEPARATE PEAKS CASO 70 HBC - 11.2PI- P,N,CTE C. 5/70
W (130.0) CASO 70 HBC - 11.2PI- P,N,CTE C. 5/70
W C SEEN IN RHO- PI+ PI- (OMEGA AND ETA ANTISELECTED IN 4 PI SYSTEM) 5/70
W 68.0 22.0 KRAMER 70 HBC + 13.1 PI+ P,2PI 11/70
W 50 (160.) TAKAHASHI 72 HBC 8. PI-P,N 2PI 12/75*
W AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

32 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON
CS A (5.5) ABRAMS 70 CNTR S CHANNEL PBAR N 1/71
CS A FER I=1 NEAR N 0.13 0.08 ALSPECTOR 73 CNTR S CHANNEL PBAR N 1/74
REFERENCES FOR T
CHIKOVAN 66 PL 22 233 CERN MISSING MASS SPECTROMETER GROUP (CERN)
FDCACCI 66 PRL 17 890 CERN MISSING MASS SPECTROMETER GROUP (CERN)
ABRAMS 67 PRL 19 1209 *COOL,GIACOMELLI,KYCIA,LENTIC,LI,+ (BNL)
ALLES-BOR 67 NC 50 A 776 ALLES-BORELLI,FRENCH,FRISK,+ (CERN+BOONN)G=
CLAYTON 67 HEIDBG.CONF.-P.57 *MASON,MURHEAD,FILIPPAS+(LIVERPOOL+ATFENS)
COOPER 68 PRL 20 1059 *LYMAN,MANNER,MUSGRAVE,VOYVODIC (ANL)

REFERENCES FOR T
ANDERSON 69 PRL 22 1390 *COLLINS,BLIEDEN+ (BNL+CERN)
BRICMAN 69 PL 29 B 451 *FERRO-LUZZI,BIZARD,+ (CERN+CAEN+SACL)
CASO 69 NC 62 A 755 *CCNTE,BENZ,+ (GENO+DESY+HAMB+MILA+SACL)
KALBFLEI 69 PL 29 B 259 G.KALBFLEISCH AND D.MILLER REVUES (BNL)
MONTAKET 69 LUND CONF.-P.189 L.MONTAKET, RAPORTEUR (CERN)
ABRAMS 70 PR D 1 1917 *COOL,GIACOMELLI,KYCIA,LENTIC,LI,+ (BNL)
CASO 70 LNC 3 707 *CONTI,TOMASINI,CORCS+(GENO+HAMB+MILA+SACL)
KALBFLEI 70 PHILA.CONF.-P.409 G.KALBFLEISCH AND D.MILLER REVUES (BNL)
KRAMER 70 PRL 25 396 *BARTON,GUTAY,LICHTMAN,MILLER,+ (PURDUE)
BACCN 71 NP B 32 66 *BUTTERWORTH,MILLER,PHELAN,+ (RHEL+LIVP)
FIELDS 71 PRL 27 1749 *COOPER,RHINES-ALLISON
YJH 71 PRL 20 922 *BARISH,CARELLI,LGBKVCIZ+ (J+RNI+ROCH)

Mesons

Data Card Listings

For notation, see key at front of Listings.

T(2200), U(2360), $\bar{N}N_{I=0}(2375)$, X(2500-3600)

ALEXANDE 72 NP B 45 29
BERTANZA 72 CHEXBRES (CERN 72-10)
BUGG 72 PR D 6 3047
CLAYTEN 72 NP B 47 81
DIEBOLD 72 BATAV,CONF.
DONALD 72 PL 40 B 586
MING MA 72 NP B 51 77
TAKAHASHI 72 PR D 6 1266

ALEXANDER, BAR-VIR, BEVARY, CARAN,+ (TEL)A
L.BERTANZA, REVIEW TALK (PISA)
+CONDO,HART,COHN,ENDCRF,+ (TENN+ORNL+CINC)
+MASON,MUIRHEAD,RIGOPoulos,+ (LIVP+PATR)
R.DIEBOLD RAPPORTEUR TALK (ANL)
+GALLETTI,EDWARDS,DE BILLY,+ (LIVP+LPNP)
+EASTMAN,OH,PARKER,SMITH,SPRAFKA (MSU)
TAKAHASHI, BARISH,+ (TG+O+PENN+NDAM+ANL)

ALPLECTOR 73 PRL 30 511
BACON 73 PR D 7 577
BETTINI 73 NC 15 A 563
BOWEN 73 PRL 30 332
DONALD 73 NP B 61 333
KIENZLE 73 PR D 7 3520

A.ASTBURY REVIEW AT PRAGUE 74 (RHEL)
+BIGI,CASALI,LARICCIA,+ (PISA+PADO+TORI)

+GARNJOST,BIGI,+ (PADO+LBL+PIA+TORI)
+EARLES,FAISSLER,BLIEDEN,+ (NEAS+STON)
+EDWARDS,GIBBINS,BRIAND,DUBOC,+ (LIVP+LPNP)
W.KIENZLE (CERN)

A.ASTBURY REVIEW AT PRAGUE 74 (RHEL)
+BIGI,CASALI,LARICCIA,+ (PISA+PADO+TORI)

+JACQUES,JONES,PANDOLAS,+ (RUTG+UPNJ+ALBA)
+GARNJOST,ROSS,+ (LBL+PADO+PIA+TORI)
+DEMARZO,GUERRIERO,+ (CANB+PARI+BROW+MIT)

U(2360)
REGION

33 U(2360, JPG=) I=1
THIS ENTRY CONTAINS THE BROAD BUMP OBSERVED IN THE S CHANNEL NBAR N, AND VARIOUS OTHER PEAKS, MOSTLY CONTROVERSIAL. OMITTED FROM TABLE. FOR A REVIEW SEE ASTBURY 74.

33 U(2360) MASS (MEV)

M	A	S CHANNEL NBAR N	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/73
M	A	FOR I=1 NBAR N				
M	N	(2360.0) (25.0)	OH	70 HBC	-OPBAR(P,N),K*2PI	1/73
M	N	NO EVIDENCE FOR THIS BUMP SEEN IN THE PEAR,P DATA OF CHAPMAN 71				1/73
M	N	NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.				12/75*
M	I	(2359.1) (42.1)	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74
M	I	ISOSPINS 0 AND 1 NOT SEPARATED				

M PEAKS FROM PRODUCTION EXPERIMENTS
M M (2382.0) (24.0) CHIKOVANI 66 MMS - 12.0 PI-P 12/72
M N NOT SEEN BY BOWEN 73.
M M 2370. 17. ANDERSON 69 ASPK - 16 PI- BKSCAT 11/69
M B 126 2340. 20. BALTAY 75 HBC + 15 PI+P,PSPI 12/75*M B DOMINANT DECAY INTO RHO RHO PI
M M
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)

33 U(2360) WIDTH (MEV)

W	M	S CHANNEL NBAR N	ABRAMS	67 CNTR	S CHANNEL PBAR N	1/73
W	N	(140.0) (24.0)	OH	70 HBC	-OPBAR(P,N),K*2PI	11/71
W	N	NO EVIDENCE FOR THIS BUMP SEEN IN THE PEAR,P DATA OF CHAPMAN 71				11/71
W	N	NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.				12/75*
W	I	(165.0) (18.0) (8.0)	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74
W	I	ISOSPINS 0 AND 1 NOT SEPARATED				

M PEAKS FROM PRODUCTION EXPERIMENTS
M M (30.0) OR LESS CHIKOVANI 66 MMS - 12.0 PI-P 8/66
M N NOT SEEN BY BOWEN 73.
M M (57.0) ANDERSON 69 ASPK - 16 PI- BKSCAT 11/69
M B 126 180. 60. BALTAY 75 HBC + 15 PI+P,PSPI 12/75*M B DOMINANT DECAY INTO RHO RHO PI

33 SIGMA (MB) FOR FORMATION BY NUCLEON ANTINUCLEON

CS	A	(3.2)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
CS	I	(2.1) (0.2)	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74
CS	I	ISOSPINS 0 AND 1 NOT SEPARATED				

REFERENCES FOR U(2360)

CHIKOVAN 66 PL 22 233 CERN MISSING MASS SPECTROMETER GROUP (CERN)
FOCACCI 66 PRL 17 890 CERN MISSING MASS SPECTROMETER GROUP (CERN)

ABRAMS 67 PRL 18 1209 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)

ANDERSON 69 PRL 22 1390 +BLESER,BIRNBAUM,EDELSTEIN,+ (BNL+CERN)
BRICMAN 69 PL 29 B 451 +FERRO-LUZZI,BIZARD,+ (CERN+CAEN+SACL)

CASO 69 LNC 3 707 +CONTE,BENZ,+ (GENO+DESY+HAMB+MILA+SACL)

ABRAMS 70 PR D 1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)
OH 70 PRL 24 1257 +PARKER,EASTMAN,SMITH,SPRAFKA,MA (MSU)

CHAPMAN 71 PR D4 1275 +GREEN,LYS,MURPHY,RING,+ (MICH)
FIELDS 71 PRL 27 1749 +COOPER,RHINES,ALL TSON (ANL+OXF)
YOH 71 PRL 26 922 +BARISH,CAROLL,LOBKOVICZ+ (CIT+BNL+RGCH)

ASTBURY 72 CERN 72-10 A.ASTBURY REVIEW AT CHEXBRES 72 (RHEL)
DIEBOLD 72 BATAV,CONF. R.DIEBOLD RAPPORTEUR TALK (ANL)
EASTMAN 72 NP B 51 29 +MING MA,OH,PARKER,SMITH,SPRAFKA (MSU)
MING MA 72 NP B 51 77 +EASTMAN,OH,PARKER,SMITH,SPRAFKA (MSU)

ALPLECTOR 73 PRL 30 511 ALPLECTOR,COHEN,CVIJANOVICH,+ (RUTG+UPNJ)
BOWEN 73 PRL 30 332 +EARLES,FAISSLER,BLIEDEN,+ (NEAS+STON)
DONNACHI 73 LNC 7 285 A.DONNACHE, P.R. THOMAS (MANCHESTER)
KIENZLE 73 PR D 7 3520 W.KIENZLE (CERN)
OH 73 NP B 51 77 +EASTMAN,MING MA,PARKER,SMITH,+ (MSU)

ASTBURY 74 CERN 74-18 A.ASTBURY REVIEW AT PRAGUE 74 (RHEL)
MING MA 74 NP B68 214 +MOUNTZ,ZEMANY,SMITH (MICH)

BALTAY 75 PRL 35 891 +CAUTIS,COHEN,KALELKA,PISELLO,+ (COLU+BING)
KEMP 75 NC 27 A 155 +LOTTIS,CONTRI,TEODORO+(DURH+GENO+MILA+LPNP)

$\bar{N}N_{I=0}(2375)$

41 N NBAR(2375, JPG=) I=0
EVIDENCE FOR RESONANCE PRELIMINARY. OMITTED FROM TABLE. FOR A REVIEW SEE ASTBURY 74.

41 N NBAR(2375) MASS

M	I	2375. (2359.1)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
M	I	ISOSPINS 0 AND 1 NOT SEPARATED	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74

41 N NBAR(2375) WIDTH

W	I	(190.0) (165.0)	ABRAMS	70 CNTR	S CHANNEL NBAR N	1/71
W	I	(18.0) (8.0)	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74
W	I	ISOSPINS 0 AND 1 NOT SEPARATED				

41 N NBAR(2375) SIGMA (MB) FOR FORMATION BN

CS	I	(2.5) (2.1)	ABRAMS	70 CNTR	S CHANNEL PBAR P	1/71
CS	I	(0.2) (0.1)	ALPLECTOR	73 CNTR	S CHANNEL PBAR P	1/74
CS	I	ISOSPINS 0 AND 1 NOT SEPARATED				

REFERENCES FOR N NBAR (2375)

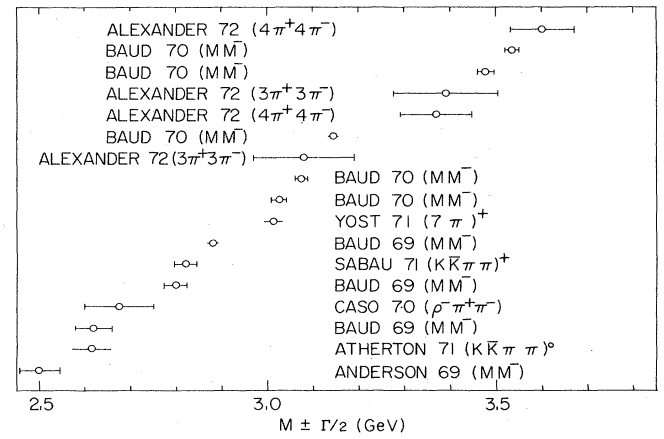
BRICMAN 69 PL 29 B 451 +FERRO-LUZZI,BIZARD,+ (CERN+CAEN+SACL)
ABRAMS 70 PR D 1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)
EASTMAN 72 NP B 51 29 +MING MA,OH,PARKER,SMITH,SPRAFKA (MSU)
MING MA 72 NP B 51 77 +EASTMAN,OH,PARKER,SMITH,SPRAFKA (MSU)
ALPLECTOR 73 PRL 30 511 ALPLECTOR,COHEN,CVIJANOVICH,+ (RUTG+UPNJ)

ASTBURY 74 CERN 74-18 A.ASTBURY REVIEW AT PRAGUE 74 (RHEL)
MING MA 74 NP B68 214 +MOUNTZ,ZEMANY,SMITH (MICH)

X(2500-3600)

46 X(2500-3600)
THIS ENTRY CONTAINS VARIOUS HIGH MASS NON-STRANGE PEAKS. OMITTED FROM TABLE.

The high mass region is covered nearly continuously by evidence for peaks of various widths and decay modes (see figure). As a satisfactory grouping into particles is not yet possible, we list all the Y=0 bumps with M > 2400 MeV together, ordered by increasing mass. Note that ANTIPOV 72 ($\pi^- p \rightarrow p \pi^-$) at 5 and 40 GeV/c see no narrow bumps.



Masses and widths of reported enhancements with Y=0, M > 2400 MeV. (—o— indicates that upper limit only was reported for the width.)

Mesons

X(2500-3600), K±, K0, K*(892)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for mass, width, and decay modes for mesons X(2500-3600). Includes entries for ANDERSON, BAUD, CASO, SABAU, BRAUN, ALEXANDER, and YOST.

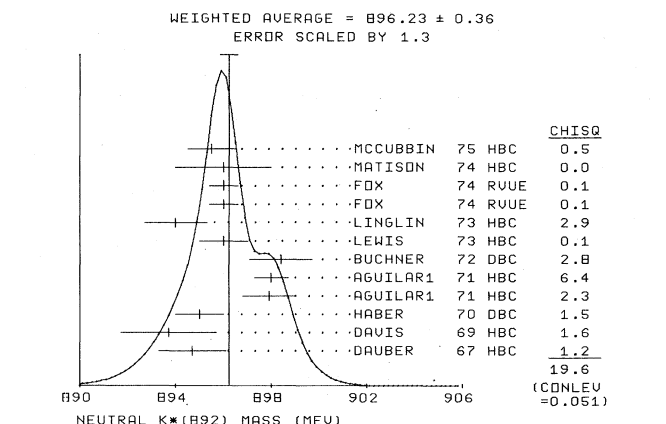
REFERENCES FOR X(2500-3600)
+COLLINS,+ (BNL+CERN)
CERN B050N SPECTROMETER GROUP (CERN)
+BAR-NIR,DAGAN,GIDIDAL,GRUNHAUS+(TEL-AWIV) (CERN)
CERN B050N SPECTROMETER GROUP (CERN)
+CENTE,TOMASINI,CORDS+(GENO+HAMB+MLLA+SACL) (CERN)
+URETSKY (BUCH+ANL) (CERN)
+MORRIS,AL BRIGHT,BRUCKER,LANNUTTI (FSU)
ALEXANDER, BAR-NIR, BEVARY,DAGAN,+ (TEL-AWIV) (CERN)
ANTIPOV 72 PL 408 147 +BRAU,BUSNELLO,DAMGAARD,+ (IHEP+CERN)
BRAUN 75 PREPRINT CRN 24 +BRICK,FRIDMAN,GERBER,JUILLOT,+ (STRB)
KALELKAR 75 THESIS(INEVIS 207) M.S.KALELKAR (COLU)

S=±1 MESON STATES

K±
10 CHARGED K(494,JP=0-) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS
K0
11 NEUTRAL K(498,JP=0-) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

K*(892) 18 K*(892,JP=1-) I=1/2
18 K*(892) MASS (MEV)
CHARGED ONLY. THIS IS WHAT APPEARS ON MESON TABLE
M S 100 (898.0) (5.0) CHADWICK 63 HBC + 1.5 K*(K PI) 12/75*

A INCLUDED IN LINGLIN 73 WORLD K+P DST
C FROM POLE EXTRAPOLATION.
D ERRORS ENLARGED BY US TO GAMMA/SQRT(N). SEE TYPED NOTE.
I INCLUSIVE REACTION. COMPLICATED BACKGROUND AND PHASE-SPACE EFFECTS
S DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED
W NUMBER OF EVENTS IN PEAK REEVALUATED BY US



Data Card Listings

For notation, see key at front of Listings.

Mesons

K*(892)

Note on K*(892) Masses and Mass Differences

Unrealistically small errors are reported by some experiments. We use simple "realistic" tests for the minimum errors on the determination of mass and width from a sample of N events:

delta_min(m) = Gamma / sqrt(N), delta_min(Gamma) = 4 * Gamma / sqrt(N)

(For detailed discussion see the April 1971 edition of this note.) We consistently increase unrealistic errors before averaging.

Although in the past we have argued against taking the mass difference m(K*0) - m(K*pm) from the separate averages of m(K*0) and m(K*pm), we no longer see any reason for such caution. In fact, the difference between the separate averages agrees with direct measurements of the mass difference, and it has a smaller statistical error.

18 K*(892) - K*(pm) MASS DIFF. (MEV)

Table with columns for experiment (D, W, S), mass (m), width (Gamma), and average values. Includes entries for BARASH, FICENEC1, FICENEC2, AGUILARI, and averages.

ERRORS ENLARGED BY US TO GAMMA/SQRT(N). SEE TYPED NOTE. DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED.

18 K*(892) WIDTH (MEV)

Table with columns for experiment (W, S), width (Gamma), and average values. Includes entries for CHADWICK, WJOJCICKI, ADELMAN, FERRO-LUZZI, GELSEMA, WANGLER, BARASH, BARLOW, CONFORTO, DAUBER, DE BAERE, KANG, SALLSTRM, DE WIT, BARNHAM, CRENNELL, DAVIS, DE BAERE, FRIEDMAN, JAHALCA, LIND, ATHERTON, HADER, AGUILAR, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA, DE BAERE, BARNHAM, CRENNELL, DAVIS, DE BAERE, FRIEDMAN, JAHALCA, LIND, ATHERTON, HADER, AGUILAR, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA.

Table listing various experiments and their results for K*(892) properties, including columns for experiment, mass, width, and average values.

A INCLUDED IN LINGLIN 73 WORLD K*P DST. C FROM POLE EXTRAPOLATION. D ERRORS ENLARGED BY US TO 4*GAMMA/SQRT(N). SEE TYPED NOTE. I INCLUSIVE REACTION. COMPLICATED BACKGROUND AND PHASE-SPACE EFFECTS. S DATA WITH MASS ERROR OF 3 MEV OR MORE NOT AVERAGED. W NUMBER OF EVENTS IN PEAK REEVALUATED BY US.

18 K*(892) PARTIAL DECAY MODES

Table showing decay modes for K*(892) into K pi, K pi pi, and K gamma, with associated decay masses.

18 K*(892) BRANCHING RATIOS

Table showing branching ratios for K*(892) into various decay channels, including K pi pi, K pi pi pi, and K gamma.

REFERENCES FOR K*(892)

List of references for K*(892) decays and properties, including authors like ALSTON, ALEXANDE, COLLEY, CHADWICK, WJOJCICKI, STUART LEE, FERRO-LUZZI, GELSEMA, WANGLER, BARASH, BARLOW, CONFORTO, DAUBER, DE BAERE, KANG, SALLSTRM, DE WIT, BARNHAM, CRENNELL, DAVIS, DE BAERE, FRIEDMAN, JAHALCA, LIND, ATHERTON, HADER, AGUILAR, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA, DE BAERE, BARNHAM, CRENNELL, DAVIS, DE BAERE, FRIEDMAN, JAHALCA, LIND, ATHERTON, HADER, AGUILAR, ABRAMOVI, BINGHAM, BEMPORAD, BRUNET, BUCHNER, CORDS, MERCER, YUTA.

Mesons

$K^*(892)$, $\kappa(1250)$

BERTHCN 73 NP B 63 54	+MONTANET, PAUL, BERTRANET, * (CERN+SACL)
CLARK 73 NP B 54 432	+LYONS, RADOJICIC (OXFORD)
LEWIS 73 NP B 60 283	+ALLEN, JACOBS, DANYSZ, BORG, + (LOWC+LOIC+CDEF)
LINGLIN 73 NP B 55 408	D. LINGLIN (CERN)
WALUCH 73 PR D 8 2837	+FLATTE, FRIEDMAN (LBL)
FOX 74 NP B80 403	G.C. FOX, M. L. GRIS (CIT)
MATISCN 74 PR D9 1872	+GALTIERI, GARNJUST, FLATTE, FRIEDMAN, + (LBL)
BRANDENB 75 PL 59 B 405	BRANDENBURG, CARNEGIE, CASHMERE, DAVIER+ (SLAC)
CARITHER 75 PRL 35 349	CARITHERS, MUHLEMANN, UNDERWOOD, + (ROCH+MCGI)
MCCUBBIN 75 NP B86 13	N. A. MCCUBBIN, L. LYONS (OXF)
PALER 75 NP B96 1	+TOVEY, SHAH, SPIRO, CHAURAND+ (RHEL+SACL+EPCL)

$\kappa(1250)$

19 K PI S WAVE, CALLED KAPPA

S-Wave $K\pi$ Interactions in the Region 750-1700 MeV

$K\pi$ interactions in the $I(J^P)=1/2(0^+)$ wave can be described by the elastic phase shift δ_0^1 from the $K\pi$ threshold (~ 630 MeV) up to at least 1100 MeV (BINGHAM 72). The first inelastic S-wave thresholds are $K\pi\pi$ and $K\eta$, neither of which is known to be important below 1400 MeV. Apart from the inelastic thresholds, the S-wave $\pi\pi$ and $K\pi$ interactions are reminiscent of each other.

All phase-shift solutions (MERCER 71, BINGHAM 72, FIRESTONE 71, 72, MATISON 72, 74, BAKER 73, GALTIERI 73, YUTA 73, FOX 74, LAUSCHER 75) share the following intrinsic ambiguities:

- 1) The standard modulo-180° ambiguity.
- 2) If one amplitude is dominant [the P wave near $K^*(892)$ or the D wave near $K^*(1420)$], then the observed S-P or S-D interference can be explained by two ambiguous S-wave solutions, known as "up" and "down". For an illustration see our 1972 edition.

The resulting "up-down" ambiguity of the S-wave phase shift in the $K^*(892)$ region (BINGHAM 72) was restricted by an analysis of $12 \text{ GeV}/c \text{ } K^+p \rightarrow K^+\pi^-\Delta^{++}$ (MATISON 72, 74, GALTIERI 73) to only two points at 890 and 900 MeV, where the P wave goes through 90°. The "up" solution is therefore excluded, except for a possibility of a very narrow ($\Gamma < 7$ MeV) S-wave resonance. Moreover, CHUNG 72 imposes positivity on physical-region $K\pi$ moments, and finds a narrow resonance most unlikely.

FIRESTONE 71 and 72 have continued $K\pi$ partial-wave analysis up to 1700 MeV. They find that δ_0^1 crosses 90° near 1300 MeV, and indeed shows the "up-down" ambiguity near the $K^*(1420)$.

Meanwhile several groups have attempted to clarify the situation around 1300 MeV. CORDS 72, FRATI 72, ROUGE 72, and LAUSCHER 75

Data Card Listings

For notation, see key at front of Listings.

give support to the resonant S-wave interpretation of FIRESTONE 71. The other groups (AGUILAR 72, BUCHNER 72, CRENNELL 72, ENGELMANN 72, BAKER 73) agree that the S wave is important but not necessarily resonant.

In a review of the status of the 0^+ nonet, MORGAN 75 shows that the resonant interpretation, $\kappa(1250)$, is consistent with the information on the other S-wave systems, $\pi\pi$, $K\bar{K}$, and $\eta\pi$ (see the corresponding mini-reviews), suggesting a non-ideally mixed nonet.

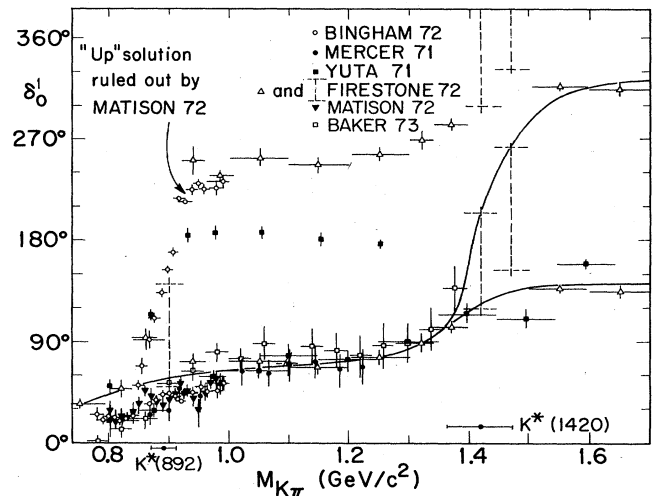


Fig. 1. The S-wave $K\pi$ phase shift. An "up-down" ambiguity occurs at the mass of $K^*(892)$, $K^*(1420)$, ..., which can be resolved by precise measurement of $\sigma(K\pi)$. At $K^*(892)$, the "up" solution was ruled out by MATISON 72; earlier "up" solutions are plotted only to show historical progress.

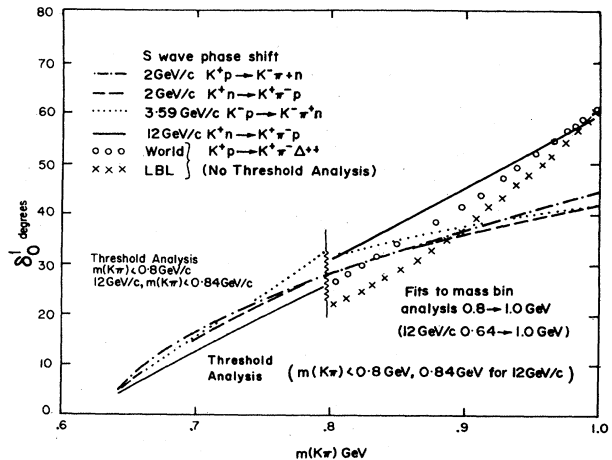


Fig. 2. Summary of six different S-wave $K\pi$ phase-shift analyses, including four solutions of FOX 74. (Figure from FOX 74.)

Data Card Listings

For notation, see key at front of Listings.

Mesons
 $\kappa(1250)$, Q

19 KAPPA MASS (MEV)				
M	1250.	100.	MORGAN	75 RVUE

19 KAPPA WIDTH (MEV)				
M	(450.)	APPROX.	MORGAN	75 RVUE

REFERENCES FOR KAPPA				
TRIPPE	68 PL 28 B 203		+CHIEN, MALAMUD, NELLEMA, SCHLEIN, +	(UCLA)
CRENNELL	69 PRL 22 487		+KARSHON, LAI, O. NEALL, SCARR	(BNL)
DODD	69 PR 177 1994		+JOLDERSMA, PALMER, SAMIOS	(BNL)
GOLDBERG	69 PL 30 B 434		SABRE COLLABOR. (SACL+AMST+BGNA+REHO+EPOL)	
SCHLEIN	69 ARGONNE CONF. 446		P. SCHLEIN	(UCLA)
FIRESTONE	71 PRL 26 1460		A. FIRESTONE, G. GOLDBERGER, D. LISSAUER	(LRL)
MERCER	71 NP B32 381		+ANTICH, CALLAHAN, CHIEN, COX, + (JOHN HOPKINS)	
YUTA	71 PRL 26 1502		+DERRICK, ENGELMANN, MUSGRAVE	(ANL+EFI)
AGUILAR	72 PR D 6 11		AGUILAR-BENITEZ, CHUNG, ETSNER	(BNL)
BINGHAM	72 NP B 41 1		+ (INTERNATIONAL K+ COLLABORATION)	
BUCHNER	72 NP B 45 333		+DEHM, CHARRIERE, CORNET, + (MPIN+CERN+BRUX)	
CHUNG	72 PRL 29 1570		+ETSNER, AGUILAR-BENITEZ	(BNL)
CRENNELL	72 PR D 6 1220		+GORDON, KWAN, WU LAI, SCARR	(BNL)
DIEBOLD	72 BATAV. CONF.		R. DIEBOLD RAPPORTEUR TALK	(ANL)
ENGELMANN	72 PR D 5 2162		ENGELMANN, MUSGRAVE, FORMAN, +	(ANL+EFI)
FIRESTONE	72 PR D 5 2188		+GOLDBERGER, LISSAUER, TRILLING	(LBL) (PMA)
FRATI	72 PR D 6 2361		+HALPERN, HARGIS, SHAPIRO, CARNAHAN, + (PENN+UNC)	
ROUGE	72 NP B 46 29		+VIDEAU, VOLTE, DE BRION, +	(EPOL+SACL)
MATISON	72 LBL 1537 (THESIS)		REVISED VERSION WILL GO TO PHYS. REV.	LBL
BAKER	73 IC/HEP/73/12		+BANERJEE, CAMPBELL, HALL, ISLAM, + (LOIC+LOWC)	
CORDS	73 NP B 54 109		+CARMONY, LANDER, MEIERE, + (PURD+UC+IUPUI)	(LBL)
GALTIERI	73 PREP. LBL 1772		+MATISON, GARNJOST, FLATTE, FRIEDMAN, +	(LBL)
LINGLIN	73 NP B 55 408		D. LINGLIN	(CERN)
YUTA	73 NP B 52 70		+ENGELMANN, MUSGRAVE, FORMAN, +	(ANL+EFI)
FOX	74 NP B80 403		G.C. FOX, M.L. GRISS	(CIT)
MATISON	74 PR D9 1872		+GALTIERI, GARNJOST, FLATTE, FRIEDMAN, +	(LBL)
ALSO	72 MATISON, 73 GALTIERI			
MORGAN	74 PL 518 71		D. MORGAN	(RHEL)
BAKER	75 NP B99 211		+BANERJEE, CAMPBELL, ALLEN, MARCH, + (LOIC+LOWC)	
LAUSCHER	75 NP B86 189		+OTTER, WIECZOREK, + (ABCLV COLLABORATION)	
MORGAN	75 PREPRINT RL75133		D. MORGAN	(RHEL)

Q REGION, $K\pi\pi(1240-1400)$

28 Q REGION (1200-1400) MEV I=1/2

The main effect in the Q region is a broad bump in the $K\pi\pi$ spectrum between 1200 and 1400 MeV (not far above the $K^*(892)\pi$ threshold), produced by K beams without charge exchange. In particular, it has been observed in coherent K^+d interactions (FIRESTONE 72) and in coherent interactions on heavy nuclei (BINGHAM 73). Throughout the entire region, $J^P = 1^+$ and $I = 1/2$.

Evidence for narrower states in the Q region has been reported from reactions with incident π^- and \bar{p} (ASTIER 69, CRENNELL 67, 72, DAVIDSON 74, DORE 75). FIRESTONE 72 observe a bump in the backward direction with a shape similar to that of the Q. WERNER 73 find no evidence for narrow states in the charge-exchange reaction $K^-p \rightarrow (K\pi\pi)^0 n$, but the data can accommodate some broad Q production. The $(K\pi\pi)^0$ system appears to have an important $J^P=1^+$ contribution (OTTER 75).

The broad Q peak does not have a simple Breit-Wigner shape. It can be fitted at all

energies by a superposition of two Breit-Wigner amplitudes (FIRESTONE 70, BARNHAM 71, BOWLER 71, BARLOUTAUD 73).

Dalitz plot analyses of the interference between the $K^*\pi$ and $K\rho$ modes show the relative magnitude and relative phase of the two decay amplitudes varying with $K\pi\pi$ mass. The $K\rho$ mode has a maximum intensity below that of $K^*\pi$. This is suggestive of the presence of two $J^P=1^+$ states, possibly mixtures of the strange members of the $J^P=1^{++}$ (" A_1 ") and 1^{+-} (B) nonets (GOLDBERGER 67, BARNHAM 71, BOWLER 71, GARFINKEL 71, BINGHAM 72, FIRESTONE 72, BOWLER 74). In addition to the dominant modes $K^*\pi$ and $K\rho$, there is some evidence for a $K\pi\pi$ mode, with the $\pi\pi$ or the $K\pi$ system in an S wave (ALEXANDER 69, BARNHAM 71, DAVIS 72, BARLOUTAUD 73).

Recent partial-wave analyses have confirmed the rather complex situation in the Q region (DEUTSCHMANN 74, ANTIPOV 75, OTTER2 75, OTTER3 75, TOVEY 75, BRANDENBURG 76). Although $J^P=1^+$ is the dominant contribution, other spin-parity states are necessary to describe the whole Q enhancement. Moreover, the $K^*\pi$ and $K\rho$ modes are not produced coherently and have different polarization properties (OTTER2 75, OTTER3 75, TOVEY 75, BRANDENBURG 76). Whereas the $K\rho$ mode approximately conserves s-channel helicity, the $K^*\pi$ mode is approximately t-channel helicity-conserving.

ANTIPOV 75 do not require the $K\rho$ decay mode to be present at 40 GeV/c.

ANTIPOV 75 and OTTER2 75 find no variation in the phase of the $1^+ K^*\pi$ wave. On the other hand BRANDENBURG 76, with high-statistics spectrometer data, observe sufficient phase variation of both the $1^+ K\rho$ and $K^*\pi$ waves to warrant proposing the existence of two 1^+ mesons, superposed on a large 1^+ , mainly $K^*\pi$, Deck background. The first state has $M \sim 1300$ MeV, $\Gamma \sim 200$ MeV and decays mainly to $K\rho$; the second has $M \sim 1400$ MeV, $\Gamma \sim 160$ MeV and decays mainly to $K^*\pi$.

Mesons

Q

Data Card Listings

For notation, see key at front of Listings.

28 Q REGION MASS (MEV)

Table listing meson production by beams other than K mesons and by K beams. Columns include particle ID, mass (MEV), and experimental parameters like beam type and energy.

28 Q LCW (CA) MASS (MEV)

Table listing meson production from experiments splitting the Q region into two peaks. Columns include particle ID, mass (MEV), and experimental parameters.

28 Q HIGH (QB) MASS (MEV)

Table listing meson production from experiments splitting the Q region into two peaks. Columns include particle ID, mass (MEV), and experimental parameters.

28 Q REGION WIDTH (MEV)

Table listing meson production by beams other than K mesons and by K beams. Columns include particle ID, mass (MEV), and experimental parameters.

28 Q LCW (CA) WIDTH (MEV)

Table listing meson production from experiments splitting the Q region into two peaks. Columns include particle ID, mass (MEV), and experimental parameters.

28 Q HIGH (QB) WIDTH (MEV)

Table listing meson production from experiments splitting the Q region into two peaks. Columns include particle ID, mass (MEV), and experimental parameters.

28 Q REGION PARTIAL DECAY MODES

Table listing decay masses for various meson decay modes. Columns include mode ID, description, and decay masses.

28 Q REGION BRANCHING RATIOS

Table listing branching ratios for various meson decay modes. Columns include mode ID, description, and branching ratios.

REFERENCES FOR Q REGION

Table listing references for the Q region, including author names and publication details.

PRODUCED BY BEAMS OTHER THAN K MESONS

Table listing meson production by beams other than K mesons. Columns include particle ID, mass (MEV), and experimental parameters.

Data Card Listings

For notation, see key at front of Listings.

Mesons Q, K*(1420)

Table listing meson data cards including names (ALEXANDE, ANDREWS, BARBARO, etc.), PRL numbers, and various physical parameters like mass and width.

Table with 22 K*(1420) WIDTH (MEV) header, listing charged and neutral decay modes with branching ratios and scale factors.

Table with 22 K*(1420) PARTIAL DECAY MODES header, listing partial decay modes (P1, P2, etc.) and their corresponding masses.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS. The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i +/- delta P_i, where delta P_i = sqrt(delta P_i delta P_i), while the off-diagonal elements are the normalized correlation coefficients (delta P_i delta P_j) / (delta P_i delta P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 22 K*(1420) MASS (MEV) header, listing mass measurements and their uncertainties.

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R1, R2, etc.).

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R3, R4, etc.).

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R5, R6, etc.).

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R7, R8, etc.).

Table with 22 K*(1420) BRANCHING RATIOS header, listing branching ratios for various decay channels (R9, R10, etc.).

Mesons

K*(1420), KN(1700), L(1770)

Data Card Listings

For notation, see key at front of Listings.

R7 K*(1420) INTO (K OMEGA) / K PI (P4)/(P1)
R7 (0.08) OR LESS SHEN 66 HBC 4.6 K+P 8/66
R7 (0.21) OR LESS BASSOMPIE 69 HBC + 5 K+P 9/69
R7 0.13 0.07 BASSOMPIE 69 HBC 0 5 K+P 9/69
R7 0.05 0.04 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R7 (0.21) OR LESS CL=.95 CHUNG 74 HBC - 7.3 K-P, K- P 12/75*

R8 K*(1420) INTO (K RHO) / (K PI) (P3)/(P1)
R8 (0.09) OR LESS CHUNG 65 HBC + 0 3.9-4.2 PI- P 8/66
R8 0.26 0.16 SCHWEINER 68 HBC 0 4.1+5.5 K- P 10/67
R8 (0.23) OR LESS BASSOMPIE 69 HBC + 5 K+P 9/69
R8 (0.31) OR LESS BASSOMPIE 69 HBC 0 5 K+P 9/69
R8 Q 15 (0.11) (0.06) BISHOP 69 HBC 3.5 K+P 9/69
R8 0.16 0.05 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R8 0.02 0.10 0.02 DEHM 74 DBC 0 4.6 K+ N 12/75*

R8 AVG 0.120 0.046 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
R8 STUDENT 0.126 0.045 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R8 FIT 0.117 0.032 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)

R9 K*(1420) INTO (K RHO) / (K*(892) PI) (P3)/(P2)
R9 (0.39) OR LESS BASSOMPIE 67 HBC + 5 K+ P 9/67
R9 (0.40) OR LESS CL=.90 FIELD 67 HBC - 3.8 K- P 6/67
R9 P 130 .13 .09 OTTER 75 HBC 08.10,16 K-P,K* N 12/75*

R10 K*(1420) INTO (K OMEGA) / (K*(892) PI) (P4)/(P2)
R10 Q (0.10) (0.04) FIELD 67 HBC - 3.8 K- P 6/67
R10 FIT 0.145 0.059 FROM FIT

R11 K*(1420) INTO (K ETA) / (K*(892) PI) (P5)/(P2)
R11 Q (0.07) (0.04) FIELD 67 HBC - 3.8 K- P 6/67
R11 FIT 0.064 0.066 FROM FIT

R12 K*(1420) INTO (K ETA) / (K PI) (P5)/(P1)
R12 (0.02) OR LESS BISHOP 69 HBC 3.5 K+ P 9/69
R12 (0.04) OR LESS CL=.95 AGUILAR 71 HBC 3.9-4.6 K- P 11/71
R12 FIT 0.035 0.036 FROM FIT

R Q FOLLOWING SUGGESTION BY AGUILAR 70, WE DO NOT MAKE USE OF MEASURE-
MENTS WHERE THE (K PI PI) BACKGROUND SUBTRACTION IS DIFFICULT DUE
TO THE NEARBY C REGION.

REFERENCES FOR K*(1420)

BADIER 65 PL 19 612 BADIER, DEMOULIN, GOLDBERG, (EPOL+SACL+ZEEHAN)
CHUNG 65 PRL 15 325 *DAHL, HARDY, HESS, JACOBS, KIRZ, MILLER (LRL)
FOCARDI 65 PL 16 351 FOCARDI, MINGUZZI, RANZI, SERRA+ (BOLOGNA+SACL)
SHEN 66 PRL 17 726 *BUTTERWORTH, FU, GOLDBERG, TRILLING (LRL)
ALSO 66 (PRIVATE COMMUN) GERSON, GOLDBER (LRL)
BASSANO 67 PRL 19 968 *GOLDBERG, GDZ, BARNES, LEITNER+ (BNL+SYRACUSE)
BASSOMPIE 67 PL 268 30 BASSOMPIERE, GOLDSCHMIDT+ (CERN+BRUX+BIRM) J P
CRENELL 67 PRL 19 66 *KALBFELDSCH, LAI, SCARR, SCHUMANN (BNL)
DAHL 67 PR 163 1377 *HARDY+HESS+KIRZ+MILLER (LRL)
ALSO 65 PRL 14 401 HARDY, CHUNG, DAHL, HESS, KIRZ, MILLER (LRL)
DE BAERE 67 NC 51 A 401 *GOLDSCHMIDT-CLERMONT, HENRI+ (BRUX+CERN)
FIELD 67 PL 248 638 *HENDRICKS+PICCIONI+YAGER (LAJOLLA)
GOLDBERG 67 PRL 19 972 G. GOLDBERGER, FIRESTONE, SHEN (LRL)
ADERHOLZ 68 NP B 5 567 *DEUTSCHMANN+ (AACH+BERL+CERN+LOIC+VIENNA)
ALSO 66 PL 22 357 BARTSCH, DEUTSCHMANN, MORRISON+ (ABCL (ICIV)
ANTICH 68 PRL 21 1862 *CALLAHAN, CARSON, COX, DENEGRI, (JHU)
DUBAL 68 THESIS 1456 L. DUBAL (GENEVE)
KANG 68 PR 176 1587 Y. W. KANG (IOWA)
SCHWEINER 68 PR 166 1317 S. SCHWEINGRUBER, DERRICK, FIELDS+ (ANL+NWES)
ALSO 67 THESIS F. L. SCHWEINGRUBER (NORTH-WESTERN, EVANSTON)
BASSOMPIE 69 NP B 13 189 BASSOMPIERE, GOLDSCHMIDT-CLERM.+ (CERN+BRUX) J P
BISHOP 69 NP B 9 403 *GOSHAW, ERWIN, WALKER (WISC)
CRENELL 69 PRL 22 487 *KARSHON, LAI, ONEALL, SCARR (BNL)
DAVIS 69 PRL 23 1071 *DERENZO, FLATTE, ALSTON, LYNCH, SOLMITZ (LRL)
DE BAERE 69 NC 61 A 397 *GOLDSCHMIDT-CLERMONT, HENRI, + (BELG+CERN)
FRIEDMAN 69 UCRL-18860 J. FRIEDMAN, PH. D. THESIS (LRL)
LIND 69 NP B 14 1 *ALEXANDER, FIRESTONE, FU, GOLDBERGER (LRL) J P
ABRAMS 70 PR D 1 2433 *EISENBERG, KIM, MARSHALL, C. HALLORAN, + (ILL)
AGUILAR 70 PRL 25 1362 AGUILAR-BENITZ, BASSANO, EISNER, + (BNL+GLAS)
BIRMINGH 70 KIEV CONF. *ASTIER, RAPPORTEURS TALK (BIRM+PLAS+OXF)
AGUILAR 71 PR D 4 2583 *EISENER, KINSCH (BNL)
BARNHAM 71 NP B 28 171 *COLLEY, JOBS, GLAS, (BNL)
CORDS 71 PR D 4 1974 *CARMONY, ERWIN, MEIERE, + (PURD+UCD+IUPUI)
SLATTERY 71 UR-875-332 (PREP) P. SLATTERY, A REVIEW OF STRANGE MESONS (ROCH)
BUCHNER 72 NP B 45 333 *DEHM, CHARRIERE, CORNET, + (MPIM+CERN+BRUX)
CRENELL 72 PR D 6 1220 *GORDON, KWAN-WU LAI, SCARR (BNL)
DEUTSCHMANN 72 NP B 36 373 DEUTSCHMANN, + (ABCLV COLLABORATION)
ENGELMANN 72 PR D 5 2162 ENGELMANN, MUSGRAVE, FORMAN, + (ANL+EFI)
FRATI 72 NP B 5 2361 *HALPERN, HARGIS, SNAPE, CARNAHAN, + (PENN+CINC)
ROUGE 72 NP B 46 29 *VIDEAU, VOLTE, DE BRION, + (EPCL+SACL)
TIECKE 72 NP B 39 596 *GRIJNS, HEINEN, DE GROOT, + (NIJN+ZEEM)
CHARRIERE 73 NP B 51 317 CHARRIERE, ORJARD, DE BAERE, + (CERN+BELG)
ALSO 73 (PRIVATE COMMUNICAT) GOLDSCHMIDT-CLERMONT (CERN)
CLARK 73 NP B 54 432 *LYONS, RADJICIC (OXFORD)
DE JONGH 73 NP B 58 110 *CORNET, CHARRIERE, + (BRUX+MONS+CERN+MPIM)
LINGLIN 73 NP B 5 408 D. LINGLIN (CERN)
WALUCH 73 PR D 8 2837 *FLATTE, FRIEDMAN (LRL)
DEHM 74 NP B 75 47 *GOEBEL, WITTEK, WOLF, + (MPIM+BRUX+MONS+CERN)
CHUNG 74 PL 518 413 *EISNER, PROTOPESCU, SAMIOS, STRAND (BNL)
ANTIPOV 75 NP B 86 381 *ASCOLI, BUSNELL, KIENZLE+ (SERP+CERN+ILL)
LAUSCHER 75 NP B 86 189 *OTTER, WIECZOREK, + (ABCLV COLLABORATION) J P
MCUBBIN 75 NP B 86 13 N. A. MCCUBBIN, L. LYONS (OXF)
OTTER 75 NP B 84 333 * (AACH+BERL+CERN+LOIC+VIEN+ATHU+AATEN+IUPUI)
BRANDENB 76 PL BRANDENBURG, CARNEGIE, CASHMORE, DAVERI+ (SLAC) J P

KN(1700)

27 KN(1700, JP= 1) I = 1/2
THIS ENTRY CONTAINS VARIOUS PEAKS IN STRANGE MESON
SYSTEMS REPORTED IN THE 1700 MEV REGION.
EVIDENCE NOT COMPELLING, OMITTED FROM TABLE.

27 KN(1700) MASS (MEV)
M (1660.0) CARMONY 67 HBC + 3.8 K-P, OMEGA K 11/67
M J (1660.0) (10.0) JOBS 67 HBC + 5 K+ P 12/75*
M J CLAIMED BY JOBS IN (K PI), (K*(892) PI), AND (K*(1420) PI)
M J MODES, K PI BUMP INTERFERES MOSTLY WITH DELTA(1236).
M (1660.) CHARRIERE 73 HBC 0 5 K+ P, K P 3PI 1/73
M 60(1710.) (15.) CHUNG 74 HBC - 7.3K-P, K- OMEGA P 12/75*

27 KN(1700) WIDTH (MEV)
W (60.0) (20.0) JOBS 67 HBC + 5 K+ P 12/75*
W (60.) CHARRIERE 73 HBC 0 5 K+ P, K P 3PI 1/73
W 60 (110.) (50.) CHUNG 74 HBC - 7.3K-P, K- OMEGA P 12/75*

27 KN(1700) PARTIAL DECAY MODES
P1 KN(1700) INTO K PI 493+ 139
P2 KN(1700) INTO K PI PI 493+ 139+ 139
P3 KN(1700) INTO K*(892) PI 892+ 139
P4 KN(1700) INTO K RHO 493+ 773
P5 KN(1700) INTO K*(1420) PI 1421+ 139
P6 KN(1700) INTO K OMEGA 493+ 783

27 KN(1700) BRANCHING RATIOS
R1 KN(1700) INTO (K PI)/(K OMEGA) (P1)/(P6)
R1 N (0.5) (0.5) CHUNG 74 HBC - 7.3 K-P, K- P 12/75*
R1 N NO K PI SIGNAL SEEN IN THIS EXPERIMENT

REFERENCES FOR KN(1700)

CARMONY 67 PRL 18 615 D. CARMONY, T. HENDRICKS, L. LANDER (LA JOLLA)
JOBS 67 PL 268 49 *BASSOMPIERE, DE BAERE + (BIRM+CERN+BRUX)
CHARRIERE 73 NP B 51 317 CHARRIERE, ORJARD, DE BAERE, + (CERN+BELG)
CHUNG 74 PL 518 412 *EISNER, PROTOPESCU, SAMIOS, STRAND (BNL)

L(1770)

23 L(1770, JP= 1) I = 1/2

The L(1770) is seen as a bump in the diffractive-like process KN -> (KPI)N. BARBARO 69 and FIRESTONE 72 find the decay is consistent with being entirely K*(1420)pi, whereas AGUILAR 70, BARTSCH 70, COLLEY 71, and DENEGRI 71 present evidence for alternate decay modes. For a review see EISNER 74.

Recent partial-wave analyses (DEUTSCHMANN 74, ANTIPOV 75, OTTER 75) have shown that the situation in the L region is complicated, with many waves contributing. The 2- K*(1420)pi S wave is important, but cannot explain the whole L enhancement (DEUTSCHMANN 74). The phase variation of the 2- wave shows no evidence for any significant resonance contribution with Gamma < 300 MeV (OTTER 75).

23 L MASS (MEV)
M 20(1780.) BERLINGHI 67 HBC + 12.7 K+P 7/67
M (1760.0) (15.0) JOBS 67 HBC + 5 K+ P 1/73
M 1745.0 20.0 AGUILAR 70 HBC - 4.6 K- P 6/70
M 1780.0 15.0 BARTSCH 70 HBC - 10.1 K- P 1/71
M (1760.0) (15.0) LUDLAM 70 HBC - 12.6 K- P 1/73
M X 1765.0 40.0 COLLEY 71 HBC + 10.4 K+P, K 2PI 1/73
M X SYSTEMATIC ERRORS ADDED CORRESP. TO SPREAD OF DIFFERENT FITS.
M (1740.0) DENEGRI 71 DBC - 12.6 K-D, K 2PI D 5/71
M 1767. 6. BLIEDEN 72 MMS - 11. -16. K- P 12/72
M P 306 1730. 20. FIRESTONE 72 DBC + 12. K+ D 1/73
M P PRODUCED IN CONJUNCTION WITH D*
M AVG 1764.6 6.7 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)
M STUDENT1765.0 5.8 AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT

Data Card Listings

For notation, see key at front of Listings.

Mesons

L(1770), K_N(1800), K*(2200)

Table with 5 columns: W, L WIDTH (MEV), BERLINGHI, JORES, AGUILAR, BARTSCH, LUDLAM, COLLEY, DENEGRI, FIREDSTONE. Includes average values and error scales.

Table with 3 columns: P1, L INTO K PI PI, DECAY MASSES. Lists partial decay modes and associated masses.

Table with 3 columns: R1, L INTO K*(1420) PI / (K PI PI), (P2)/(P1). Lists branching ratios and experimental evidence.

Table with 3 columns: Author, PRL, PL, REFERENCE FOR L(1770). Lists various experimental references for the L(1770) meson.

K_N(1800) 60 KN(1800, JP=3-)

This entry includes peaks of low statistical significance seen in the K_N and K_N spectra in the region of 1800 MeV (CARMONY 71, AGUILAR 73, SPIRO 76). FIRESTONE 71 and BRANDENBURG 76 observe structure in the K_N angular distribution, the simplest explanation for which is a rapid rise of the F-wave amplitude around 1800 MeV, interfering strongly with other waves. BRANDENBURG 76 propose the existence of a J^P=3⁻ resonance, but we prefer to wait for further confirmation before including this entry in the Table.

Table with 5 columns: M, C, DISAGREEMENT BETWEEN THE FIT AND DATA ON BOTH SIDES OF THE SIGNAL, APPROX., APPARENT INTERFERENCE WITH OTHER AMPLITUDES PRECLUDES, PRECISE DETERMINATION. Includes CARMONY 71, AGUILAR 73, BRANDENB 76, SPIRO 76.

Table with 5 columns: W, C, DISAGREEMENT BETWEEN THE FIT AND DATA ON BOTH SIDES OF THE SIGNAL, APPROX., OR LESS, SEE NOTE A ABOVE. Includes CARMONY 71, AGUILAR 73, BRANDENB 76, SPIRO 76.

Table with 3 columns: P1, KN(1800) INTO K PI, DECAY MASSES. Lists partial decay modes and masses for KN(1800).

Table with 3 columns: R1, KN(1800) INTO (K PI)/(K*(892) PI + K RHO), (P1)/(P2+P3). Lists branching ratios and statistical compatibility.

Table with 3 columns: Author, PRL, PL, REFERENCE FOR KN(1800). Lists references for the KN(1800) meson.

K*(2200) 40 K*(2200, JP=) ENHANCEMENT SEEN IN (ANTI)HYPERON-NUCLEON MASS NEAR THRESHOLD. INTERPRETATION UNCERTAIN. OMITTED FROM TABLE.

Table with 5 columns: M, C, 20, 20, LISSAUER 70 HBC, SLATTERY 71 RVUE. Includes CARMONY 71, LISSAUER 70, SLATTERY 71, SPIRO 76.

Table with 5 columns: W, C, 20, 20, LISSAUER 70 HBC, SLATTERY 71 RVUE. Includes CARMONY 71, LISSAUER 70, SLATTERY 71, SPIRO 76.

Table with 3 columns: Author, PRL, PL, REFERENCE FOR K*(2200). Lists references for the K*(2200) meson.

Mesons

EXOTICS, NEW HEAVY MESONS

Data Card Listings

For notation, see key at front of Listings.

EXOTIC MESONS

EXOTICS

70 EXOTICS

THE PURPOSE OF THIS ENTRY IS TO PROVIDE A LIST OF REFERENCES FOR EXOTIC MESON SEARCHES (SEE MAIN TEXT, SEC. III AND TABLE I), AS WELL AS THEORETICALLY BASED SUGGESTIONS FOR EXPERIMENTS. NOTE THAT LIPKIN 73 PROPOSES EXPERIMENTS WHICH ARE CONCLUSIVE EVEN IF NEGATIVE RESULTS ARE OBTAINED.

REFERENCES FOR EXOTICS

REPORTS ON SEARCHES

ROSENFEL	68 PHILA.CONF.P.455	A.H.ROSENFELD	(LRL)
ODD	69 PR 177 1991	+JGLDERSMA, PALMER, SAMIOS	(BNL)
CHO	70 PL 32 B 409	+DERRICK, JOHNSON, MUSGRAVE, +	(ANL+NWES+KANS)
GIACOMEL	70 PL 33 B 373	G.GIACOMELLI +	(BGNA+SACL+ZEEP+REHO+EPOL)
LYS	70 PR D 2 2525	J.LYS +	(MICH)
ROSNER	70 EXP.MESON SPECTROSCOPY, ED. C. BALTAY AND A.H.ROSENFELD, P.499		
BUHL	72 NP B 37 421	+CLINE, TERRELL	(WISCONSIN)
COHEN	73 NP B 53 1	+FERBEL, SLATTERY, WERNER	(ROCHESTER)
ALAM	74 PL 53 B 207	+BRABSON, GALLOWAY, +	(IND+PURD+SLAC+VAND)
COHEN	74 BOSTON	D.COHEN REVIEW TALK	(COLU)
OREN	74 NP B 71 189	+COOPER, FIELDS, RHINES, WHITMORE, +	(ANL+OXF)
BALTAY	75 PL 57 B 293	+CAUTIS, COHEN, KALEL KAR, PISELLO, +	(COLU+BING)
DAVIS	75 NP B 96 426	+AMAR, KRUPAC, YARGER, +	(KANS+CCAC+ANL)

SUGGESTIONS FOR SEARCHES

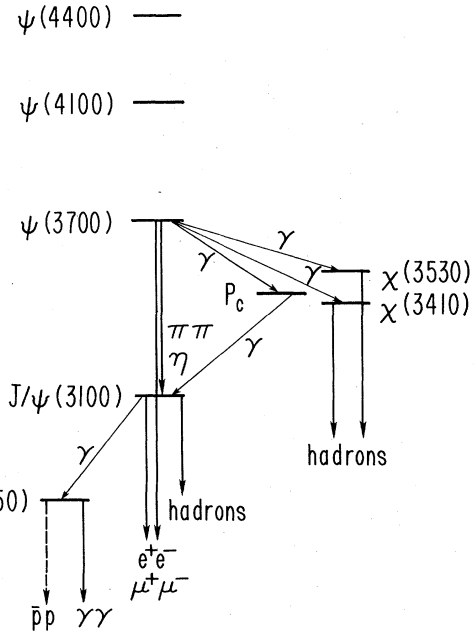
ROSNER	68 PRL 21 950+1468	J.L.ROSNER	(TEL-AVIV)
ROSNER	70 EXP.MESON SPECTROSCOPY, ED. C.BALTAY AND A.H.ROSENFELD, P.499		
FATMAN	73 PL 43 B 307	D.FATMAN, G.GOLDBERGER, Y.ZARMI	(CERN)
LIPKIN	73 PR D 7 2262	H.J.LIPKIN	(ARGONNE+FNAL)

NEW HEAVY MESONS

The New Heavy Mesons

Since the discovery of the $J/\psi(3100)$ particle in p-Be collisions at the AGS (AUBERT 74), and in e^+e^- collisions at SPEAR (AUGUSTIN 74), the size of the family of "new particles" has increased to at least 8 members, with a complex set of decay modes and mutual transitions; the main features of the experimental situation are illustrated in the figure below.

In the following entries we compile the experimental data as of January 1976, with only a minimum of mini-reviews and footnotes. For detailed discussions of the implications and possible interpretations of the properties of the new particles, see, e.g., the Proceedings of the 1975 International Symposium on Lepton and Photon Interactions at High Energies, Stanford, 1975 [W. T. Kirk (SLAC), editor].



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The New Heavy Mesons

Extracting Resonance Widths from e^+e^- Colliding Beam Formation Experiments

In an e^+e^- colliding beam formation experiment, the true shape of an observed resonance is distorted primarily by the effects of (1) soft-photon processes, and (2) beam energy spread due to processes such as quantum fluctuations in emission of synchrotron radiation. The spread in energy due to (2) may usually be approximated by a Gaussian distribution, the effect of which vanishes rapidly at energies sufficiently removed from resonance. The major effect of (1) is a decrease in the effective c.m. energy for some fraction of the collisions, because of the emission of bremsstrahlung by the electron or positron before annihilation. Hence, though the nominal beam energy may be well above the resonance region, a certain fraction of the collisions occur at or near resonance. This gives rise to the well-known high-mass radiative tails of the $J/\psi(3100)$ and $\psi(3700)$ resonances.

Data Card Listings

For notation, see key at front of Listings.

Mesons

NEW HEAVY MESONS, J/ψ(3100)

Because of these effects, the most reliable means for determining resonance widths is to use a method based on the area under the line shape [see, e.g., "Notes from the SLAC Theory Workshop on the ψ", ed. R. Pearson, SLAC-PUB-1515 (1974)]. This method, familiar in nuclear physics, minimizes the sensitivity to the details of the beam energy spread. Corrections for the radiative processes, which depend on the limits of integration of the areas, still need to be made. This discussion assumes the resolution is adequate to allow a reasonable separation of signal from background (which itself is subject to radiative processes).

For formation of a resonance of mass M in e⁺e⁻ collisions, with subsequent decay via channel i,

$$\sigma_i(W) = \sigma_0 \frac{\Gamma_e \Gamma_i / 4}{(M-W)^2 + \Gamma^2 / 4}$$

where W is the total center of mass energy; Γ, Γ_e, and Γ_i are the total width and partial widths for coupling to e⁺e⁻ and channel i, respectively; and a Breit-Wigner line shape with energy-independent partial widths is assumed. The quantity σ₀ is given by

$$\sigma_0(W) = \frac{4\pi(2J+1)}{W^2}$$

where J is the spin of the resonance. For a narrow resonance, the area under the resonant line is given by

$$A_i = \frac{\pi}{2} \frac{\Gamma_e \Gamma_i}{\Gamma} \sigma_0(M)$$

independent of the energy resolution of the apparatus. Determination of the mass, spin, and integrated channel cross sections, A_i, allows, then, determination of the quantities Γ_eΓ_i/Γ, and, assuming there are no undetected decay modes (i.e., in our case Γ = Γ_{e⁺e⁻} + Γ_{μ⁺μ⁻} + Γ_{hadrons}), determination of the total and partial widths.

In the Data Card Listings, we tabulate and average the quantity Γ_eΓ_i/Γ only when it was not used to determine the partial widths and/or branching ratios.

J/ψ(3100)

70 J/PSI(3100,JPG=1--1) 1=0

The J/ψ(3100) was discovered in p-Be collisions at the AGS (AUBERT 74), and in e⁺e⁻ collisions at SPEAR (AUGUSTIN 74), and since then has been extensively investigated in e⁺e⁻ collisions at SPEAR, DORIS, and ADONE, as well as at conventional accelerators with photon and hadron beams.

Evidence bearing on the hadronic nature of the J/ψ(3100) comes from the observation of a large forward photoproduction cross section. Using vector dominance arguments one derives a total cross section for J/ψ(3100) on nucleons over a large energy range of the order of 1 mb (ANDREWS 75, DAKIN 75, KNAPP1 75, MARTIN 75). Although this cross section is an order of magnitude smaller than the corresponding cross sections of the well known vector mesons, it is large enough to suggest that the J/ψ(3100) is probably a hadron. The hadronic interpretation of the J/ψ(3100) is also supported by the apparent conservation of isospin and G-parity in direct hadronic decays (JEAN-MARIE 75). For a detailed discussion of the hadronic nature of J/ψ(3100), see HARARI 75.

The J^{PC} = 1⁻⁻ assignment, suggested by production in s-channel e⁺e⁻ collisions, is confirmed by the observation of an interference between the resonant and QED amplitudes, and by the angular distribution of the lepton pairs in the final state (BOYARSKI 75). The I^G = 0⁻ assignment was determined by a study of multipion decays (JEAN-MARIE 75). These quantum numbers are compatible with all the data on partial decay widths (non-observation of the γγ mode, observation of the ΛΛ̄ mode, etc.).

70 J/PSI(3100) MASS (MEV)

M	(3100.)		AUBERT	74	SPEC	28.	PP(E+E-)	2/75*
M	L	(3105.)	(3.)	AUGUSTIN	74	SPEAR	E+E-	2/75*
M	D	3095.	4.	BOYARSKI	75	SPEAR	E+E-	3/75*
M	S	(3089.5)		CRIGGEE	75	DORIS	E+E-	2/75*
M		3098.	6.	PREPOST	75	SPEC	13.-21.GAMMA D	1/76*
M		3103.	6.	BEMPORAD	75	ADONE	E+E-	1/76*
M	L	BOYARSKI 75 IS A REEVALUATION OF AUGUSTIN 74 BASED ON A RECALIBRATION OF THE SPEAR BEAM ENERGY.						
M	D	MASS, WIDTH, PARTIAL WIDTHS, AND BRANCHING RATIOS ALL OBTAINED FROM ONE OVERALL FIT TO DATA OF THIS EXPERIMENT.						
M	S	ERROR OF ABOUT 1 PER CENT FROM THE UNCERTAINTY IN CALIBRATION OF THE BEAM ENERGY.						
M	AVG	3097.6	2.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)				3/75*
M	STUDENT	3097.6	3.3	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT				2/75*

70 J/PSI(3100) WIDTH (KEV)

SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.									
W	69.	15.	BOYARSKI	75	SMAG	E+E-	3/75*		
W	68.	26.	BALDINI1	75	FRAG	E+E-	1/76*		
W	60.	25.	ESPOSITO	75	FRAM	E+E-	1/76*		
W	66.9	11.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)						
W	STUDENT	66.9	12.4	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					

Mesons

J/ψ(3100)

Data Card Listings

For notation, see key at front of Listings.

70 J/PSI(3100) PARTIAL DECAY MODES

		DECAY MASSES		
P1	J/PSI(3100) INTO E+ E-	.5+ .5		
P2	J/PSI(3100) INTO MU+ MU-	105+ 105		
P3	J/PSI(3100) INTO HADRONS			
P4	J/PSI(3100) INTO GAMMA INTO HADRONS			
HADRONIC DECAYS				
P	HADRONIC DECAYS			
P11	J/PSI(3100) INTO PI+ PI-			
P12	J/PSI(3100) INTO PI+ PI- PI0			
P13	J/PSI(3100) INTO 2(PI+ PI-)			
P14	J/PSI(3100) INTO 2(PI+ PI-) PI0			
P15	J/PSI(3100) INTO 3(PI+ PI-)			
P16	J/PSI(3100) INTO 3(PI+ PI-) PI0			
P17	J/PSI(3100) INTO 4(PI+ PI-)			
P18	J/PSI(3100) INTO 4(PI+ PI-) PI0			
P19	J/PSI(3100) INTO K0S K0L			
P20	J/PSI(3100) INTO K+ K-			
P21	J/PSI(3100) INTO PI+ PI- K+ K-			
P22	J/PSI(3100) INTO 2(PI+ PI-) K+ K-			
P23	J/PSI(3100) INTO RHO PI			
P24	J/PSI(3100) INTO RHO PI PI			
P25	J/PSI(3100) INTO OMEGA PI PI			
P26	J/PSI(3100) INTO PHI PI PI			
P27	J/PSI(3100) INTO K K*(892)			
P28	J/PSI(3100) INTO K K*(1420)			
P29	J/PSI(3100) INTO K*(892) K*(892)			
P30	J/PSI(3100) INTO K*(1420) K*(1420)			
P31	J/PSI(3100) INTO K*(892) K*(1420)			
P32	J/PSI(3100) INTO PEAR P			
P33	J/PSI(3100) INTO LAMBDA ANTILAMBDA			
P34	J/PSI(3100) INTO NUCLEON ANTINUCLEON PI			
P35	J/PSI(3100) INTO P PEAR PI+ PI-			
P36	J/PSI(3100) INTO P PBAR PI+ PI- NEUTRALS			
RADIATIVE DECAYS				
P	RADIATIVE DECAYS			
P	RADIATIVE DECAYS			
P51	J/PSI(3100) INTO GAMMA GAMMA			
P52	J/PSI(3100) INTO PI0 GAMMA			
P53	J/PSI(3100) INTO ETA GAMMA			
P54	J/PSI(3100) INTO ETA PRIME GAMMA			
P55	J/PSI(3100) INTO X(2750) GAMMA	2750+ 0		

70 J/PSI(3100) PARTIAL WIDTHS (KEV)

SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.

W1	J/PSI(3100) INTO E+ E-	(G1)		
W1	4.8	0.6	BOYARSKI 75 SMAG	E+E- 2/75*
W1 B	(4.6)	(.8)	BALDINI 75 FRAG	E+E- 3/75*
W1 B	ASSUMING EQUAL PARTIAL WIDTHS FOR (E+E-) AND (MU+MU-)			
W1	4.6	1.0	ESPOSITO 75 FRAM	E+E- 1/76*
W1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W1	4.75	0.51	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
W1	4.75	0.55	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
G2				
W2	J/PSI(3100) INTO MU+ MU-			
W2	4.8	0.6	BOYARSKI 75 SMAG	E+E- 2/75*
W2	5.0	1.0	ESPOSITO 75 FRAM	E+E- 3/75*
W2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W2	4.85	0.51	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
W2	4.85	0.55	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
G3				
W3	J/PSI(3100) INTO HADRONS			
W3	59	14.	BOYARSKI 75 SMAG	E+E- 2/75*
W3	59.	25.	BALDINI 75 FRAG	E+E- 3/75*
W3	50.	25.	ESPOSITO 75 FRAM	E+E- 1/76*
W3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
W3	57.3	10.9	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
W3	57.3	11.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
G4				
W4	J/PSI(3100) INTO GAMMA INTO HADRONS			
W4	12.	2.	BOYARSKI 75 SMAG	E+E- 1/76*

70 J/PSI(3100) BRANCHING RATIOS

FOR THE BRANCHING RATIOS R1 - R4, SEE ALSO THE PARTIAL WIDTHS ABOVE, AND (PARTIAL WIDTHS)*R1 BELOW.

R1	J/PSI(3100) INTO (E+ E-)/TOTAL	(P1)		
R1	0.069	0.009	BOYARSKI 75 SMAG	E+E- 3/75*
R2	J/PSI(3100) INTO (MU+ MU-)/TOTAL			
R2	0.069	0.009	BOYARSKI 75 SMAG	E+E- 3/75*
R3	J/PSI(3100) INTO (HADRONS)/TOTAL			
R3	0.86	0.02	BOYARSKI 75 SMAG	E+E- 3/75*
R4	J/PSI(3100) INTO (E+ E-)/(MU+ MU-)			
R4	1.00	0.05	BOYARSKI 75 SMAG	E+E- 2/75*
R4	0.93	0.10	FORD 75 SPEC	E+E- 3/75*
R4	.91	.15	ESPOSITO 75 FRAM	E+E- 1/76*
R4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
R4	0.980	0.043	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R4	0.980	0.047	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT	
R5	J/PSI(3100) INTO (GAMMA INTO HADRONS)/TOTAL			
R5	.17	.02	BOYARSKI 75 SMAG	E+E- 1/76*

HADRONIC DECAYS

R8	J/PSI(3100) INTO (PI+ PI-)/TOTAL			
R8	(.0003)OR LESS CL=.90	WIJK 75 DASP	E+E-	1/76*
R9	J/PSI(3100) INTO (2(PI+ PI-))/TOTAL			
R9	76	.004	JEAN-MARI 75 SMAG	E+E- 1/76*
R10	J/PSI(3100) INTO (2(PI+ PI-) PI0)/TOTAL			
R10	675	.04	JEAN-MARI 75 SMAG	E+E- 1/76*
R11	J/PSI(3100) INTO (3(PI+ PI-))/TOTAL			
R11	32	.004	JEAN-MARI 75 SMAG	E+E- 1/76*

R12	J/PSI(3100) INTO (3(PI+ PI-) PI0)/TOTAL			
R12	181	.029	JEAN-MARI 75 SMAG	E+E- 1/76*
R13	J/PSI(3100) INTO (4(PI+ PI-) PI0)/TOTAL			
R13	13	.009	JEAN-MARI 75 SMAG	E+E- 1/76*
R14	J/PSI(3100) INTO (PI+ PI- K+ K-)/TOTAL			
R14	A 83	.005	ABRAMS4 75 SMAG	E+E- 1/76*
R14	A	INCLUDING CONTRIBUTION FROM K*(892)K*(1420)		
R15	J/PSI(3100) INTO (2(PI+ PI-) K+ K-)/TOTAL			
R15		.003	ABRAMS4 75 SMAG	E+E- 1/76*
R16	J/PSI(3100) INTO (RHO PI)/(PI+ PI- PI0)			
R16	(.7)	OR MORE	JEAN-MARI 75 SMAG	E+E- 1/76*
R17	J/PSI(3100) INTO (RHOO PI0)/(RHOO+ PI++)			
R17		.59	JEAN-MARI 75 SMAG	E+E- 1/76*
R18	J/PSI(3100) INTO (RHO PI)/TOTAL			
R18	153	.013	JEAN-MARI 75 SMAG	E+E- 1/76*
R19	J/PSI(3100) INTO (OMEGA PI PI)/(2(PI+ PI-) PI0)			
R19	J	(.2)	JEAN-MARI 75 SMAG	E+E- 1/76*
R20	J/PSI(3100) INTO (RHO PI PI PI)/(2(PI+ PI-) PI0)			
R20	J	(.3)	JEAN-MARI 75 SMAG	E+E- 1/76*
R20	FINAL STATE 2(PI+PI-)PI0			
R21	J/PSI(3100) INTO (PHI PI+ PI-)/(OMEGA PI+ PI-)			
R21		.2	ABRAMS4 75 SMAG	E+E- 1/76*
R22	J/PSI(3100) INTO (K0S K0L)/TOTAL			
R22		(.0002)OR LESS CL=.90	ABRAMS4 75 SMAG	E+E- 1/76*
R23	J/PSI(3100) INTO (K+ K-)/TOTAL			
R23		(.0006)OR LESS CL=.90	WIJK 75 DASP	E+E- 1/76*
R24	J/PSI(3100) INTO (K0 K*(892))/TOTAL			
R24	57	.0024	ABRAMS4 75 SMAG	E+E- 1/76*
R25	J/PSI(3100) INTO (K+ K*(892))/TOTAL			
R25	87	.0031	ABRAMS4 75 SMAG	E+E- 1/76*
R25		.0007	WIJK 75 DASP	E+E- 1/76*
R25	POSSIBLY SEEN			
R26	J/PSI(3100) INTO (K0 K*(1420))/TOTAL			
R26		(.0019)OR LESS CL=.90	ABRAMS4 75 SMAG	E+E- 1/76*
R27	J/PSI(3100) INTO (K+ K*(1420))/TOTAL			
R27		(.0019)OR LESS CL=.90	ABRAMS4 75 SMAG	E+E- 1/76*
R28	J/PSI(3100) INTO (K*(892)0 K*(892)0)/TOTAL			
R28		(.0006)OR LESS CL=.90	ABRAMS4 75 SMAG	E+E- 1/76*
R29	J/PSI(3100) INTO (K*(1420)0 K*(1420)0)/TOTAL			
R29		(.0018)OR LESS CL=.90	ABRAMS4 75 SMAG	E+E- 1/76*
R30	J/PSI(3100) INTO (K*(892)0 K*(1420)0)/TOTAL			
R30	30	.0037	ABRAMS4 75 SMAG	E+E- 1/76*
R31	J/PSI(3100) INTO (PEAR PI)/TOTAL			
R31	A 105	.0021	ABRAMS4 75 SMAG	E+ E- 1/76*
R31	A 40	.0023	WIJK 75 DASP	E+ E- 1/76*
R31	A	ASSUMING ANGULAR DISTRIBUTION (1.+COS(THETA)**2)		
R31	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)			
R31	STUDENT	0.00216	0.00036	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT
R32	J/PSI(3100) INTO (PEAR PI)/(MU+ MU-)			
R32	A 20	.051	CRTEGEE 75 PLUT	E+ E- 1/76*
R32	A	ASSUMING ANGULAR DISTRIBUTION (1.+COS(THETA)**2)		
R33	J/PSI(3100) INTO (LAMBDA ANTILAMBDA)/TOTAL			
R33	19	.0016	ABRAMS4 75 SMAG	E+ E- 1/76*
R34	J/PSI(3100) INTO (NUCLEON ANTINUCLEON PI)/TOTAL			
R34	A 87	.0037	ABRAMS4 75 SMAG	E+ E- 1/76*
R34	A	INCLUDES P PBAR PI0, NBAR P PI- AND N PBAR PI+		
R35	J/PSI(3100) INTO (P PBAR PI+PI-)/TOTAL			
R35	G 125	SEEN	GOLDBABER 75 SMAG	E+E- 2/76*
R35	G	INCLUDES LAMBDA ANTILAMBDA		
R36	J/PSI(3100) INTO (PEAR PI+ PI- NEUTRALS)/TOTAL			
R36	I 91	SEEN	GOLDBABER 75 SMAG	E+E- 2/76*
R36	I	INCLUDES SIGMA ANTISIGMA ETC.		
R37	J/PSI(3100) INTO (LAMBDA ANTISIGMA)/(LAMBDA ANTILAMBDA)			
R37		(.22) OR LESS	GOLDBABER 75 SMAG	E+E- 2/76*
RADIATIVE DECAYS				
R	RADIATIVE DECAYS			
R	RADIATIVE DECAYS			
R51	J/PSI(3100) INTO (2 GAMMA)/(E+ E-)			
R51		(.050)OR LESS CL=.90	BRAUNSCH 75 DASP	(P4)/(P1) E+E- 2/75*
R52	J/PSI(3100) INTO (PI0 GAMMA)/(E+ E-)			
R52		(.0.13) OR LESS CL=.90	BRAUNSCH 75 DASP	(P5)/(P1) E+E- 2/75*
R52	B	(.06) OR LESS CL=.90	BACCI 75 FRAG	E+E- 1/76*
R52	B	RE-STAT'ED BY US USING (E+E-)/HADRONS = .08		
R53	J/PSI(3100) INTO (ETA GAMMA)/TOTAL			
R53	U	.0014	WIJK 75 DASP	E+ E- 1/76*
R53	U	USING TOTAL WIDTH (69+-15) KEV		
R53		(.016)OR LESS CL=.90	BACCI 75 FRAG	E+E- 1/76*
R54	J/PSI(3100) INTO (ETA PRIME GAMMA)/(ETA GAMMA)			
R54		(5.) OR LESS CL=.90	WIJK 75 DASP	E+E- 1/76*
R54	B	4.0	HEINTZE 75 DESY	E+E- 1/76*
R54	B	(10.) OR LESS CL=.90	BALDINI2 75 FRAG	E+E- 1/76*
R54	B	RE-STAT'ED BY US USING (ETA GAMMA)/TOTAL = .0014		
R56	J/PSI(3100) INTO (X(2750) GAMMA)/(ETA GAMMA) FINAL STATE (3 GAMMA)			
R56	16	(1.)	HEINTZE 75 DESY	E+E- 1/76*

Data Card Listings

For notation, see key at front of Listings.

Mesons

J/ψ(3100), ψ(3700)

Table with columns for particle name, G-parity, JPC, and various experimental references. Includes entries for J/ψ(3100) and ψ(3700) with associated decay widths and branching ratios.

ψ(3700) 71 PSI(3700, JPC=1--1) I=0
The ψ(3700) was discovered in e+e- collisions at SPEAR (ABRAMS 74), and since then has been extensively investigated in e+e- collisions at SPEAR and DORIS, as well as at conventional accelerators.
The hadronic interpretation of the ψ(3700) is closely related to that of J/ψ(3100), due to the decays ψ(3700) → J/ψ(3100)ππ and ψ(3700) → J/ψ(3100)η. These decays also determine the IG = 0- assignment for the ψ(3700) (TANENBAUM 75).
The JPC assignment, suggested by production in s-channel e+e- collisions, is confirmed by the observation of an interference between the reso-

nant and QED amplitudes, and by the angular distribution of the lepton pairs in the final state (LÜTH 75). These quantum numbers are compatible with all the data on partial decay widths (the multiplication decays mentioned above, the non-observation of J/ψ(3100)π0 and γγ modes, etc.).

71 PSI(3700) MASS (MEV)
Table listing mass measurements from various experiments including ABRAMS 74, CRIEGEE 75, LUTH 75, and PREPOST 75. Includes error bars and average values.

71 PSI(3700) WIDTH (KEV)
Table listing width measurements from experiments like LUTH 75 and SMAG.

71 PSI(3700) PARTIAL DECAY MODES
Table listing various decay channels such as INTO E+ E-, INTO MU+ MU-, INTO GAMMA INTO HADRONS, etc.

71 PSI(3700) PARTIAL DECAY MODES (continued)
Table listing radiative decays and other partial decay modes with branching fractions and correlation coefficients.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS
The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, Pij, as follows: The diagonal elements are Pij ± δPij, where δPij = sqrt(δPij δPij), while the off-diagonal elements are the normalized correlation coefficients (δPij δPkl) / (δPij δPkl). For the definitions of the individual Pij, see the listings above; only those Pij appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.
Table with columns P13, P14, P15, P58, POTHER and corresponding values.

71 PSI(3700) PARTIAL WIDTHS (KEV)
Table listing partial widths for decays into e+e-, SMAG, and HAERONS with associated G-parity and I-spin values.

71 PSI(3700) BRANCHING RATIOS
Table listing branching ratios for decays into e+e- and SMAG, including a note about the overall fit.

Mesons

$\psi(3700)$

Data Card Listings

For notation, see key at front of Listings.

R2	PSI(3700) INTO (MU+ MU-)/TOTAL					
R2 H	.0084 .0018	HILGER	75 SPEC	E+E-	1/76*	
R2 H	RE-STAT'D BY US USING (J/PSI(3100)+ANYTHING)/TOTAL = .6					
R3	PSI(3700) INTO (HADRONS)/TOTAL					
R3	.981 .003	LUTH	75 SMAG	E+E-	1/76*	
R4	PSI(3700) INTO (MU+ MU-)/(E+ E-)					
R4	.89 .16	LUTH1	75 SMAG	E+E-	1/76*	
R5	PSI(3700) INTO (GAMMA INTO HADRONS)/TOTAL					
R5	.029 .004	LUTH	75 SMAG	E+E-	1/76*	
R	DECAYS INTO J/PSI(3100) + ANYTHING					
R10	PSI(3700) INTO (J/PSI(3100) + ANYTHING)/TOTAL					
R10	.57 .08	ABRAMS4	75 SMAG	E+E-	1/76*	
R10	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R10	0.581 0.045					
R11	PSI(3700) INTO (J/PSI(3100) + NEUTRALS)/(J/PSI(3100) + ANYTHING)					
R11 S	(.44) (.03)	ABRAMS1	75 SMAG	E+E-	1/76*	
R11 S	SUPERSEDED BY TANENBAUM 76 (FOOTNOTE 13 OF TANENBAUM 75, AND					
R11 S	PRIV. COMMUNICATION).					
R11	.41 .02	TANENBAUM	76 SMAG	E+E-	2/76*	
R11	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R11	0.410 0.019					
R12	PSI(3700) INTO (J/PSI(3100) P+ PI-)/TOTAL					
R12	.32 .04	TANENBAUM	75 SMAG	E+E-	1/76*	
R12	.36 .06	WIJK	75 DASP	E+E-	1/76*	
R12	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R12	0.332 0.033					
R12	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
R12	0.331 0.027					
R12	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R13	PSI(3700) INTO (J/PSI(3100) P0 P0)/TOTAL					
R13	.18 .06	WIJK	75 DASP	E+E-	1/76*	
R13	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R13	0.172 0.024					
R14	PSI(3700) INTO (J/PSI(3100) P0 P0)/(J/PSI(3100) P+ P-)					
R14 H	(.64) (.15)	HILGER	75 SPEC	E+E-	1/76*	
R14 H	IGNORING THE (J/PSI ETA) AND (J/PSI GAMMA GAMMA) DECAYS					
R15	PSI(3700) INTO (J/PSI(3100) ETA)/TOTAL					
R15	.043 .008	TANENBAUM	75 SMAG	E+E-	1/76*	
R15	.037 .015	WIJK	75 DASP	E+E-	1/76*	
R15	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R15	0.0417 0.0071					
R15	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
R15	0.0416 0.0070					
R15	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R16	PSI(3700) INTO (J/PSI(3100) GAMMA OR J/PSI(3100) P0)/TOTAL					
R16	(.0015) OR LESS CL=.90	TANENBAUM	75 SMAG	E+E-	2/76*	
R50	PSI(3700) INTO (J/PSI(3100) GAMMA GAMMA)/TOTAL					
R50 A	10 SEEN	CASP3	75 DASP	E+E-	1/76*	
R50 B	5 SEEN	DASP3	75 DASP	E+E-	1/76*	
R50 A	15 SEEN	HEINTZE	75 SPEC	E+E-	1/76*	
R50 B	150 SEEN	TANENBAUM	75 SMAG	E+E-	1/76*	
R50 AB	DECAY PROCEEDS VIA INTERMEDIATE STATE PC (SEE R51)					
R50 A	WITH J/PSI(3100) INTO (E+ E-)					
R50 B	WITH J/PSI(3100) INTO (MU+ MU-)					
R	HADRONIC DECAYS					
R20	PSI(3700) INTO (P+ PI-)/TOTAL					
R20	(.0009) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R20	(.00019) OR LESS CL=.90	FELDMAN	75 SMAG	E+E-	1/76*	
R21	PSI(3700) INTO (RHOO P0)/TOTAL					
R21	(.001) OR LESS CL=.90	ABRAMS4	75 SMAG	E+E-	1/76*	
R22	PSI(3700) INTO (2(P+ PI-) P0)/TOTAL					
R22	.0035 .0015	ABRAMS4	75 SMAG	E+E-	1/76*	
R23	PSI(3700) INTO (K+ K-)/TOTAL					
R23	(.0016) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R23	(.00023) OR LESS CL=.90	FELDMAN	75 SMAG	E+E-	1/76*	
R24	PSI(3700) INTO (P+ PI- K+ K-)/TOTAL					
R24	(.0005)	ABRAMS4	75 SMAG	E+E-	1/76*	
R25	PSI(3700) INTO (PBAR P)/TOTAL					
R25	(.001) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R25	.0004 .0002	ABRAMS4	75 SMAG	E+E-	1/76*	
R	RADIATIVE DECAYS					
R41	PSI(3700) INTO (GAMMA GAMMA)/TOTAL					
R41	(.008) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R41	(.005) OR LESS CL=.95	HUGHES	75 SPEC	E+E-	1/76*	
R42	PSI(3700) INTO (P0 GAMMA)/TOTAL					
R42	(.01) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R42	(.007) OR LESS CL=.95	HUGHES	75 SPEC	E+E-	1/76*	
R43	PSI(3700) INTO (ETA GAMMA)/TOTAL					
R43	(.0013) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R43	(.024) OR LESS CL=.95	HUGHES	75 SPEC	E+E-	1/76*	
R43	U RE-STAT'D BY US USING (MU+ MU-)/TOTAL = .01					
R44	PSI(3700) INTO (ETA PRIME GAMMA)/TOTAL					
R44	(.014) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R51	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (J/PSI(3100) GAMMA)					
R51	.04 .02	WIJK	75 DASP	E+E-	1/76*	
R51	.036 .007	TANENBAUM	75 SMAG	E+E-	1/76*	
R51	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R51	0.0364 0.0066					
R51	STUDENT AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
R51	0.0363 0.0066					
R51	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)					
R52	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (GAMMA GAMMA)					
R52	(.000013) OR LESS CL=.90	WIJK	75 CASP	E+E-	1/76*	
R53	PSI(3700) INTO (X(2750) GAMMA)/TOTAL, X(2750) INTO (GAMMA GAMMA)					
R53	(.00037) OR LESS CL=.90	WIJK	75 DASP	E+E-	1/76*	
R53	(.003) OR LESS CL=.95	HUGHES	75 SPEC	E+E-	1/76*	

R54	PSI(3700) INTO (CHI(3410) GAMMA)/TOTAL					
R54	*10**2					
R54 F	FOR DECAY MODES OF THE CHI INDICATED					
R54 F	(.13) (.05)	FELDMAN	75 CHI INTO (P+PI-JOR(K+K			
R54 F	(.07)	FELDMAN	75 CHI INTO (P+PI-K+K-)			
R54 F	(.14) (.07)	FELDMAN	75 CHI INTO (2(P+PI-J)			
R54 F	(.1)	FELDMAN	75 CHI INTO (3(P+PI-J)			
R54 F	(.5) OR LESS CL=.90	FELDMAN	75 CHI INTO (J/PSI GAMMA)			
R54 F	FOR DECAY MODES OF THE CHI INDICATED					
R55	PSI(3700) INTO (CHI(3530) GAMMA)/TOTAL					
R55	*10**2					
R55 F	FOR DECAY MODES OF THE CHI INDICATED					
R55 F	(.027) OR LESS CL=.90	FELDMAN	75 CHI INTO (P+PI-JOR(K+K			
R55 F	(.2)	FELDMAN	75 CHI INTO (P+PI-K+K-)			
R55 F	(.2)	FELDMAN	75 CHI INTO (2(P+PI-J)			
R55 F	(.2)	FELDMAN	75 CHI INTO (3(P+PI-J)			
R55 F	FOR DECAY MODES OF THE CHI INDICATED					
R56	PSI(3700) INTO (PC GAMMA)/TOTAL, PC INTO (P+PI-) OR (K+K-)					
R56	(.00027) OR LESS CL=.90	FELDMAN	75 SMAG	E+E-	1/76*	

71 PSI(3700) G(I)*G(E+E-)/G(TOTAL) (KEV)

THIS COMBINATION OF A PARTIAL WIDTH WITH THE PARTIAL WIDTH INTO E+E- AND WITH THE TOTAL WIDTH IS OBTAINED FROM THE INTEGRATED CROSS-SECTION INTO CHANNEL(I) IN THE E+E- ANNIHILATION. SEE THE MINI-REVIEW ON EXTRACTING RESONANCE WIDTHS.

G3	G(HADRONIC)*G(E+E-)/G(TOTAL)					
G3	2.2 .4	ABRAMS4	75 SMAG	E+E-	1/76*	

***** REFERENCES FOR PSI(3700) *****

ABRAMS	74 PRL 33 1453	*BRIGGS,AUGUSTIN,BOYARSKI+	(LBL+SLAC)	V71
ABRAMS1	75 PRL 34 1181	*BRIGGS,CHINCENSKY,FRIEDBERG,+	(LBL+SLAC)	V71
ABRAMS4	75 STANFORD SYMP.	G.S.ABRAMS	(LBL)	V71
AUBERT	75 PRL 33 1624	*BECKER,BIGGS,BURGER,GLENN+	(MIT+BNL)	V71
CAMERINI	75 PRL 35 483	*EARNED,PREFPOST,ASH,ANDERSON,+	(WISC+SLAC)	V71
CRIGEE	75 PL 93B 489	*BEHNE,FRANKE,HORLITZ,KRECHLOCK+	(DESY)	V71
DASP3	75 PL 57B 407	BRAUNSCHWIEG,KONIGS,+	(AACH+DESY+MIM+TKY)	V71
FELDMAN	75 PRL 35 821	*JEAN-MARIE,SADOULET,VANNUCCI,+	(LBL+SLAC)	V71
FELDMAN	75 STANFORD SYMP.	G.J.FELDMAN	(SLAC)	V71
GREGO	75 PL 56B 367	*PANCHERI-SRIVASTAVA,SRIVASTAVA	(FRAS)	V71
HEINTZE	75 STANFORD SYMP.	J.FEINTZE	(HEIDELBERG)	V71
JACKSON	75 NIM 128 13	J.D.JACKSON,D.SCHARRE	(LBL)	V71
HILGER	75 PRL 35 625	*BERON,FORD,HOFSTADTER,HOWELL,+	(STAN+PENN)	V71
HUGHES	75 PREP-HEPL 765	*BERON,CARRINGTON,FORD,HILGER,+	(STAN+PENN)	V71
LUTH	75 PRL 35 1124	*BOYARSKI,LYNCH,BREIBENBACH,+	(SLAC+LBL JJC)	V71
LUTH1	75 SLAC-PUB-1599	V.LUTH, PALERMIC CNF.PROC.+	(SLAC+LBL)	V71
LIBERMAN	75 STANFORD SYMP.	A.D.LIBERMAN	(STANFORD)	V71
PREFPOST	75 STANFORD SYMP.	R.PREFPOST	(WISCONSIN)	V71
SIMPSON	75 PRL 35 699	*BERON,FORD,HILGER,HOFSTADTER,+	(STAN+PENN)	V71
TANENBAU	75 PRL 35 1323	TANENBAUM,WHITAKER,ABRAMS,+	(LBL+SLAC)	V71
TANENBAU	75 PRL SLAC-PUB1696	TANENBAUM,ABRAMS,BOYARSKI,+	(SLAC+LBL)	V71
WIJK	75 STANFORD SYMP.	B.F.WIJK	(DESY)	V71

Additional States in S-channel e+e- Collisions

No evidence for narrow states has been found in the mass intervals 1910-2545 MeV and 2970-3090 MeV (BACCI 75, ESPOSITO 75). Except for the $\psi(3700)$, no additional narrow states have been found in the mass interval 3200-7600 MeV (BOYARSKI 75, SCHWITTERS 75).

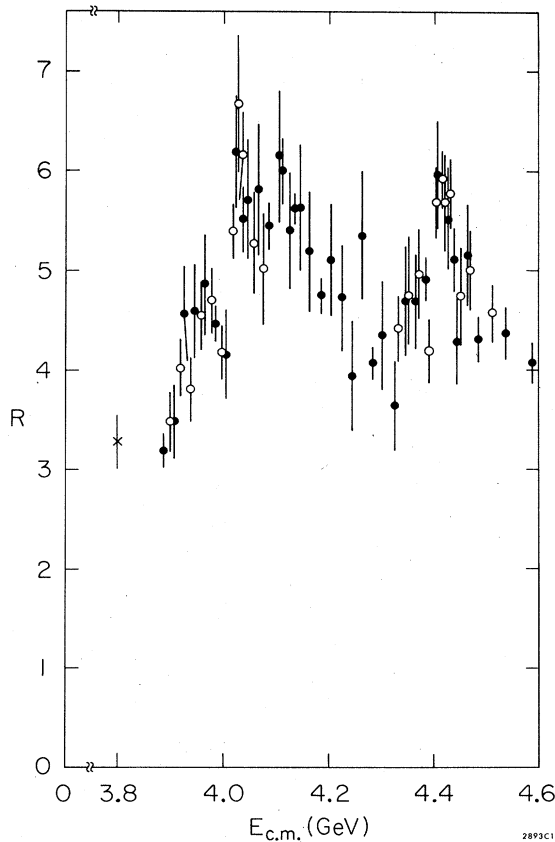
The broad enhancement observed in the e^+e^- total cross section at SPEAR (AUGUSTIN 75) is now resolved into two distinct objects at 4100 and 4400 MeV (SIEGRIST 76), with indications of even more detailed substructure especially at 4100 MeV (see figure). Due to there being very limited information available on the properties of $\psi(4100)$ and $\psi(4400)$, their resonant interpretation is not yet established. If in fact they are resonances, then their relatively large widths suggest that they are hadronic, with $J^{PC} = 1^{--}$.

Data Card Listings

For notation, see key at front of Listings.

Mesons

$\psi(4100)$, $\psi(4400)$, $X(2750)$



The ratio of cross sections $\frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$ in the $\psi(4100)$ - $\psi(4400)$ region (from SIEGRIST 76).

$\psi(4100)$

72 PSI(4100, JPG=) I=

72 PSI(4100) MASS (MEV)

M	(4150.)	AUGUSTIN 75 SPEAR	2.4-5.0 E+E-	2/75*
M	(4100.)	SCHWITTER 75 SMAG	E+E-	1/76*

REFERENCES FOR PSI(4100)

AUGUSTIN 75 PRL 34 764	+BOYARSKI, ABRAMS, BRIGGS+	(SLAC+LBL)	V72
BACCI 75 PL 58B 481	+BIDDOLI, PENSO, STELLA, +	(ROMA+FRAS)	V72
BOYARSKI 75 PRL 34 762	+BREIDENBACH, ABRAMS, BRIGGS, +	(SLAC+LBL)	V72
ESPOSITO 75 PL 58B 478	+FELICETTI, PERUZZI, +	(FRAS+NAPO+PADO+ROMA)	V72
SCHWITTE 75 STANFORD SYMP.	R.F.SCHWITTERS	(SLAC)	V72
SIEGRIST 76 SUBM. TO PRL	+ABRAMS, BOYARSKI, BREIDENBACH, +	(LBL+SLAC)	V72

$\psi(4400)$

73 PSI(4400, JPG=) I=

73 PSI(4400) MASS (MEV)

M	4414.	7.	SIEGRIST 76 SMAG	E+E-	2/76*
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73 PSI(4400) WIDTH (MEV)

M	33.	10.	SIEGRIST 76 SMAG	E+E-	2/76*
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73 PSI(4400) BRANCHING RATIOS

R1	PSI(4400) INTO (E+ E-)/TCTAL *10**5				
R1	1.3	.3	SIEGRIST 76 SMAG	E+E-	2/76*

REFERENCES FOR PSI(4400)

SCHWITTE 75 STANFORD SYMP.	R.F.SCHWITTERS	(SLAC)	V73
SIEGRIST 76 SUBM. TO PRL	+ABRAMS, BOYARSKI, BREIDENBACH, +	(LBL+SLAC)	V73

$X(2750)$

54 X(2750, JPG=) I=

OBSERVED IN THE SEQUENTIAL RADIATIVE DECAY OF THE $J/\psi(3100)$ INTO $X(2750)$ GAMMA, $X(2750)$ INTO GAMMA GAMMA, AND POSSIBLY $P\bar{B}AR P$ (HEINTZE 75). THIS SUGGESTS QUANTUM NUMBER ASSIGNMENTS $C=+$, $IG=0+$ OR $1-$. NEEDS CONFIRMATION. OMITTED FROM TABLE.

54 X(2750) MASS (MEV)

M	(2750.)	HEINTZE 75 DESY	E+E-	1/76*
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54 X(2750) PARTIAL DECAY MODES

P1	X(2750) INTO GAMMA GAMMA	DECAY MASSES
P2	X(2750) INTO $P\bar{B}AR P$	

54 X(2750) BRANCHING RATIOS

R1	X(2750) INTO ($P\bar{B}AR P$)/TOTAL	WIJK 75 DASP	E+E-	1/76*
R1	2 POSSIBLY SEEN			

REFERENCES FOR X(2750)

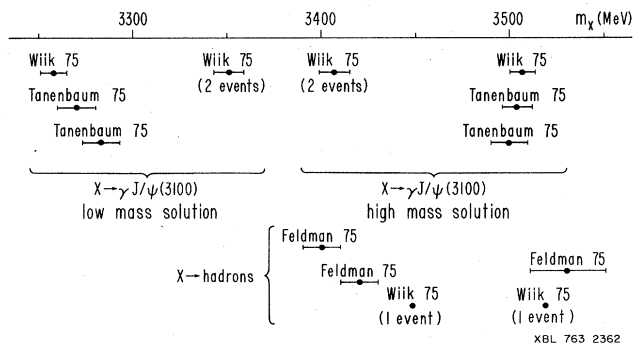
HEINTZE 75 STANFORD SYMP.	J.FEINTZE	(HEIDELBERG)	V54
WIJK 75 STANFORD SYMP.	B.H.WIJK	(DESY)	V54

States Observed in Radiative Decays of $\psi(3700)$

The properties, and even the number, of intermediate states in radiative decay of $\psi(3700)$ are not well known. The main problems are:

- 1) small statistics in most experiments;
- 2) ambiguity in the determination of the mass of P_c ;
- 3) it is not known whether the $X(3530)$ is a single state, or whether the relatively large observed width is due to a superposition of several narrow states.

The data seem to require the presence of at least 3 states in the mass interval 3250-3550 MeV; the figure summarizes all the present claims.



Claims for intermediate states X in the radiative decay $\psi(3700) \rightarrow \gamma X$.

Mesons

P_c(3300 or 3500), χ(3410), χ(3530)

Data Card Listings

For notation, see key at front of Listings.

P_c(3300 or 3500)

55 PC(3300 OR 3500, JPC=) I=

OBSERVED IN THE RADIATIVE SEQUENTIAL DECAY OF THE PSI(3700) INTO PC GAMMA, PC INTO J/PSI(3100) GAMMA (BRAUNSWIG 75, CONFIRMED BY TANENBAUM 75), THEREFORE C=+. MASS DETERMINATION AMBIGUOUS DUE TO TWO POSSIBLE (GAMMA J/PSI) COMBINATIONS IN THE FINAL STATE (J/PSI GAMMA GAMMA). NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

55 PC MASS (MEV)

Table with columns for mass solutions (LW and HIGH MASS), particle names (DASP3, TANENBAUM, WIJK), and branching ratios (E+E-, 1/76*). Includes average values and error factors.

55 PC PARTIAL DECAY MODES

Table listing decay modes (P1-P4) and masses (DECAY MASSES).

55 PC BRANCHING RATIOS

Table listing branching ratios (R1-R3) and dominant states.

REFERENCES FOR PC

Bibliography table listing references for PC (DASP3, FELDMAN, HEINTZE, SIMPSON, TANENBAUM, WIJK).

χ(3410)

56 CHI(3410, JPC=) I=

OBSERVED IN THE RADIATIVE DECAY OF PSI(3700) INTO CHI(3410) GAMMA (FELDMAN 75), THEREFORE C=+. THE OBSERVED DECAY INTO (PI+PI-) OR (K+K-) IMPLIES G=+, JP=0+, 2+, CONFIRMATION OF BOTH (PI+PI-) AND (K+K-) MODES WOULD ESTABLISH I=0. NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

56 CHI(3410) MASS (MEV)

Table with columns for mass solutions (M, M AVG, M STUDENT) and branching ratios (E+E-, 1/76*).

56 CHI(3410) PARTIAL DECAY MODES

Table listing decay modes (P1-P6) and masses (DECAY MASSES).

56 CHI(3410) BRANCHING RATIOS

Table listing branching ratios (R) and references for CHI(3410).

χ(3530)

57 CHI(3530, JPC=) I=

OBSERVED IN RADIATIVE DECAY OF PSI(3700) INTO CHI(3530) GAMMA (FELDMAN 75), THEREFORE C=+. CAN BE INTERPRETED AS A SINGLE (BROAD) STATE, OR AS A SUPERPOSITION OF SEVERAL NARROW STATES. NEEDS CLARIFICATION (SEE THE MINIREVIEW ON STATES OBSERVED IN RADIATIVE DECAYS OF PSI(3700). OMITTED FROM TABLE.

57 CHI(3530) MASS (MEV)

Table with columns for mass solutions (M) and branching ratios (E+E-, 1/76*).

57 CHI(3530) PARTIAL DECAY MODES

Table listing decay modes (P1-P6) and masses (DECAY MASSES).

57 CHI(3530) BRANCHING RATIOS

Table listing branching ratios (R) and references for CHI(3530).

Data Card Listings

For notation, see key at front of Listings.

Baryons
N's and Δ 'sNote on N's and Δ 'sI. Determination of Resonance Parameters

Values of masses, widths, and branching ratios are obtained mainly from phase-shift analyses. In production experiments, in fact, it is seldom clear which of the many states at similar masses is being observed. In addition to a few complete phase-shift analyses, we have other analyses, done by using somewhat incomplete data, by various different groups, but we are quite far from having reliable masses and widths derived therefrom.

There are essentially two problems in obtaining reliable resonance parameters. First there is often disagreement as to just what the values of the partial-wave amplitudes are. This problem is obviously related to the quality and quantity of the data and to the procedures used to determine the amplitudes. Secondly, even if smooth curves were available for the amplitudes, there would still be some parametrization-dependent ambiguity in deciding what the resonance parameters should be. From a theoretical standpoint the most unambiguously defined resonance parameters are the pole position and residue, and it has been found in practice that, given sufficiently precise partial-wave amplitudes, these quantities can be extracted in a stable and parametrization-independent way, in spite of the fact that they require an extrapolation away from the physical region. This point has been discussed in detail with regard to the $\Delta(1232)$ in previous editions of this review^{1,2}. We list available pole parameter determinations in the Data Card Listings, and many further discussions can be found in the corresponding references, e.g., NOGOVA 73, SPEARMAN 74, BALL 75, LICHTENBERG 75, LONGACRE 75, and VASAN 75.

At the beginning of the Data Card Listings for N's and Δ 's, we present a table giving our evaluation of the N and Δ resonances based on information contained in the Listings. In the Table of Particle Properties, we do not quote values and errors for most parameters, but give only ranges for masses and widths in order to

emphasize that in some cases these parameters are quite poorly determined.

References for Section I

1. Particle Data Group, Rev. Mod. Phys. **43**, S114 (1971).
2. Particle Data Group, Phys. Lett. **39B**, 103 (1972).

For other references see the Data Card Listings.

II. $\pi N \rightarrow \pi N, \eta N, K\Sigma, K\Lambda$

The most recent available $\pi N \rightarrow \pi N$ amplitudes are from the analysis of the Saclay group, AYED 74. Preliminary results of more recent analyses have been discussed at meetings^{1,2}, but resonance parameters are not yet available. The results of AYED 74 are shown in Figs. II.1-II.6. The figure captions summarize the known resonances in the various partial waves. We restrict our discussion here to mentioning the possible new effects seen in recent analyses. All resonance parameters quoted under AYED 74 derive from energy-dependent fits (Breit-Wigner with background) to the partial-wave amplitudes which were determined in an energy-independent shortest-path analysis.

S₃₁^{''}: In Fig. II.1, the energy-dependent fit of AYED 74 suggests that the broad structure around 2000 MeV is associated with a second resonance with

$$M = 2001 \text{ MeV}, \Gamma_{\text{tot}} = 307 \text{ MeV}, x_{\text{el}} = 0.08$$

There is some additional evidence for such an effect from LANGBEIN 73 and DEANS 75, which are energy-independent and energy-dependent analyses, respectively, of $K\Sigma$ associated production.

P₃₃^{''}: The small dip in the imaginary part of P_{33} and the zero in the real part around 1900 MeV are interpreted in the energy-dependent fit of AYED 74 as a second P_{33} resonance,

$$M = 1904 \text{ MeV}, \Gamma_{\text{tot}} = 204 \text{ MeV}, x_{\text{el}} = 0.19$$

ALMEHED 72 found two effects, one at ~ 1680 MeV and another at ~ 2150 MeV, as did DEANS 75.

However, LANGBEIN 73 found evidence for the higher mass effect only. The existence of some sort of effect at ~ 2200 MeV seems definitely established

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

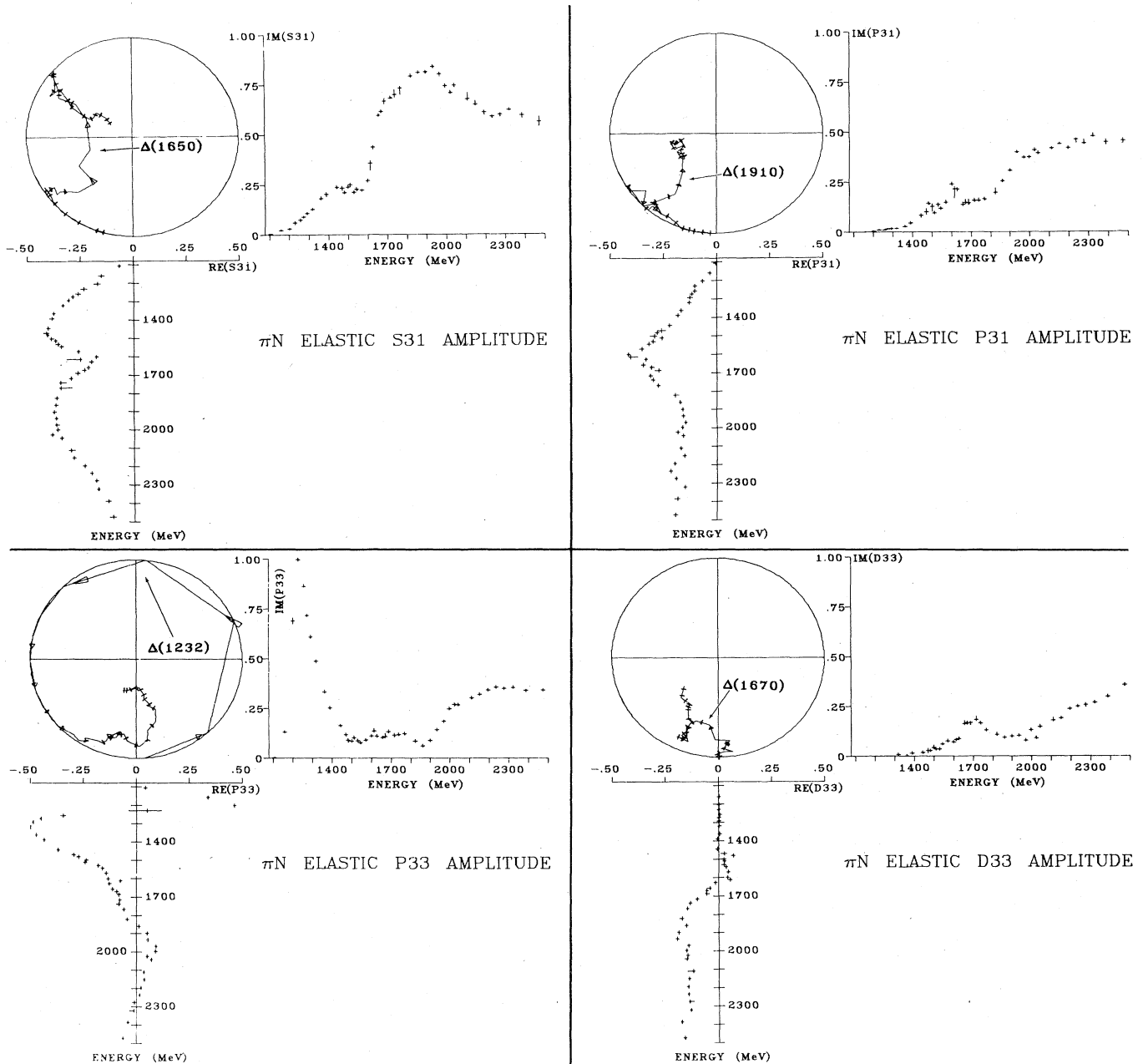


Fig. II.1. Amplitudes for $I = 3/2$ π N elastic scattering in the $J = 1/2$ and $J = 3/2$ waves from AYED '74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $\Delta(1650)$, $\Delta(1910)$, $\Delta(1232)$, and $\Delta(1670)$ in the S₃₁, P₃₁, P₃₃, and D₃₃ waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for other possible resonances in the S₃₁ and P₃₃ waves.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N's and Δ 's

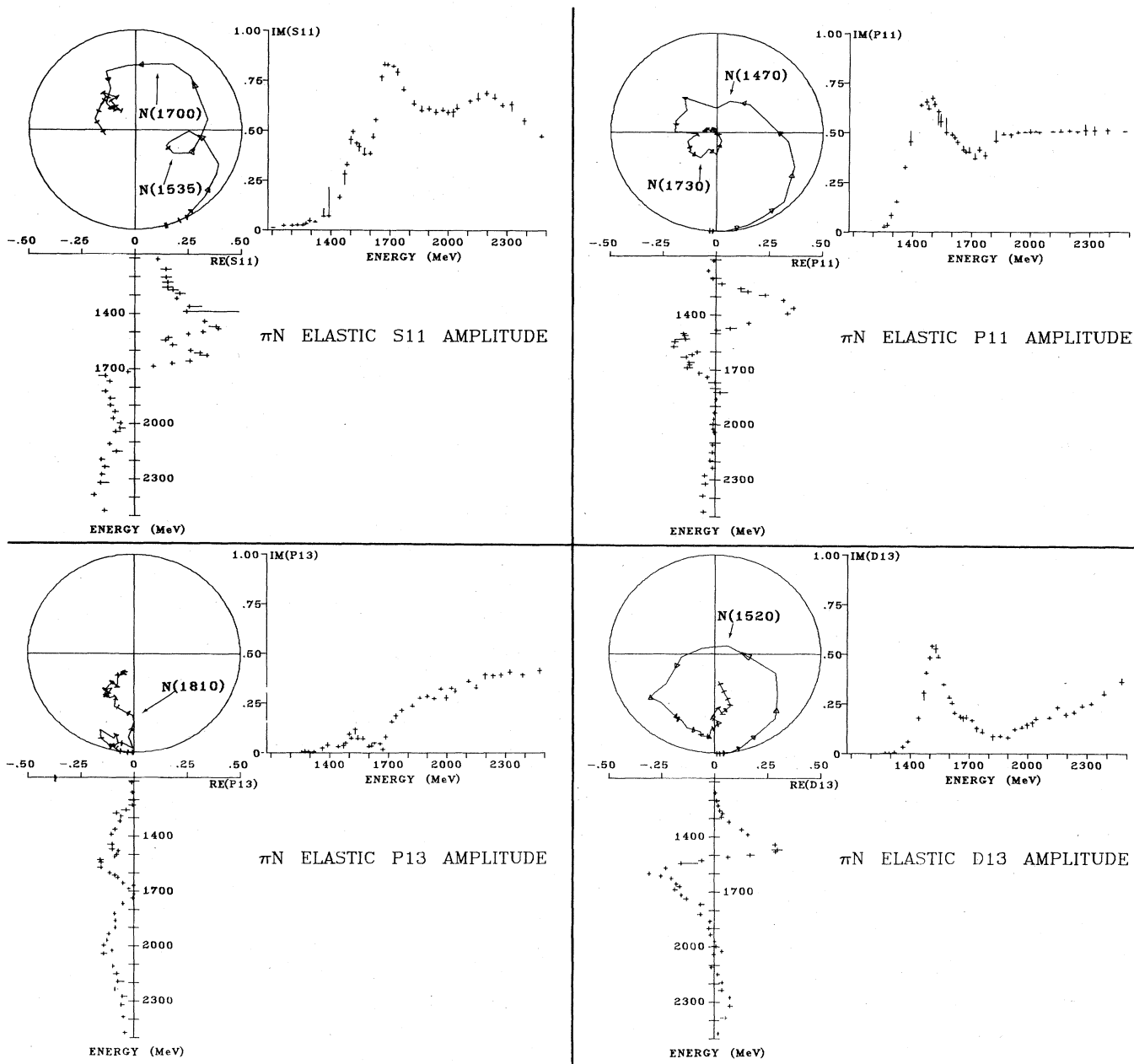


Fig. II.2. Amplitudes for $I = 1/2$ πN elastic scattering in the $J = 1/2$ and $J = 3/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the N(1535) and N(1700) in the S₁₁ wave, the N(1470) and N(1730) in the P₁₁ wave, the N(1810) in the P₁₃ wave, and the N(1520) in the D₁₃ wave; these are indicated on the above Argand plots. The P₁₁ wave also contains the nucleon pole, 138 MeV below threshold. See the Data Card Listings and the accompanying mini-review for other possible resonances in these waves.

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

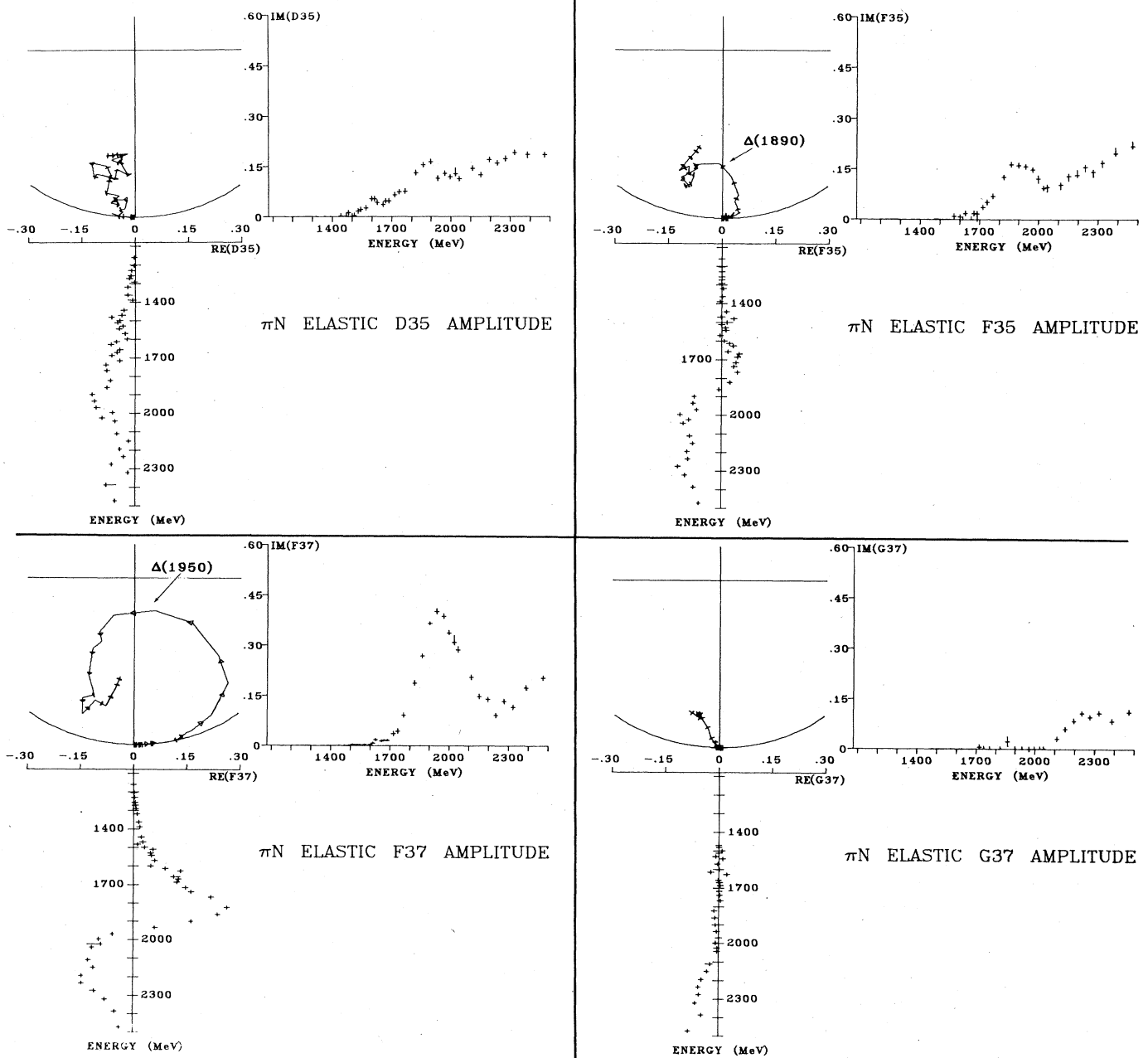


Fig. II.3. Amplitudes for $I = 3/2$ π N elastic scattering in the $J = 5/2$ and $J = 7/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $\Delta(1890)$ and $\Delta(1950)$ in the F35 and F37 waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for another possible resonance in the D35 wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

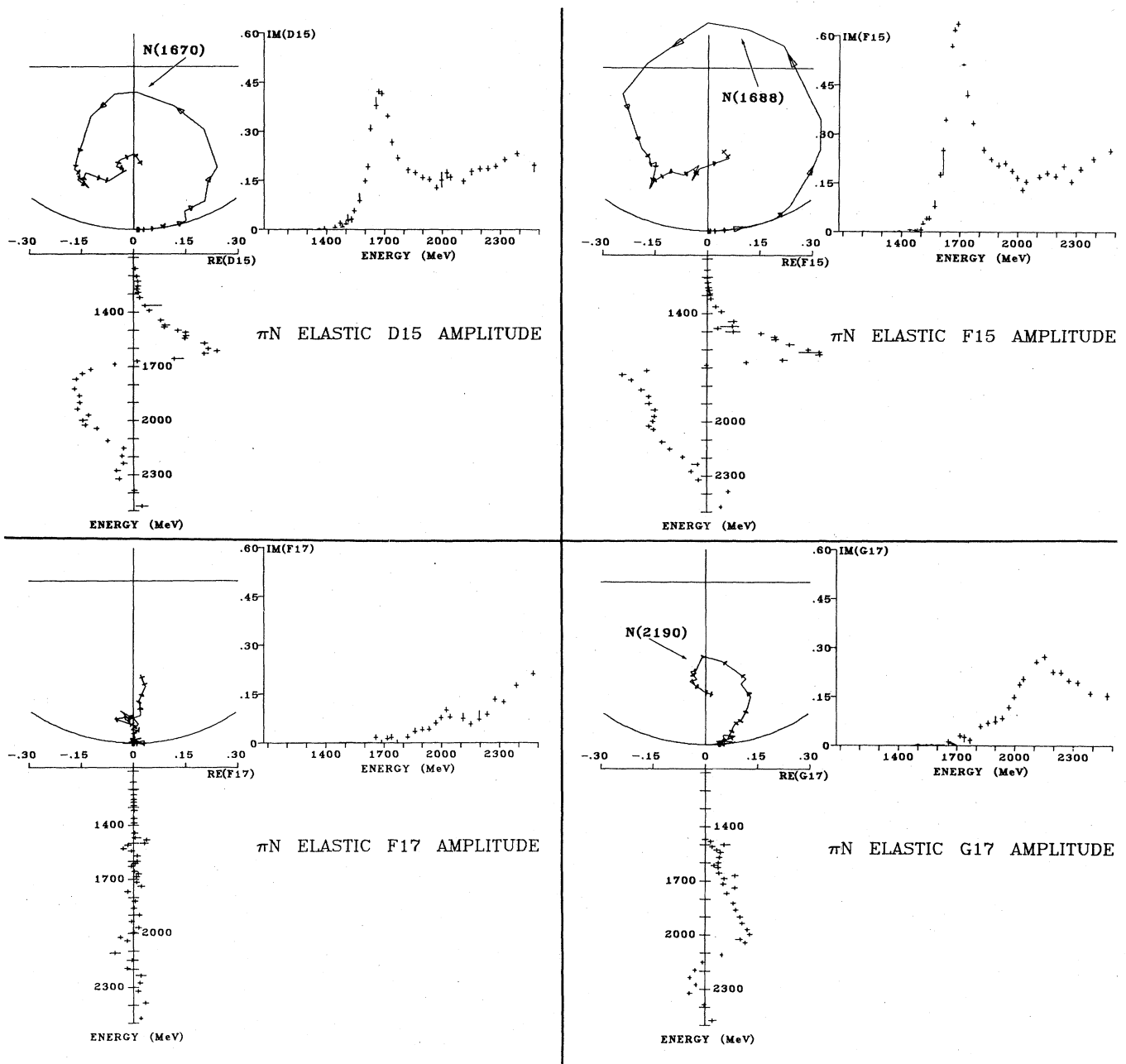
N's and Δ 's

Fig. II.4. Amplitudes for $I = 1/2$ πN elastic scattering in the $J = 5/2$ and $J = 7/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from elastic threshold to 2500 MeV. Established resonances in these waves are the $N(1670)$, $N(1688)$, and $N(2190)$ in the D_{15} , F_{15} , and G_{17} waves, respectively; these are indicated on the above Argand plots. See the Data Card Listings and the accompanying mini-review for other possible resonances in the D_{15} , F_{15} , and F_{17} waves.

Baryons

N's and Δ 's

Data Card Listings

For notation, see key at front of Listings.

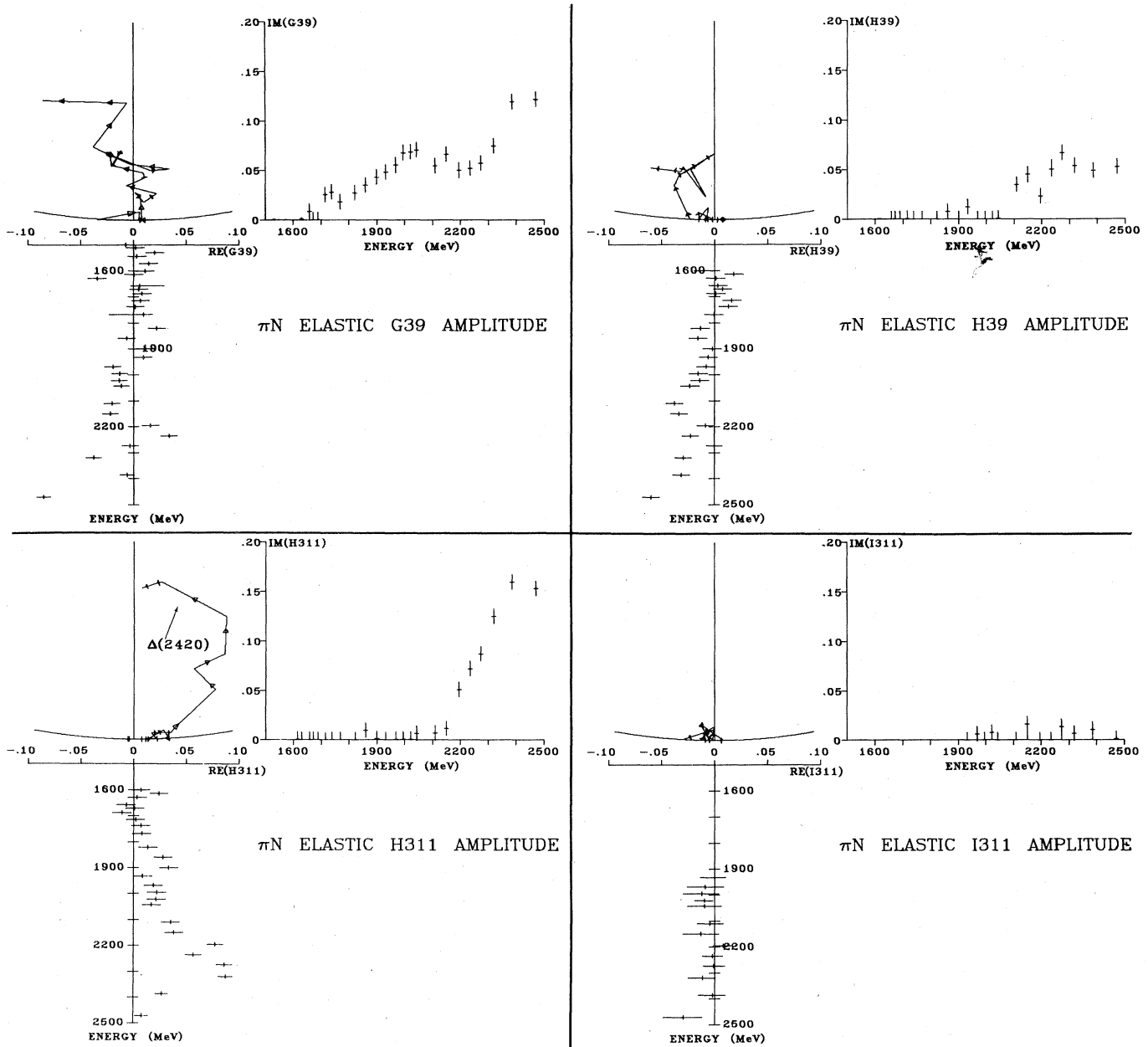


Fig. II.5. Amplitudes for $I = 3/2$ πN elastic scattering in the $J = 9/2$ and $J = 11/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All energy axes run from 1500 to 2500 MeV. The only established resonance in these waves is the $\Delta(2420)$ in the $H_3 11$ wave; it is indicated on the above $H_3 11$ Argand plot. See the Data Card Listings and the accompanying mini-review for another possible resonance in the G_{39} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

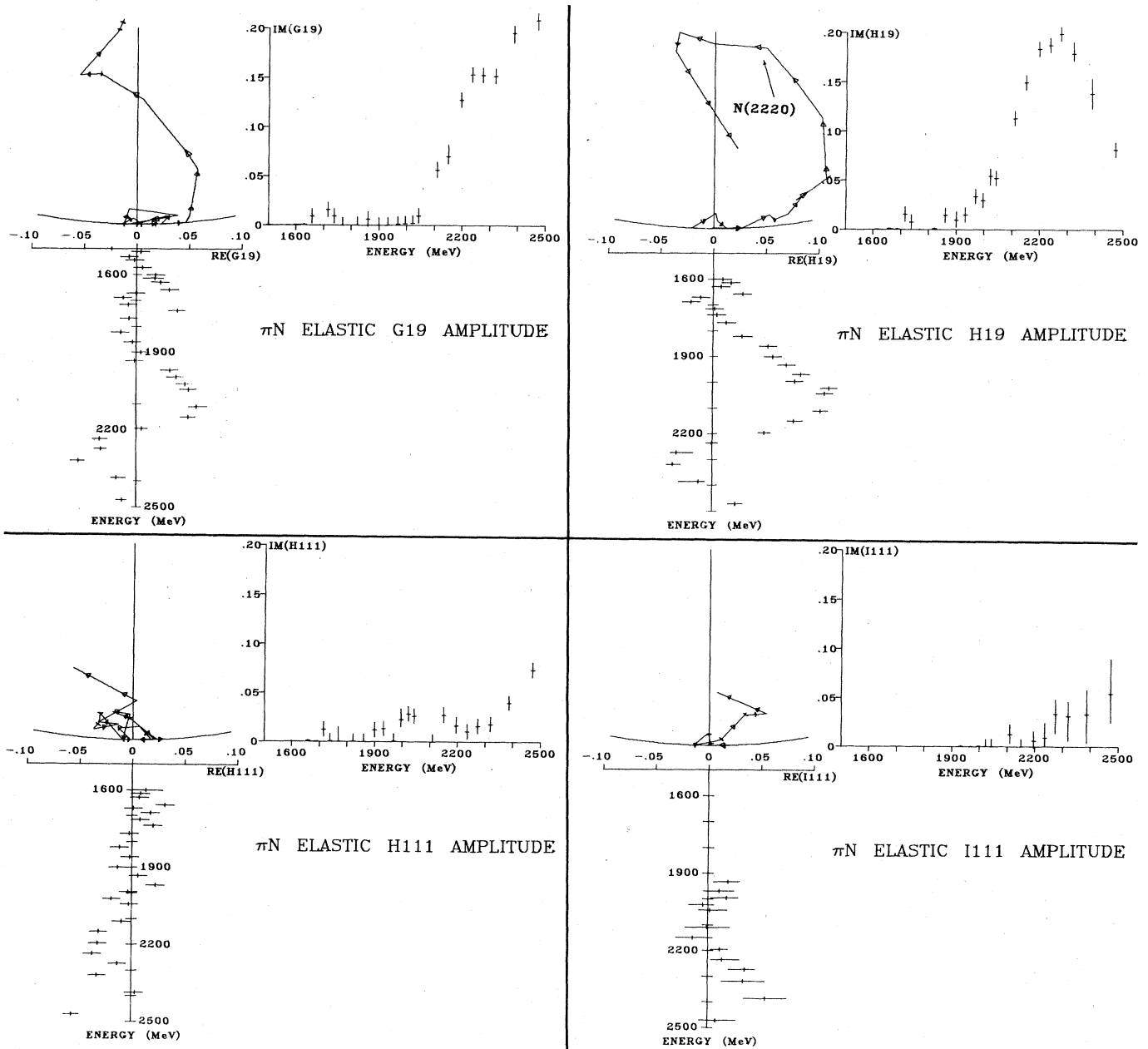
N's and Δ 's

Fig. II.6. Amplitudes for $I = 1/2$ π N elastic scattering in the $J = 9/2$ and $J = 11/2$ waves from AYED 74. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts vs. energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 50 MeV and a base-to-tip length of 5 MeV. All the energy axes run from 1500 to 2500 MeV. The only established resonance in these waves is the $N(2220)$ in the H_{19} wave; it is indicated on the above H_{19} Argand plot. See the Data Card Listings and the accompanying mini-review for another possible resonance in the G_{19} wave.

Baryons

N's and Δ 's

by the backward π^+p scattering experiment of REY 74, but the quantum numbers are not yet established [see the Listings for $\Delta(2160)$].

P_{31} : The P_{31} resonance mass quoted by AYED 74 ($M = 1786$ MeV) is considerably lower than in previous analyses (except AYED 70) which find M in the region 1900 to 1950 MeV.

D_{35} : AYED 74, LANGBEIN 73, and DEANS 75 all find evidence for a D_{35} state somewhere in the mass range 1900-2000 MeV, ALMEHED 72 found an effect at 2200 MeV, but with such a large width, 600 MeV, that a resonance interpretation is questionable.

S_{11}''' : In Fig. II.2 are shown the low spin $I = 1/2$ waves of AYED 74. They associate the broad effect at ~ 2300 MeV with a third S_{11} resonance:

$$M = 2283 \text{ MeV}, \Gamma_{\text{tot}} = 310 \text{ MeV}, x_{\text{el}} = 0.14 .$$

ALMEHED 72 find a similar effect at 2100 MeV.

P_{11}' : AYED 74 claim the "Roper" to be split:

$$M_{\text{low}} = 1413 \text{ MeV}, \Gamma_{\text{tot}} = 187 \text{ MeV}, x_{\text{el}} = 0.55 ;$$

$$M_{\text{high}} = 1532 \text{ MeV}, \Gamma_{\text{tot}} = 89 \text{ MeV}, x_{\text{el}} = 0.12 .$$

While such a splitting might be a natural explanation for the relatively low mass and broad width of the "Roper" as seen in production experiments, the evidence in Fig. II.2, especially the projections, seems only tentative. The new Saclay ηn data and partial-wave analysis of FELTESSE 75 are not inconsistent with the parameters of the high mass state. However, an equally acceptable fit to their data is obtained by introducing a new P_{13} resonance with $M = 1530$ MeV, $\Gamma = 79$ MeV, and $\sqrt{x_{\text{el}}} = 0.33$.

D_{13}'' : A somewhat larger effect in this wave at ~ 1700 MeV is identified by AYED 74 as a second D_{13} resonance:

$$M = 1710 \text{ MeV}, \Gamma_{\text{tot}} = 100 \text{ MeV}, x_{\text{el}} = 0.09 .$$

The projections reveal a shoulder in the imaginary part and a small dip in the real part. A similar effect exists in ALMEHED 72, but was not claimed as a resonance. Further evidence for this state is found by LANGBEIN 73, DEANS 75, and KNAESEL 75 ($K\bar{A}$ associated production), as well as recent pion

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photoproduction analyses and the LBL-SLAC $\pi\pi\pi N$ analysis (see III and IV below).

D_{13}''' : Along with ALMEHED 72, AYED 74 find evidence for a D_{13} state at ~ 2000 MeV:

$$M = 2029 \text{ MeV}, \Gamma_{\text{tot}} = 116 \text{ MeV}, x_{\text{el}} = 0.10 .$$

The evidence in the projections seems less than for the effect at 1700 MeV. Evidence for coupling of this state to the $K\bar{\Sigma}$ channel is found by DEANS 75, but not by LANGBEIN 73.

D_{15}'' : The rapid decrease in the real part of this amplitude around 2100 MeV (see Fig. II.4) is associated with a second D_{15} resonance by AYED 74:

$$M = 2100 \text{ MeV}, \Gamma_{\text{tot}} = 220 \text{ MeV}, x_{\text{el}} = 0.08 .$$

These parameters are in reasonable agreement with those of ALMEHED 72.

F_{15}'' : The broad shoulder in the real part and dip in the imaginary part are identified with a second F_{15} resonance by AYED 74:

$$M = 1989 \text{ MeV}, \Gamma_{\text{tot}} = 179 \text{ MeV}, x_{\text{el}} = 0.08 .$$

ALMEHED 72 find an effect at ~ 2200 MeV, and LANGBEIN 73 and DEANS 75 find some evidence for coupling to the $K\bar{\Sigma}$ channel at ~ 2000 MeV.

Higher Waves. Besides confirming the already accepted $H_{3,11}$ and H_{19} , AYED 74 present evidence for two additional states in the 2200 MeV region:

$$G_{39}: M = 2174 \text{ MeV}, \Gamma_{\text{tot}} = 205 \text{ MeV}, x_{\text{el}} = 0.04 ;$$

$$G_{19}: M = 2133 \text{ MeV}, \Gamma_{\text{tot}} = 193 \text{ MeV}, x_{\text{el}} = 0.10 .$$

The G_{39} effect in Fig. II.5 is not too convincing. Notice that ALMEHED 72 claimed a P_{33} resonance in this same region; see $\Delta(2160)$ for further possibilities in this mass region. In contrast, the evidence for G_{19} , Fig. II.6, looks quite good; note, in particular, the behavior of the real part projection in the region 1900-2500 MeV.

References for Section II

1. R. L. Kelly, in New Directions in Hadron Spectroscopy (ANL-HEP-CP-75-58), eds. S. L. Kramer and E. L. Berger, 1975.
2. E. Pietarinen, Universität Karlsruhe Reports TKP 4/75 and TKP 15/75.

For other references see the Data Card Listings.

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Baryons
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In the isobar model, the amplitude for the reaction $\pi N \rightarrow N\pi\pi$ is written

$$T = \sum \left\{ \begin{array}{l} T_{\Delta\pi}^{\text{JILL}'}(E) BW_{\Delta}(E_{\pi\pi}) X_{\Delta\pi}^{\text{JILL}'} \\ + T_{\rho N}^{\text{JILL}'}(E) BW_{\rho}(E_{\pi\pi}) X_{\rho N}^{\text{JILL}'} \\ + T_{\epsilon N}^{\text{JILL}'}(E) BW_{\epsilon}(E_{\pi\pi}) X_{\epsilon N}^{\text{JILL}'} \end{array} \right\}.$$

In this expression the BW's denote either appropriate Breit-Wigner's or the corresponding two-body amplitude. The functions X contain all the angular information; these are well-defined functions depending only on which isobars are used in the model. The partial-wave amplitudes, $T_{\Delta\pi}^{\text{JILL}'}$, etc., may be indicated by

$$\Delta\pi, LL'_{2I2J} \quad \rho N, LL'_{2I2J} \quad \epsilon N, LL'_{2I2J}$$

where L is the incoming (πN) angular momentum, and L' is the outgoing angular momentum between the isobar [$\rho, \epsilon (\pi\pi I = 0, S \text{ wave}), \Delta$] and the remaining hadron (N or π); as usual I and J are the isospin and total spin ($\vec{J} = \vec{L} + \vec{S} = \vec{L}' + \vec{S}'$), respectively. Often the ρ has a subscript 1 or 3 which denotes twice the ρN total spin. The generalization of the model to include more two-body final-state interactions is obvious, but only the three included above, Δ , ρ , and ϵ , have actually been used in presently published analyses. It should be noted that one assumes the partial-wave amplitudes depend only on the c.m. energy E and not on the diparticle subenergies, $E_{\pi\pi}$ and $E_{\pi N}$.

The largest isobar analysis to date is that of the LBL-SLAC collaboration, LONGACRE 75 and Ref. 1, which analyzes 200K events of $\pi^- p \rightarrow \pi^- \pi^0 p$, $\pi^+ \pi^- n$ and $\pi^+ p \rightarrow \pi^+ \pi^0 p$, $\pi^+ \pi^+ n$ in the c.m. energy range 1300-2000 MeV. Details of this analysis, particularly the evolution of the preferred solution B, have been discussed in the previous edition of this review². Results from this solution are shown in Fig. III.1. In this edition we list resonance parameters, including pole positions, from LONGACRE 75 in the Data Card Listings. Breit-Wigner masses, widths, and resonance couplings were extracted in two ways by

LONGACRE 75, and we list the results of both methods, in addition to the pole positions. Method 1 used Breit-Wigner fitting to individual partial waves, and method 2 used a coupled-channel K-matrix fit to all waves with the same IJP. Among the most interesting results are further evidence for the existence of the $D_{13}(1700)$ and the $P_{33}(1690)$, further evidence for dominantly F-wave decay of the $F_{37}(1950)$ into $\Delta\pi$, and good agreement with quark-model predictions³.

Some information on isobar couplings is also available from earlier analysis. CHAVANON 74⁴ analyzed $\pi^+ p \rightarrow \pi^+ \pi^0 p$ in the range 1580-1970 MeV and obtained information on Δ -resonance decays to $\Delta\pi$ and $N\rho$. MEHTANI 72 analyzed the $\Delta\pi$ channel in the region 1820-2090 MeV and found the first evidence for strong F-wave decay of the $F_{37}(1950)$. The analysis of DIEM 70 covered the region 1550-1650 MeV.

In connection with all of these analyses it should be kept in mind that the isobar model is really a model and is not an exact representation of the $\pi N \rightarrow \pi\pi N$ amplitude. There are a number of ways in which the model could seriously fail to represent the physical amplitude. The problem that has received the most attention recently is the assumption that the isobar partial-wave amplitudes do not depend on the diparticle sub-energies, an assumption which is inconsistent with unitarity. For discussion of this problem see Refs. 2, 5, 6, and other references quoted therein.

References for Section III

1. D. J. Herndon et al., Phys. Rev. D11, 3183 (1975).
2. Particle Data Group, Phys. Lett. 50B, No. 1 (1974).
3. R. J. Cashmore et al., Nucl. Phys. B92, 37 (1975).
4. P. Chavanon et al., Nucl. Phys. B76, 157 (1974).
5. R. Aaron et al., Phys. Rev. D12, 1984 (1975).
6. Y. Goradia and T. A. Lasinski, LBL-3626 (1975); submitted to Phys. Rev. D.

For other references see the Data Card Listings.

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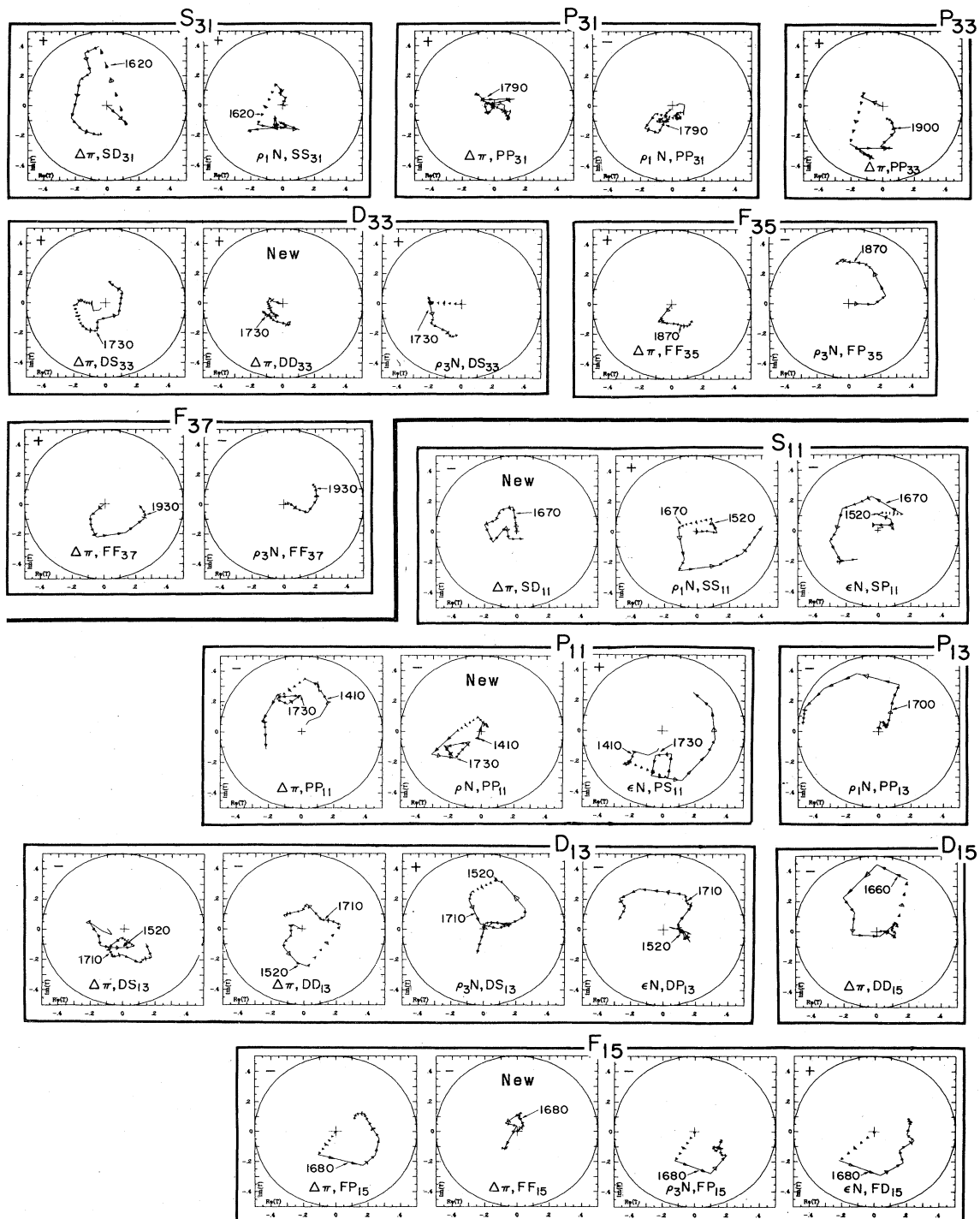


Fig. III.1. Solution B of the LBL-SLAC isobar analysis of $\pi N \rightarrow N\pi\pi$. See mini-review for notation. The four waves labeled "New" were not used in earlier solutions. The sign in the upper left-hand corner of each plot indicates how to transform from the sign conventions of this group to that of the "baryon-first" convention. (See main text, Sec. IV B.)

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IV. Photon Couplings

(R. L. Crawford and R. G. Moorhouse, Glasgow Univ.)

Photon couplings can be studied in formation reactions like

$$\gamma N \rightarrow N^* \rightarrow \pi N, \eta N, K\Lambda, \pi\Delta, \dots$$

A partial-wave analysis of these processes is the standard technique to determine the coupling strengths $g(N^* N \gamma)$. Up to now, almost all results are derived from the analysis of single-pion photoproduction. In the following we, therefore, define the conventions of pion photoproduction in which the results will be quoted.

The process $\gamma N \rightarrow N^* \rightarrow \pi N$ for a specific intermediate resonance can be symbolically described as

$$\langle \pi N | H_{\pi} | N^* \rangle \langle N^* | H_{\gamma} | \gamma N \rangle$$

The first term is measured in strong interactions, e.g., by partial-wave analysis of elastic πN scattering. A common feature of almost all analyses of pion photoproduction is a strong reliance on the knowledge of resonance parameters from πN partial-wave analyses. Pion photoproduction does not give a particularly accurate method of determining the resonance masses and widths, due partly to a lack of sufficiently precise data at many energies and partly because photoproduction is complicated by the fact that the photon has helicity states ± 1 and can react as an isoscalar and isovector particle. Consequently, several couplings for $N^* \rightarrow \gamma N$ have to be determined (2 for Δ and 4 for N).

Isospin Decomposition. Ignoring possible isotensor components, the reactions leading to the four possible final charge states are described by three isospin amplitudes. One set of these consists of the amplitudes A^{Δ} , A^P , and A^N which are respectively the amplitudes for the reaction to proceed by an $I = 3/2$ state, an $I = 1/2$ state with charge = +1, and an $I = 1/2$ state with charge = 0:

$$\begin{aligned} A(\gamma p \rightarrow \pi N) &= C_{\pi N}^{3/2} A^{\Delta} + C_{\pi N}^{1/2} A^P, \\ A(\gamma n \rightarrow \pi N) &= C_{\pi N}^{3/2} A^{\Delta} + C_{\pi N}^{1/2} A^N. \end{aligned} \quad (1)$$

The C-G coefficients, $C_{\pi N}^I$, for the coupling to a specific πN state are given explicitly by

$$\begin{aligned} A(\gamma p \rightarrow \pi^+ n) &= -\sqrt{1/3} A^{\Delta} - \sqrt{2/3} A^P, \\ A(\gamma p \rightarrow \pi^0 p) &= \sqrt{2/3} A^{\Delta} - \sqrt{1/3} A^P, \\ A(\gamma n \rightarrow \pi^- p) &= \sqrt{1/3} A^{\Delta} - \sqrt{2/3} A^N, \\ A(\gamma n \rightarrow \pi^0 n) &= \sqrt{2/3} A^{\Delta} + \sqrt{1/3} A^N. \end{aligned} \quad (2)$$

Walker's¹ amplitudes, A^{V3} , A^{V1} , and A^S for the production of an isospin eigenstate by the isovector and isoscalar parts of the electromagnetic current and the corresponding CGLN² amplitudes $A^{(3)}$, $A^{(1)}$, and $A^{(0)}$ are given by

$$\begin{aligned} A^{V3} &= A^{\Delta} = \sqrt{3} A^{(3)}, \\ A^{V1} &= 1/2(A^N - A^P) = \sqrt{1/3} A^{(1)}, \\ A^S &= 1/2(A^N + A^P) = \sqrt{2/3} A^{(0)}. \end{aligned} \quad (3)$$

Resonance Couplings. Since we are interested in intermediate resonances, we can approximate the energy dependence of the photoproduction partial waves by a Breit-Wigner form with added background. Using the helicity and parity eigenstates,³ $C_{\lambda}^{\ell\pm}(W)$, this gives

$$C_{\lambda}^{\ell\pm}(W) = \epsilon \left\{ \frac{\Gamma^{\lambda}(N^* \rightarrow \gamma N) \Gamma(N^* \rightarrow \pi N)}{kq} \right\}^{1/2} \frac{W}{W^2 - m_R^2 - iW\Gamma} + \text{background} \quad (4)$$

where ϵ is the sign of the amplitude and k, q are the c.m. momenta of the initial, final states. The resonance energy is m_R and Γ is its total width. In the following discussion of resonance couplings, we use the notation \tilde{A} for the imaginary part of the resonance coupling to an amplitude $A(W)$ evaluated at resonance ($W = m_R$). Thus

$$\tilde{C}_{\lambda}^{\ell\pm} = \epsilon \left\{ \frac{\Gamma^{\lambda} \Gamma_{\pi}}{kq |1|^2} \right\}^{1/2} \quad (5)$$

A dominant feature in pion photoproduction is the Born approximation which contains the nucleon pole in the s- and u-channels and the pion pole in the t-channel. It reproduces, for example, the experimentally observed forward peak in charged-pion photoproduction. In partial-wave analyses, the sign factor ϵ is well determined relative to the Born terms.

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Introducing helicity amplitudes A_{λ}^{jP} for the decay $N^*(j^P) \rightarrow (\gamma N)_{\lambda}$, where j^P labels the spin and parity of the N^* , we can calculate the radiative width⁴, $\Gamma_{\gamma}^{\lambda}$,

$$\Gamma_{\gamma}^{\lambda} = \frac{k^2}{\pi} \frac{m_N}{m_R} \frac{1}{2j+1} \left(\tilde{A}_{\lambda}^{jP} \right)^2, \quad (6)$$

where m_N is the nucleon mass. Introducing this expression into Eq. (5) we find

$$\tilde{C}_{\lambda}^{\ell\pm} = \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} \tilde{A}_{\lambda}^{jP}. \quad (7)$$

We quote the results of partial-wave analyses in terms of the amplitudes \tilde{A}_{λ}^{jP} in units of $\text{GeV}^{-1/2}$.

The total radiative width Γ_{γ} and the corresponding contribution $\sigma_T^{\ell\pm}$ of the partial waves $C_{\lambda}^{\ell\pm}$ to the total cross section are given by

$$\Gamma_{\gamma} = \sum_{\lambda=-3/2}^{3/2} \Gamma_{\gamma}^{\lambda} = \frac{k^2}{\pi} \frac{m_N}{m_R} \frac{2}{2j+1} \left\{ \left(\tilde{A}_{1/2}^{jP} \right)^2 + \left(\tilde{A}_{3/2}^{jP} \right)^2 \right\}, \quad (8)$$

$$\sigma_T^{jP} = 2 \left(C_{\pi N}^I \right)^2 \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \left\{ \left(\tilde{A}_{1/2}^{jP} \right)^2 + \left(\tilde{A}_{3/2}^{jP} \right)^2 \right\}. \quad (9)$$

The Hebb-Walker amplitudes¹, $\tilde{A}_{\ell\pm}^1$, $\tilde{B}_{\ell\pm}^1$, are related to the \tilde{A}_{λ}^{jP} by

$$\tilde{A}_{\ell\pm}^1 = \mp \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} C_{\pi N}^I \tilde{A}_{1/2}^{jP}, \quad (10)$$

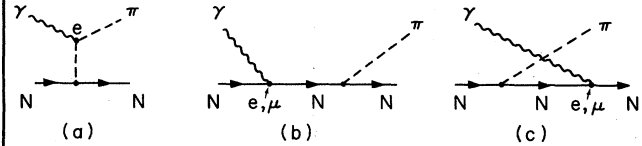
$$\tilde{B}_{\ell\pm}^1 = \pm \left\{ \frac{1}{(2j+1)\pi} \frac{k}{q} \frac{m_N}{m_R} \frac{\Gamma_{\pi}}{\Gamma^2} \right\}^{1/2} \times \left\{ \frac{16}{(2j-1)(2j+3)} \right\}^{1/2} C_{\pi N}^I \tilde{A}_{3/2}^{jP}. \quad (11)$$

We see from (2), (4), (5) and (7) that, apart from the factor $[16/(2j-1)(2j+3)]^{1/2}$, these are just the imaginary part of the resonance part of the amplitude (2) at resonance. Consequently, they are more directly related to experiment than the \tilde{A}_{λ}^{jP} , which incorporate an additional uncertainty in the partial width $\Gamma(N^* \rightarrow \pi N)$.

Methods of Partial-Wave Analysis. (a) Simple Isobar Model: In this method, the partial-wave amplitudes are formulated as in (4) (or similarly), with the resonance or resonances built partly

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Fig. IV.1. Feynman diagrams for the Born terms in pion photoproduction. The couplings of the photon to the pion and nucleon are indicated by an e to represent a coupling via charge and by a μ for a coupling via anomalous magnetic moment.

using knowledge of m_R , Γ , and $\Gamma(N^* \rightarrow \pi N)$ from πN elastic partial-wave analysis. The parametrization of the background is generally author-dependent, as is the detailed form of the resonance term, but the background is generally assumed to vary more slowly than the resonances. When the partial waves of Eq. (4) have been summed to form the angular-dependent helicity amplitudes, the Born amplitudes given by the Feynman diagrams of Fig. IV.1 must be added. [Alternatively, these Born amplitudes may previously have been projected into partial waves and included in the background terms of Eq. (4).]

As indicated in Fig. IV.1, the photon couples to the nucleon through both the electric charge, e , and the anomalous magnetic moment, μ . The minimum gauge invariant form, which includes the pion pole, is given by the pion pole plus the nucleon pole terms with the electric charge interaction only. Generally, only these terms, which we can refer to as the electric Born terms, are included explicitly because they reproduce the forward peak in charged-pion production. The magnetic terms can be assumed to be included in the real parts of the background and resonance terms in Eq. (4). However, this certainly raises problems of principle since some of the electric Born terms might also be subsumed in the background and resonances.

The best values of the resonance couplings, $\Gamma(N^* \rightarrow \pi N)$, and background parameters are found as those which give minimum χ^2 for all data within the energy range, in the usual manner for DPWA.

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(b) Fixed-t Dispersion Relations (FTDR): The difficulties of principle with the Born approximation in the simple isobar model are among the reasons which have led to the use of FTDR:

$$\text{Re}A_i^{(\pm,0)}(s,t) = B_i^{(\pm,0)}(s,t) + \frac{1}{\pi} P \int_{(m_N+m_\pi)^2}^{\infty} ds' \times \left\{ \frac{\text{Im}A_i^{(\pm,0)}(s',t)}{s'-s} + \xi_i \frac{\text{Im}A_i^{(\mp,0)}(s',t)}{s'-u} \right\} \quad (12)$$

for the invariant amplitudes $A_i^{(\pm,0)}(s,t)$ ($i=1, 2, 3, 4$), where $B_i^{(\pm,0)}(s,t)$ are the Born amplitudes⁵ of Fig. IV.1, $\xi_1 = \xi_2 = \xi_4 = 1$, $\xi_3 = -1$ and $\pm, 0$ denote $\gamma p \rightarrow \pi^+ n$, $\gamma n \rightarrow \pi^- p$, $\gamma p \rightarrow \pi^0 p$ respectively. $A_i^{(\pm,0)}(s,t)$ can be expressed in terms of the helicity amplitudes, and vice versa, by standard formulae⁵. $B_i^{(\pm,0)}(s,t)$ includes both the electric Born and anomalous magnetic moment terms, and thus no ambiguity arises as in the simple isobar model because all the real part (including the "resonant" real part) is manufactured by the unique prescription in Eq. (12).

The $\text{Im}A_i^{(\pm,0)}(s,t)$ are constructed using only the imaginary part of Eq. (4), or alternative formulae⁵. The $\text{Re}A_i^{(\pm,0)}(s,t)$ are obtained from Eq. (12) and minimization of χ^2 in fitting to experimental data to find the unknown parameters of Eq. (4) as is usual in DPWA.

A second advantage of FTDR over the simple isobar model is that the imaginary parts appear to be resonance-dominated^{1,5} (but not resonance-saturated), and this is not true of the real parts. The FTDR method only parametrizes the imaginary parts, and the smaller influence therein of background terms, relative to the size of the background in the real parts, is helpful. This advantage is diminished by the need, apparent from Eq. (12), to parametrize $\text{Im}A_i^{(\pm,0)}(s,t)$ outside the energy region of fitting for many analyses, thus effectively adding extra parameters to describe the resulting real background in the energy region being studied. This has been treated in two recent analyses^{6,7} by fitting data simultaneously in the resonance region and in the high energy (Regge) region.

Another, and in principle serious, disadvantage is that $\text{Im}A_i^{(\pm,0)}(s,t)$ in Eq. (12) is required for values of t outside the physical range corresponding to s' . Formally, $\text{Im}A_i^{(\pm,0)}(s',t)$ is given by the partial-wave expansion in $\cos\theta$ (θ is the production angle in the center-of-mass system), and this is the method used in practice to calculate the dispersion integrals. While the convergence of this expansion has not been proved for important parts of the unphysical region, Devenish, Lyth, and Rankin⁸ have surmised on the basis of the Mandelstam representation that the method is good up to $-t \leq 1$ (GeV/c)² in π^\pm production and to $-t \leq 1.5$ (GeV/c)² in π^0 production.

At low energies, the unitarity of the S-matrix imposes a phase condition on the partial-wave amplitudes known as Watson's theorem. It states that in the elastic region the complex phase of each $C_\lambda^{\ell\pm}(W)$ is equal to the scattering phase of the corresponding πN partial wave. This is important across the first resonance region in photoproduction since it imposes a strict phase condition on the two P_{33} partial waves and results in important non-resonant S-wave background in the imaginary parts contributing to the dispersion integrals. However, while the effects of Watson's theorem are significant in the first resonance region, the consequences in a FTDR calculation of ignoring it appear to be slight at higher energies.

The use of FTDR in pion photoproduction has a long history in the first resonance region⁹ where their role was envisaged as largely predictive or synthetic. The independently motivated extension by the Berkeley^{5,10,11,12}, Lancaster^{6,13}, and Glasgow^{7,14} based groups to a larger energy region has changed their role to one that is almost purely analytic.

(c) Energy-Independent Partial-Wave Analyses: Recent improvements in experimental data have allowed Berends and Donnachie¹⁵ to extend the energy-independent analysis of pion photoproduction to beyond the second resonance region. Previous analyses¹⁶ were restricted by the data then available to the first resonance region where Watson's theorem can be used to get a unique

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solution for the partial waves. In general, the energy-independent analyses are in good qualitative agreement with the energy-dependent analyses but there are some qualitative discrepancies. The significance of these is not yet clear.

Definitions of Resonance Parameters and Errors: From analyses which use a form like that of Eq. (4) there are reported in the Data Card Listings and Table IV.1, photonic resonance couplings which are obtained through Eq. (4) from a Breit-Wigner partial width. The FTDR method fits only the imaginary part of Eq. (4) to experimental data. The FTDR method of the Berkeley group^{5,10,11,12} uses a K-matrix ansatz for the imaginary part and in these cases it is the corresponding K-matrix pole quantity that is reported. (Regretably, no authors give also partial widths corresponding to T-matrix residues.)

If no errors are assigned, the authors have given a unique result without quoting an error. The Berkeley^{5,10,11,12} and Glasgow⁷ analyses quote as an error the spread around a central value of a number of solutions. The Lancaster group^{6,13} estimate a "real error" for each parameter of a given solution that corresponds to the change in value required to increase "the best possible χ^2 " by 1%.

The variation between the central values of the various papers are generally about the same size as or greater than the quoted uncertainties in individual papers. This indicates that the systematic effects of different parametrizations and choices of experimental data are at least as important as statistical effects when estimating the uncertainty in the couplings. This is reflected in the errors given in Table IV.1.

Recent Partial-Wave Analyses: We describe only the most recent analyses by the various groups involved. A description of earlier analyses may be found in the previous edition of this Review.³

(a) Simple Isobar Model: The most important of these is still METCALF 74¹⁷ which uses the methods of Walker¹ to analyze data for $\gamma p \rightarrow \pi^+ n$,

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TABLE IV.1. Photon resonance couplings.

State	λ	Helicity Couplings (GeV) ^{-1/2} 10 ⁻³		Helicity Couplings (GeV) ^{-1/2} 10 ⁻³	
		Analyses Average ⁽¹⁾		Quarks ⁽²⁾	
		\tilde{A}_λ^p	\tilde{A}_λ^n	\tilde{A}_λ^p	\tilde{A}_λ^n
P ₁₁ ⁺ (1470)	1/2	-74±15	34±35	27	-18
D ₁₃ ⁺ (1520)	1/2	-10±15	-75±15	-34	-31
	3/2	171±15	-129±10	109	-109
S ₁₁ ⁺ (1535)	1/2	63±25	-49±35	156	-108
D ₁₅ ⁺ (1670)	1/2	16±10	-32±36	0	-38
	3/2	21±12	-62±40	0	-53
F ₁₅ ⁺ (1688)	1/2	-5±30	25±10	-10	30
	3/2	127±35	-16±20	60	0
S ₁₁ ⁰ (1700)	1/2	43±30	-37±40	0	30
D ₁₃ ⁰ (1700)	1/2	-20±45	18±55	0	-10
	3/2	19±45	18±80	0	-40
P ₁₁ ⁰ (1780)	1/2	18±40	18±50	-40	10
P ₁₃ ⁰ (1810)	1/2	-25±50	24±70	100	-30
	3/2	-31±50	-4±60	-30	0
		$\tilde{A}_\lambda^\Delta = \tilde{A}_\lambda^{\sqrt{3}}$	$\tilde{A}_\lambda^\Delta = \tilde{A}_\lambda^{\sqrt{3}}$		
P ₃₃ ⁺ (1232)	1/2	-139±5		-108	
	3/2	-256±5		-187	
S ₃₁ ⁺ (1650)	1/2	46±36		47	
D ₃₃ ⁺ (1670)	1/2	72±26		88	
	3/2	72±45		84	
P ₃₃ ⁰ (1690)	1/2	-2±40		23	
	3/2	-12±50		39	
F ₃₅ ⁺ (1890)	1/2	21±30		-20	
	3/2	-10±60		-90	
P ₃₁ ⁺ (1910)	1/2	-11±20		-30	
F ₃₇ ⁺ (1950)	1/2	-69±16		-50	
	3/2	-76±20		-70	

(1) Average of the couplings from MOORHOUSE 73, DEVENISH 73, MOORHOUSE 74, METCALF 74, KNIES 74, DEVENISH2 74, CRAWFORD 75, BARBOUR 76, and Ref. 12. The errors given are an estimate that takes into account the errors quoted in these analyses and the differences between the analyses.

(2) The naive, l -excitation, quark model is used in the 4-dimensional oscillator form of Feynman, Kislinger, and Ravndal. The non-relativistic quark model with recoil gives generally very similar results.

$\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ from the first through the fourth resonance regions. The partial waves are parametrized as in Eq. (4) with the background taken to be an independent number at each energy fitted. The electric Born terms are added explicitly.

Data Card Listings

For notation, see key at front of Listings.

Baryons N's and Δ's

Other recent isobar analyses have been on a much smaller scale. KRIVETS 74¹⁸ is an analysis of $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ across the first, second, and third resonance regions using a rather small data set consisting only of differential cross-section measurements. Several solutions are given with a large spread in the values of the couplings and those for the resonance couplings to γn are stated to be preliminary. HEMM1 73¹⁹ is a fit to forward $\gamma p \rightarrow \pi^0 p$ differential cross-section data across the second and third resonance regions and evaluates the $\lambda = 1/2$ couplings in this region. HEMM2 73²⁰ is an isobar analysis of $\gamma n \rightarrow \pi^0 n$ across the second resonance region. ROSSI 73²¹ and BENEVENTANO 73²² are analyses of data for $\gamma n \rightarrow \pi^- p$ across the second resonance region. Earlier analyses^{1,23} are not included in the Data Card Listings.

(b) Fixed-t Dispersion Relations:

MOORHOUSE 73¹⁰ and MOORHOUSE 74⁵ are FTDR analyses from threshold through the third resonance region ($1160 < W < 1780$ MeV) using data for $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$. The first uses the results of seven fits with a χ^2 per data point ranging from 9.7 with 52 variables to 5.7 with 74. The second uses the average of three solutions with a χ^2 per data point ranging from 4.0 (56 variables) to 3.0 (74 variables). KNIES 74¹¹ and Ref. 12 extend the energy range of the Berkeley analyses up to a center-of-mass total energy of 2 GeV, and thus include the energy region of the $\Delta(1950)F_{37}$. Several solutions are described with typical values of the χ^2 per data point of between 3.5 and 4.5. In all these analyses, the imaginary parts of the partial waves are parametrized using a K-matrix formalism.

DEVENISH 73¹³ analyze data from the first through the third resonance regions and parametrize the imaginary parts of the partial waves as Breit-Wigner forms without background except for the π^0 -production S wave. The dispersion integral is cut off at $W = 1.9$ GeV, but parametrized real background is added to allow for the integrals above 2 GeV. Three solutions are given with χ^2 per data point of 10.5, 3.0, and 4.7 for

37, 68, and 68 free parameters, respectively. DEVENISH2 74⁶ treats the problem of the high energy contributions to the dispersion integral by simultaneously fitting data in the resonance region and at high energies and evaluates the couplings for all resonances for $W \leq 2$ GeV. It uses data for $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$ at all available energies and for $|t| < 1$ (GeV/c)². A χ^2 per data point of 4.8 is obtained with 48 resonance parameters and 50 high energy parameters.

CRAWFORD 75¹⁴ is an FTDR analysis from the first through the fourth resonance regions of $\gamma p \rightarrow \pi^+ n$, $\pi^0 p$ and $\gamma n \rightarrow \pi^- p$. Data are fitted for $W \leq 2$ GeV and $|t| \leq 1.5$ (GeV/c)². Three solutions are given with a best χ^2 per data point of 3.1 with 58 free parameters. The masses and widths of all resonances below 1.7 GeV are evaluated. Stable values are obtained that agree well with the values from elastic πN partial-wave analyses. BARBOUR 76⁷ is an extension of the previous analysis to include the fitting of high energy data for $W > 2.6$ GeV. One fit is presented with a χ^2 per data point of 2.6 for the resonance region data and 1.8 for the high energy data. No data are fitted for $2 \text{ GeV} < W < 2.6 \text{ GeV}$, thus omitting the region of the $H_{3,11}$.

(c) Energy-Independent Analyses: These have now been extended to the second resonance region by Berends and Donnachie¹⁵ as described above. Earlier analyses¹⁶ fitted only in the first resonance region. Resonance couplings have not been quoted in these analyses.

Electroproduction in the Resonance Region:

Although single-pion electroproduction can be considered in many ways to be similar to photoproduction²⁴, a partial-wave analysis of the process is considerably more difficult than in photoproduction. This is due to the less accurate and complete experimental data that are available and because the scalar part of the virtual photon in electroproduction contributes an extra (scalar) coupling for each resonance. The scalar amplitudes uncouple from the T-matrix in the photoproduction limit. Also, it is necessary, in principle, to make an analysis at each value of λ^2 , the (virtual photon mass)², at which there are

Baryons

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Data Card Listings

For notation, see key at front of Listings.

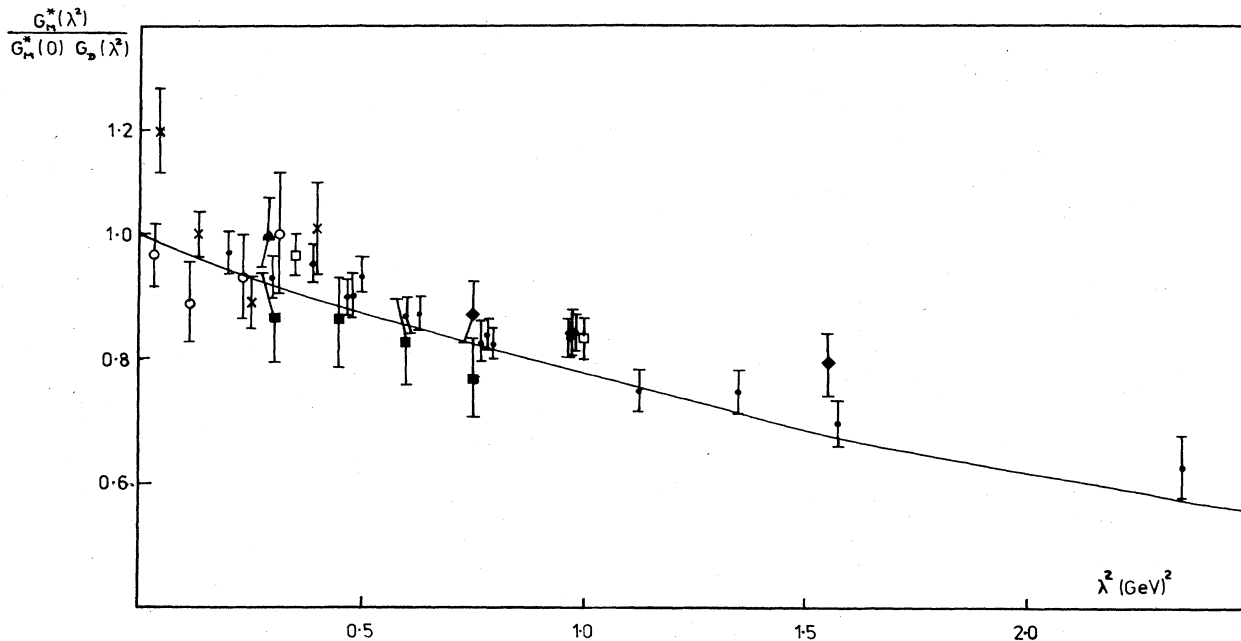


Fig. IV.2. The $\Delta(1232)$ magnetic moment form factor, $G_M^*(\lambda^2)$, normalized to $G_M^*(0) = 3$, compared to the nucleon dipole form factor, $G_D(\lambda^2) = (1 + \lambda^2/0.71)^{-2}$. The sources of the data points are \bullet ref. 25, \circ ref. 27, \times ref. 28, \blacksquare ref. 29, \square ref. 30, \blacklozenge ref. 31, \blacktriangle ref. 32, and \blacklozenge ref. 33. The solid line is from Devenish and Lyth³⁴.

data. The information to be obtained from such an analysis is the size of the scalar couplings to the resonance and the form factors of the scalar transverse couplings, i.e., their λ^2 dependence. This information has been extracted in a number of ways, depending on the experimental data available and the energy range being studied.

(a) Total Cross Sections: In a single-arm experiment, in which only the final electron is detected, it is possible to measure only σ_T and σ_S , the total cross sections for the absorption of transverse and scalar photons. In the first resonance region, due to the dominance of the $M_{1+}^{(3)}$ multipole, σ_T can give the form factor for the magnetic coupling to the $P_{33}(1232)$. There have been many analyses of this type, both the form factors for γp and γn coupling to the $P_{33}(1232)$ having been measured. We quote as examples of these, the analyses of Bartel et al.²⁵ and of Köbberling et al.²⁶ Due to the unknown non-resonant scalar background and because the scalar coupling is relatively small, σ_S gives

only qualitative information about the scalar coupling to this resonance. Above the second resonance region, due to the large number of resonances, the total cross sections again give little more than qualitative information.

(b) Energy-Independent Analyses: Differential cross-section measurements have been made for π^0 and π^+ electroproduction in a number of coincidence experiments. Although the data are not good enough for a full partial-wave analysis, it is possible, due to the small number of partial waves involved, to use the dominance of the $M_{1+}^{(3)}$ multipole to measure $|M_{1+}|$, $\text{Re } E_{1+}^* M_{1+} / |M_{1+}|^2$, and $\text{Re } S_{1+}^* M_{1+} / |M_{1+}|^2$ for π^0 production across the first resonance region. This has been done by a number of groups²⁷⁻³². Because $M_{1+}^{(3)}$ is so large and is purely imaginary at the P_{33} mass, these terms can be taken to be respectively $\text{Im} M_{1+}^{(3)}$, $E_{1+}^{(3)} / M_{1+}^{(3)}$, and $S_{1+}^{(3)} / M_{1+}^{(3)}$ at the resonance.

(c) Energy-Dependent Analyses: Energy-dependent partial-wave analyses using FTDR have been made for pion electroproduction by Crawford³³

Data Card Listings

For notation, see key at front of Listings.

Baryons N's and Δ's

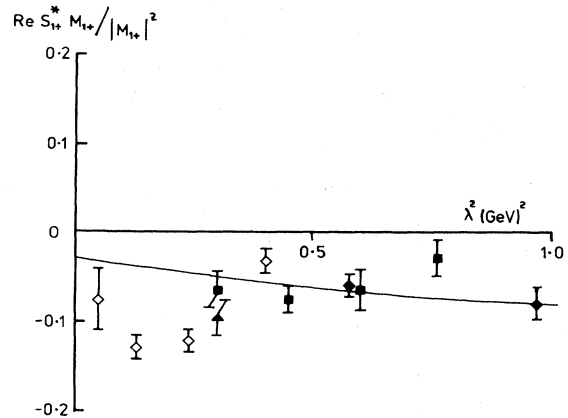
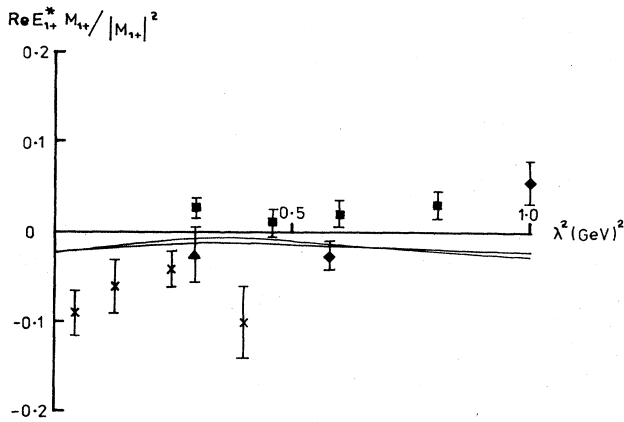


Fig. IV.3. The ratios $\text{Re}(E_{1+}^* M_{1+} / |M_{1+}|^2)$ and $\text{Re}(S_{1+}^* M_{1+} / |M_{1+}|^2)$ for π^0 electroproduction at the $\Delta(1232)$. The data points and curves are indicated as in Fig. IV.2. Crawford³³ and Devenish and Lyth³⁴ give the ratios $E_{1+}^{(3)} / M_{1+}^{(3)}$ and $S_{1+}^{(3)} / M_{1+}^{(3)}$.

and by Devenish and Lyth³⁴. The first fitted only in the first resonance region and over a relatively small range of λ^2 . The second is an extensive analysis from the first to the fourth resonance regions. Using the results of photo-production analyses to give a starting point, the form factors for the resonance couplings are given algebraic forms that reflect their analytic structure in the complex λ^2 plane. Data for π^0 and π^+ production and for a range of values of λ^2 are fitted simultaneously and the results of a number of fits are presented. The analysis evaluates the couplings for the $P_{33}(1232)$, $P_{11}(1470)$, $D_{13}(1520)$, $S_{11}(1535)$, $D_{33}(1670)$, $F_{15}(1688)$, $S_{11}(1700)$, $D_{13}(1700)$, $P_{13}(1780)$, $F_{37}(1950)$, and $D_{13}(2040)$ resonances. The results are discussed below.

The couplings and form factors for the $P_{33}(1232)$ are now established for $\lambda^2 < 1$ (GeV)² and there is agreement on all general features between the basic methods of analysis (Figs. IV.2 and IV.3). The λ^2 dependence of $M_{1+}^{(3)}$ is usually expressed in terms of the form factor $G_M^*(\lambda^2)$,²⁷

$$G_M^*(\lambda^2) = 2M \left(\frac{3}{2\alpha} \frac{|\vec{k}|}{\sin^2 \delta_{33}} \frac{\Gamma |M_{1+}^{(3)}|^2}{q^2} \right)^{1/2},$$

where \vec{k} and \vec{q} are respectively the center-of-mass three-momenta of the pion and the virtual photon.

The behavior of G_M^* for the charge = +1 state, normalized to $G_M^*(0) = 3$, is compared with the nucleon dipole form, $G_D(\lambda^2) = (1 + \lambda^2/0.71)^{-2}$, in Fig. IV.2. It is clearly seen that G_M^* falls off significantly faster with λ^2 than the nucleon dipole. Measurements of G_M^* for the charge = 0 state^{26,35} show that it has the same form as for the charge = +1 state. The measurements of the ratios of $E_{1+}^{(3)}$ and $S_{1+}^{(3)}$ to $M_{1+}^{(3)}$ are all consistent with these being small, although the $S_{1+}^{(3)}$ is clearly established as being non-zero with a value of about $-0.06 M_{1+}^{(3)}$. This is a small violation of the quark model selection rule that predicts that there should be no scalar excitation of the $P_{33}(1232)$ ³⁶.

In the second to the fourth resonance regions, Devenish and Lyth obtain definite results for the couplings to the charge = +1 states of the $P_{11}(1470)$, $D_{13}(1520)$, $S_{11}(1535)$, $F_{15}(1688)$, and $F_{37}(1950)$ resonances. These are expressed in terms of the multipole partial waves:

$P_{11}(1470)$ The form factor for the magnetic multipole rapidly vanishes and there is no evidence for a scalar coupling.

$D_{13}(1520)$ The electric multipole decreases rapidly with λ^2 in an approximate dipole form, but the electric dipole decreases significantly less quickly. There is therefore a rapid change in

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the ratio of the transverse helicity partial waves. The scalar coupling is non-zero at $\lambda^2 = 0$ but decreases rapidly.

S₁₁(1535) The electric multipole decreases slowly with λ^2 . The scalar coupling is comparable to the electric one at $\lambda^2 = 0$ but may decrease in an approximately dipole form.

F₁₅(1688) The form factors for the transverse multipoles are similar to those for D₁₃(1520). The scalar coupling decreases rapidly.

F₃₇(1950) The magnetic multipole probably decreases more slowly than the dipole form. The electric multipole is small and its λ^2 behavior is not clearly determined. The scalar multipole is small and decreases rapidly.

The form factors for the S₁₁(1535) can be checked in η electroproduction. This process proceeds mainly through S-wave excitation in the second resonance region and is thus dominated by the S₁₁(1535). The results for the cross section of this process³⁷ again require that at least one of the form factors should fall off slowly with λ^2 .

The rapid change in the ratio of the transverse helicity amplitudes of the D₁₃(1520) and the F₁₅(1688), along with the almost constant value of the ratio for the P₃₃(1232), gives qualitative agreement with quark model calculations^{4,38}. However, some quark models find that the form factor for the S₁₁(1535) should decrease rapidly and that the one for the P₁₁(1470) should stay up, and are therefore in conflict with the results above.

Information in this Edition: The Baryon Table contains the branching fractions, Γ_Y/Γ for 14 resonances.

The Data Card Listings contain the photon resonance couplings (\tilde{A}_λ^j)^P for p, n, and Δ from the analyses of Moorhouse and Oberlack¹⁰ (MOORHOUSE 73), Moorhouse, Oberlack, and Rosenfeld⁵ (MOORHOUSE 74), Knies, Moorhouse, and Oberlack¹¹ (KNIES 74), Devenish, Lyth, and Rankin^{13,6} (DEVENISH 73 and DEVENISH2 74),

Data Card Listings

For notation, see key at front of Listings.

Crawford¹⁴ (CRAWFORD 75), Barbour and Crawford⁷ (BARBOUR 76), Metcalf and Walker¹⁷ (METCALF 74), Hemmi et al.¹⁹ (HEMMI1 73), Hemmi et al.²⁰ (HEMMI2 73), Rossi et al.²¹ (ROSSI 73), and Beneventano et al.²² (BENEVENTANO 73).

The average of the listed results along with an estimate of the uncertainty in each from the spread of results is given in Table IV.1. We give results from a naive (λ -excitation) quark model for comparison.

References for Section IV

1. R. L. Walker, Phys. Rev. **182**, 1729 (1969).
2. G. F. Chew, M. L. Goldberger, F. E. Low and Y. Nambu, Phys. Rev. **106**, 1345 (1956).
3. These amplitudes have been described in the previous edition of this Review, Phys. Lett. **50B**, No. 1 (1974). See also reference 5.
4. L. A. Copley, G. Karl and E. Obryk, Nucl. Phys. **B13**, 303 (1969).
5. R. G. Moorhouse, H. Oberlack and A. H. Rosenfeld, Phys. Rev. **D9**, 1 (1974).
6. R. C. E. Devenish, D. H. Lyth and W. A. Rankin, Phys. Lett. **B52**, 227 (1974).
7. I. M. Barbour and R. L. Crawford, Glasgow University preprint (1975). Results reported by A. Donnachie, rapporteur's talk, Stanford Conference (1975).
8. R. C. E. Devenish, D. H. Lyth and W. A. Rankin, Daresbury report DNPL/P 109 (1972).
9. F. A. Berends, A. Donnachie and D. L. Weaver, Nucl. Phys. **B4**, 1 and 54 (1967). J. Engels, A. Müllensiefen and W. Schmidt, Phys. Rev. **175**, 1951 (1968).
10. R. G. Moorhouse and H. Oberlack, Phys. Lett. **43B**, 44 (1973).
11. G. Knies, R. G. Moorhouse and H. Oberlack, Phys. Rev. **D9**, 2680 (1974).
12. G. Knies, R. G. Moorhouse, H. Oberlack, A. Rittenberg and A. H. Rosenfeld, Proc. of the XVII Int. Conf. on High Energy Physics, London 1974, edited by J. Smith (Science Research Council, Chilton), reported by D. H. Lyth, p. II, 150.
13. R. C. E. Devenish, D. H. Lyth and W. A. Rankin, Phys. Lett. **47B**, 53 (1973).
14. R. L. Crawford, Nucl. Phys. **B97**, 125 (1975).
15. F. A. Berends and A. Donnachie, paper submitted to Stanford Conference (1975).

Data Card Listings

Baryons

For notation, see key at front of Listings.

N's and Δ's, p, n

16. P. Noelle, W. Pfeil and D. Schwela, Nucl. Phys. B26, 461 (1971). W. Pfeil and D. Schwela, Nucl. Phys. B45, 379 (1972). F. A. Berends and D. L. Weaver, Nucl. Phys. B30, 575 (1971). Yu. M. Aleksandrov, V. F. Grushin, E. M. Leiken and A. Ya. Rotvain, Nucl. Phys. B45, 589 (1972).
17. W. J. Metcalf and R. L. Walker, Nucl. Phys. B76, 253 (1974).
18. A. G. Krivets et al. Yad. Fiz. 20, 804 (1974), Sov. J. Nucl. Phys. 20, 430 (1975).
19. Y. Hemmi et al. Phys. Lett. 43B, 79 (1973).
20. Y. Hemmi et al., Nucl. Phys. B55, 333 (1973).
21. V. Rossi et al. Nuovo Cimento 13A, 59 (1973).
22. M. Beneventano et al., Nuovo Cimento 19A, 529 (1973).
23. Y. C. Chau, N. Dombey and R. G. Moorhouse, Phys. Rev. 163, 1632 (1967). R. G. Moorhouse and W. A. Rankin, Nucl. Phys. B23, 181 (1970). R. L. Walker, Proc. of the Daresbury Conf. on Electron and Photon Interactions at High Energy (edited by D. Braben, Daresbury Laboratory, 1969).
24. See for example Berends, Donnachie and Weaver, Reference 9.
25. W. Bartel et al., Phys. Lett. 28B, 148 (1968).
26. M. Köbberling et al., Nucl. Phys. B82, 201 (1974).
27. W. W. Ash et al., Phys. Lett. 24B, 165 (1967).
28. C. Mistretta et al., Phys. Rev. 184, 1487 (1969).
29. R. Siddle et al., Nucl. Phys. B35, 93 (1971).
30. S. Golster et al., Phys. Rev. D5, 519 (1972).
31. J. C. Alder et al., Nucl. Phys. B46, 573 (1972).
32. K. Bätzer et al., Bonn University preprint (1973).
33. R. L. Crawford, Nucl. Phys. B28, 573 (1971).
34. R. C. E. Devenish and D. H. Lyth, Nucl. Phys. B93, 109 (1975).
35. R. V. Akhmerov et al., Sov. J. Nucl. Phys. 21, 57 (1975).
36. C. Becchi and G. Morpurgo, Phys. Lett. 17, 352 (1965).

37. P. S. Kummer, et al., Phys. Rev. Lett. 30, 873 (1973). U. Beck et al., Phys. Lett. 51B, 103 (1974). J. C. Alder et al., Nucl. Phys. B91, 386 (1975).
38. F. Ravndal, Phys. Rev. D4, 1466 (1971). L. A. Copley, G. Karl and E. Obryk, Phys. Rev. D4, 2844 (1971). R. G. Lipes, Phys. Rev. D5, 2849 (1972). T. Abdullah and F. E. Close, Phys. Rev. D5, 2332 (1972).

STATUS OF N* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --										OTHER CHANN.			
			TOTAL* CR.S.	PI	N	ETA	N	K	LAM	K	SIG	PI		DE		
N*(940)	P11	****														
N*(1470)	P11	****	****													EPS N
N*(1520)	D13	****	****	*												RHC N
N*(1535)	S11	****	****	***												RHO N
N*(1670)	D15	****	****	***												RHC N
N*(1688)	F15	****	****	*												RHC N
N*(1700)	S11	****	****	*												EPS N
N*(1700)	D13	**	**	*						*						EPS N
N*(1780)	P11	**	**	*						*						EPS N
N*(1810)	P13	**	**	*						*						RHO N
N*(1990)	F17	**	**	*						*						
N*(2000)	F15	*	*	*						*						
N*(2040)	D13	**	**	*						*						
N*(2100)	S11	*	*	*						*						
N*(2100)	D15	*	*	*						*						
N*(2190)	G17	**	**	***						*						
N*(2220)	H19	**	**	***						*						
N*(2650)		**	**	*						*						
N*(3030)		***	***	*						*						
N*(3245)		*	*	*						*						
N*(3690)		*	*	*						*						
N*(3755)		*	*	*						*						
DE(1232)	P33	****	****	****						F						****
DE(1650)	S31	****	**	****						O						RHO N
DE(1670)	D33	***	**	***						R						****
DE(1690)	P33	*	*	*						B						****
DE(1890)	F35	**	*	***						I						RHC N
DE(1900)	S31	*	*	*												****
DE(1910)	F31	**	*	**						D						RHC N
DE(1950)	F37	***	***	****						D						RHO N
DE(1960)	D35	**	*	*						E						****
DE(2160)	**	**	***	***						N						****
DE(2420)	F311	**	**	***						F						****
DE(2850)	**	**	***	**												****
DE(3230)	**	**	*	*						C						****

*** GOOD, CLEAR, AND UNMISTAKABLE.
** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
* NEEDS CONFIRMATION.
* WEAK.
* ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

S=0 I=1/2 NUCLEON STATES (N)

p

16 PROTON(938, J=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

n

17 NEUTRON(939, J=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

Baryons
N(1470)

Data Card Listings

For notation, see key at front of Listings.

N(1470)

61 N*1/2(1470, JP=1/2+) I=1/2
MASS AND WIDTH ARE BEST DETERMINED FROM PARTIAL WAVE ANALYSES. WE LIST PRODUCTION EXPERIMENTS SEPARATELY -- SEE BELOW.
AYED 74 CLAIM TWO P11 STATES IN THE 1500 MEV REGION. WE TENTATIVELY LIST BOTH HERE.

P11

Table with columns for mass (MEV) and various experimental data points for N(1470).

Table with columns for width (MEV) and various experimental data points for N(1470).

Table with columns for real part of pole position (MEV) and various experimental data points for N(1470).

Table with columns for 2*imag part of pole position (MEV) and various experimental data points for N(1470).

Table with columns for absolute value of pole residue (MEV) and various experimental data points for N(1470).

Table with columns for phase of pole residue (radians) and various experimental data points for N(1470).

Table with columns for partial decay modes and various experimental data points for N(1470).

Table with columns for branching ratios and various experimental data points for N(1470).

Table with columns for branching ratios and various experimental data points for N(1470).

Table with columns for branching ratios and various experimental data points for N(1470).

Table with columns for various experimental data points for N(1470) including decay widths and theoretical estimates.

Table with columns for photon decay amplitudes and various experimental data points for N(1470).

Table with columns for references for N(1470) listing various scientific publications.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1470), N(1520)

Table listing authors and their affiliations for papers not referred to in data cards. Includes names like BRICMAN, VALLADAS, VILLET, etc.

THE FOLLOWING ARE THEORETICAL PAPERS CONCERNING THE N(1470) -- RESONANT STATE, IF ANY, THE HBC EXPERIMENTS SEE ENHANCEMENTS MAINLY IN THE P PI MASS PLOT. FOR ZERO CHARGE EXCHANGE, SUCH FINAL STATES ARE KNOWN TO HAVE LARGE DECK-TYPE BACKGROUND. THIS FACT COMPLICATES THE INTERPRETATION OF THIS BUMP AS A RESONANCE.

1470 MEV REGION - PRODUCTION EXPERIMENTS

91 N(1470, JP=) I=1/2 PRODUCTION EXPERIMENTS

THE BUMP SEEN IN PRODUCTION EXPERIMENTS AT LOW INVARIANT MASS MOST LIKELY CORRESPONDS TO THE P11 (SEE ABOVE) ENHANCEMENTS MAINLY IN THE P PI MASS PLOT. FOR ZERO CHARGE EXCHANGE, SUCH FINAL STATES ARE KNOWN TO HAVE LARGE DECK-TYPE BACKGROUND. THIS FACT COMPLICATES THE INTERPRETATION OF THIS BUMP AS A RESONANCE.

Main data table for N(1470) production experiments. Columns include mass (MEV), production method (e.g., APPROX, CERN), and various parameters like p, m, s, t, w, x, y, z.

91 N(1470) WIDTH (MEV) (PRCD. EXP.)

Table listing width measurements for N(1470) in MeV, including authors like BELL, SHAPIRA, etc.

91 N(1470) PARTIAL DECAY MODES (PRCD. EXP.)

Table listing partial decay modes for N(1470), such as N(1470) to N(1470) + pi, N(1470) to N(1470) + pi + pi, etc.

Table listing production experiments for N(1470) with columns for author, mass, production method, and various parameters.

REFERENCES FOR N(1470) (PRCD. EXP.)

Table listing references for N(1470) production experiments, including authors like COCCONI, ADELMAN, etc.

N(1520)

62 N(1520, JP=3/2-) I=1/2 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

Table listing production experiments for N(1520) with columns for author, mass, production method, and various parameters.

62 N(1520) WIDTH (MEV)

Table listing width measurements for N(1520) in MeV, including authors like BAREYRE, DONNACHI, etc.

62 N(1520) REAL PART OF POLE POSITION (MEV)

Table listing real part of pole position for N(1520) in MeV.

62 N(1520) 2*IMAG PART OF POLE POSITION (MEV)

Table listing imaginary part of pole position for N(1520) in MeV.

Baryons
N(1520)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle ID, decay mode, and decay masses. Includes entries for N(1520) partial decay modes.

Table with columns for particle ID, branching ratios, and other parameters. Includes entries for N(1520) branching ratios.

Text block providing additional information for the branching ratios table, including a note about inelasticity.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into N(1232) pi pi total.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into gamma N.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into gamma N.

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into gamma N.

AVERAGE MEANINGLESS (SCALE FACTOR = 3.9)

Table with columns for particle ID, decay mode, and other parameters. Includes entries for N(1520) into gamma N.

AVERAGE MEANINGLESS (SCALE FACTOR = 3.1)

REFERENCES FOR N(1520)

Table with columns for author names, publication details, and references. Includes entries for BRANDSEN, KIRZ, NAMYSLOM, etc.

Table with columns for author names, publication details, and references. Includes entries for KIRZ, BAREYRE, CROUCH, etc.

Data Card Listings
For notation, see key at front of Listings.

Baryons
N(1535)

N(1535)

63 N*1/2(1535, JP=1/2-) I=1/2 S11
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED. IT IS STRONGLY ASSOCIATED WITH THE ETA N CHANNEL.

Table with columns for mass (MEV) and various parameters. Includes entries for HENDRY, MICHAEL, UCHIYAMA, BAREYRE, DONNACHI, DELCOURT, etc.

63 N*1/2(1535) WIDTH (MEV)

Table with columns for width (MEV) and various parameters. Includes entries for HENDRY, MICHAEL, UCHIYAMA, BAREYRE, DONNACHI, DELCOURT, etc.

63 N*1/2(1535) REAL PART OF POLE POSITION (MEV)

Table with columns for real part of pole position (MEV) and various parameters. Includes entry for LONGACRE.

63 N*1/2(1535) 2*IMAG PART OF POLE POSITION (MEV)

Table with columns for 2*imag part of pole position (MEV) and various parameters. Includes entry for LONGACRE.

63 N*1/2(1535) PARTIAL DECAY MODES

Table with columns for partial decay modes and various parameters. Includes entries for INTO PI N, INTO N ETA, INTO N PI PI, etc.

63 N*1/2(1535) BRANCHING RATIOS

Table with columns for branching ratios and various parameters. Includes entries for INTO (PI N)/TOTAL, DOMINANT INEL DECAY, INTO (N ETA)/TOTAL, etc.

Table with columns for various parameters and values. Includes entries for INTO GAMMA PROTON TO ETA PROTON, FROM PI N INTO ETA N, FROM PI N TO K LAMBDA, etc.

63 N*1/2(1535) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for photon decay amplitudes and various parameters. Includes entries for INTO GAM P, HELICITY=1/2, INTO GAM N, HELICITY=1/2, etc.

REFERENCES FOR N*1/2(1535)

Table with columns for references and various parameters. Includes entries for HENDRY, MICHAEL, UCHIYAMA, BAREYRE, DONNACHI, DELCOURT, etc.

Baryons

N(1535), N(1670)

THE FOLLOWING ARTICLES DEAL WITH THE REACTIONS PI- P TO ETA N AND GAMMA P TO ETA P NEAR THRESHOLD. THE DATA AND THE THEORETICAL ARTICLES ARE USEFUL IN UNDERSTANDING THE BEHAVIOR OF THE S11 AMPLITUDE AS DETERMINED IN PI P PHASE-SHIFT ANALYSES. FURTHER REFERENCES MAY BE FOUND IN THEM.

Table listing experimental and theoretical references for N(1535) and N(1670) baryons, including authors like BULOS, BACCI, JONES, RICHARDS, PREPOST, BLOCM, HEUSCH, BINNIE, BALL, DORSEN, MINAMI, DEANS, LOGAN, MENCUCCI, MIAMI, MOSS, DEANS, PAL, BALL, LEFTVIRE.

1520 MEV REGION - PRODUCTION EXPERIMENTS

8 N*1/2(1520, JP=) I=1/2 PRODUCTION EXPERIMENTS

THIS INFORMATION REFERS TO EITHER THE D13 OR THE S11 STATE SEEN AT THIS MASS. FOR SPIN-PARITY ANALYSIS OF THIS MASS REGION, SEE JOHNSTAD 72.

8 N*1/2(1520) MASS (MEV) (PROD. EXP.)

Table of mass measurements for N*1/2(1520) with columns for mass, width, and references.

8 N*1/2(1520) WIDTH (MEV) (PROD. EXP.)

Table of width measurements for N*1/2(1520) with columns for width, mass, and references.

8 N*1/2(1520) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for N*1/2(1520) listing decay channels and branching ratios.

8 N*1/2(1520) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for N*1/2(1520) listing various decay channels and their relative probabilities.

REFERENCES FOR N*1/2(1520) (PROD. EXP.)

Table of references for N*1/2(1520) production experiments, listing authors and their respective publications.

Data Card Listings

For notation, see key at front of Listings.

Table listing data card entries for N(1670) resonance, including authors like EDLSTEIN, COOPER, BRAUN, MUSGRAVE, STRACHMA, WEBB.

N(1670) 64 N*1/2(1670, JP=5/2-) I=1/2 D15 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

Table of mass and width measurements for N*1/2(1670) with columns for mass, width, and references.

Table of real and imaginary parts of pole positions for N*1/2(1670) with columns for real part, imaginary part, and references.

Table of partial decay modes for N*1/2(1670) listing decay channels and branching ratios.

Table of branching ratios for N*1/2(1670) listing various decay channels and their relative probabilities.

Table of partial decay modes for N*1/2(1670) listing decay channels and branching ratios.

Table of branching ratios for N*1/2(1670) listing various decay channels and their relative probabilities.

Table of references for N*1/2(1670) production experiments, listing authors and their respective publications.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1670), N(1688)

R8 N*1/2(1670) FROM PI N TO N*3/2(1232) PI, D-WAVE SORT(PI*P11) 11/75*
R8 L (-.45)OR -.50 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

SEE NOTE PRECEDING THE N*1/2(1688) INELASTIC DECAY MODE MEASUREMENTS.

64 N*1/2(1670) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns for particle ID (A1, A2, A3, A4), helicity, and various decay amplitudes and phase shifts for N*1/2(1670) and N*1/2(1688).

Table containing references for N*1/2(1670) and N*1/2(1688), listing authors and publication details.

N(1688)

65 N*1/2(1688, JP=5/2+) I=1/2 F'15

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

65 N*1/2(1688) MASS (MEV)

Table listing mass measurements for N(1688) from various experiments, including Bransden, Barbour, and others.

65 N*1/2(1688) WIDTH (MEV)

Table listing width measurements for N(1688) from various experiments, including Barbour, Davies, and others.

65 N*1/2(1688) REAL PART OF POLE POSITION (MEV)

Table listing real part of pole position for N(1688) from Longacre et al.

65 N*1/2(1688) 2*IMAG PART OF POLE POSITION (MEV)

Table listing imaginary part of pole position for N(1688) from Longacre et al.

65 N*1/2(1688) PARTIAL DECAY MODES

Table listing partial decay modes for N(1688) into various particles like pi N, eta N, lambda K, etc.

65 N*1/2(1688) BRANCHING RATIOS

Table listing branching ratios for N(1688) into various decay channels.

MORE INFORMATION ON THE INELASTIC DECAY MODES OF THE 1690 MEV BUMP, AS SEEN IN PRODUCTION EXPERIMENTS, MAY BE FOUND BELOW

Table listing inelastic decay modes and branching ratios for N(1688) into various channels like eta N, lambda K, etc.

Baryons

N(1688), N(1700)

Data Card Listings

For notation, see key at front of Listings.

R10 N*1/2(1688) FROM PI N TC N RHO,S=3/2,P-WAVE SQRT(P1*P14) 11/75*
R10 L (+.27)OR +.30 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

65 N*1/2(1688) PHOTON DECAY AMPL(GEV**1/2)
FCR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*1/2(1688) INTO GAM P, HELICITY=1/2 (GEV**1/2)
A1 +.015 +.023 DEVENISH 73 DPWA PI N PHOTO PROD 2/74

A3 N*1/2(1688) INTO GAM N, HELICITY=1/2 (GEV**1/2)
A3 +.035 +.049 DEVENISH 73 DPWA PI N PHOTO PROD 2/74

A4 N*1/2(1688) INTO GAM N, HELICITY=3/2 (GEV**1/2)
A4 -.018 -.039 DEVENISH 73 DPWA PI N PHOTO PROD 2/74

REFERENCES FOR N*1/2(1688)
SEE A PREVIOUS EDITION (RMP 37, 633, 1965) FOR EARLIER REFERENCES.

BRANDSEN 65 PL 19 420 +ODDONNELL, MCOORHOUSE (DURHAM, RHEL IJJP
HEUSCH 66 PRL 17 1019 C A HEUSCH, C Y PRESCOTT, R F DASHEN (CIT)
TRIPP 67 NP B3 10 + LEITH, + (LRL,SLAC,CERN,HEID,SLACLAY)

N(1700) 66 N*1/2(1700, JP=1/2-1) I=1/2 S11
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

66 N*1/2(1700) MASS (MEV)
M (1695.0) BRANDSEN 65 RVUE PHASE-SHIFT ANAL 9/66
M 1 (1700.0) MICHAEL 66 RVUE FITS BAREYRE S11 7/66

66 N*1/2(1700) WIDTH (MEV)
W 1 (240.0) MICHAEL 66 RVUE 7/66
W 3 (300.0) DGNNACH1 68 RVUE 8/69

66 N*1/2(1700) REAL PART OF PCLE POSITION (MEV)
RE (1648.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

66 N*1/2(1700) 2*IMAG PART OF POLE POSITION (MEV)
IM (117.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

66 N*1/2(1700) PARTIAL DECAY MODES
P1 N*1/2(1700) INTO PI N DECAY MASSES 139+ 938
P2 N*1/2(1700) INTO N ETA 939+ 548

66 N*1/2(1700) BRANCHING RATIOS
R1 N*1/2(1700) INTO (PI N)/TOTAL (P1)
R1 (1.0) APPROX MICHAEL 66 RVUE 7/66

R2 N*1/2(1700) FROM PI N TO K LAMBDA SQRT(P1*P3) 4/75*
R2 (+.20) (.05) ORITO 69 RVUE 4/75*
R2 A (-.21)OR .23 WAGNER 71 IPWA PI- P TO K LAMB 4/75*
R2 -.179 .033 DEVENISH 74 IPWA O FIXED T DISP REL 4/75*
R2 (.12) KNASEL 75 DPWA O PI- P TO KO LAM 11/75*

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1700)

R7 N*1/2(1700) FROM PI N TO K SIGMA ... R7 1 (-.11) LANGBEIN 73 IPWA ... R7 5 (.066)TO .137 DEANS 75 DPWA ...

66 N*1/2(1700) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*1/2(1700) INTO GAM P, HELICITY=1/2 (GEV**1/2) ... A1 .024 .033 DEVENISH 73 DPWA ... A1 +.066 .042 MOORHOUSE 73 DPWA ...

***** REFERENCES FOR N*1/2(1700) *****

BRANDSEN 65 PL 19 420 +ODONNELL, MOORHOUSE (DURHAM, RHELIJIP) ... MICHAEL 66 PL 21 93 C MICHAEL (OXF) ... BAREYRE 68 PR 165 1731 P BAREYRE, C BRICMAN, G VILLET (SACLAY)IJP ...

N(1700)

18 N*1/2(1700, JP=3/2-1) I=1/2 D13 AYED 74 AND LONGACRE 75 BOTH FIND EVIDENCE FOR THIS STATE. THERE IS ADDITIONAL EVIDENCE FROM PHOTOPRODUCTION AND ASSOCIATED PRODUCTION.

18 N*1/2(1700) MASS (MEV)

M 3 (1730.) DONNACH2 68 RVUE PHAS.SHIPT-CERN1 10/69 ... M 3 (1690.) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69 ...

M 1 (1790.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73 ... M 1 NCT SEEN IN SOLUTION 2 OF LANGBEIN73 ...

18 N*1/2(1700) WIDTH (MEV)

W 1 (90.) DEANS 72 MPWA GAM P-K LM,SOL D 9/73 ... W 1 (120.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73 ...

18 N*1/2(1700) REAL PART OF POLE POSITION (MEV)

RE (1710.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

18 N*1/2(1700) 2*IMAG PART OF POLE POSITION (MEV)

IM (607.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

18 N*1/2(1700) PARTIAL DECAY MODES

P1 N*1/2(1700) INTO PI N 139+ 938 ... P2 N*1/2(1700) INTO LAMBDA K 1115+ 497 ... P3 N*1/2(1700) INTO GAM P,HELICITY=3/2 0+ 938 ...

18 N*1/2(1700) BRANCHING RATIOS

R1 N*1/2(1700) FROM GAMMA PROTON TO K LAMBDA DEANS 72 MPWA SQRTP((P3+P4)*P2) 9/73 ... R1 (-.0077) GAM P-K LM,SOL D 9/73 ...

18 N*1/2(1700) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*1/2(1700) INTO GAM P, HELICITY=1/2 (GEV**1/2) ... A1 -.103 .130 DEVENISH 73 DPWA ... A1 -.048 .050 DEVENIS2 74 DPWA ...

A3 N*1/2(1700) INTO GAM N, HELICITY=1/2 (GEV**1/2)

A3 .013 .222 DEVENISH 73 DPWA PI N PHOTO PROD 2/74 ... A3 -.021 .058 DEVENIS2 74 DPWA PI N PHOTO-PRCD 4/75*

A4 N*1/2(1700) INTO GAM N, HELICITY=3/2 (GEV**1/2)

A4 -.088 .087 DEVENISH 73 DPWA PI N PHOTO PROD 2/74 ... A4 -.026 .067 DEVENIS2 74 DPWA PI N PHOTO-PRCD 4/75*

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1700), N(1780)

Table of data cards for N(1700) and N(1780) baryons, including authors like AMALDI, BALLAM, BEKETOV, BOESEBEC, ELLIS, MA, MORSE, RUSHBROD, EDLSTEIN, JOHNSTAD, KARSHON, LAMSA, OH, RONAT, ABE, DAVIDSON, LICHTMAN, BLOBEL, BRAUNI, CAVALLI, MUSGRAVE, STRACHMA, WEBB, MERLO, ANTIPOV, LEDNICKY.

N(1780)

14 N*1/2(1780, JP=1/2+) I=1/2 P=1 THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

14 N*1/2(1780) MASS (MEV)

Table listing mass measurements for N(1780) with columns for author, mass value, and reference.

14 N*1/2(1780) WIDTH (MEV)

Table listing width measurements for N(1780) with columns for author, width value, and reference.

14 N*1/2(1780) REAL PART OF POLE POSITION (MEV)

Table listing real part of pole position for N(1780).

14 N*1/2(1780) 2*IMAG PART OF POLE POSITION (MEV)

Table listing imaginary part of pole position for N(1780).

14 N*1/2(1780) PARTIAL DECAY MODES

Table listing partial decay modes for N(1780) with columns for mode, branching ratio, and reference.

14 N*1/2(1780) BRANCHING RATIOS

Table listing branching ratios for N(1780) with columns for mode, ratio, and reference.

Table of data cards for N(1780) baryons, including authors like R2, R2 A, R2 B, R3, R3 B, R4, R4 B, R4 B, R4 B, R5, R5, R5, R6, R6 2, R7, R7, R7 1, R7 1, R8, R8 L, R9, R9 L, R10, R10 L, R11, R11 5, R11 5.

14 N*1/2(1780) PHOTON DECAY AMPL(1GEV**1/2)

Table listing photon decay amplitudes for N(1780) with columns for author, amplitude, and reference.

REFERENCES FOR N*1/2(1780)

Table listing references for N(1780) with columns for author, reference number, and journal.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1780), N(1810), N(1990)

PAPERS NOT REFERRED TO IN DATA CARDS.

DEANS 69 PR 177 2623 S R DEANS (UNIV S FLORIDA)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

N(1810)

15 N*1/2(1810, JP=3/2+) I=1/2

P13

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

15 N*1/2(1810) MASS (MEV)

Table with columns: M, X, W, A, M, M, M, M, M, M, M, L, L. Rows contain mass values and associated parameters like DGNACH1, LEA, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) WIDTH (MEV)

Table with columns: W, W, W, W, W, W, W, W, W, W, L. Rows contain width values and associated parameters like DGNACH1, LEA, WAGNER, ALMEHED, HICKS, KNASEL, LONGACRE.

15 N*1/2(1810) REAL PART OF POLE POSITION (MEV)

Table with columns: RE, W, W, W, W, W, L. Rows contain real part values and associated parameters like LONGACRE.

15 N*1/2(1810) 2*IMAG PART OF POLE POSITION (MEV)

Table with columns: IM, W, W, W, W, W, L. Rows contain imaginary part values and associated parameters like LONGACRE.

15 N*1/2(1810) PARTIAL DECAY MODES

Table with columns: P1, P2, P3, P4, P5, P6, P7, P8, P9, P10. Rows contain decay mode descriptions and associated parameters.

15 N*1/2(1810) BRANCHING RATIOS

Table with columns: R1, R2, R3, R4, R5, R6, R7, R8. Rows contain branching ratio values and associated parameters like DGNACH1, LEA, WAGNER, DEVENISH, KNASEL, DEANS, HICKS, DEANS, LONGACRE.

15 N*1/2(1810) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with columns: A1, A2, A3, A4. Rows contain photon decay amplitude values and associated parameters like DEVENISH, KNIES, METCALF, CRAWFORD, BARBOUR.

REFERENCES FOR N*1/2(1810)

Table with columns: DGNACH1, RUSH, BOTKE, DEANS, LEA, AYED, CARRERAS, DAVIES, WAGNER, ALMEHED, DEANS, DEVENISH, HICKS, AYED, DEVENISH, DEVENISH, METCALF, CRAWFORD, DEANS, KNASEL, LONGACRE, BARBOUR, DEANS, DONNACHI, AYED, APLIN. Rows contain references for various parameters.

N(1990)

17 N*1/2(1990, JP=7/2+) I=1/2

F17

THE MOST RECENT PI N PARTIAL WAVE ANALYSIS, AYED 74, FINDS EVIDENCE FOR THIS STATE. THERE IS ALSO SOME INDICATION IN ASSOCIATED PRODUCTION.

17 N*1/2(1990) MASS (MEV)

Table with columns: M, M, M, M, M, M, M, M, M, M, M, L. Rows contain mass values and associated parameters like DGNACH1, KIRSOPP, WAGNER, LEA, ALMEHED, HICKS, DEANS, LONGACRE.

Data Card Listings

For notation, see key at front of Listings.

Baryons

N(1990), N(2000), N(2040)

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Includes entries for N(1990) and N(2000).

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Includes entries for N(2040).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists various decay channels for N(1990).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists various decay channels for N(2000).

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(1990).

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(2000).

17 N(1990) PHOTON DECAY AMPL(GEV***-1/2) FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists photon decay amplitudes for N(1990).

REFERENCES FOR N(1990) ALMEHED 72 NP B40 157, DEANS 72 PRD 6 1906, LANGBEIN 73 NP B53 251, AYED 74 PRIVATE COMMCTN., ALSO 73 AIX CONFERENCE, DEANS 75 NP B96 90.

N(2040) THIS STATE IS NOW SEEN BY THE SACLAY GROUP, AYED 74.

REFERENCES FOR N(1990) DONNACHI 68 PL 268 161, KIRSOPP 68 THESIS, DEANS 69 PR 185 1797, LEA 69 PL 298 584, ALMEHED 72 NP B40 157, DEANS 72 PRD 6 1906, HICKS 73 PRD 7 2614, LANGBEIN 73 NP 853 251, AYED 74 PRIVATE COMMCTN., ALSO 73 AIX CONFERENCE, DEVENISH 74 NP 881 330, DEANS 75 NP B96 90, BARBOUR 76 SBMTD. TO NP, DEANS 69 PR 177 2623, AYED 70 PL 318 598, APLIN 71 NP B22 253.

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Lists mass and width for N(2040).

Table with 4 columns: Particle ID, Width (MeV), and other properties. Lists width for N(2040).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists partial decay modes for N(2040).

N(2000) THE MOST RECENT PI N PARTIAL WAVE ANALYSIS, AYED 74, FINDS EVIDENCE FOR THIS STATE. THERE IS ALSO SOME INDICATION IN ASSOCIATED PRODUCTION.

Table with 4 columns: Particle ID, Branching Ratios, and other properties. Lists branching ratios for N(2040).

Table with 4 columns: Particle ID, Mass (MeV), Width (MeV), and other properties. Lists mass and width for N(2000).

Table with 4 columns: Particle ID, Decay Mode, Decay Masses, and other properties. Lists partial decay modes for N(2040).

Baryons

N(2040), N(2100), N(2190)

Data Card Listings

For notation, see key at front of Listings.

16 N*1/2(2040) PHOTON DECAY AMPL(GEV**--1/2)
 FER DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-
 REVIEW PRECEDING THE BARYON LISTINGS.

A1	N*1/2(2040) INTO GAM P, HELICITY=1/2 (GEV**--1/2)	PI N PHOTO-PROD	4/75*
A1	.026 .052 DEVENISZ 74 DPWA		
A2	N*1/2(2040) INTO GAM P, HELICITY=3/2 (GEV**--1/2)	PI N PHOTO-PRCD	4/75*
A2	.128 .057 DEVENISZ 74 DPWA		
A3	N*1/2(2040) INTO GAM N, HELICITY=1/2 (GEV**--1/2)	PI N PHOTO-PRCD	4/75*
A3	.053 .083 DEVENISZ 74 DPWA		
A4	N*1/2(2040) INTO GAM N, HELICITY=3/2 (GEV**--1/2)	PI N PHCTC-PRCD	4/75*
A4	.100 .141 DEVENISZ 74 DPWA		

 REFERENCES FOR N*1/2(2040)
 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
 DONNACHIE RAPPORTEUR.S TALK (GLAS)
 R G KIRSOPP (EDIN)
 LEA, GADES, WARD, COWAN, + (RHEL, BRISTOL, DARE)
 B CARRERAS, A DONNACHIE (DARE, MCHS)
 +LOVELACE (LUND, RUTG)IJP
 DEANS, JACOBS, LYONS, MONTGOMERY (SOUTH FLA, IJJP
 +DEANS, JACOBS, LYONS+ (CARN+ORNL+SOUTH FLA, IJJP
 A YED, BAREYRE (SACL)IJP
 A YED, BAREYRE (SACL)IJP
 DEVENISH, LYTH, RANKIN (DESY, LANC, BCNN)IJP
 +MITCHELL, MONTGOMERY, + (SFLA, ALABAMA)IJP
 PAPERS NOT REFERRED TO IN DATA CARDS.
 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
 +BAREYRE, VILLET (SACLAY)
 +COWAN, GIBSON, GILMCRE++ (RHEL, BRISTOL)

N(2100) 04 N*1/2(2100, JP=1/2-) I=1/2 **S¹¹**
 NOW ALSO SEEN BY SACLAY, AYED 74.

04 N*1/2(2100) MASS (MEV)

M	(2070.)	ROYCHOUD 71 DPWA	3/72
M	(2100.)	ALMEHED 72 IPWA	2/72
M	(2283.)	AYED 74 IPWA	2/74

04 N*1/2(2100) WIDTH (MEV)

W	(200.)	ALMEHED 72 IPWA	2/72
W	(310.)	AYED 74 IPWA	2/74

04 N*1/2(2100) PARTIAL DECAY MODES

P1	N*1/2(2100) INTO PI N	DECAY MASSES	139+ 938
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04 N*1/2(2100) BRANCHING RATIOS

R1	N*1/2(2100) INTO (PI N)/TOTAL	(P1)	
R1	7 (0.5)	ALMEHED 72 IPWA	2/72
R1	(.14)	AYED 74 IPWA	2/74

 REFERENCES FOR N*1/2(2100)
 ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSEN (DURH)IJP
 ALMEHED 72 NP 840 157 +LOVELACE (LUND, RUTG)IJP
 AYED 74 PRIVATE COMMCTN. A YED, BAREYRE (SACL)IJP
 ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)IJP

N(2100) 05 N*1/2(2100, JP=5/2-) I=1/2 **D¹⁵**
 NOW ALSO SEEN BY SACLAY, AYED 74.

05 N*1/2(2100) MASS (MEV)

M	(2100.)	ALMEHED 72 IPWA	2/72
M	(2100.)	AYED 74 IPWA	2/74

05 N*1/2(2100) WIDTH (MEV)

W	(150.)	ALMEHED 72 IPWA	2/72
W	(220.)	AYED 74 IPWA	2/74

05 N*1/2(2100) PARTIAL DECAY MODES

P1	N*1/2(2100) INTO PI N	DECAY MASSES	139+ 938
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05 N*1/2(2100) BRANCHING RATIOS

R1	N*1/2(2100) INTO (PI N)/TOTAL	(P1)	
R1	7 (0.2)	ALMEHED 72 IPWA	2/72
R1	(.08)	AYED 74 IPWA	2/74

 REFERENCES FOR N*1/2(2100)
 ALMEHED 72 NP 840 157 +LOVELACE (LUND, RUTG)IJP
 AYED 74 PRIVATE COMMCTN. A YED, BAREYRE (SACL)IJP
 ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)IJP

N(2190) 71 N*1/2(2190, JP=7/2-) I=1/2 **G¹⁷**
 ROYCHOUDHURY 71 FIND SOME INDICATION OF P11 AND FIT IN
 THIS REGION. BRANSEN 71 ALSO FIND P11, F15, AND G19 RESO-
 NANT NEAR THIS MASS. AYED 74 FIND A G19.

71 N*1/2(2190) MASS (MEV)

M	(2190.0)	DIDDENS 63 CNTR	PI+ P TOTAL
M	(2210.0)	HOHLER 64 RVUE	DATA + DISP REL
M	(2190.0)	APPROX YOKOSAWA 66 CNTR	PI- P DSIG + POL 7/66
M	(2265.0)	DONNACHIE 68 RVUE	PHASE-SHIFT ANAL 6/68
M	(2000.0)	LEA 69 CNTR	PI-P ELASTIC 8/69
M	2180. 25. APPROX ANDERSON 70 MMS		PI- P TO PI- MMS 2/71
M	6 (2158.0)	AYED 70 IPWA	1/71
M	6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM		
M	(2260.0)	HULL 70 MPWA	SMALL ANGLE PI-P 1/71
M	(2160.0) (50.0)	AMALDI 71 CNTR	P P AT 24 GEV 10/71
M	(2160.)	BRANSDEN 71 DPWA	3/72
M	(2200.)	ROYCHOUD 71 DPWA	3/72
M	(2225.)	ALMEHED 72 IPWA	2/72
M	(2190.)	OTT 72 MPWA	0 PI-P BKWD ELSTC 2/73
M	1 (2208.)	HICKS 73 MPWA	GAM P-ETA P 9/73
M	1 ONLY STATES FROM TABLE VII OF HICKS73 ARE INCLUDED IN LISTINGS.		9/73
M	1 M AND W ARE FROM SOLUTION C2, BR=SQRT(G)/W WITH G FROM TABLE VII.		4/75*
M	(2141.) (20.)	ABE 74 + P+P->P+X, JCBN PK	2/74
M	(2133.)	AYED 74 IPWA	2/74
M	2 (2133.)	AYED 74 IPWA	2/74
M	2 THIS IS A G19 RESONANCE LISTED HERE UNTIL CONFIRMED.		2/74

71 N*1/2(2190) WIDTH (MEV)

W	(200.0)	DIDDENS 63 CNTR	7/66
W	(200.0)	HOHLER 64 RVUE	7/66
W	(2200.0)	APPROX YOKOSAWA 66 CNTR	6/68
W	3 (298.0)	DONNACHIE 68 RVUE	2/71
W	275. 70. APPROX ANDERSON 70 MMS		PI- P TO PI- MMS 1/71
W	6 (325.0)	AYED 70 IPWA	1/71
W	(239.0)	HULL 70 MPWA	SMALL ANGLE PI-P 1/71
W	(150.)	ALMEHED 72 IPWA	2/72
W	1 (193.)	HICKS 73 MPWA	GAM P-ETA P 9/73
W	(243.)	AYED 74 IPWA	2/74
W	(193.)	AYED 74 IPWA	2/74

SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

71 N*1/2(2190) PARTIAL DECAY MODES

P1	N*1/2(2190) INTO PI N	DECAY MASSES	139+ 938
P2	N*1/2(2190) INTO LAMBDA K		1115+1765
P3	N*1/2(2190) INTO N PI PI		938+ 139+ 139
P4	N*1/2(2190) INTO GAM P, HELICITY=3/2		0+ 938
P5	N*1/2(2190) INTO GAM P, HELICITY=1/2		0+ 938
P6	N*1/2(2190) INTO GAM N, HELICITY=3/2		0+ 939
P7	N*1/2(2190) INTO GAM N, HELICITY=1/2		0+ 939
P8	N*1/2(2190) INTO ETA N		548+ 938
P9	N*1/2(2190) INTO SIGMA K		493+1189

71 N*1/2(2190) BRANCHING RATIOS

R1	N*1/2(2190) INTO (PI N)/TOTAL	(P1)	
R1	(0.3)	APPROX DIDDENS 63 CNTR	7/66
R1	(0.3)	APPROX YOKOSAWA 66 CNTR	7/66
R1	3 (0.349)	DONNACHIE 68 RVUE	6/68
R1	6 (0.150)	AYED 70 IPWA	1/71
R1	(0.09)	HULL 70 MPWA	SMALL ANGLE PI-P 1/71
R1	7 (0.35)	ALMEHED 72 IPWA	2/72
R1	(.25)	OTT 72 MPWA	0 PI-P BKWD ELSTC 2/73
R1	(.161)	AYED 74 IPWA	2/74
R1	(.095)	AYED 74 IPWA	2/74

R2 N*1/2(2190) FROM GAMMA PROTON TO K LAMBDA SORT((P4+P5)*P2) 9/73
 R2 (.0161) DEANS 72 MPWA GAM P-K LM, SOL D 9/73

R3 N*1/2(2190) FROM GAMMA PROTON TO ETA PROTON SQRT((P4+P5)*P8) 9/73
 R3 (.0094) HICKS 73 MPWA GAM P-ETA P 9/73

R4 N*1/2(2190) FROM PI N TO K SIGMA SQRT(P1*P9) 11/75*
 R4 4 (.014) TO .019 DEANS 75 DPWA PI N TO K SIGMA 11/75*
 R4 4 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS. 11/75*

 REFERENCES FOR N*1/2(2190)
 DIDDENS 63 PRL 10 262 +JENKINS, KYCIA, RILEY (BNL) I
 HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
 YOKOSAWA 66 PRL 16 714 +SUGA, FILL, ESTERLING, BOOTH (ANL, CHIC) JP
 DONNACHIE 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
 ALSO 68 THESIS DONNACHIE RAPPORTEUR.S TALK (GLAS)
 R G KIRSOPP (EDIN)
 LEA 69 PL 298 584 LEA, GADES, WARD, COWAN, + (RHEL, BRISTOL, DARE)
 ANDERSON 70 PRL 25, 699 +BLESER, BLIEDEN, COLLINS++ (BNL, CARN)
 AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACL)IJP
 HULL 70 PR D2 1783 J HULL, R LEACOCK (ISU)

Data Card Listings

Baryons

For notation, see key at front of Listings. N(2190), N(2220), N(2650), N(3030), N(3245)

AMALDI 71 PL 348 435 +BIANCASTELLI, BOSIO, + (I SANITA ROMA+GERN)
BRANDSEN 71 NP 826 511 ,OGDEN (DURHI)JP
ALSC 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANDSEN (DURH)JP
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANDSEN (DURH)JP

ALME-ED 72 NP 840 157 +LOVELACE (LUND, RUTG)JP
DEANS 72 PRD 6 1906 DEANS, JACOBS, LYONS, MCNTGOMERY (SOUTH FLA.)JP
DTT 72 PL 428 133 +TRISCHUK, VAVRA, RICHARDS, + (MCGI, STLO, IOWA)JP
ALSO 72 MCGILL THESIS J. VAVRA (MCGI) JP
HICKS 73 PRD 7 2614 +DEANS, JACOBS, LYONS+ (CARA+ORNL+SOUTH FLA.)JP

ABE 74 PL 538 114 +ALSPECTOR, BCMBEROWITZ+ (RLTG, UPNJ, FSU)
AYED 74 PRIVATE COMMCTN. AYED, BAREYRE (SACL)JP
ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)JP
DEANS 75 NP 896 90 +MITCHELL, MCNTGOMERY, + (SFLA, ALABAMA)JP

AYED 70 PL 318 598 +BAREYRE, VILLET (SACLAY)
QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN DATA CARDS.

BARGER 66 PRL 16 913 V BARGER, D CLINE (NISC) P
CARROLL 66 PRL 16 288 +CORBETT, DAMERELL, MIDDLEMAS, + (RHEL, OXF)J-L
CARROLL 66 PRL 17 1274 +CORBETT, DAMERELL, MIDDLEMAS, + (RHEL, OXF)J-L
ERRATUM CHANGING THE RATHER WEAK DETERMINATION OF J-L TO +1 (2.)
KORMANYO 66 PRL 16 709 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P
BUSZA 67 NC 52A 331 +DAVIS, DUFF, HEYMANN, + (LOUC, WESTFIELD)

N(2220)

H19

90 N*1/2(2220, JP=9/2+) I=1/2
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

90 N*1/2(2220) MASS (MEV)
M (2200.) APPROX. BUSZA 67 DSPK LEG. POLYN. ANAL. 2/71
M 6 (2221.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM HULL 70 MPWA SMALL ANGLE PI-P 1/71
M (2245.0) HULL 70 MPWA 2/74
M (2249.) AYED 74 IPWA

90 N*1/2(2220) WIDTH (MEV)
W 6 (258.0) AYED 70 IPWA 1/71
W (329.0) HULL 70 MPWA 1/71
W (347.) AYED 74 IPWA 2/74

90 N*1/2(2220) PARTIAL DECAY MODES
P1 N*1/2(2220) INTO PI N DECAPY MASSES 139+ 938
P2 N*1/2(2220) INTO N PI 939+ 548

90 N*1/2(2220) BRANCHING RATIOS
R1 N*1/2(2220) INTO (PI N)/TOTAL (P1) 1/71
R1 6 (0.140) AYED 70 IPWA 1/71
R1 (0.15) HULL 70 MPWA SMALL ANGLE PI-P 2/74
R1 (.204) AYED 74 IPWA

REFERENCES FOR N*1/2(2220)
BUSZA 67 NC 52A 331 +DAVIS, DUFF, HEYMANN, NIMMON + (LOUC+LOWC)
AYED 70 KIEV CONF R AYED+P BAREYRE, G VILLET (SACL)JP
HULL 70 PR D2 1783 J HULL, R LEACOCK (ISU)
AYED 74 PRIVATE COMMCTN. AYEC, BAREYRE (SACL)JP
ALSO 73 AIX CONFERENCE AYED, BAREYRE (SACL)JP

N(2650) BUMPS

72 N*1/2(2650, JP=) I=1/2 PRODUCTION EXPERIMENTS
ROYCHOUDHURY 71 CLAIM F15(2400) AND G19(2400) TO BE POSSIBLE RESONANCES. BRANDSEN 71 FIND THE POSSIBLE RESONANT CANDIDATES S11(2520) AND H19(2590).

72 N*1/2(2650) MASS (MEV) (PROD. EXP.)
M (2700.0) ALVAREZ 64 CNTR PI PHCTOPROD
M (2660.0) HOHLER 64 RVUE DATA + DISP REL
M (2600.0) WAHLIG 64 DSPK 0 PI-P CH EX
M (2633.0) BARGER 66 FIT TOTAL + CH EX 11/67
M 2649.0 10.0 CITRGN 66 CNTR PI+- P TOTAL 7/66

72 N*1/2(2650) WIDTH (MEV) (PROD. EXP.)
W (100.0) ALVAREZ 64 CNTR 7/66
W (200.0) HOHLER 64 RVUE 11/67
W (425.0) BARGER 66 FIT TOTAL + CH EX 7/66
W 360.0 20.0 CITRGN 66 CNTR

72 N*1/2(2650) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*1/2(2650) INTO PI N DECAPY MASSES 139+ 938
P2 N*1/2(2650) INTO LAMBDA K 1115+ 497
P3 N*1/2(2650) INTO N PI PI 938+ 139+ 139

72 N*1/2(2650) BRANCHING RATIOS (PROD. EXP.)
R1 N*1/2(2650) INTO (PI N)/TOTAL (P1)
R1 ONLY (J+1/2)*(PI N/TOTAL) MEASURED FOR THIS STATE
R1 B (0.456) (0.018) BARGER 66 RVUE TOTAL + CH EXC. 11/67
R1 0.436 0.028 CITRGN 66 CNTR TOTAL CROSS-SEC. 11/67
R1 B (0.30) BARGER 67 RVUE USES KORMANYOS67 11/67
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 D (0.24) DIKMN 67 RVUE USES KORMANYOS66 11/67
R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (0.06) KORMANYOS 67 CNTR PI-P AT 180 DEG. 11/67

REFERENCES FOR N*1/2(2650) (PROD. EXP.)
ALVAREZ 64 PL 12 710 +BAR-YAM, KERN, LUCKEY, OSBORNE, + (MIT, CEA)
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
WAHLIG 64 PRL 13 103 +MANNELLI, SODICKSON, FACKLER, WARD, + (MIT)
BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)
CITRGN 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMN 67 PRL 18 798 F N DIKMN (MICH)
KORMANYO 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P

PAPERS NOT REFERRED TO IN DATA CARDS.
BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE, ORSAY)J-L
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)
WAHLIG 68 PR 168 1515 M A WAHLIG, I MANNELLI (MIT, PISA)
FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH CITRGN 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES CCMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

BRANDSEN 71 NP 826 511 ,OGDEN (DURHI)JP
ALSO 70 NP 816 461 ROYCHOUDHURY, PERRIN, BRANDSEN (DURH)JP
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANDSEN (DURH)JP

N(3030) BUMPS

73 N*1/2(3030, JP=) I=1/2 PRODUCTION EXPERIMENTS

73 N*1/2(3030) MASS (MEV) (PROD. EXP.)
M (3080.0) HOHLER 64 RVUE DATA + DISP REL 7/66
M (3030.0) CITRGN 66 CNTR PI+- P TOTAL 7/66

73 N*1/2(3030) WIDTH (MEV) (PROD. EXP.)
W (400.0) CITRGN 66 CNTR 7/66

73 N*1/2(3030) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*1/2(3030) INTO PI N DECAPY MASSES 139+ 938
P2 N*1/2(3030) INTO N PI 938+ 139+ 139

73 N*1/2(3030) BRANCHING RATIOS (PROD. EXP.)
R1 N*1/2(3030) INTO (PI N)/TOTAL (P1)
R1 ONLY (J+1/2)*(PI N/TOTAL) MEASURED FOR THIS STATE
R1 B (0.088) (0.016) BARGER 66 RVUE TOTAL + CH EXC. 11/67
R1 (0.048) CITRGN 66 CNTR TOTAL CROSS-SEC. 11/67
R1 B (0.12) BARGER 67 CNTR USES KORMANYOS66 11/67
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 D (0.016) DIKMN 67 RVUE USES KORMANYOS67 11/67
R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES

REFERENCES FOR N*1/2(3030) (PROD. EXP.)
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
BARGER 66 PR 151 1123 V BARGER, M OLSSON (WISC)
CITRGN 66 PR 144 1101 +GALBRAITH, KYCIA, LEONTIC, PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMN 67 PRL 18 798 F N DIKMN (MICH)

PAPERS NOT REFERRED TO IN DATA CARDS
KORMANYO 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH, ANL) P
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)

N(3245) BUMPS

74 N*1/2(3245, JP=) PRODUCTION EXPERIMENTS
EXISTENCE NOT CONCLUSIVELY ESTABLISHED. I-SPIN NOT DETERMINED, BUT THE NARROW WIDTH PRECLUDES IDENTIFICATION WITH THE N*3/2(3230). OMITTED FROM TABLE.

74 N*1/2(3245) MASS (MEV) (PROD. EXP.)
M 3245.0 10.0 KORMANYOS 67 CNTR PI-P 180 DEG EL 6/68

74 N*1/2(3245) WIDTH (MEV) (PROD. EXP.)
W (35.0) OR LESS KORMANYOS 67 CNTR 6/68

Baryons

N₇(3245), N(3690), N₇(3755), Δ(1232)

Data Card Listings

For notation, see key at front of Listings.

74 N* /2(3245) PARTIAL DECAY MODES (PROD. EXP.)
P1 N* /2(3245) INTO PI N
DECAY MASSES
139+ 938

74 N* /2(3245) BRANCHING RATIOS (PROD. EXP.)
R1 N* /2(3245) INTO (PI N)/TOTAL (P1)
R1 J IS NOT KNOWN. FOLLOWING IS (J+1/2)*(PI N)/TOTAL
R1 (0.37) KORMANYOS 67 CNTR 6/68

REFERENCES FOR N* /2(3245) (PROD. EXP.)
KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P

N(3690) BUMPS
75 N*1/2(3690, JP=) I=1/2 PRODUCTION EXPERIMENTS
A BUMP SEEN IN THE INVARIANT MASS OF A VERY COMPLICATED STATE IN + SEVEN PIS, SO AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. NOT INCLUDED IN TABLE.

75 N*1/2(3690) MASS (MEV) (PROD. EXP.)
M 3690.0 10.0 BARTKE 67 HBC + PI+P 8 PRONGS 8/67

75 N*1/2(3690) WIDTH (MEV) (PROD. EXP.)
W 50.0 30.0 BARTKE 67 HBC + 8/67

75 N*1/2(3690) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*1/2(3690) INTO N + 7 PIS
DECAY MASSES

REFERENCES FOR N*1/2(3690) (PROD. EXP.)
BARTKE 67 PL 248 118 +CZYZEWSKI, GANYSZ, + (CRACOW, ORSAY) I

N7(3755) BUMPS
76 N* /2(3755, JP=) PRODUCTION EXPERIMENTS
A SMALL PEAK IN THE (P P PBAR) INVARIANT MASS FROM 8.4 BEV/C PI+ P TO PI+ P P PBAR EVENTS. AS EVIDENCE FOR A NEW RESONANCE IT IS NOT CONCLUSIVE. OMITTED FROM TABLE.

76 N* /2(3755) MASS (MEV) (PROD. EXP.)
M 3755.0 8.0 EHRlich 68 HBC + PI+ P P PBAR 6/68

76 N* /2(3755) WIDTH (MEV) (PROD. EXP.)
W 40.0 20.0 EHRlich 68 HBC + 6/68

76 N* /2(3755) PARTIAL DECAY MODES (PROD. EXP.)
P1 N* /2(3755) INTO PI+ P P PBAR
DECAY MASSES
139+ 938+ 938+ 938

REFERENCES FOR N* /2(3755) (PROD. EXP.)
EHRlich 68 PRL 20 686 R EHRlich, R J PLANC, J B WHITTAKER (RUTGERS)

S=0 I=3/2 NUCLEON STATES (Δ)

Δ(1232)
33 N*3/2(1232, JP=3/2+) I=3/2 P33

THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED. SEE CARTER 71 AND CARTER 73 FOR PI N CROSS-SECTION DATA IN THIS REGION.

33 N*3/2(1232) MASS (MEV)

M (1234.) ROPER 65 DPWA ++0 PHASE SHIFT AN. 2/72
M (1235.) ALMEHD 72 IPWA 2/74
M 3 (1243.3) (1241.7) CHENG 73 FIT CARTER 71 2/74

M++ 1236.0 0.55 OLSSON 65 RVUE ++ TOTAL-SIGMA DATA 1/74
M++ 2 1231.0 1.5 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
M++ 1 1231.1 +2 CARTER 73 IPWA ++ PI N 88-310 MEV 9/73
M++ 1 EXPERIMENTAL QUANTITY-SEE CARTER73 FOR COULOMB BARRIER CORRECTIONS 9/73
M++ 2 EXPERIMENTAL QUANTITY-SEE CARTER71 FOR COULOMB BARRIER CORRECTIONS 1/74
M++ AVERAGE MEANINGLESS (SCALE FACTOR = 8.4)

M+ (1231.8) BERENDS 75 IPWA + GA P TO PI NUC 4/75*
M+ 1230.6 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
M+ (1231.7) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*
MO 1236.45 0.65 OLSSON 65 RVUE 0
MO 1232.9 0.6 CARTER 71 MPWA 0 PI-P SIG. TOTAL 1/71
MO AVERAGE MEANINGLESS (SCALE FACTOR = 4.0)

33 N*3/2(1232) WIDTH (MEV)
W (120.) ROPER 65 DPWA ++0 PHASE SHIFT AN. 2/72
W (129.) ALMEHD 72 IPWA 2/74
W 3 (152.2) (145.8) CHENG 73 FIT CARTER 71 2/74
W (120.) TSCHANG 73 FIT CARTER71 P33 1/74
W (120.) AYED 74 IPWA 2/74

M++ 120.0 2.0 OLSSON 65 RVUE ++
M++ 2 111.1 1.8 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
M++ 1 111.5 .4 CARTER 73 IPWA ++ PI N 88-310 MEV 9/73
M++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

M+ 120.2 3.9 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
M+ (117.4) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*
MO 119.6 2.4 OLSSON 65 RVUE 0
MO 114.7 3.0 CARTER 71 MPWA 0 PI-P SIG TOT. 1/71
MO AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)

33 (N*0) - (N**+) MASS DIFFERENCE (MEV)
D R (0.45) (0.85) OLSSON 65 RVUE 1/74
D 2 1.3 1.9 CARTER 71 MPWA ++ PI+P SIG. TOTAL 2/73
D 1 1.4 .4 CARTER 73 IPWA PI N 88-310 MEV 9/73
D R REDUNDANT WITH DATA IN MASS LISTING.
D AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 (N*0) - (N**+) WIDTH DIFFERENCE (MEV) 9/73
WD 2 6.5 2.3 CARTER 71 MPWA ++ PI+P SIG. TOTAL 1/74
WD 1 10.3 1.2 CARTER 73 IPWA PI N 88-310 MEV 9/73
WD AVERAGE MEANINGLESS (SCALE FACTOR = 1.5)

33 N*3/2(1232) REAL PART OF PCLE POSITION(MEV)
REE M (1214.) MICHAEL 67 2/74
REE P (1211.) BALL 72 2/73
REE 3 (1210.7) (1210.7) CHENG 73 FIT DELTA 33 2/73
REE 3 (1210.7) (1210.7) CHENG 73 FIT CARTER 71 2/74
REE 1214.5 10. NOGOVA 73 FIT ALMEHD72 2/74
REE (1213.) SPEARMAN 74 FIT ZERO TRJCTRY 4/75*
REE M FIT INCLUDES OLSSON 65 PARAMETERS PLUS SCATTERING LENGTH PLUS 6
REE M PHASE SHIFT VALUES FOR TPI=120 TO 492 MEV
REE P ERROR EST. FROM FITS WITH SOMEWHAT VARYING ASSUMPTIONS
REE AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

R++ U 1211.5 .6 BALL 75 ++ FIT CARTER 73 11/75*
R++ U 1210.9 .8 LICHTENB 75 ++ FIT CARTER 73 11/75*
R++ C 1209.6 .5 VASAN 75 ++ FIT CARTER 73 1/76*
R++ C FROM FITS TO COULOMB CORRECTED CARTER 73 PHASE SHIFT. 1/76*
R++ U (1210.5) TO (1210.8) VASAN 75 ++ FIT CARTER 73 1/76*
R++ U FROM FITS TO UNCORRECTED CARTER 73 PHASE SHIFT. 1/76*
R++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)

REO U (1211.6) BALL 75 0 FIT CARTER 73 11/75*
REO U 1210.9 1.4 LICHTENB 75 0 FIT CARTER 73 11/75*
REO C 1210.75 .6 VASAN 75 ++ FIT CARTER 73 1/76*
REO U (1210.2) VASAN 75 ++ FIT CARTER 73 1/76*
REO AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 N*3/2(1232) IMAG PART OF PCLE POSITION(MEV)
IME M (52.) MICHAEL 67 2/74
IME (50.) BALL 72 2/73
IME P 49.5 1.8 PDG 72 FIT DELTA 33 2/73
IME 3 (50.7) (50.6) CHENG 73 FIT CARTER 71 2/74
IME (48.6) 5. NOGOVA 73 FIT ALMEHD72 2/74
IME (49.) SPEARMAN 74 FIT ZERO TRJCTRY 4/75*
IME AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

I++ U 50.1 .6 BALL 75 ++ FIT CARTER 73 11/75*
I++ U 49.6 .75 LICHTENB 75 ++ FIT CARTER 73 11/75*
I++ C 50.4 .5 VASAN 75 ++ FIT CARTER 73 1/76*
I++ U (49.9) TO (50.0) VASAN 75 ++ FIT CARTER 73 1/76*
I++ AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

IMO U (53.0) BALL 75 0 FIT CARTER 73 11/75*
IMO U 53.25 1.75 LICHTENB 75 0 FIT CARTER 73 11/75*
IMO C 52.8 .75 VASAN 75 ++ FIT CARTER 73 1/76*
IMO U (52.9) TO (53.1) VASAN 75 ++ FIT CARTER 73 1/76*
IMO AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

33 N*3/2(1232) ABSOLUTE VALUE OF PCLE RESIDUE (MEV)

ABS (53.) BALL 73 FIT DELTA 33 9/73
A++ C (52.4) TO (53.2) VASAN 75 ++ FIT CARTER 73 1/76*
A++ U (52.1) TO (52.4) VASAN 75 ++ FIT CARTER 73 1/76*

ABO C (54.8) TO (55.0) VASAN 75 ++ FIT CARTER 73 1/76*
ABO U (55.2) TO (55.3) VASAN 75 ++ FIT CARTER 73 1/76*

Baryons

$\Delta(1650)$, $\Delta(1670)$

Data Card Listings
For notation, see key at front of Listings.

$\Delta(1650)$

82 N*3/2(1650, JP=1/2-) I=3/2
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

S₃₁

Table with columns for mass (MEV), width (MEV), and various parameters for the Delta(1650) resonance. Includes entries for DEVLIN, BAREYRE, DONNACHI, and others.

Table showing the real part of the pole position (MEV) and 2*imag part of the pole position (MEV) for the Delta(1650) resonance.

Table showing the partial decay modes for the Delta(1650) resonance, including decay masses and branching ratios.

Table showing the branching ratios for the Delta(1650) resonance into various channels.

Table showing the partial decay modes for the Delta(1650) resonance, including decay masses and branching ratios.

Table showing the branching ratios for the Delta(1650) resonance into various channels.

Table showing the partial decay modes for the Delta(1650) resonance, including decay masses and branching ratios.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table showing the branching ratios for the Delta(1650) resonance into various channels.

REFERENCES FOR N*3/2(1650)

Table listing references for the Delta(1650) resonance, including authors like DEVLIN, BAREYRE, DONNACHI, etc.

CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
LONGACRE 75 PL 558 415 +ROSENFELD,LASINSKI,SMADJA+ (LEL,SLAC)IJP

BARBOUR 76 BMTD. TO NP I. M. BARBUR,R. L. CRAWFCRD (GLAS)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.

CAPRUTHE 60 PRL 4 303 P CARRUTHERS (CORNELL) I
DEVLIN 62 PR 125 690 T J DEVLIN, B J MOYER, V PEREZ-MENDEZ (LRL) I
HELLAND 64 PR 134 81062 +DEVLIN,HAGGE,LCNGCI,MOYER,WOOD (LRL) I

Table with columns for mass (MEV), width (MEV), and various parameters for the Delta(1670) resonance.

Table showing the real part of the pole position (MEV) and 2*imag part of the pole position (MEV) for the Delta(1670) resonance.

Table showing the partial decay modes for the Delta(1670) resonance, including decay masses and branching ratios.

Table showing the branching ratios for the Delta(1670) resonance into various channels.

Table showing the partial decay modes for the Delta(1670) resonance, including decay masses and branching ratios.

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table showing the branching ratios for the Delta(1670) resonance into various channels.

REFERENCES FOR N*3/2(1670)

Table listing references for the Delta(1670) resonance, including authors like CRAWFORD, BARBOUR, CAPRUTHE, etc.

Data Card Listings
For notation, see key at front of Listings.

Baryons
Δ(1670), Δ(1690), Δ(1890)

A1 +.101 .011 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
A1 (+.120) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.4)
A2 N*3/2(1670) INTO GAM NUCLEON, FELICITY=3/2 (GEV**=1/2)
A2 .110 .039 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A2 +.022 .052 MOORHOUS 73 DPWA PI N PHOTO-PROD 2/73
A2 .072 .014 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*

Δ(1690) 19 N*3/2(1690, JP=3/2+) I=3/2 P33
RECENT ANALYSES INDICATE A POSSIBLE P33 RESONANCE SOMEWHERE IN THE 1600-2000 MEV REGION. SEE ALSO THE N*3/2(1690) LISTINGS.

19 N*3/2(1690) MASS (MEV)
M 3 (1690.) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
M 3 (1690.) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
M 3 WHERE MAX. ABSORPTION IS -DONNACH1, 2, KIRSOPP EYEBALL FIT CERN 1 10/69

19 N*3/2(1690) WIDTH (MEV)
W 3 (281.) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
W 3 (240.) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
W 6 (598.0) AYED 70 IPWA 1/71
W 7 (220.) ALMEHED 72 IPWA 2/72
W 7 (204.) AYED 74 IPWA 2/74
W L 1900. OR 1640. LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) REAL PART OF POLE POSITION (MEV)
RE (1609.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) 2*IMAG PART OF POLE POSITION (MEV)
IM (323.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

19 N*3/2(1690) PARTIAL DECAY MODES
P1 N*3/2(1690) INTO PI N 139+ 938
P2 N*3/2(1690) INTO K SIGMA 493+1189
P3 N*3/2(1690) INTO N*3/2(1232) PI,P-WAVE 1232+ 139

19 N*3/2(1690) BRANCHING RATIOS
R1 N*3/2(1690) INTO (PI N)/TOTAL (P1)
R1 3 (.10) DONNACH2 68 RVUE PHAS-SHIFT-CERN1 10/69
R1 3 (.08) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
R1 6 (0.135) AYED 70 IPWA 1/71
R1 7 (0.1) ALMEHED 72 IPWA 2/72
R1 1 (.194) AYED 74 IPWA 2/74

R3 N*3/2(1690) FROM PI N TO K SIGMA SQRT(PI*P2) 11/75*
R3 2 (.006) TO .042 DEANS 75 DPWA PI N TO K SIGMA 11/75*
R3 2 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS. 11/75*

19 N*3/2(1690) PHOTON DECAY AMPL(GEV**=1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*3/2(1690) INTO GAM NUCLEON, HELICITY=1/2 (GEV**=1/2)
A1 .016 .055 DEVENISH 73 DPWA PI N PHOTO PROD 2/74
A1 -.033 .037 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*
A1 .003 .015 KNIES 74 DPWA PI N PHOTO PRCD 2/74
A1 .0 .038 METCALF 74 DPWA PI N PHOTO-PROD 2/74
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)
A2 N*3/2(1690) INTO GAM NUCLEON, HELICITY=3/2 (GEV**=1/2)
A2 .074 .064 DEVENISH 73 DPWA PI N PHOTO PRCD 2/74
A2 .308 .046 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*

REFERENCES FOR N*3/2(1690)
DONNACH2 68 VIENNA 139 DONNACHIE RAPPORTEUR'S TALK (GLAS)
KIRSOPP 68 THESIS R G KIRSOPP (EDIN)
AYED 70 KIEV CNF R AYED,P BAREYRE, G VILLET (SACL)IJP
FEUERBACH 70 NP 168 85 FEUERBACHER+HOLLADAY (VANDERBILT)
ALMEHED 72 NP 840 157 +LOVELACE (LUND,RUTG)IJP
DEVENISH 73 PL 478 53 MOORHOUS 73 PL 438 44 MOORHOUSE, OBERLACK (LOUC+BNN+LANC)IJP (GLAS+LBL)IJP

Δ(1890) 11 N*3/2(1890, JP=5/2+) I=3/2 F35
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

11 N*3/2(1890) MASS (MEV)
M 3 (1913.0) DONNACH1 68 RVUE PHASE-SHIFT ANAL 8/69
M 6 (1837.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM DAVIDES 70 RVUE P-S ANAL SCL A 8/69
M 4 (1841.0) ALMEHED 72 IPWA 2/72
M 7 (1875.) MEHTANI 72 DPWA PI+P TO PI+PIOP 9/73
M 1 1890. TO 1986. LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
M (1890.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M (1869.) AYED 74 IPWA 2/74
M L 1870. OR 1830. LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) WIDTH (MEV)
W 3 (350.0) DONNACH1 68 RVUE 8/69
W 6 (198.0) AYED 70 IPWA 1/71
W 4 (136.0) DAVIDES 70 RVUE SOL A 8/69
W 7 (250.) ALMEHED 72 IPWA 2/72
W 1 273. TO 322. MEHTANI 72 DPWA PI+P TO PI+PIOP 9/73
W (180.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
W (140.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (255.) AYED 74 IPWA 2/74
W L 255. OR 220. LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) REAL PART OF POLE POSITION (MEV)
RE (1813.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) 2*IMAG PART OF POLE POSITION (MEV)
IM (193.) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

11 N*3/2(1890) PARTIAL DECAY MODES
P1 N*3/2(1890) INTO PI N 139+ 938
P2 N*3/2(1890) INTO K SIGMA 938+ 139+ 139
P3 N*3/2(1890) INTO K SIGMA 493+1189
P4 N*3/2(1890) INTO N*3/2(1232) PI 1232+ 139
P5 N*3/2(1890) INTO GAM NUCLEON, HELICITY=1/2 0+ 938
P6 N*3/2(1890) INTO GAM NUCLEON, HELICITY=3/2 0+ 938
P7 N*3/2(1890) INTO N RHO 938+ 773
P8 N*3/2(1890) INTO N*3/2(1232) PI,P-WAVE 1232+ 139
P9 N*3/2(1890) INTO N*3/2(1232) PI,P-WAVE 1232+ 139
P10 N*3/2(1890) INTO N RHO,S=3/2,P-WAVE 938+ 773

Baryons

$\Delta(1890)$, $\Delta(1900)$, $\Delta(1910)$

11 N*3/2(1890) BRANCHING RATIOS
R1 N*3/2(1890) INTO (PI N)/TOTAL
R2 N*3/2(1890) INTO (K SIGMA)/TOTAL
R3 N*3/2(1890) INTO (K SIGMA)*(PI N)/TOTAL**2
R4 N*3/2(1890) FROM PI N TO K SIGMA
R5 N*3/2(1890) FROM PI N TO K SIGMA
R6 N*3/2(1890) FROM PI N TO N RHO,S=3/2,P-WAVE
R7 N*3/2(1890) FROM PI N TO N RHO,S=3/2,P-WAVE

11 N*3/2(1890) PHOTON DECAY AMPL(GEV**-1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.
A1 N*3/2(1890) INTO GAM NUCLEON, HELICITY=1/2 (GEV**-1/2)
A2 N*3/2(1890) INTO GAM NUCLEON, HELICITY=3/2 (GEV**-1/2)

REFERENCES FOR N*3/2(1890)
DCNNACHI 68 PL 268 161
ALSO 68 VIENNA 139
ALSO 68 THESIS
AYED 70 KIEV CCFN
DAVIES 70 NP 821 359
FEUERBACH 70 NP 168 85
KALMUS 70 PR D2 1824
ALMEHED 72 NP 840 157
MEHTANI 72 PRL 29 1634
LANGBEIN 73 NP 853 251
AYED 74 PRIVATE COMMCTN.
ALSO 73 AIX CONFERENCE
DEVENIS2 74 PL 528 227
KNIES 74 PRD 9 2680
METCALF 74 NP 876 253
CRAWFORD 75 NP 897 125
DEANS 75 NP 896 90
LCNGACRE 75 PL 558 415
BARBCUR 76 SBMTD. TO NP
AYED 70 PL 318 598

$\Delta(1900)$ 30 N*3/2(1900, JP=1/2-) I=3/2 S31
THIS EFFECT IS SEEN IN TWO CHANNELS.

30 N*3/2(1900) MASS (MEV)
M (1920.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
M (1870.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M (2001.) AYED 74 IPWA 2/74
30 N*3/2(1900) WIDTH (MEV)
W (140.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
W (160.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (307.) AYED 74 IPWA 2/74
30 N*3/2(1900) PARTIAL DECAY MODES
P1 N*3/2(1900) INTO PI N 139+ 938
P2 N*3/2(1900) INTO K SIGMA 493+1189

Data Card Listings

For notation, see key at front of Listings.

30 N*3/2(1900) BRANCHING RATIOS 9/73
R1 N*3/2(1900) FROM PI N TO K SIGMA SQRT(PI*P2) 9/73
R1 (.11) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
R1 (.12) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
R1 1 (.076) DEANS 75 DPWA PI N TO K SIGMA 11/75*
R2 1 VALUE GIVEN IS FROM SOLUTION 1, NOT PRESENT IN SOLUTIONS 2,3,4. 11/75*

REFERENCES FOR N*3/2(1900)
LANGBEIN 73 NP 853 251 LANGBEIN,WAGNER (MUNICH)IJP
AYED 74 PRIVATE COMMCTN. AYED,BAREYRE (SACL)IJP
ALSC 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP
DEANS 75 NP 896 90 +MITCHELL,MCATGMERY,+ (SFLA,ALABAMA)IJP

$\Delta(1910)$ 12 N*3/2(1910, JP=1/2+) I=3/2 P31
THE EXISTENCE OF THIS RESONANCE IS WELL ESTABLISHED.

12 N*3/2(1910) MASS (MEV)
M 3 (1934.0) DCNNACHI 68 RVUE PHASE-SHIFT ANAL 8/69
M 6 (1783.0) AYED 70 IPWA 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM (GLAS) 8/69
M 4 (194.0) ALMEHED 72 IPWA P-S ANAL SCL A 2/72
M 7 (1900.) ALMEHED 72 IPWA PI N-K SIG,SOL 1 9/73
M (1980.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M (1950.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M (1786.) AYED 74 IPWA 2/74

12 N*3/2(1910) WIDTH (MEV)
W 3 (339.0) DCNNACHI 68 RVUE 8/69
W 6 (308.0) AYED 70 IPWA 1/71
W 4 (290.) DAVIES 70 RVUE SCL A 8/69
W 7 (200.) ALMEHED 72 IPWA 2/72
W (190.) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
W (170.) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (222.) AYED 74 IPWA 2/74
SEE NOTES ACCOMPANYING MASSES QUOTED

12 N*3/2(1910) PARTIAL DECAY MODES
P1 N*3/2(1910) INTO PI N 139+ 938
P2 N*3/2(1910) INTO N PI PI 938+ 139+ 139
P3 N*3/2(1910) INTO K SIGMA 493+1189
P4 N*3/2(1910) INTO N*3/2(1232) PI 1232+ 139
P5 N*3/2(1910) INTO GAM NUCLEON, HELICITY=1/2 0+ 938
P6 N*3/2(1910) INTO N RHO 938+ 773

12 N*3/2(1910) BRANCHING RATIOS
R1 N*3/2(1910) INTO (PI N)/TOTAL (P1)
R1 3 (0.30) DCNNACHI 68 RVUE 8/69
R1 6 (0.128) AYED 70 IPWA 1/71
R1 4 (0.18) DAVIES 70 RVUE SOL A 8/69
R1 7 (0.33) ALMEHED 72 IPWA 2/72
R1 1 (.158) AYED 74 IPWA 2/74
R2 N*3/2(1910) INTO (K SIGMA)/TOTAL (P3)
R2 1 (0.08)OR LESS FEUERBACH 70 RVUE PI P TO K+ SIG+ 7/70
R2 1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68
R2 1 MODEL USED MAY DOUBLE COUNT.

12 N*3/2(1910) PHOTON DECAY AMPL(GEV**-1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.
A1 N*3/2(1910) INTO GAM NUCLEON, HELICITY=1/2 (GEV**-1/2)
A1 .000 .025 DEVENIS2 74 DPWA PI N PHOTO-PROD 4/75*
A1 .010 .012 KNIES 74 DPWA PI N PHOTO-PROD 2/74
A1 -.032 .065 METCALF 74 DPWA PI N PHOTO-PROD 2/74
A1 -.009 .013 CRAWFORD 75 DPWA PI N PHOTO-PROD 1/76*
A1 (-.031) BARBOUR 76 DPWA PI N PHOTO-PROD 1/76*

REFERENCES FOR N*3/2(1910)
DCNNACHI 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
ALSO 68 VIENNA 139 DCNNACHIE RAPPORTEUR'S TALK (GLAS)
ALSO 68 THESIS R G KIRSOPP (EDIN)
AYED 70 KIEV CCFN R AYED,P BAREYRE, G VILLET (SACL)IJP
DAVIES 70 NP 821 359 A DAVIES (GLAS)
FEUERBACH 70 NP 168 85 FEUERBACHER+HOLLADAY (VANDERBILT)
ALMEHED 72 NP 840 157 +LOVELACE (LUND,RUTG)IJP
LANGBEIN 73 NP 853 251 LANGBEIN,WAGNER (MUNICH)IJP
AYED 74 PRIVATE COMMCTN. AYED,BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP
DEVENIS2 74 PL 528 227 DEVENISH,LYTH,RANKIN (DESY,LANC,BONN)IJP
KNIES 74 PRD 9 2680 KNIES,MOORHOUSE,OBERLACK (LBL,GLAS)IJP
METCALF 74 NP 876 253 W J METCALF,R L WALKER (CIT)IJP

Data Card Listings

For notation, see key at front of Listings.

Baryons

Δ(1910), Δ(1950), Δ(1960)

CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
DEANS 75 NP 896 90 +MITCHELL,MCNTGCMERY,+ (SFLA,ALABAMA)IJP
BARBOUR 76 SBMTD. TO NP I. M. BARBOUR,R. L. CRAWFCRD (GLAS)IJP
PAPERS NOT REFERREC TO IN DATA CARDS.
CARYANN 65 PR 138 8433 CARAYANNOPOULOS,TAUTFEST,WILLMANN (PURD)
A PARTIAL WAVE ANALYSIS OF PI+P TO SIGMA+ K+ (PURD)
AYED 70 PL 318 958 +BAREYRE+VILLET (SACLAY)

Δ(1950)

83 N*3/2(1950, JP=7/2+) I=3/2 F37
THE EXISTENCE OF THIS RESNANCE IS WELL ESTABLISHED.

83 N*3/2(1950) MASS (MEV)
M (1920.0) DUKE 65 CNTR PI-P EL + PCL 7/66
M (1950.0) APPROX YOKOSAWA 66 CNTR PI- P DSIG + PCL 7/66
M 1 (1975.0) BAREYRE 68 RVUE PHASE-SHIFT ANAL 11/67
M 3 (1946.0) WHERE CROSS SECTION IS GREATEST - EYEBALL FIT 6/68
M 6 (1931.0) DONNACHI 68 RVUE PHASE-SHIFT ANAL 1/71
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM
M 4 (1935.0) DAVIES 70 RVUE P-S ANAL SOL A 8/69
M (1950.0) (30.0) KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
M (1920.0) ROYCHOUD 71 DPWA 3/72
M 7 (1925.0) ALMEHED 72 IPWA 2/72
M 2 1920 TO 1951 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
M (1928.0) AYED 74 IPWA 2/74
M L 1930. OR 1925. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
M L THE 2 SETS OF PARAMETERS ARE FROM METHODS 1 AND 2 OF LONGACRE 75. 11/75*

83 N*3/2(1950) WIDTH (MEV)
W (170.0) DUKE 65 CNTR 7/66
W (200.0) APPROX YOKOSAWA 66 CNTR 7/66
W 1 (190.0) BAREYRE 68 RVUE 11/67
W 3 (221.0) DONNACHI 68 RVUE 6/68
W 6 (197.0) AYED 70 IPWA 1/71
W 4 (221.0) DAVIES 70 RVUE SOL A 8/69
W (300.0) (60.0) KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
W 7 (200.0) ALMEHED 72 IPWA 2/72
W 2 234 TO 269 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
W (237.0) AYED 74 IPWA 2/74
W L 235. OR 240. LONGACRE 75 IPWA PI N TO 2PI N 11/75*
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

83 N*3/2(1950) REAL PART OF POLE POSITION (MEV) 11/75*
RE (1924.0) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

83 N*3/2(1950) 2*IMAG PART OF POLE POSITION (MEV) 11/75*
IM (258.0) LONGACRE 75 IPWA PI N TO 2PI N 11/75*

83 N*3/2(1950) PARTIAL DECAY MODES
P1 N*3/2(1950) INTO PI N DECAY MASSES 139+ 938
P2 N*3/2(1950) INTO SIGMA K 1189+ 493
P3 N*3/2(1950) INTO N*3/2(1232) PI 1232+ 139
P4 N*3/2(1950) INTO Y*1(1365) K 1384+ 493
P5 N*3/2(1950) INTO N*3/2(1232) RHO 1232+ 773
P6 N*3/2(1950) INTO NEUTRON PI+ PI* 939+ 139+ 139
P7 N*3/2(1950) INTO N*3/2(1232) PI PI (NOT RHO) 1232+ 938
P8 N*3/2(1950) INTO GAM NUCLEON, HELICITY=1/2 0+ 938
P9 N*3/2(1950) INTO GAM NUCLEON, HELICITY=3/2 0+ 938
P10 N*3/2(1950) INTO N RHO 938+ 773
P11 N*3/2(1950) INTO N*3/2(1232) PI, F=-WAVE 1232+ 139
P12 N*3/2(1950) INTO N*3/2(1232) PI, H=-WAVE 1232+ 139
P13 N*3/2(1950) INTO N RHO, S=3/2, F=-WAVE 938+ 773

83 N*3/2(1950) BRANCHING RATIOS
R1 N*3/2(1950) INTO (PI N)/TOTAL (P1) VERY ENERGY DEP 7/66
R1 (0.41) DUKE 65 CNTR 7/66
R1 (0.4) APPROX YOKOSAWA 66 CNTR 11/67
R1 1 (0.57) BAREYRE 68 RVUE 6/68
R1 3 (0.386) DONNACHI 68 RVUE 1/71
R1 6 (0.496) AYED 70 IPWA 8/69
R1 4 (0.51) DAVIES 70 RVUE SOL A 2/72
R1 7 (0.4) ALMEHED 72 IPWA 2/74
R1 (0.405) AYED 74 IPWA 2/74
SEE THE NOTES ACCOMPANYING THE MASSES QUOTED.

R2 N*3/2(1950) INTO (SIGMA K)+(PI N)/TOTAL**2 (P2*P1) 10/69
R2 (0.004) (0.008) BORREANI 68 HBC PI+P 1.35-1.68 7/70
R2 1 ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68 PI P TO K+ SIG+
R2 1 MODEL USED MAY DOUBLE COUNT. FEUERBACH 70 RVUE
R2 0.0081 0.0013 KALMUS 70 DPWA PI+P TO K+ SIG+ 1/71
R3 N*3/2(1950) FROM PI N TO SIGMA K SQRT(PI*P3) 9/73
R3 2 .37 TO .48 MEHTANI 72 MPWA PI+P TO PI+PIOP 9/73
R3 2 MOSTLY F WAVE DECAY
R4 N*3/2(1950) FROM PI N TO SIGMA K SQRT(PI*P2) 9/73
R4 (.04) LANGBEIN 73 IPWA PI N-K SIG,SOL 1 9/73
R4 (.05) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
R4 5 (1.02) DEANS 75 DPWA PI N TO K SIGMA 11/75*
R4 5 RANGE GIVEN IS FROM FOUR BEST SOLUTIONS..
R5 N*3/2(1950) FROM PI N TO N*3/2(1232) PI, F=-WAVE SQRT(PI*P11) 11/75*
R5 L (-.2510R -.32 LONGACRE 75 IPWA PI N TO 2PI N 11/75*

R6 N*3/2(1950) FROM PI N TO N RHO, S=3/2, F=-WAVE SQRT(PI*P13) 11/75*
R6 L (.1810R -.24 LONGACRE 75 IPWA PI N TO 2PI N 11/75*
MORE INFORMATION ON INELASTIC DECAY MODES OF BUMPS, SEEN IN PRODUCTION EXPERIMENTS AROUND 1950 MEV, MAY BE FOUND IN THE NEXT ENTRY

83 N*3/2(1950) PHOTON DECAY AMPL(GEV**1/2)
FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1 N*3/2(1950) INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)
A1 -.088 .025 DEVENIS2 74 DPWA PI N PHOTO-PRD 4/75*
A1 -.070 .012 KNIES 74 DPWA PI N PHOTO-PRD 2/74
A1 -.055 .029 METCALF 74 DPWA PI N PHOTO-PRD 2/74
A1 (-.080) .005 MCCRHOUSE 74 DPWA PI N PHOTO-PRD 2/74
A1 -.038 .014 CRAWFORD 75 DPWA PI N PHOTO-PRD 1/76*
A1 (-.076) BARBOUR 76 DPWA PI N PHOTO-PRD 1/76*
A1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.3)
A2 N*3/2(1950) INTO GAM NUCLEON, HELICITY=3/2 (GEV**1/2)
A2 -.080 .021 DEVENIS2 74 DPWA PI N PHOTO-PRD 4/75*
A2 -.078 .010 KNIES 74 DPWA PI N PHOTO-PRD 2/74
A2 -.093 .024 METCALF 74 DPWA PI N PHOTO-PRD 2/74
A2 (-.180) MCCRHOUSE 74 DPWA PI N PHOTO-PRD 2/74
A2 -.038 .014 CRAWFORD 75 DPWA PI N PHOTO-PRD 1/76*
A2 (-.065) BARBOUR 76 DPWA PI N PHOTO-PRD 1/76*
A2 AVERAGE MEANINGLESS (SCALE FACTOR = 1.5)

REFERENCES FOR N*3/2(1950)

DUKE 65 PRL 15 468 +JONES,KEMP,MURPHY,PRENTICE, + (RHEL,OXF)IJP
YOKOSAWA 66 PRL 16 714 +SUWA, FILL, ESTERLING, BOOTH (ANL,CHIC)IJP
BAREYRE 68 PR 165 1731 P BAREYRE, C BRICMAN, G VILLET (SACLAY)IJP
BORREANI 68 UCLR 18350 BORREANI, KALMUS (LRL)
DONNACHI 68 PL 268 161 A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP
ALSO 68 VIENNA 139 DONNACHIE RAPPORTEUR'S TALK (GLAS)
ALSO 68 THESIS R G KIRSOPP (EDIN)
AYED 70 KIEV CONF R AYED, P BAREYRE, G VILLET (SACL)IJP
DAVIES 70 NP 821 359 A DAVIES (GLAS)
FEUERBACH 70 NP 168 85 FEUERBACHER+HCLLDAY (VANDERBILT)
KALMUS 70 PR D2 1824 G KALMUS, G BORREANI, J LOUIE (LRL)
ROYCHOUD 71 NP 827 125 R K ROYCHOUDHURY, B H BRANSDEN (DURH)IJP
ALMEHED 72 NP 840 157 +LCVELACE (LUND,RUTG)IJP
MEHTANI 72 PRL 29 1634 +FUNG, KERNAN, SCHALK, + (UCR +LBL)
LANGBEIN 73 NP 853 251 LANGBEIN,WAGNER (MUNICH)IJP
AYED 74 PRIVATE COMMCTN. AYED,BAREYRE (SACL)IJP
ALSO 73 AIX CONFERENCE AYED,BAREYRE (SACL)IJP
DEVENIS2 74 PL 528 227 DEVENIS, LYTH, RANKIN (DESY,LANC,BCN)IJP
KNIES 74 PRD 9 2680 KNIES, MCCRHOUSE, OBERLACK (LBL, GLAS)IJP
METCALF 74 NP 874 253 W J METCALF, R L WALKER (CIT)IJP
MCCRHOUSE 74 PRD 9 1 MCCRHOUSE, OBERLACK, ROSENFELD (GLAS+LBL)IJP
CRAWFORD 75 NP 897 125 R L CRAWFORD (GLAS)IJP
DEANS 75 NP 896 90 +MITCHELL,MCNTGCMERY,+ (SFLA,ALABAMA)IJP
LONGACRE 75 PL 558 415 +ROSENFELD,LASINSKI,SMADJA+ (LBL,SLAC)IJP
BARBOUR 76 SBMTD. TO NP I. M. BARBOUR,R. L. CRAWFCRD (GLAS)IJP
PAPERS NOT REFERREC TO IN DATA CARDS.

HCHLER 63 NP 48 470 G HOHLER, G EBEL (KARLSRUHE) I
LAYSON 63 NC 27 724 W H LAYSON (CERN) IJ
AUVIL 64 NC 33 473 P AUVIL, C LOVELACE (LOIC)IJP
HELLAND 64 PR 134 81062 +DEVLIN,HAGGE,LCNGO,MOYER,WOOD (LRL) IJ
HOHLER 64 PL 12 149 G HOHLER, J GIESECKE (KARLSRUHE) I
HOLLADAY 65 PR 139 81348 W G HOLLADAY (VANDERBILT)
JOHNSON 67 UCLR-17483 THESIS C H JOHNSON (LRL)
DONNACHI 69 NP 108 433 A DONNACHIE, R KIRSOPP (GLAS+EDIN)
AYED 70 PL 318 598 +BAREYRE+VILLET (SACLAY)

Δ(1960) 13 N*3/2(1960, JP=5/2-) I=3/2 D35
THIS EFFECT IS SEEN IN TWO CHANNELS.

13 N*3/2(1960) MASS (MEV)
M 3 (1954.0) DONNACHI 68 RVUE PHASE-SHIFT ANAL 6/68
M 3 (1970.0) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
M X (1950.0) APPROX LEA 69 CNTR PI-P ELASTIC 8/69
M X SEE ALSO AFLIN 70
M 3 WHERE MAX. ABSORPTION IS -DONNACHI, 2, KIRSOPP EYEBALL FIT CERN 1 10/69
M 7 (2200.0) ALMEHED 72 IPWA 2/72
M 1 (1560.0) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
M 1 NOT SEEN IN SOLUTION 1 OF LANGBEIN 73 9/73
M (1889.0) AYED 74 IPWA 2/74

13 N*3/2(1960) WIDTH (MEV)
W 3 (311.00) DONNACHI 68 RVUE 8/69
W 3 (400.0) KIRSOPP 68 RVUE PHASE SHIFT ANAL 10/69
W 7 (600.0) ALMEHED 72 IPWA 2/72
W 1 (150.0) LANGBEIN 73 IPWA PI N-K SIG,SOL 2 9/73
W (1121.0) AYED 74 IPWA 2/74
SEE THE NOTES ACCOMPANYING MASSES QUOTED

Baryons

$\Delta(1960)$, $\Delta(2160)$

Data Card Listings

For notation, see key at front of Listings.

13 N*3/2(1960) PARTIAL DECAY MODES

P1	N*3/2(1960)	INTG PI N	DECAY MASSES	
P2	N*3/2(1960)	INTO K SIGMA	139+ 938	
			493+1189	

13 N*3/2(1960) BRANCHING RATIOS

R1	N*3/2(1960)	INTO (PI N)/TOTAL	(P1)	
R1 3	(.154)	DONNACHI 68 RVUE	PHASE SHIFT ANAL.	10/69
R1 3	(.12)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL.	10/69
R1 7	(0.25)	ALMEHD 72 IPWA		2/72
R1	(.081)	AYED 74 IPWA		2/74

R2 N*3/2(1960) INTO (K SIGMA)/TOTAL (P2)

R2 1	(0.013)	(0.01)	FEUERBACH 70 RVUE	PI P TO K+ SIG+	7/70
R2 1			ASSUME MASS, WIDTH, X(ELAST) OF DONNACHIE 68		
R2 1			MODEL USED MAY DOUBLE COUNT.		

R3 N*3/2(1960) FROM PI N TO K SIGMA

R3 1	(.08)	LANGBEIN 73 IPWA	PI N-K SIG,SOL 2	9/73
R3 2	(.018) TO .035	DEANS 75 DPWA	PI N TO K SIGMA	11/75*
R3 2		RANGE GIVEN IS FROM FOUR BEST SOLUTIONS.		11/75*

13 N*3/2(1960) PHOTON DECAY AMPL(GEV**1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

A1	N*3/2(1960)	INTO GAM NUCLEON, HELICITY=1/2 (GEV**1/2)	1/76*
A1	+0.03	-CRAWFORD 75 DPWA	PI N PHOTO-PRCD
A1	(-.085)	BARBOUR 76 DPWA	PI N PHOTO-PRCD

A2	N*3/2(1960)	INTO GAM NUCLEON, HELICITY=3/2 (GEV**1/2)	1/76*
A2	-.010	+CRAWFORD 75 DPWA	PI N PHOTO-PRCD
A2	(+.066)	BARBOUR 76 DPWA	PI N PHOTO-PRCD

REFERENCES FOR N*3/2(1960)

DONNACHI 68 PL 268 161	A DONNACHIE, R G KIRSOPP, C LOVELACE (CERN)IJP	
KIRSOPP 68 THESIS	R G KIRSOPP (EDIN)	
LEA 69 PL 298 584	LEA, DADES, WARD, COWAN, + (RH,EL, BRISTOL, DARE)	
FEUERBACH 70 NP 168 85	FEUERBACHER+HOLLADAY (VANDERBILT)	
ALMEHD 72 NP 840 157	+LOVELACE (LUND, RUTGIIJP)	
LANGBEIN 73 NP 853 251	LANGBEIN,WAGNER (MUNICH)IJP	
AYED 74 PRIVATE COMMCTN.	AYED,BAREYRE (SACLIIJP)	
ALSC 73 AIX CONFERENCE	AYED,BAREYRE (SACLIIJP)	
CRAWFORD 75 NP 897 125	R L CRAWFORD (GLASIIJP)	
DEANS 75 NP 896 90	+MITCHELL, MCNTGCHERY,+ (SFLA-ALABAMA)IJP	
BARBOUR 76 SBMTD. TO NP	I. M. BARBOUR, R. L. CRAWFORD (GLASIIJP)	
	PAPERS NOT REFERRED TO IN DATA CARDS.	
DONNACHI 69 NP 108 453	A DONNACHIE, R KIRSOPP (GLAS+EDIN)	
AYED 70 PL 318 598	+BAREYRE+VILLET (SACLAY)	
APLIN 71 NP 832 253	+COWAN,GIBSON,GILMORE++ (RH,EL, BRISTOL)	

1950 MEV REGION - PRODUCTION AND σ_{TOTAL} EXP'TS

70 N*3/2(1950, JP= 1 I=3/2) PRODUCTION EXPERIMENTS

70 N*3/2(1950) MASS (MEV) (PROD. EXP.)

M	(1922.0)	APPROX	COOL 56 CNTR	PI+ P TOTAL	7/66
M	(192.0)	(15.0)	BRISSON 61 CNTR	PI+ P TOTAL	7/66
M	(1900.0)	(9.0)	DEVLIN 65 CNTR	PI+ P TOTAL	
M	(2080.0)	(12.0)	YDGN 67 HBC	+ 3 BEV/C PI-P	8/67
M	N		THIS BUMP IS NOT SEEN BY CHUNG 68 AT 3.2 GEV/C		
M	(1860.0)		COLTCN 72 HBC	++ PP TO PI+PN 7GEV	1/73
M	(1895.0)	(15.0)	COLLEY 74 HBC	++ K+P TO K+PI+PI-	10/74*
M	(1890.0)	5. TO 10.	BRAUNZ 75 BC	PBAR P AND D,5-7	11/75*
M	C	(1881.0)	CHUNG 75 HBC	++ PI+ P AND K+ P	1/76*
M	C		MOST PROBABLE JP ASSIGNMENT IS 7/2+		1/76*
M	(1880.0)	(10.0)	GAIDOS 75 HBC	++ PI+P TO N* 2PI	1/76*

70 N*3/2(1950) WIDTH (MEV) (PROD. EXP.)

M	(256.0)	(39.0)	DEVLIN 65 CNTR		
M	(40.0)	(20.0)	YDGN 67 HBC	+	8/67
M	(180.0)		COLTCN 72 HBC	++ PP TO PI+PN 7GEV	1/73
M	(230.0)	(50.0)	COLLEY 74 HBC	++ K+P TO K+PI+PI-	10/74*
M	(120.0)	10. TO 20.	BRAUNZ 75 BC	PBAR P AND D,5-7	11/75*
M	C	(219.0)	CHUNG 75 HBC	++ PI+ P AND K+ P	1/76*
M	(180.0)	(30.0)	GAIDOS 75 HBC	++ PI+P TO N* 2PI	1/76*

70 N*3/2(1950) PARTIAL DECAY MODES (PROD. EXP.)

P1	N*3/2(1950)	INTO PI N	DECAY MASSES	
P2	N*3/2(1950)	INTO SIGMA K	139+ 938	
P3	N*3/2(1950)	INTO N*3/2(1232) PI	1185+ 493	
P4	N*3/2(1950)	INTO Y*(11385) K	1232+ 139	
P5	N*3/2(1950)	INTO N*3/2(1232) RHO	1384+ 493	
P6	N*3/2(1950)	INTO NEUTRON PI+ PI+	1232+ 773	
P7	N*3/2(1950)	INTO N*3/2(1232) PI PI (NOT RHO)	939+ 139+ 139	
			1232+ 139+ 139	

70 N*3/2(1950) BRANCHING RATIOS (PROD. EXP.)

R1	N*3/2(1950)	INTO (PI N)/TOTAL	(P1)	
R1	(0.57)	(0.12)	DEVLIN 65 CNTR	
R2	N*3/2(1950)	INTO (SIGMA K)/(PI N)	(P2)/(P1)	
R2	0.059	0.024	CHINOWSKY 68 HBC	++ PP TO P SIG K
				11/68

R3	N*3/2(1950)	INTO N*3/2(1232) PI PI (NOT RHO)	(P7)	
R3	SEEN	CHINOWSKY 68 HBC	++ PP TO (P 3PI) N	11/68
R3	SEEN	BOGGILD 70 HBC	PP TO N3PI(NTRL)	6/70

R4	N 3/2(1950)	INTO (PI N)/(N*3/2(1232) PI)	(P1)/(P3)	
R4	(0.55) OR LESS	LEE 67 HBC	PI-P 3.63 BEV/C	11/67

R5	N*3/2(1950)	INTO ((PI N)*(NEUTRON PI+ PI+))/TOTAL**2	(P1*P6)	
R5	0.05	0.013	GALLOWAY 68 RVUE	++ PI+P TO N 2PI+
				6/68

R6	N*3/2(1950)	INTO (Y*(11385) K)/(PI N)	(P4)/(P1)	
R6	0.035	0.015	CHINOWSKY 68 HBC	++ PP TO P LAM K PI
				11/68

R7	N*3/2(1950)	INTO (N*3/2(1232) RHO)/(PI N)	(P5)/(P1)	
R7	(0.45)	APPROX	CHINOWSKY 68 HBC	++ PP TO (P 3PI) N
R7			+MGBS+ROE, SINCLAIR-VANDER VELDE (MICH)	11/68
			THIS INCLUDES CORRECTION FOR UNSEEN DECAY (ISPIN FACTOR 5/3).	

R8	N*3/2(1950)	INTO (N*3/2(1232) RHO)/TOTAL	(P5)	
R8	SEEN	YDGN 67 HBC	+	8/67
R8	NOT SEEN	BOGGILD 70 HBC	PP TO N3PI(NTRL)	6/70

REFERENCES FOR N*3/2(1950) (PRCD. EXP.)

COOL 56 PR 103 1082	R CCCL, D PICCINI, D CLARK (BNL) I
BRISSON 61 NC 19 210	+DETUEF,FALK-VAIRANT,VAN ROSSUM,+ (SACLAY) I
DEVLIN 65 PRL 14 1031	T J DEVLIN, J SOLOMON, G BERTSCH (PRINCETON) I
LEE 67 PR 159 1156	+MGBS+ROE, SINCLAIR-VANDER VELDE (MICH)
YDGN 67 PL 248 307	+BERENYI, KEV, PRENTICE, + (TORONTG,WISC)
CHINOWSKY 68 PR 171 1421	CHINOWSKY, CNDON, KINSEY, KLEIN, + (LRL, SLAC)
CHUNG 68 PR 165 1491	S U CHUNG, DAHL, KIRZ, MILLER (LRL)
GALLOWAY 68 PL 268 334	K F GALLOWAY (INDIANA) I
BOGGILD 70 NP 816 503	+KOREA-AHO+JACOBSEN+ (BOHR+ HELS+OSLO-STOH)
COLTCN 72 PR 08 95	E COLTCN, A KIRSCHBAUM (LBL)
COLLEY 74 NP 865 205	COLLEY+HODGINS, KINSEY, MILNE, + (BIRM+GLAS) I
BRAUNZ 75 NP 895 503	+GERBER, MAUREK, MICHALON, SCHIBY+ STBRB, LPNP I
CHUNG 75 PL 578 384	+PROTOPOESCU, EISNER+ (BNL+CASE+LBL+UCSC)IJP
ALSO 75 PRD 12 693	CHUNG, PROTOPOESCU, EISNER+ (BNL+CASE+UCSC)IJP
GAIDOS 75 PRD 12 2565	GAIDOS, MILLER (PURDUE) IJP

PAPERS NOT REFERRED TO IN DATA CARDS.

DEUTSCHM 75 NP 899 397 DEUTSCHMANN+(AACH+BONN+BERL+CERN+CRAC+EID)

$\Delta(2160)$ \rightarrow 9 N*3/2(2160,) I=3/2

EARLY ANALYSES FOUND EVIDENCE FOR A RESONANCE NEAR THIS MASS IN THE P33 PARTIAL WAVE, AND UNLESS STATED OTHERWISE, ALL DATA CARDS BELOW APPLY TO THIS WAVE. IN ADDITION, ROYCHOUDHURY 71 FIND POSSIBLE EVIDENCE FOR P31, D33, AND D35 RESONANCES IN THIS MASS REGION. IN A SIMILAR ANALYSIS BRANSDEN 71 FOUND SOME EVIDENCE FOR S31, D33, AND D35 RESONANCES IN THIS REGION. VCN SCHLIPPE 72 SUGGESTS A G39. THE MOST RECENT PI N ANALYSIS FINDS A G39 RESONANCE IN THIS MASS REGION. THE EFFECT SEEN IN K SIG PROD. IS 100 MEV LOWER IN MASS. A PRONOUNCED SHARP DIP IS OBSERVED IN PI+ P BACKWARD SCATTERING AT 2200 MEV BY REY 74.

DUAL INTERFERENCE MODEL ANALYSIS OF MA 75 FINDS SIGNAL FOR P33, P31, AND D35, BUT NOT FOR G39.

9 N*3/2(2160) MASS (MEV)

M 3	(2160.)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69
M	(2120.)	ROYCHOUD 71 DPWA		3/72
M 7	(2150.)	ALMEHD 72 IPWA		2/72
M 1	(1980.)	LANGBEIN 73 IPWA	PI N-K SIG,SOL 1	9/73
M 2		NOT SEEN IN SOLUTION 2 OF LANGBEIN73		9/73
M 2	(2174.)	AYED 74 IPWA		2/74
M 2		AYED 74 RESULT IS A G39 RESONANCE.		2/74
M 4	(2196.)	(46.) (41.) REY 74 MPWA	++ PI+ P 180 DEG CS	10/74*
M 4		BAKER 74 AND REY 74 FIND NEGATIVE PARITY (SPIN UNDETERMINED).		

9 N*3/2(2160) WIDTH (MEV)

W 3	(260.)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69
W 7	(200.)	ALMEHD 72 IPWA		2/72
W 1	(190.)	LANGBEIN 73 IPWA	PI N-K SIG,SOL 1	9/73
W 2	(205.)	AYED 74 IPWA		2/74
W 4	(302.)	(143.) REY 74 MPWA	++ PI+ P 180 DEG CS	10/74*

SEE THE NOTES ACCOMPANYING MASSES QUOTED

9 N*3/2(2160) PARTIAL DECAY MODES

P1	N*3/2(2160)	INTO PI N	DECAY MASSES	
P2	N*3/2(2160)	INTO K SIGMA	139+ 938	
			493+1189	

9 N*3/2(2160) BRANCHING RATIOS

R1	N*3/2(2160)	INTO (PI N)/TOTAL	(P1)	
R1 3	(.25)	KIRSOPP 68 RVUE	PHASE SHIFT ANAL	10/69
R1 7	(0.3)	ALMEHD 72 IPWA		2/72
R1 2	(0.45)	AYED 74 IPWA		2/74
R1 4	REY74 FINDS (J+1/2)X=.81+/-(.54/.35)			10/74*

R2	N*3/2(2160)	FROM PI N TO K SIGMA	SORT(P1*P2)	9/73
R2 1	(.08)	LANGBEIN 73 IPWA	PI N-K SIG,SOL 1	9/73
R2 5	(.048) TO .120	DEANS 75 DPWA	PI N TO K SIGMA	11/75*
R2 5		RANGE GIVEN IS FROM FOUR BEST SOLUTIONS.		11/75*

REFERENCES FOR N*3/2(2160)

KIRSOPP 68 THESIS	R G KIRSOPP (EDIN)
ROYCHOUD 71 NP 827 125	R K ROYCHOUDHURY, B H BRANSDEN (DURH)IJP
ALMEHD 72 NP 840 157	+LOVELACE (LUND, RUTGIIJP)
LANGBEIN 73 NP 853 251	LANGBEIN,WAGNER (MUNICH)IJP

Data Card Listings

For notation, see key at front of Listings.

Baryons

Δ(2160), Δ(2420), Δ(2850)

AYED 74 PRIVATE COMMCTN.
ALSO 73 AIX CONFERENCE
REY 74 PRL 32 908
ALSO 74 PRL 33 250
ALSO 75 PRD 11 1777
DEANS 75 NP 896 90

BRANDEN 71 NP 826 511
ALSO 70 NP 816 461
VON SCHL 72 LNC 4 767
BAKER 74 PRL 32 251
MA 75 PRD 11 1852

Δ(2420) 84 N*3/2(2420, JP=11/2+) I=3/2 H3 11

BOTH ROYCHOUDHURY 71 AND BRANDEN 71 SEE A POSSIBLE RESONANT F35 IN THIS MASS REGION. IN ADDITION BRANDEN 71 FIND A RESONANT P33 AT 2600 MEV.

84 N*3/2(2420) MASS (MEV)
M 6 (2312.0)
M 6 FROM ENER. DEP. FIT OF ARGAND DIAGRAM
M (2400.)
M (2400.)
M (2440.)
M (2392.)
M 1 (2404.) (63.)

84 N*3/2(2420) WIDTH (MEV)
W 6 (347.0)
W (289.)
W 1 (484.) (75.)

84 N*3/2(2420) PARTIAL DECAY MODES
P1 N*3/2(2420) INTO PI N
P2 N*3/2(2420) INTO SIGMA K

84 N*3/2(2420) BRANCHING RATIOS
R1 N*3/2(2420) INTO (PI N)/TOTAL
R1 6 (0.113)
R1 7 (4.)
R1 (1.09)
R1 1 (.157) (.070) (.035)REY
R1 1 REY 74 DETERMINES (J+1/2)X ONLY, WE HAVE DIVIDED BY 6.

AYED 70 KIEV CONF
BRANDEN 71 NP 826 511
ALSO 70 NP 816 461
ROYCHOUD 71 NP 827 125
OTT 72 PL 428 133
ALSO 72 MCGILL THESIS
AYED 74 PRIVATE COMMCTN.
ALSO 73 AIX CONFERENCE
REY 74 PRL 32 908
ALSO 74 PRL 33 250
ALSO 75 PRD 11 1777

BELLAMY 67 PRL 19 476
AYED 70 PL 318 598

2420 MEV REGION - PRODUCTION AND σTOTAL EXP'TS

69 N*3/2(2420, JP=) I=3/2 PRODUCTION EXPERIMENTS

69 N*3/2(2420) MASS (MEV) (PROD. EXP.)
M (2360.0)
M (2520.0) (40.0)
M (2440.0)
M (2400.0) APPROX
M B (2452.0)
M B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE
M B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
M 2423.0 10.0

69 N*3/2(2420) WIDTH (MEV) (PROD. EXP.)
W (200.0)
W (245.0)
W B (275.0)
W 310.0 20.0

69 N*3/2(2420) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*3/2(2420) INTO PI N
P2 N*3/2(2420) INTO SIGMA K
P3 N*3/2(2420) INTO N*3/2(1232) PI
P4 N*3/2(2420) INTO NEUTRON PI+ PI+

69 N*3/2(2420) BRANCHING RATIOS (PROD. EXP.)
R1 N*3/2(2420) INTO (PI N)/TOTAL
R1 (0.067) APPROX
R1 0.113 0.0036
R1 B (0.12)
R1 D (0.163)
R1 D USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (0.06) KORMANYOS 67 CNTR ASSUMING J=11/2 7/66

R2 N*3/2(2420) INTO (PI N)*(NEUTRON PI+ PI+)/(TOTAL**2)
R2 0.0195 0.0048 GALLOWAY 68 RVUE 6/68

REFERENCES FOR N*3/2(2420) (PROD. EXP.)
DIDDENS 63 PRL 10 262
ALVAREZ 64 PRL 12 710
HCHLER 64 PL 12 149
WAHLIG 64 PR 13 103
BARGER 66 PR 151 1123
CITRON 66 PR 144 1101
BARGER 67 PR 155 1792
DIKMEN 67 PR 18 798
KORMANYOS 67 PR 164 1661
GALLOWAY 68 PL 268 334

BAACKE 67 NC 51A 761
DOBROWOL 67 PL 248 203
DOLEN 68 PR 166 1768
WAHLIG 68 PR 168 1515
J BAACKE, M YVERT
DOBROWOLSKI, GUSKOV, LIKHACHEV, +
R DOLEN, D HORN, C SCHMID
M A WAHLIG, I MANNELLI
FINAL VERSION OF DATA USED IN WAHLIG 64-. IN CONJUNCTION WITH CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

Δ(2850) BUMPS

85 N*3/2(2850, JP=) I=3/2 PRODUCTION EXPERIMENTS

85 N*3/2(2850) MASS (MEV) (PROD. EXP.)
M (2870.0)
M (2700.0) APPROX
M (2850.0)
M 2850.0 12.0
M (2883.1) (26.) (28.)

85 N*3/2(2850) WIDTH (MEV) (PROD. EXP.)
W (150.0)
W 400.0 40.0
W (380.) (141.) (201.)

85 N*3/2(2850) PARTIAL DECAY MODES (PROD. EXP.)
P1 N*3/2(2850) INTO PI N
P2 N*3/2(2850) INTO P PI PI
P3 N*3/2(2850) INTO N PI PI

85 N*3/2(2850) BRANCHING RATIOS (PROD. EXP.)
R1 N*3/2(2850) INTO (PI N)/TOTAL
R1 ONLY (J+1/2)*(PI N)/TOTAL MEASURED FOR THIS STATE
R1 B (0.224) (0.016)
R1 0.261 0.048
R1 B (0.40)
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 C (0.49)
R1 C USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (0.10)
R1 D (0.06) OR LESS CL=-.95 HALDORSE 72 HC PP 19 GEV/C 12/72
R1 D UPPER LIMIT ON ELASTICITY ALSO FIND J=9/2 OR MORE.
R1 (.28) (.13) (.19) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74**

REFERENCES FOR N*3/2(2850) (PROD. EXP.)
HCHLER 64 PL 12 149
WAHLIG 64 PRL 13 103
BARDADIN 66 PL 21 357
BARGER 66 PR 151 1123
CITRON 66 PR 144 1101
G HOHLER, J GIESECKE
MANNELLI, SODICKSON, FACKLER, WARD, +
BARDADIN-OTWINSKA, CANYSZ, +
V BARGER, M CLSSEN
GALBRAITH, KYCIA, LEONTIC, PHILLIPS, +
BARGER 67 PR 155 1792
DIKMEN 67 PR 18 798
DOBROWOL 67 PL 248 203
KORMANYOS 67 PR 164 1661
HALDORSE 72 NC 10A 468
REY 74 PRL 32 908
ALSO 74 PRL 33 250
ALSO 75 PRD 11 1777

Baryons

$\Delta(2850)$, $\Delta(3230)$, EXOTIC NUCLEONS, Z^* 's

Data Card Listings

For notation, see key at front of Listings.

PAPERS NOT REFERRED TO IN DATA CARDS.

BAACKE 67 NC 51A 761 J BAACKE, M YVERT (KARLSRUHE,ORSAYIJ-L
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)
WAHLIG 68 PR 168 1515 M A WAHLIG, I MANNELLI (MIT,PISA)
FINAL VERSION OF DATA USED IN WAHLIG 64. IN CONJUNCTION WITH
CITRON 66 TOTAL CROSS SECTIONS, THIS CHARGE EXCHANGE DATA GIVES
COMPLEX ELASTIC SCATTERING AMPLITUDE AT 0 DEGREES.

$\Delta(3230)$
BUMPS

86 N*3/2(3230, JP=) I=3/2 PRODUCTION EXPERIMENTS

86 N*3/2(3230) MASS (MEV) (PRCD. EXP.)

M (3230.0) CITRCN 66 CNTR PI+ P TOTAL 7/66
(3296.) (79.) (78.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

86 N*3/2(3230) WIDTH (MEV) (PROD. EXP.)

W (440.0) CITRCA 66 CNTR 7/66
W (687.) (1043.) (323.) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

86 N*3/2(3230) PARTIAL DECAY MODES (PROD. EXP.)

DECAY MASSES
P1 N*3/2(3230) INTO PI N 139+ 938
P2 N*3/2(3230) INTO N PI PI 938+ 139+ 139

86 N*3/2(3230) BRANCHING RATIOS

R1 N*3/2(3230) INTO (PI N)/TOTAL (PI)
R1 ONLY (J+1/2)* (PI N)/TOTAL MEASURED FOR THIS STATE
R1 B (0.03) (0.01) BARGER 66 RVUE TOTAL + CH EXC. 11/67
CITRON 66 CNTR TOTAL CROSS. SEC. 11/67
R1 B (0.06) TO 0.1 BARGER 67 CNTR USES KORMANYOS66 11/67
R1 B USES REGGE AMP.+RESON. TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGRE
R1 B FOR CRITICISM OF THIS METHOD, SEE DOLEN 68.
R1 D (0.25) DIKMAN 67 RVUE USES KORMANYOS67 11/67
R1 C USES ONLY RESONANCES TO CALCULATE DIF. CROSS SECTIONS AT 180 DEGREES
R1 (.45) (.09) (.13) REY 74 MPWA ++ PI+ P 180 DEG CS 10/74*

REFERENCES FOR N*3/2(3230) (PROD. EXP.)

BARGER 66 PR 151 1123 V BARGER, M CLSSON (WISC)
CITRON 66 PR 144 1101 +GALBRAITH,KYCIA,LEONTIC,PHILLIPS, + (BNL) I
BARGER 67 PR 155 1792 V BARGER, D CLINE (WISC) P
DIKMAN 67 PRL 18 798 F N DIKMAN (MICH)
REY 74 PRL 32 908 REV,LENNOX,POIRIER,PRETZL (NDAM+MPI)IP
ALSO 74 PRL 33 250 REV,LENNOX,POIRIER,PRETZL (NCAM+MPI)IP
ALSC 75 PRD 11 1777 LENNCX,POIRIER,REV,SANDER* (NDAM+FNAL+ANL)IP

PAPERS NOT REFERRED TO IN DATA CARDS

KORMANYOS 67 PR 164 1661 KORMANYOS, KRISCH, OFALLON, + (MICH,ANL) P
DOLEN 68 PR 166 1768 R DOLEN, D HORN, C SCHMID (CIT)

EXOTIC NUCLEONS - 1640 MEV REGION

THIS IS NOT A COMPLETE LIST. WE WILL TABULATE EXOTICS FROM NOW ON

EX(1640, JP=) I=5/2

THIS IS NOT A COMPLETE LIST. WE TABULATE
ONLY FROM 1970 ON.

IN A MISSING MASS EXPERIMENT, PI+ P TO PI- X+*,
BIRULEV 71 FIND NO EVIDENCE FOR EXOTIC (I=5/2) RESONANCES IN THE
MASS INTERVAL 1.2 TO 2.2 GEV.

EX(1640) MASS (MEV)

M A 29 1627. 12. PRICE 70 DBC -- K-D AT 4.91GEV/C 3/71
M A FOUR S. D. EFFECT

EX(1640) WIDTH (MEV)

W B 29 30. OR LESS CL=.90 PRICE 70 DBC -- PI-PI-N BUMP 3/71
W B CROSS SECTION 13.0+-3.9 MICROBARNS

EX(1640) CROSS SECTION LIMITS (MICROBARNS)

CS B 40. OR LESS BANNER 70 DSPK +++ PI+P,1.9 GEV/C 7/70
CS B I=5/2 LIMIT GIVEN ABOVE IS FOR MASS RANGE 1540-1750 MEV

REFERENCES FOR EX(1640)

BANNER 70 NP 815 205 +CHEZE,HAMEL,TEIGER,ZACCONE + (SACLAY)
PRICE 70 PL 338,595 +BERG,SALANT,WATERS,WEBSTER,WEINBERG (VAND)

PAPERS NOT REFERRED TO IN DATA CARDS

AMMANN 71 PL 348 533 +CARMONY,GARFINKEL,GLTAY,MILLER,YEN (PURD)
BIRULEV 71 SJNP 12 536 +VOVENKO,GUSKOV,DJBRVOVLSKII,+ (JINR)
JOHNSON 71 PL 348 428 D JOHNSON (ANL)

Note on the S = +1 Baryon System

Models based on states of three quarks have
been successful in the description of the spectrum
and decay of the known baryon resonances. Three
quarks cannot produce S = +1 baryon resonances
(Z*s), and this has probably been the primary
motivation for the great amount of experimental
effort that has gone into S = +1 baryon physics
below ~ 2 GeV/c during the last several years. In
addition, the S = +1 system offers an opportunity
to study single and double pion production in a
small number of rather distinct quasi-two-body
channels (K Δ , K*N, and K* Δ). Recent developments
in the field are summarized below.

I = 1 System

The most recent K⁺p total cross-section data
are from BNL (CARROLL 73; 0.41 to 1.06 GeV/c) and
Arizona-Michigan (BOWEN 73; 0.57 to 1.16 GeV/c).
Neither measurement exhibits the dip of approxi-
mately 1.5 mb at 0.7 GeV/c reported by BUGG 68 and
BOWEN 70. Recent differential elastic cross-section
measurements (including extensions and final results
of previously reported results) have come from
Glasgow-Bologna-Trieste (CAMERON 74; 0.13-0.76
GeV/c), Birmingham-RHEL (ADAMS 73; 0.43 to 0.94
GeV/c), Bristol-RHEL-Aarhus-Southampton (CHARLES
73; 0.7 to 1.9 GeV/c), Maryland-ANL (ABE 75; 0.86
to 2.12 GeV/c), Washington (ADAMS 75; 1.0-1.5 GeV/c
at 180°), UCL-RHEL (BARBER 73; 1.37-2.26 GeV/c),
Yale-BNL (PATTON 75; 1.7-3.0 GeV/c near 0° and
180°), and ANL (YUTA 74; 2.53, 2.76, and 3.20
GeV/c). Below 750 MeV/c (where the inelastic cross
section is less than 0.1 mb), the total elastic
cross sections from the first three above experi-
ments also fail to exhibit the dip seen in earlier
total cross-section experiments. Coulomb interfer-
ence effects are observed by CAMERON 74 and ADAMS
73; both confirm the early conclusion of GOLDBABER
62 that the low energy S-wave interaction is
repulsive. The most recent polarization measure-
ments are those of Maryland-ANL-Northwestern-FNAL
(BARNETT 73; 1.37-2.31 GeV/c) and Yale-BNL (PATTON
75; 1.7-3.0 GeV/c near 0° and 180°). Recent studies
of pion production in K⁺p interactions are those of
CERN-Saclay (BERTHON 73; KN π , 1.21-1.69 GeV/c),
CIT-UCLA (LOKEN 72; KN π and KN $\pi\pi$, 1.37-2.17 GeV/c),

Data Card Listings

For notation, see key at front of Listings.

Baryons
Z*'s

France-Saclay-IC-Westfield (BRUNET 72 and LEWIS 73; $\text{KN}\pi$, $\text{KN}\pi\pi$, and $\text{KN}\pi\pi\pi$, 2.11-2.72 GeV/c), and ANL (MUSGRAVE 75; $\text{KN}\pi$, 2.53-3.20 GeV/c). The general conclusion is that most pion production is associated with $\Delta(1232)$ and/or $\text{K}^*(892)$ production, though other resonances are also produced at the higher energies. The $I=1$ total and partial cross sections are shown as dashed lines in Fig. 1.

Five partial-wave analyses of K^+p elastic scattering have been carried out since 1973. CAMERON 74 performed an elastic energy-independent analysis of data at momenta up to 870 MeV/c, including their new high-statistics cross-section data below 755 MeV/c. S and P waves were kept up to 400 MeV/c with D waves being added at higher momenta, and solutions at different energies were linked by the shortest-path method, yielding a unique overall solution. An energy-dependent fit, with an effective-range parametrization for the S wave and a zero-effective-range scattering-length parametrization for P and D waves, was also performed and agreed well with the energy-independent analysis. The χ^2 per data point (χ^2/ND) was well below 1.0 at most momenta in both fits. Threshold parameters were extracted, the dominant threshold effect being a repulsive S wave. The energy-dependent ADAMS 73 analysis included data at momenta up to 1.5 GeV/c, but was aimed primarily at determining the amplitudes below 1.0 GeV/c. Below about 0.5 GeV/c a zero-effective-range scattering-length parametrization was used for S and P waves, with D and F waves set to zero; at higher energies the η and δ parameters of the S, P, D, and F waves were taken to be polynomials in c.m. momentum. Waves with $\ell \geq 4$ were set to the values calculated from two- and three-pion exchange by ALCOCK 73. The best χ^2 values obtained were about 2250 for 1600 degrees of freedom, of which an estimated 300 comes from inconsistent data. Two solutions were found; they are essentially identical below 1 GeV/c but differ significantly at higher energies. MARTIN 75 have performed an energy-dependent analysis of K^+p scattering up to 1.5 GeV/c as part of a simultaneous fit of $I=1$ and $I=0$ KN scattering. Below the inelastic threshold the inverse amplitudes for waves with $J \leq 5/2$ were parametrized as rational functions of the c.m. momentum which obeyed elastic

unitarity and had correct elastic threshold behavior.

To allow for inelasticity a (fitted) negative term was added to the imaginary part of the inverse amplitudes above the inelastic threshold. The F_{17} , G_{17} , and G_{19} waves were also kept, and set to the values calculated by ALCOCK 73. Two similar solutions were found with $\chi^2/\text{ND} = 1.1$. ARNDT 74

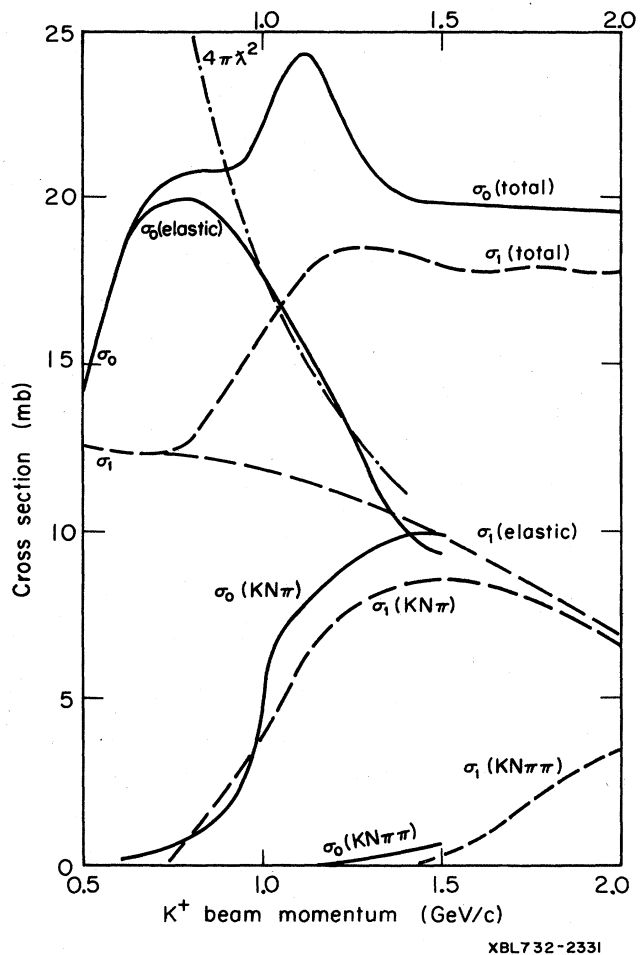


Fig. 1. KN total and partial cross sections. Subscripts indicate isospin; $I=1$ ($I=0$) cross sections are indicated by dashed (solid) curves. The total cross-section curves are from CARROLL 73, who use their recent data and that of BOWEN 73 as well as older data. The elastic $I=1$ curve is hand drawn through new and old elastic data. The inelastic $I=1$ curves are taken from LOKEN 72. The inelastic $I=0$ curves are taken from GIACOMELLI 72, and the elastic $I=0$ curve is obtained by subtracting these from the $I=0$ total cross-section curve of CARROLL 73.

Baryons

Z*'s

have analyzed data up to 2.0 GeV/c using an energy-dependent, two-channel K-matrix formalism in which the inelastic channel coupling to the S wave is taken to be K^*N , and that coupling to higher waves (with $\ell \leq 4$) is taken to be $K\Delta$. A unique solution

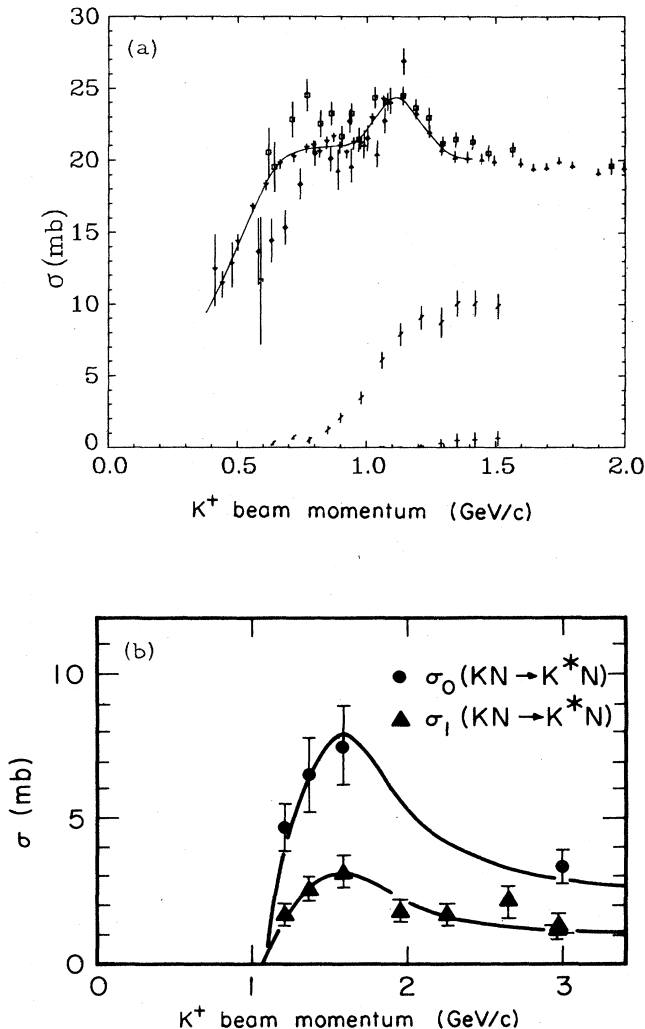


Fig. 2. (a) Unfolded $I=0$ cross sections as quoted by the various authors discussed in this mini-review:

- ◇ BOWEN 73 σ_T
- BUGG 68 σ_T (as unfolded by CARROLL 73)
- ▽ CARROLL 73 σ_T
- △ COOL 70 σ_T
- / GIACOMELLI 72 $\sigma(\pi KN)$
- GIACOMELLI 72 $\sigma(\pi\pi KN)$

(b) Energy dependence of the $I=0$ and $I=1$ cross sections for the reaction $KN \rightarrow K^*N$ (HIRATA 70).

Data Card Listings

For notation, see key at front of Listings.

with a χ^2 per degree of freedom (χ^2/DF) of 1.33 for 3822 data points was found, and it was possible to decrease this to 1.16 by dropping 70 data points with $\chi^2 > 8$ in the original fit. The parametrized P_{13} partial wave had a resonance pole at (1787 - 100i) MeV. The CUTKOSKY 76 analysis included data from 0.8 to 2.5 GeV/c. The ACE parametrization (see MILLER 72) was used along with contributions from P , ρ , A_2 , Λ , Σ , and low mass di-pion exchange, the latter being based on the results of ALCOCK 73. An iterative fitting procedure was followed in which energy-independent fitting was alternated with energy smoothing using a parametrization based on partial-wave dispersion relations which included $K\Delta$, K^*N , and $K^*\Delta$ threshold effects for all waves with $J \leq 7/2$. Typical χ^2/DF values were 1.0 to 1.1.

The question of the existence of a Z_1^* in the elastic channel cannot yet be considered settled. The most likely candidate has long been considered to be the P_{13} wave around 1900 MeV. ARNDT 74 found such an effect, but the three other analyses described above which reached high enough energy to observe the effect (ADAMS 73, MARTIN 75, and CUTKOSKY 76) failed to do so. See KELLY 75 for a discussion of the experimental implications of this discrepancy.

Results of the analyses of ADAMS 73 and CUTKOSKY 76 are shown in Figs. 3 and 4. Solution 2 of ADAMS 73 is plotted below 1.0 GeV/c with errors¹ on real and imaginary parts ranging from 3% to 50% for the various waves. The data set used for the results plotted included a preliminary version of the CAMERON 74 data, and the results are similar to those of CAMERON 74. This solution also agrees qualitatively with CUTKOSKY 76 at higher energies; solution 1 disagrees. The results² of CUTKOSKY 76 are plotted from 0.8 to 2.5 GeV/c. These results are quantitatively quite different from those of MARTIN 75 and ARNDT 74.

Two multi-channel partial-wave analyses of the $I=1$ KN , $K\Delta$, and K^*N channels at 1.21, 1.29, 1.38, and 1.69 GeV/c have been performed by LESQUOY 75. One analysis used a quasi-two-body approach, and the other the isobar model approach. It was not possible to draw any firm conclusions about the

Data Card Listings

Baryons

Z*'s

For notation, see key at front of Listings.

K*N channel or to discriminate between results of other analyses of the elastic channel alone. Results on the $K\Delta$ channel indicate four comparably large waves: SD_{11} , PP_{11} , PP_{13} , and either DD_{13} or DD_{15} depending on the method of analysis. This is in contrast to earlier analyses (e.g., GRIFFITH 72) which found this channel to be dominated by the PP_{13} wave over most of the 1.0 to 1.5 GeV/c region. BERTHON 73 find (like GRIFFITH 72) no indication in the $K\Delta$ channel of rapid phase variation signaling possible resonant effects.

I = 0 System

The experiments of CARROLL 73 and BOWEN 73 (see above) measured K^+d total cross sections in the same energy ranges as their K^+p measurements. These experiments agree rather well with each other and with older data except that the CARROLL 73 cross sections are systematically higher than those of BOWEN 70 (0.36 to 0.72 GeV/c). Unfolded I=0 cross sections, new and old, are shown in Fig. 2(a); the disagreement between CARROLL 73 and BOWEN 73 at low energies is primarily due to differences in the unfolding procedure. The I=0 cross section unfolded from the newer K^+p and K^+d data rises to a plateau in the 0.7 to 0.9 GeV/c region and has a hump of about 3 mb in the 0.9 to 1.3 GeV/c region. Older data had indicated a hump of comparable size in the plateau region. As shown in Figs. 1 and 2(b), the plateau is associated with a broad peak in the elastic cross section, while the hump is associated with a rising K^*N cross section.

In a partial-wave analysis of the I=0 elastic channel GIACOMELLI 74 fit data in the 0.38 to 1.51 GeV/c region, including the new K^+n differential cross-section data of the BGRT collaboration (GIACOMELLI 73; 0.64 to 1.51 GeV/c). Three types of partial-wave analysis (all with S, P, and D waves only) were used to analyze these data: an energy-dependent analysis (ED) using the parametrization of LEA 68 and I=1 partial waves from previous solutions of GIACOMELLI 70, ALBROW 71, and LOVELACE 71; an energy-independent analysis (EI-1) which included K^+p data and determined both the I=0 and I=1 partial waves; and an energy-independent analysis (EI-2) in which the I=1 partial waves were fixed at the values of ALBROW 71's

solution γ . In analyses EI-1 and EI-2 several variants of the LOVELACE 71 shortest-path technique were used to link solutions at different energies. The solutions found by these different methods fell into three main families denoted A, C, and D, the best χ^2/ND values being 1.4-1.9, 1.6-1.7, and 1.5-1.8, respectively. Examples of each class from an ED analysis using the I=1 partial waves of GIACOMELLI 70 are shown in Fig. 5. Some criteria for choice between these solutions are offered by comparison with particular features of the data. The single existing measurement of $K^+n \rightarrow K^0p$ polarization (RAY 69; 0.60 GeV/c) is consistent with classes C and D only. The I=0 total cross sections of CARROLL 73 and BOWEN 73 were not included in the fitting, but comparison is made with CARROLL 73 and the data are found to be consistent with class D only. LONDON 74 has calculated $K_{Lp} \rightarrow K_{Sp}$ amplitudes by combining the best (i.e., lowest χ^2) energy-dependent solution from each of GIACOMELLI 74's classes with existing I=1 KN and $\bar{K}N$ solutions and compared the results with $K_{Lp} \rightarrow K_{Sp}$ cross-section data. He finds that the class D solution is preferred. Similar calculations have been made more recently, and with much more regeneration data, by ALEXANDER 75 and CAMERON 75. The former favors class A, while the latter favors classes C and D but is unable to distinguish between them. No conclusive evidence for Z*'s emerges from the analysis of GIACOMELLI 74; the most likely candidate is the class-D P_{01} wave, but GIACOMELLI 74 point out that the resonance interpretation is questionable for some of the class-D solutions and for all of the class-C solutions. A Breit-Wigner plus quadratic background fit to the class-D P_{01} wave yields the following parameters: $M = 1740$ MeV, $\Gamma = 300$ MeV, $x = 0.85$.

A more recent effort is the energy-dependent simultaneous I=0 and I=1 analysis of MARTIN 75, which makes use of the new K^+n elastic and charge-exchange data of DAMERELL 75. The parametrization used is described above. Two qualitatively similar solutions are found. Both are quite different from all the GIACOMELLI 74 solutions, but the P_{01} wave most resembles that of the class-C solution (see Fig. 5). MARTIN 75 find no significant evidence for resonance-like energy dependence.

Baryons
Z*'s

Data Card Listings
For notation, see key at front of Listings.

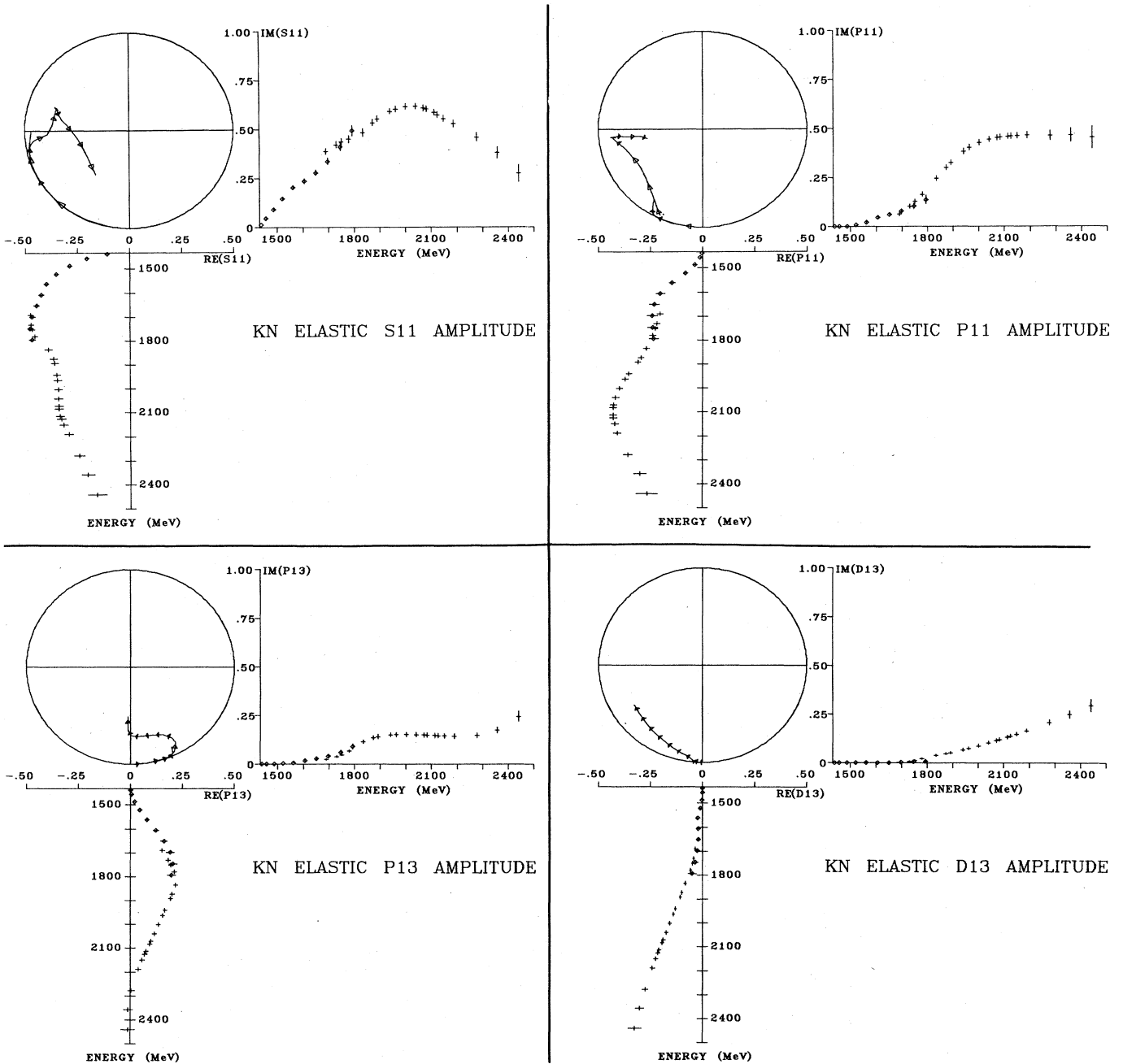


Fig. 3. Amplitudes for $I=1$ KN elastic scattering in the $J=1/2$ and $J=3/2$ waves from ADAMS 73 (ϕ) and CUTKOSKY 76 (\dagger). The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 100 MeV and a base-to-tip length of 20 MeV. All the energy axes run from elastic threshold to 2500 MeV.

Data Card Listings

For notation, see key at front of Listings.

Baryons

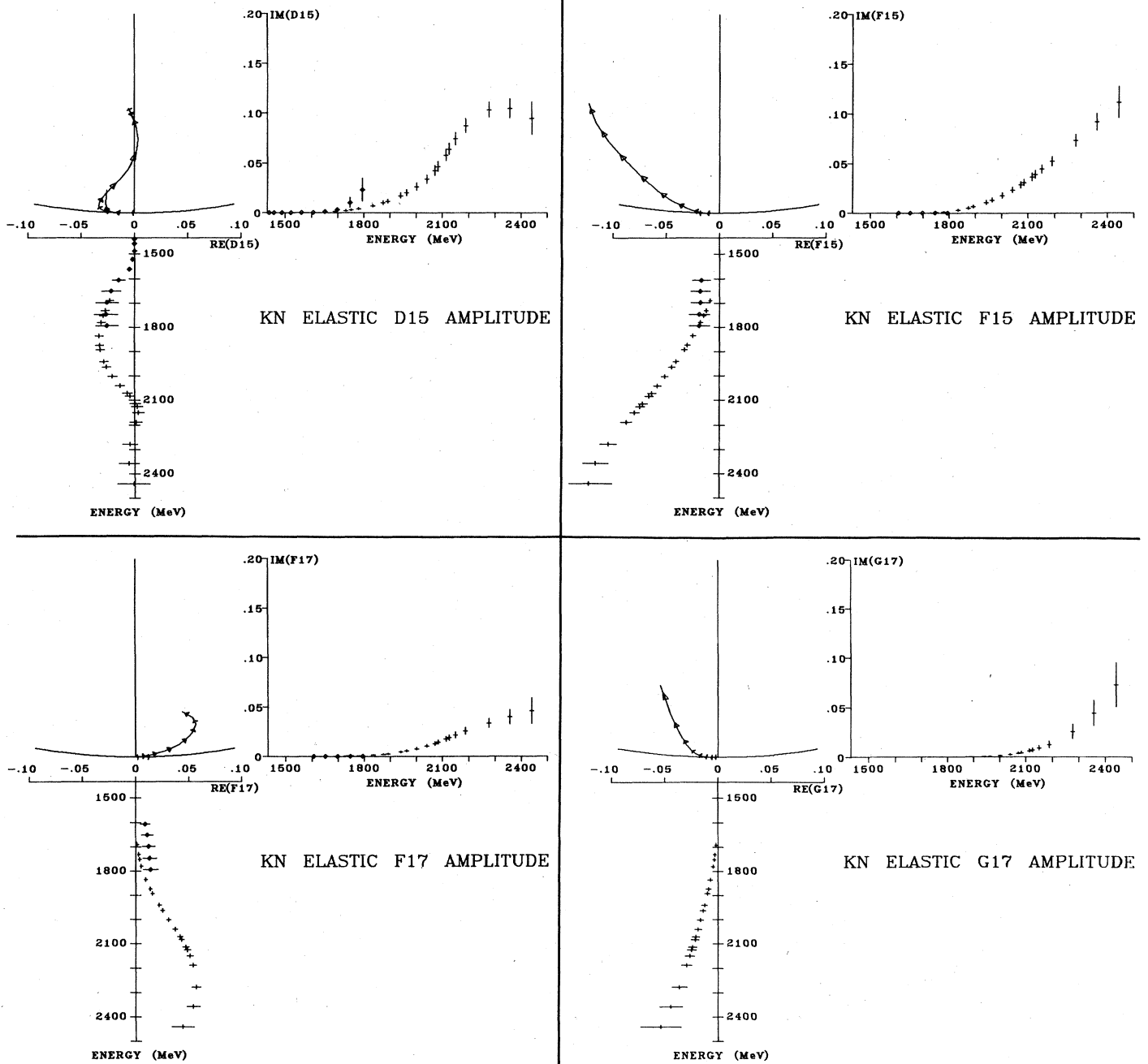
 Z^* 's

Fig. 4. Amplitudes for $I=1$ KN elastic scattering in the $J=5/2$ and $J=7/2$ waves from ADAMS 73 (ϕ) and CUTKOSKY 76 (\dagger). The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows are plotted on the Argand plots with bases positioned at integer multiples of 100 MeV and a base-to-tip length of 20 MeV. All the energy axes run from elastic threshold to 2500 MeV.

Baryons
Z*'s

Data Card Listings
For notation, see key at front of Listings.

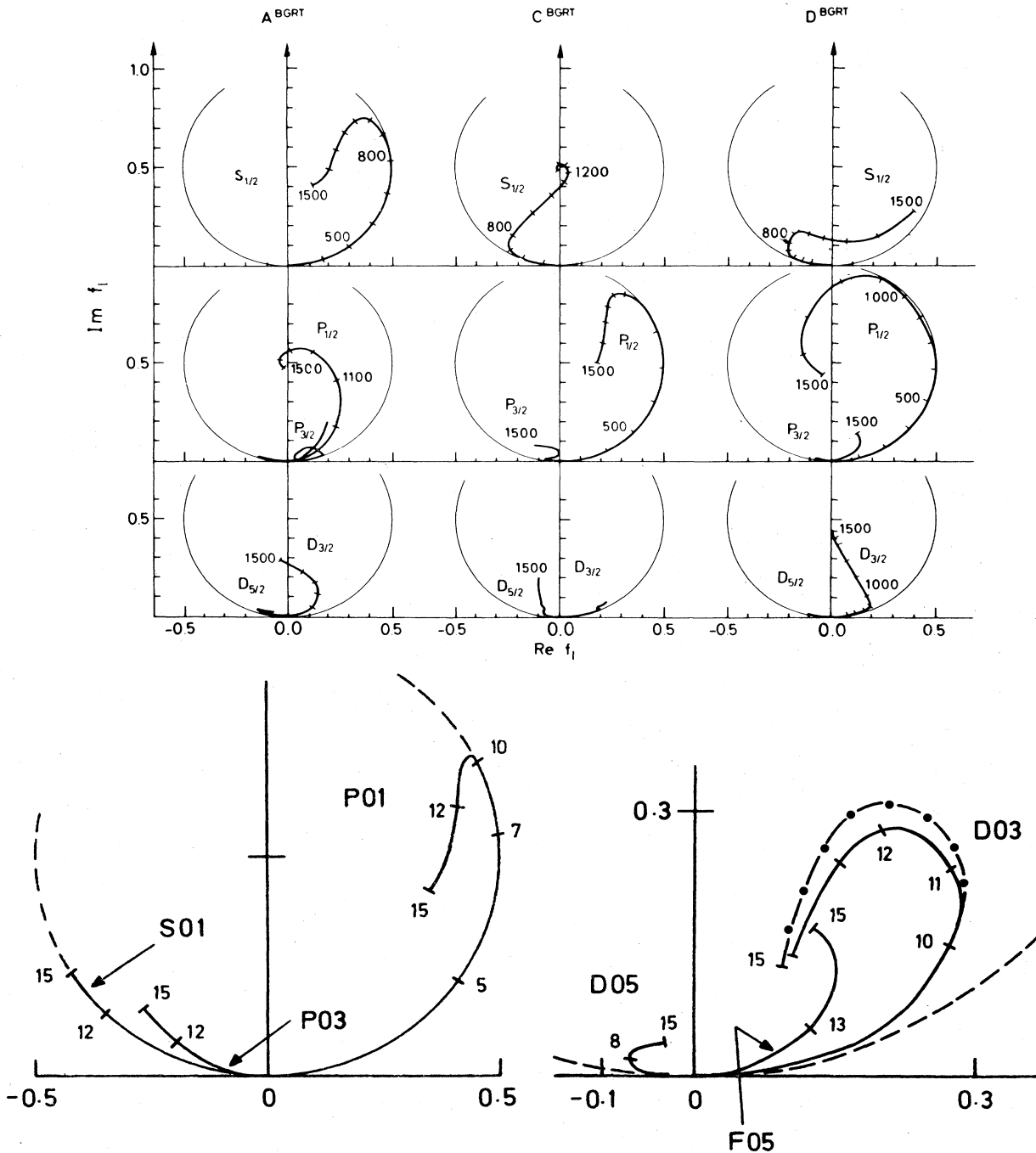


Fig. 5. Amplitudes for $I=0$ KN elastic scattering from GIACOMELLI 74 (upper plots) and MARTIN 75 (lower plots). The GIACOMELLI 74 results are from the energy-dependent analysis and exemplify the three main classes of solutions found in the complete analysis; the energy dependence of the amplitudes is displayed by tick marks at integral multiples of 100 MeV. The MARTIN 75 results are from the preferred solution 2, with the solution-1 D_{03} wave (which exhibited the largest difference between solutions 1 and 2 of any wave) shown as a broken line for comparison; the energy dependence of the amplitudes is displayed by tick marks at the indicated integral multiples of 100 MeV.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Z*'s, Z₀(1780), Z₀(1865), Z₁(1900)

Production Experiments

There are no new developments in production experiments, and the observation of ERNE 70 that present upper limits for Z* production are not small still holds. Cross-section limits for the production of broad Z*'s are comparable to the observed N* and Y* production cross sections.

References

- J. D. Dowell, private communication.
 - R. E. Cutkosky, private communication.
- See the Data Card Listings for other references.

S=1 I=0 EXOTIC STATES (Z₀)

Z₀(1780) 95 Z*0(1780, JP=1/2+) I=0 **P₀₁**

SEE THE MINI-REVIEW PRECEDING THIS LISTING.
WILSON 72 AND GIACOMELLI 74 FIND SOME SOLUTIONS WITH RESONANT-LIKE BEHAVIOR IN THE P₀₁ PARTIAL WAVE. THE EFFECT SEEN IN THE I=0 TOTAL CROSS SECTIONS, IF A RESONANCE, MUST HAVE SPIN=1/2, BECAUSE THE INELASTIC CROSS SECTION IS VERY SMALL AND THE TOTAL CROSS SECTION IS ABOUT 4*PI/K**2.

95 Z*0(1780) MASS (MEV)

M	1780.0	10.0	COOL	70 CNTR	K+P, D TOTAL	1/71
M	D	SEEN	DOWELL	70 CNTR	K+P, D TOTAL	7/70
M	D	SEE ALSO DISCUSSION OF LYNCH 70				7/70
M	W	(1800.)	WILSON	72 PWA	K+N P ₀₁ WAVE	3/72
M	W	ESTIMATE OF PARAMETERS FROM BW + QUADRATIC BACKGROUND FIT TO P ₀₁ .				3/72
M	1	(1750.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 1	9/73
M	1	(1825.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
M	1	FIT 1=FIT OF SINGLE L=1 BW+BACKGROUND TO I=0 TCS FROM .4-1.1 GEV/C				9/73
M	1	FIT 2=FIT OF L=1 AND L=2 BWS TO SAME DATA; SEE Z ₀ (1865) FOR L=2 PART				9/73
M		(1740.)	GIACMEL	74 PWA	.38-1.51 GEV/C	10/74*

95 Z*0(1780) WIDTH (MEV)

W	(565.0)		COOL	70 CNTR	K+P, D TOTAL	1/71
W	(300.)		WILSON	72 PWA	K+N P ₀₁ WAVE	3/72
W	1	(600.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 1	9/73
W	1	(845.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
W		(300.)	GIACMEL	74 PWA	.38-1.51 GEV/C	10/74*

95 Z*0(1780) PARTIAL DECAY MODES

P1	Z*0(1780) INTO K N	DECAY MASSES
		493+ 939
P2	Z*0(1865) INTO N K*(892)	938+ 892

95 Z*0(1780) BRANCHING RATIOS

R1	Z*0(1780) INTO (K N)/TOTAL	(P1)	
R1	(0.95)	COOL	70 CNTR K+P, D TOTAL 1/71
R1	(0.85)	WILSON	72 PWA K+N P ₀₁ WAVE 3/72
R1	(.75)	CARROLL	73 CNTR IF J=1/2; FIT 1 9/73
R1	(.91)	CARROLL	73 CNTR IF J=1/2; FIT 2 9/73
R1	(.85)	GIACMEL	74 PWA .38-1.51 GEV/C 10/74*

REFERENCES FOR Z*0(1780)

COOL	70 DUKE CONF 47	R L COOL	(BNL)
ALSO	69 PL 308 564	ABRAMS, COOL, GIACOMELLI, KYCIA, LI +	(BNL)
ALSC	70 PR D1 1887	COOL, GIACOMELLI, KYCIA, LEONTIC, LI +	(BNL)
DOWELL	70 DUKE 53	J. D. DOWELL	(BIRM)
WILSON	72 NP 842 445	+GRIFFITHS, HIRATA +	(BGNA+GLAS+RCMA+TRST)
CARROLL	73 PL 458 531	+KYCIA, LI, MICHAEL, MOCKETT, RAHM +	(BNL)
GIACMEL	74 NP 871 138	GIACOMELLI +	(BGNA+GLAS+ROMA+TRST) IJP

PAPERS NOT REFERRED TO IN DATA CARDS

LYNCH	70 DUKE 9	G LYNCH (REVIEWER OF CR. SEC. DATA)	(LRL)
HIRATA	71 NP B30 157	+GOLDHABER, HALL, SEEGER, THRILLING, WOHLL (LBL) IJP	
BCWEN	73 PR D7 22	+JENKINS, KALBACH, PETERSEN +	(ARIZ+MICH)
JOHNSON	74 508 343	JOHNSON, VLASSOPULOS	(CFRN, DURH)
CAMERON	75 PALERM CCNF.	+CAPLUPPI +	(BGNA+EDIN+GLAS+PISA+RHEL) IJP

EXPERIMENTS MAINLY ABOUT ELASTIC CHANNELS --
GOLDHABER 62 PRL 9 135 GOLDHABER, CHINDOSKY, GOLDHABER+ (LRL+UCLA) IJP
RAY 69 PR 183 1183 RAY, BURRIS, FISK, KRAEMER, HILL+ (*CFRN+BNL)
ARMITAGE 72 NAL PAPER 391 *ASTON, DUEROOTH, ELLISON, + (MCHS+DARE)
GIACOMEL 72 NP 842 437 GIACOMELLI + (BGNA+GLAS+RCMA+TRST)
GIACOMEL 73 NP 856 346 GIACOMELLI, + (BGNA+GLAS+RCMA+TRST)
ALSO 73 BGNA PPT. AE-73/4 GIACOMELLI, GRIFFITHS, + (BGNA+GLAS+RCMA+TRST)
LONDON 74 PRD 9 1565 LONDON (BNL) IJP
ALEXANDE 75 PL 588 484 ALEXANDER, BAR-NIR, BENARY+ (TELAF+EID) IJP
DAMERELL 75 NP 894 374 *POTCHKISS, WICKENS, BENTLEY+ (RHEL, BIRM)

EXPERIMENTS MAINLY ABOUT INELASTIC CHANNELS --
GIACOMEL 72 NP 837 577 GIACOMELLI + (BGNA+GLAS+ROMA+TRST)

Z₀(1865) 96 Z*0(1865, JP=3/2-) I=0 **D₀₃**
THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K* N THRESHOLD. SEE HIRATA 68 AND 70. WILSON 72 AND GIACOMELLI 73 REPORT PARTIAL WAVE ANALYSES. AARN 73 CLAIMS A RESONANCE IN A MODEL DEPENDENT PWA. SEE ALSO Z*0(1780).

96 Z*0(1865) MASS (MEV)

M	(1860.0)	(15.0)	CARTER	67 THEO	DISPERSION REL.	8/67
M	(1868.0)	(10.0)	COOL	70 CNTR	K+P, D TOTAL	8/67
M	(1830.)		AARON	73 MPWA	I=0 KN .6-1.6G/C	9/73
M	1	(1840.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73
M	1	FIT2=FIT OF L=1 AND L=2 BWS TO I=0 TCS FROM .4-1.1 GEV/C.				9/73
M	1	SEE Z ₀ (1780) FOR FIT 1 AND L=1 PART OF FIT 2.				9/73

96 Z*0(1865) WIDTH (MEV)

W	(200.0)	(50.0)	CARTER	67 THEO		8/67
W	(160.0)	(30.0)	COOL	70 CNTR		8/67
W	(100.)		AARON	73 MPWA	I=0 KN .6-1.6G/C	9/73
W	1	(75.)	CARROLL	73 CNTR	KN I=0 TCS; FIT 2	9/73

96 Z*0(1865) PARTIAL DECAY MODES

P1	Z*0(1865) INTO K N	DECAY MASSES
		493+ 939
P2	Z*0(1865) INTO N K*(892)	938+ 892

96 Z*0(1865) BRANCHING RATIOS

R1	Z*0(1865) INTO (K N)/TOTAL	(P1)	
R1	(.155)	(.025)	CARTER 67 THEO IF J=3/2 9/73
R1	(.115)	(.025)	COOL 70 CNTR IF J=3/2 9/73
R1	1	(.085)	CARROLL 73 CNTR IF J=3/2; FIT 2 9/73
R2	Z*0(1865) INTO N K*(892)	(P2)	
R2		MAIN INELASTIC DECAY	HIRATA .68 HBC 11/68

REFERENCES FOR Z*0(1865)

CARTER	67 PRL 18 801	A A CARTER	(CAVENDISH)
HIRATA	68 PRL 21 1485	HIRATA, WOHLL, GOLDHABER, TRILLING	(LRL)
COOL	70 PR D1 1887	COOL, GIACOMELLI, KYCIA, LEONTIC, LI +	(BNL)
ALSO	66 PRL 17 102	+GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, +	(BNL) I
ALSC	69 PL 308 564	ABRAMS, COOL, GIACOMELLI, KYCIA, LI +	(BNL)
AARON	73 PRD 7 1401	AARON, RICH, HOGAN, SRIVASTAVA	(LASL+NEAS) IJP
CARROLL	73 PL 458 531	+KYCIA, LI, MICHAEL, MOCKETT, RAHM +	(BNL) I

PAPERS NOT REFERRED TO IN DATA CARDS

HIRATA	70 DUKE 429	+GOLDHABER, SEEGER, TRILLING, WOHLL	(LRL)
AARON	71 PRL 26 407	+AMADO, SILBAR	(NEAS, PENN, LASL) IJP
HIRATA-1	71 NP B33 445	+GOLDHABER, HALL, SEEGER, TRILLING, WOHLL (LBL)	
GIACOMEL	72 NP 837 577	GIACOMELLI +	(BGNA+GLAS+ROMA+TRST)
WILSON	72 NP 842 445	+GRIFFITHS, HIRATA +	(BGNA+GLAS+ROMA+TRST)

S=1 I=1 EXOTIC STATES (Z₁)

Z₁(1900) 97 Z*1(1900, JP=3/2+) I=1 **P₁₃**

THIS EFFECT IS STRONGLY ASSOCIATED WITH THE K-DELTA THRESHOLD. SEE THE MINIREVIEW PRECEDING Z*0(1780)

97 Z*1(1900) MASS (MEV)

M	1	(1922.0)	AYED	70 IPWA	P13, SOL. I	6/70
M	1	(1899.0)	AYED	70 IPWA	P13, SOL. II	6/70
M	1	(2030.0)	AYED	70 IPWA	S11, SOL. III	6/70
M	1	THREE SCALNS IN ORDER OF DECREASING SIGNIFICANCE, THOUGH AYED 70				
M	1	GIVE PARAMETERS, THEY CONCLUDE RESONANT INTERPRETATION DOUBTFUL.				
M	2	(1830.)	BARNETT	70 IPWA	P13, SCALN III	9/73
M	2	RESONANCE SIGNAL BARELY ABOVE BACKGROUND DUE TO THE LARGE ERRORS				
M	2	IN THE AMPLITUDES RESULTING FROM THE ANALYSIS				
M		1900.0	COOL	70 CNTR	K+P TOTAL	1/71
M		(1880.)	10.0	ALSBRO	71 IPWA ** SOL. GAMMA	10/71
M	K	(1890.)	KATO	71 IPWA	SOL II (FIT BW)	10/71
M	K	(2040.)	KATO	71 IPWA	SOL II (FIT BW)	10/71
M	K	KATO 71 ESTIMATE RESONANCE PARAMETERS -- UPDATED PHASE SHIFTS				3/72
M	K	PUBLISHED IN MILLER 72.				3/72

Baryons

Z₁(1900), Z₁(2150), Z₁(2500)

Data Card Listings

For notation, see key at front of Listings.

Table with columns for particle name, mass, width, and various parameters. Includes entries for Z1(1900) WIDTH (MEV).

Table with columns for particle name, real part of pole position, and parameters. Includes entries for Z1*(1900) REAL PART OF POLE POSITION.

Table with columns for particle name, imaginary part of pole position, and parameters. Includes entries for Z1*(1900) IMAGINARY PART OF POLE POSITION.

Table with columns for particle name, partial decay modes, and decay masses. Includes entries for Z1(1900) PARTIAL DECAY MODES.

Table with columns for particle name, branching ratios, and dispersion relations. Includes entries for Z1(1900) BRANCHING RATIOS.

Table with columns for particle name, main inelastic decay, and parameters. Includes entries for Z1(1900) MAIN INELASTIC DECAY.

Table with columns for particle name, references for Z1(1900), and various experimental data points.

Table with columns for particle name, total-cross-section experiments, and various data points.

Table with columns for particle name, a k-matrix analysis of some of the early k+p data, and various data points.

Table with columns for particle name, theoretical and model dependent analyses, and various data points.

Table with columns for particle name, experiments mainly about inelastic channels, and various data points.

Table with columns for particle name, the main elastic scattering and polarization experiments, and various data points.

Table with columns for particle name, various parameters, and references. Includes entries for BARNETT, CHARLES, CAMERON, YUTA, ABE, ADAMS, PATTON.

Table with columns for particle name, partial-wave analyses, and various data points.

Table with columns for particle name, earlier analyses that do not include recent polarization data, and various data points.

Table with columns for particle name, latest review talks and papers, and various data points.

Complex block containing a box labeled 'Z1(2150) BUMPS' with an arrow pointing to the right, and text describing a small bump in total cross section at PK=1.8 GEV/C.

Table with columns for particle name, Z1(2150) mass (MEV), and various data points.

Table with columns for particle name, Z1(2150) width (MEV), and various data points.

Table with columns for particle name, Z1(2150) partial decay modes, and various data points.

Table with columns for particle name, Z1(2150) branching ratios, and various data points.

Table with columns for particle name, references for Z1(2150), and various data points.

Complex block containing a box labeled 'Z1(2500) BUMPS' with an arrow pointing to the right, and text describing a small bump in total cross section at PK=2.7 GEV/C.

Table with columns for particle name, Z1(2500) mass (MEV), and various data points.

Table with columns for particle name, Z1(2500) width (MEV), and various data points.

Table with columns for particle name, Z1(2500) partial decay modes, and various data points.

Table with columns for particle name, Z1(2500) branching ratios, and various data points.

Table with columns for particle name, references for Z1(2500), and various data points.

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's**Z₁ CROSS SECTION LIMITS**

SEE MINTREVIEW PRECEDING Z*0

CS	UNITS	MICROBARN	50.		BASSOMPIE 68 HBC	K+P TO Z** PI+	10/69
CS	A	LESS THAN	-2	+3	ANDERSON 69 ASPK +	PI-P TO K-Z**	10/69
CS	A	ABOVE LIMIT FOR	M=1.2 TO 1.4	GEV -	CL= 99 P.C.		
CS	B	LESS THAN	1.4	+1.9	ANDERSON 69 ASPK +	PI-P TO K-Z**	10/69
CS	B	ABOVE LIMIT FOR	M=1.5 TO 2.5	GEV			

REFERENCES FOR Z*1 CROSS SECTION LIMITS

BASSOMPI 68 PL 278 468
ANDERSON 69 PL 298 136

BASSOMPIERRE, + (CERN, BRUXELLES)
+BLESER, BLIEDEN, COLLINS, + (BNL, CARNEGIE)

PAPERS NOT REFERRED TO IN DATA CARDS

TYSON 67 PRL 19 255 +GREENBERG, HUGHES, LU, MINEHART, MCRI, (YALE)
MORI 68 PL 288 152 +GREENBERG, HUGHES, LU, ROTHBERG, + (YALE)
MORI 69 PR 185 1687 +GREENBERG, HUGHES, LU, MINEHART, + (YALE)
MORI 69 REPLACES TYSON 67 AND MORI 68.

Note on Λ 's and Σ 's

The number of confirmed Y^* states has not increased greatly in the last few years, although there has been a larger increase in the number of proposed, but unconfirmed, possible new resonances. Since our last edition (1974) we have made only one addition to the Y^* portion of the Baryon Table, the $\Lambda(1860)$, but have entered four new states into the Data Card Listings, and encountered several further indications that some of the states we list may really be more than one resonance. New high-statistics experiments are needed to clarify the situation.

Just as the recently discovered N^* 's are only weakly coupled to the $\pi N \rightarrow \pi N$ reaction, so also are the recently proposed Y^* 's only weakly coupled to the $\bar{K}N \rightarrow \bar{K}N$, $\bar{K}N \rightarrow \Lambda\pi$, and $\bar{K}N \rightarrow \Sigma\pi$ reactions. For this reason the newer Y^* 's are more difficult to uncover; in invariant mass distributions they usually appear as small peaks or make no appearance at all. Rather, when the 2-body reactions are partial-wave analyzed, some of the amplitudes are found to traverse resonance-like counterclockwise circles. The results of partial-wave analysis give J^P information, whereas a peak seen in an invariant mass distribution or a total cross section often cannot be analyzed for its quantum numbers. We will keep information coming from formation experiments and from production experiments separate whenever necessary.

Formation Experiments

Partial-wave analyses have been performed on $\bar{K}N$, $\Lambda\pi$, $\Sigma\pi$, EK , and $\Lambda\omega$, plus some quasi-two-body channels. Given the present accuracy of the data it is not possible to perform a completely energy-independent analysis, that is, solve for the partial-wave amplitudes at each energy in a model-independent way.

For the great majority of the analyses done so far, the energy range covered was rather limited, corresponding usually to a single bubble chamber experiment. A disturbing feature that appears when examining the partial waves obtained in a particular analysis is that they do not join smoothly with the partial waves given in analyses done for the same channel in a different energy range. One way to avoid this lack of continuity is to analyze the data in as large an energy range as possible. This has been done by six analyses which are described below and three of which are illustrated in Figs. 1-11. See the Data Card Listings for information on other analyses covering limited energy ranges.

Four of the analyses described below (RLIC 76, LEA 73, LANGBEIN 72, and KIM 71) have attempted a multi-channel approach using data on the three 2-body reactions $\bar{K}N$, $\Lambda\pi$, and $\Sigma\pi$, with a fictitious channel sometimes being introduced to account for the global effect of the remaining final states. The latter have large cross sections, so the present multi-channel analyses do not really impose any more stringent unitarity constraints than those already contained in the single-channel fits. However, there is an advantage in the determination of resonance parameters (masses, widths, and branching fractions) since they are fit simultaneously to data in all three channels. The other two analyses (BAILLON 75 and VAN HORN 75) deal only with the $\Lambda\pi$ channel which is, in principle, the easiest to handle as only one isospin has to be considered and both the differential cross section and baryon polarization are measured simultaneously in bubble chamber experiments. These two analyses cover the mass range 1540-2200 MeV and make use of the Legendre polynomial expansion of the angular distributions.

Baryons

Λ 's and Σ 's

a) The most recent analysis is the work done by the Rutherford Laboratory-Imperial College collaboration, RLIC 76. Here the $\bar{K}N$, $\Sigma\pi$, and $\Lambda\pi$ channels have been considered in the mass range 1480-2170 MeV. It is a conventional energy-dependent analysis, the presence of a resonance in a partial wave being detected by comparing the goodness of the fit when this wave is parametrized as a smooth background to the alternative fit when a Breit-Wigner is added to the background. The data used have been carefully selected in order to eliminate inconsistencies (usually the older and statistically less accurate points have been rejected). Angular distributions were directly used in the fit except when the quality of the data was such that no loss of information occurred by using Legendre coefficients (e.g., $K^-p \rightarrow \Sigma^0\pi^0$). Internal consistency is introduced by requiring that the mass and width of the resonances be the same for each of the three channels; a "weighted average" of the three values has been done when necessary and used as a fixed parameter in the final fit. Argand diagrams from the final solution are given in Figs. 1-11. Some suspected resonances are confirmed by this analysis, but many other reported "resonant effects" are not found and new possible resonances are proposed. The situation, in particular for the low partial waves, is still very confused even for masses as low as ~ 1600 MeV.

b) The most recent $\Lambda\pi$ analysis, BAILLON 75, is a semi-energy-independent analysis. The D_{15} and F_{17} partial waves were constrained to lie near the well known $D_{15}(1765)$ and $F_{17}(2030)$ resonances, with their generally accepted parameters, in the energy region where they are important. This constraint was found necessary in order to reduce the number of possible solutions and to fix the arbitrary phase. The latter is undetermined, as in all the inelastic channels, and if unconstrained would make it difficult to join amplitudes at neighboring energies. Above a mass of 1750 MeV only one solution is found. Below this mass, two possible solutions are proposed; one of them (solution 1) has a much greater number of "new resonant structures" and is somewhat preferred on the basis of the behavior of the S and P waves for very low masses

Data Card Listings

For notation, see key at front of Listings.

(around 1580 MeV). Using the method of Barrelet-zeroes^{1,2} it is found that only these 2 solutions are possible if resonances are required to exist in the D_{15} and F_{17} waves. To extract the resonance parameters, the amplitudes were fitted to a combination of Breit-Wigner's plus background. The parameters corresponding to both solutions are given in the data cards, and Argand diagrams are shown in Figs. 1-6.

c) VAN HORN 75 has done an energy-dependent fit of the $\Lambda\pi$ channel. In addition to the best fit, VAN HORN 75 also presents all the other ambiguous solutions that can be generated by the method of Barrelet-zeroes. Among the "ambiguous" solutions, seven of them are found to preserve the established resonances $D_{13}(1670)$, $D_{15}(1765)$, $F_{15}(1915)$, and $F_{17}(2030)$, but with couplings to the $\Lambda\pi$ channel which are sometimes very different from their generally accepted values (cf., BAILLON 75). Also, new resonant structures appear in all the waves, particularly in the lower spins. This analysis is instructive in so far as indicating that there could be an entire constellation of other possible resonances beyond those which appear at the primitive stage of the analysis. However, one should keep in mind that these possible resonances do not necessarily correspond to a simple or even plausible parametrization of the amplitudes. (The values we have listed on the data cards correspond to the average of the 20 best original fits of this analysis, all containing the established, plus a few probable, resonances).

d) The analysis of LEA 73 is a multi-channel energy-dependent partial-wave analysis with parametrized K-matrix elements. The momentum range covered is 440 to 1190 MeV/c, and 99 parameters were used in the fit. Established resonances were constrained to have parameters near their generally accepted values, taken from a previous edition of these tables (for this reason we do not list the results of LEA 73 for these resonances in the following data cards). New resonances are identified using poles of the K-matrix. The procedure used to get the resonance parameters is not described, the values of the fitted K-matrix elements are not given, and no explicit comparison of the results with alternative parametrizations is made.

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's

e) LANGBEIN 72 performed single-energy fits at 40 momenta between 436 and 1226 MeV/c. This work is the nearest approximation to an energy-independent partial-wave analysis existing for the $S=-1$ system. The partial waves corresponding to established resonances were, however, constrained to have Breit-Wigner forms. Approximately 90 acceptable single-energy fits per energy were generated and were used in shortest-path searches. Several candidates for acceptable shortest paths were generated and a preferred path was chosen by rejecting those that failed to reproduce known resonance behavior. Resonances in this solution were identified by loops in Argand diagrams correlated with peaks in the ≥ 3 -body final-state cross section. Resonance parameters were then extracted by fitting Breit-Wigner's with both multiplicative and additive background. A disturbing characteristic of this work is that the amplitudes which were not specifically chosen to be resonant (sometimes even those, as in the case $D_{15}(1765) \rightarrow \Lambda\pi$) show an erratic behavior which cannot be justified simply by the fluctuation of the data.

f) KIM 71 fit data from threshold to 1226 MeV/c using the Ross and Shaw³ effective-range expansion of the inverse multi-channel K-matrix. The data in each of seven energy intervals bounded by 0, 534, 658, 806, 916, 1022, 1117, and 1226 MeV/c, were fit with a constant effective-range matrix. Only the $F_{15}(1915)$ was fixed to a Breit-Wigner form, all other waves included being parametrized by the K-matrix formalism. Resonances were identified by a method involving the appearance of loops in the Argand diagram, peaks in the speed plot, and poles of the K-matrix, the exact procedure not being reported.

None of the three older multi-channel fits above extend beyond a mass of 1.9 GeV (the highest-energy data from the only high-statistics experiment existing before 1971, ARMENTEROS 68). These analyses do not join smoothly with either older (LITCHFIELD 71) or more recent (HEMINGWAY 75) partial-wave analyses at higher energies. The three more recent analyses extend to 2.2 GeV, and for masses higher than 2.2 GeV, the only existing work is an energy-dependent analysis of the $\Lambda\pi$ and $\Sigma\pi$ channels for momenta between 1.93 and

2.5 GeV/c by the Saclay-College de France collaboration (BELLEFON1 75 and BELLEFON 76).

Production Experiments

This type of experiment is often difficult to analyze. Information on $I=0$ states is possible only when there is no $I=1$ state at a similar mass. The main controversies at the present time concern resonances in the 1600 to 1700 MeV region. See the mini-reviews on $\Sigma(1620)$ and $\Sigma(1670)$ in these Listings.

Figures

Argand plots of 15 $S=-1$ partial waves are shown in Figs. 1-11. The analyses shown were picked largely for illustrative purposes rather than on the basis of our judgment of their quality; as discussed above there are a number of analyses extant and no clear choice of the "best" ones is possible.

Errors on Masses and Widths

Often the quoted errors in partial-wave analyses are only statistical, and the values of masses and widths can change by more than these errors when a new parametrization is used. For this reason we report the values of M , Γ , and x_i obtained by different authors even if they analyze the same data. The spread of these masses and widths is certainly a better estimate of the uncertainties than the statistical errors. Sometimes the errors quoted are obtained by the inspection of various fits done with different hypotheses (see, for example, BERTHON 70, GALTIERI 70, VAN HORN 75, RLIC 76). For three states, $\Lambda(1520)$, $\Lambda(1815)$, and $\Sigma(1765)$, there are enough data available to perform an overall fit of the various x_i of the type discussed in the main text (Sec. VII B). In this case we are forced to use the errors, however small they may be, but we warn the reader that the final errors are not to be taken seriously.

In conclusion, in the Baryon Table we choose not to give errors on masses and total widths determined primarily by partial-wave analyses, but, whenever necessary, to show a range of values. As for the branching ratios, we use the errors when needed to perform an overall fit, but we caution the reader.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

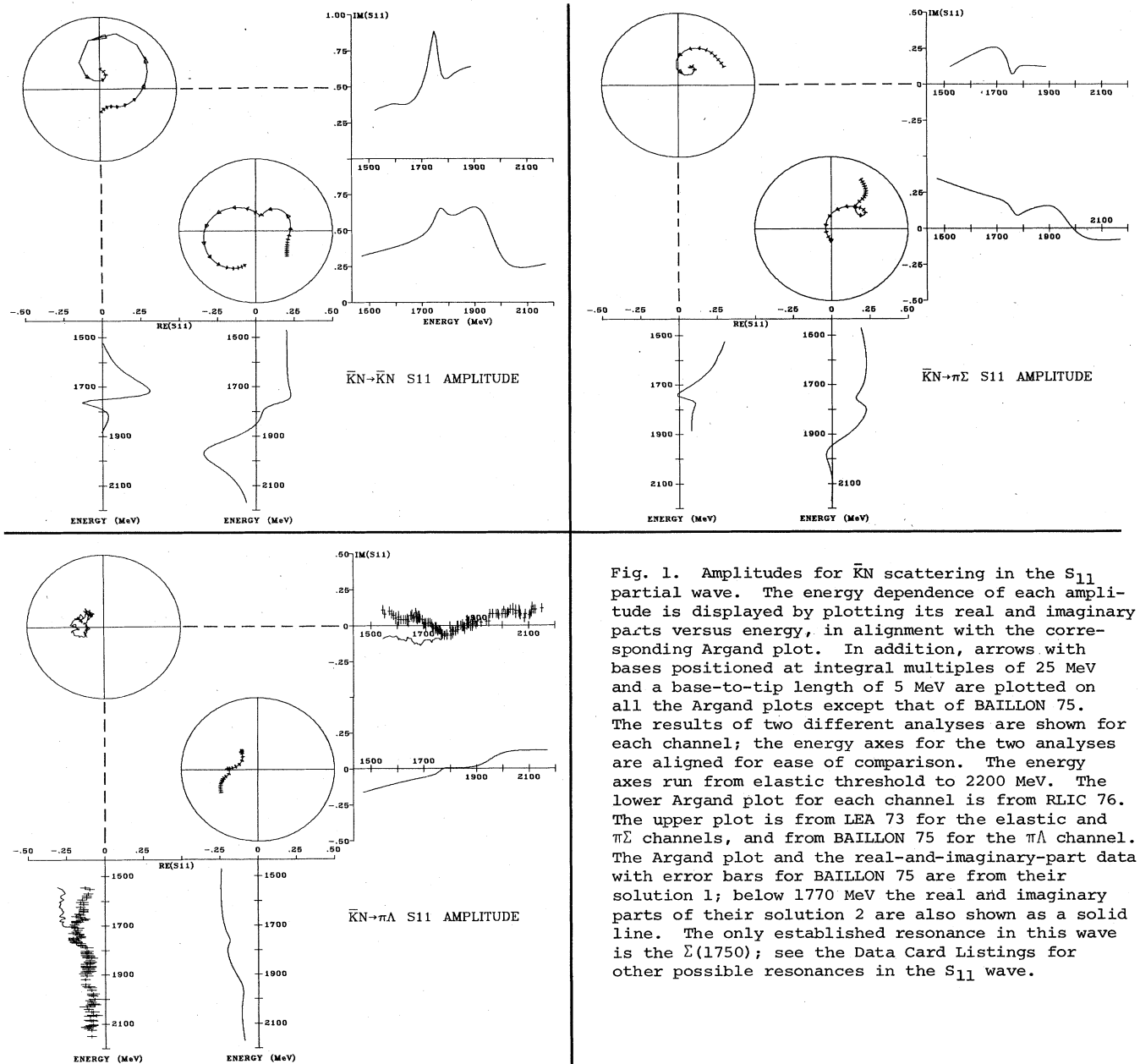


Fig. 1. Amplitudes for $\bar{K}N$ scattering in the S_{11} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. The only established resonance in this wave is the $\Sigma(1750)$; see the Data Card Listings for other possible resonances in the S_{11} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Λ 's and Σ 's

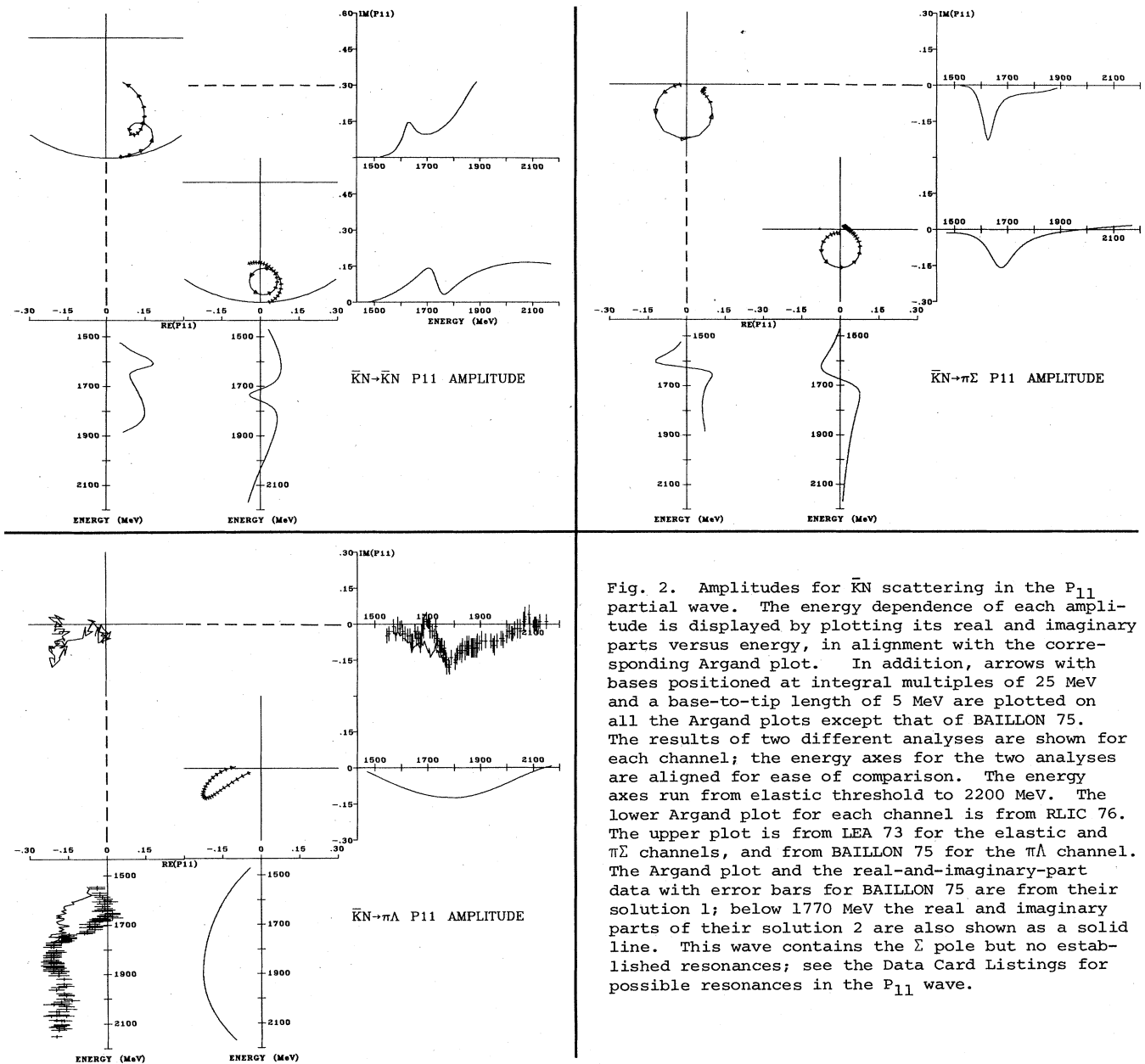


Fig. 2. Amplitudes for $\bar{K}N$ scattering in the P_{11} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. This wave contains the Σ pole but no established resonances; see the Data Card Listings for possible resonances in the P_{11} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

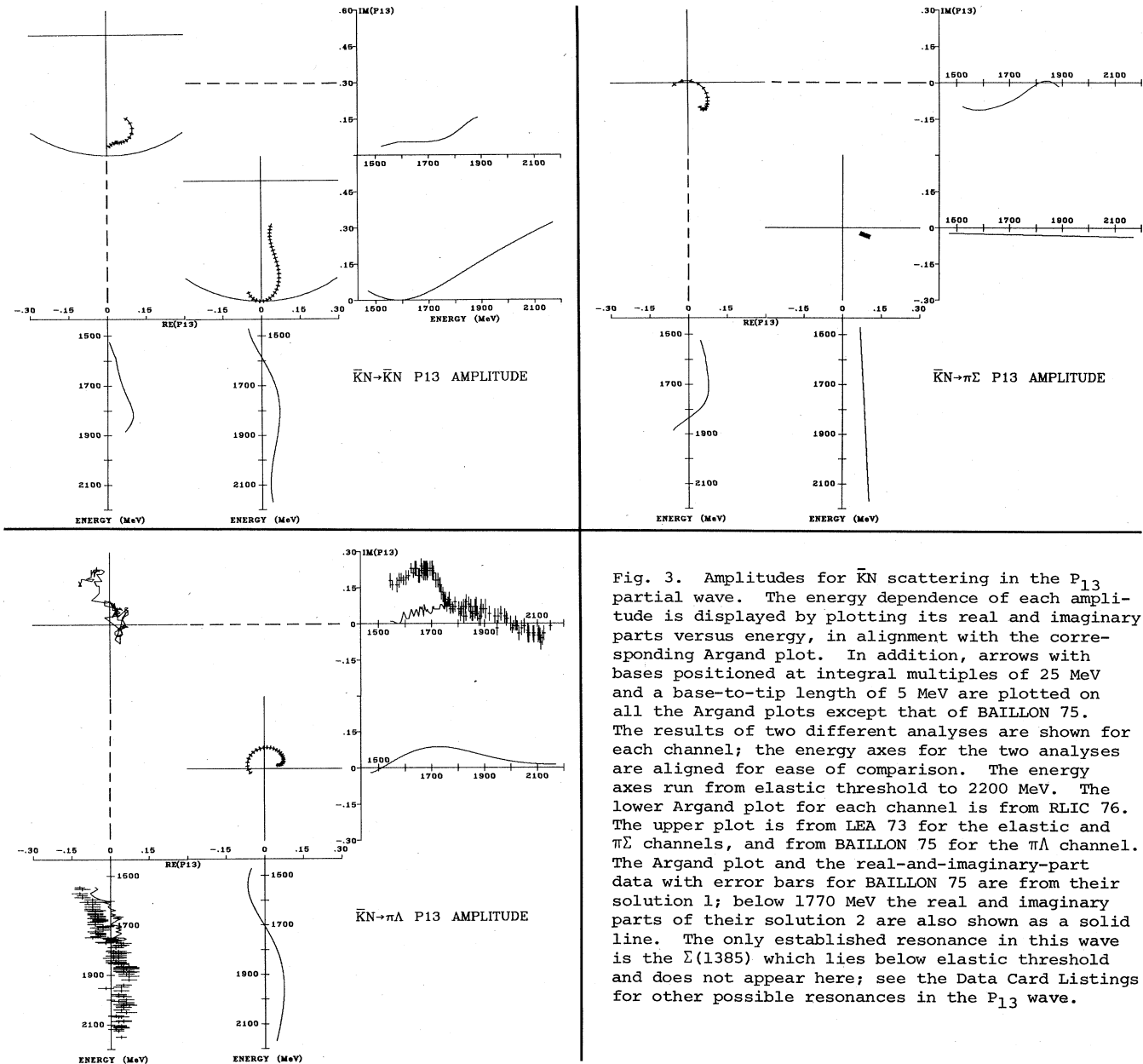


Fig. 3. Amplitudes for $\bar{K}N$ scattering in the P_{13} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line. The only established resonance in this wave is the $\Sigma(1385)$ which lies below elastic threshold and does not appear here; see the Data Card Listings for other possible resonances in the P_{13} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

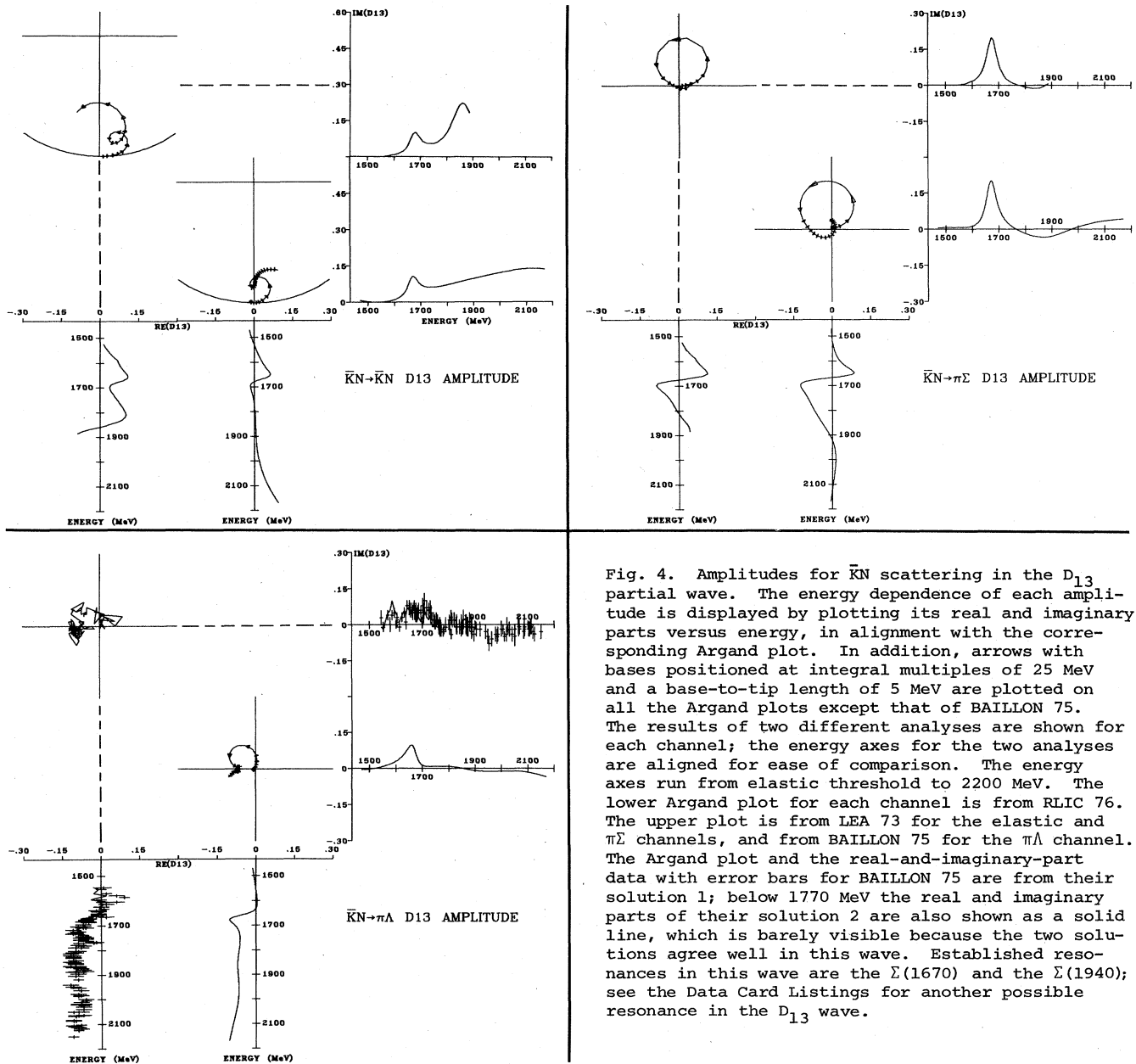
 Λ 's and Σ 's

Fig. 4. Amplitudes for $\bar{K}N$ scattering in the D_{13} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line, which is barely visible because the two solutions agree well in this wave. Established resonances in this wave are the $\Sigma(1670)$ and the $\Sigma(1940)$; see the Data Card Listings for another possible resonance in the D_{13} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

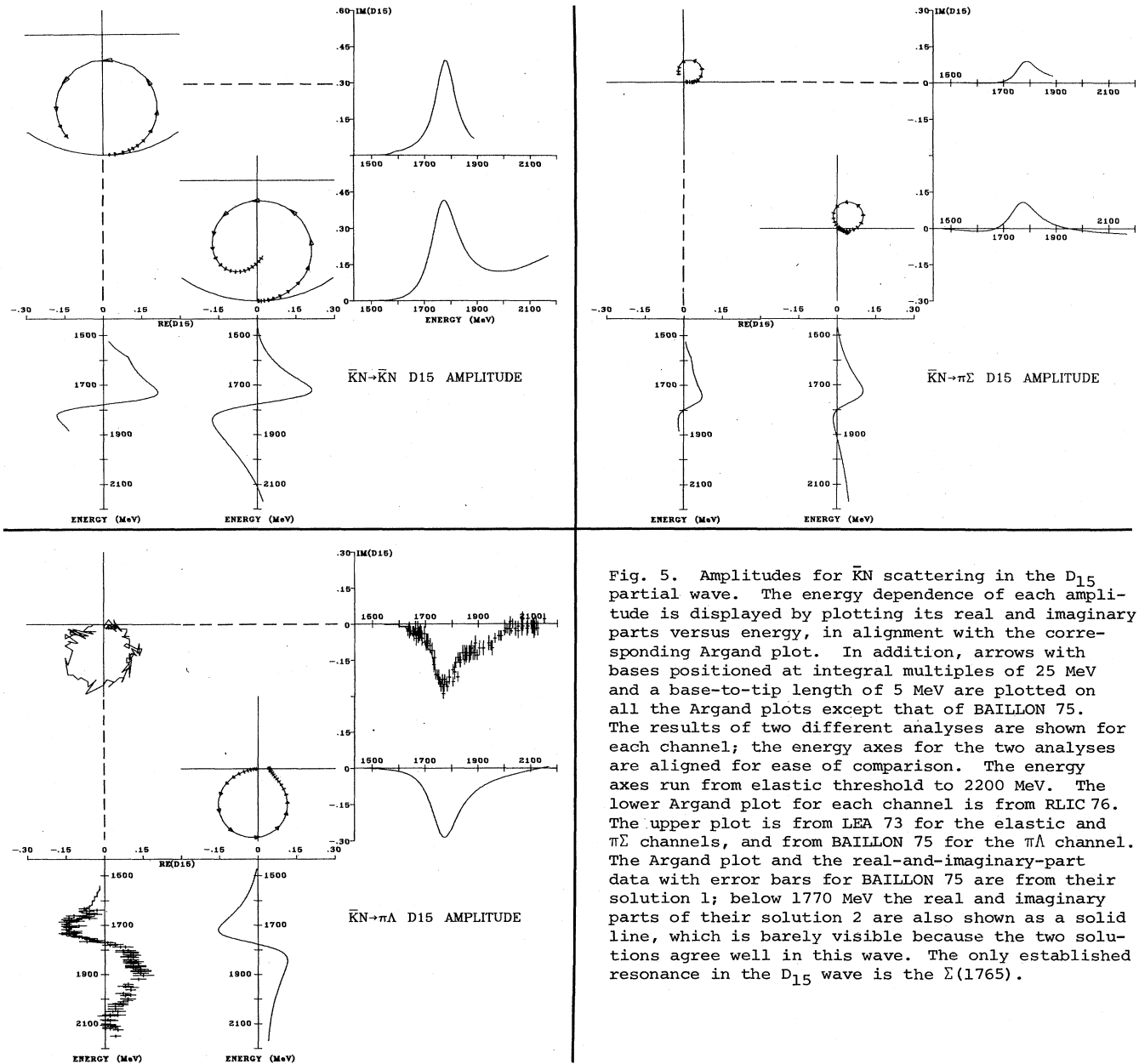


Fig. 5. Amplitudes for $\bar{K}N$ scattering in the D_{15} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The Argand plot and the real-and-imaginary-part data with error bars for BAILLON 75 are from their solution 1; below 1770 MeV the real and imaginary parts of their solution 2 are also shown as a solid line, which is barely visible because the two solutions agree well in this wave. The only established resonance in the D_{15} wave is the $\Sigma(1765)$.

Data Card Listings

For notation, see key at front of Listings.

Baryons
 Λ 's and Σ 's

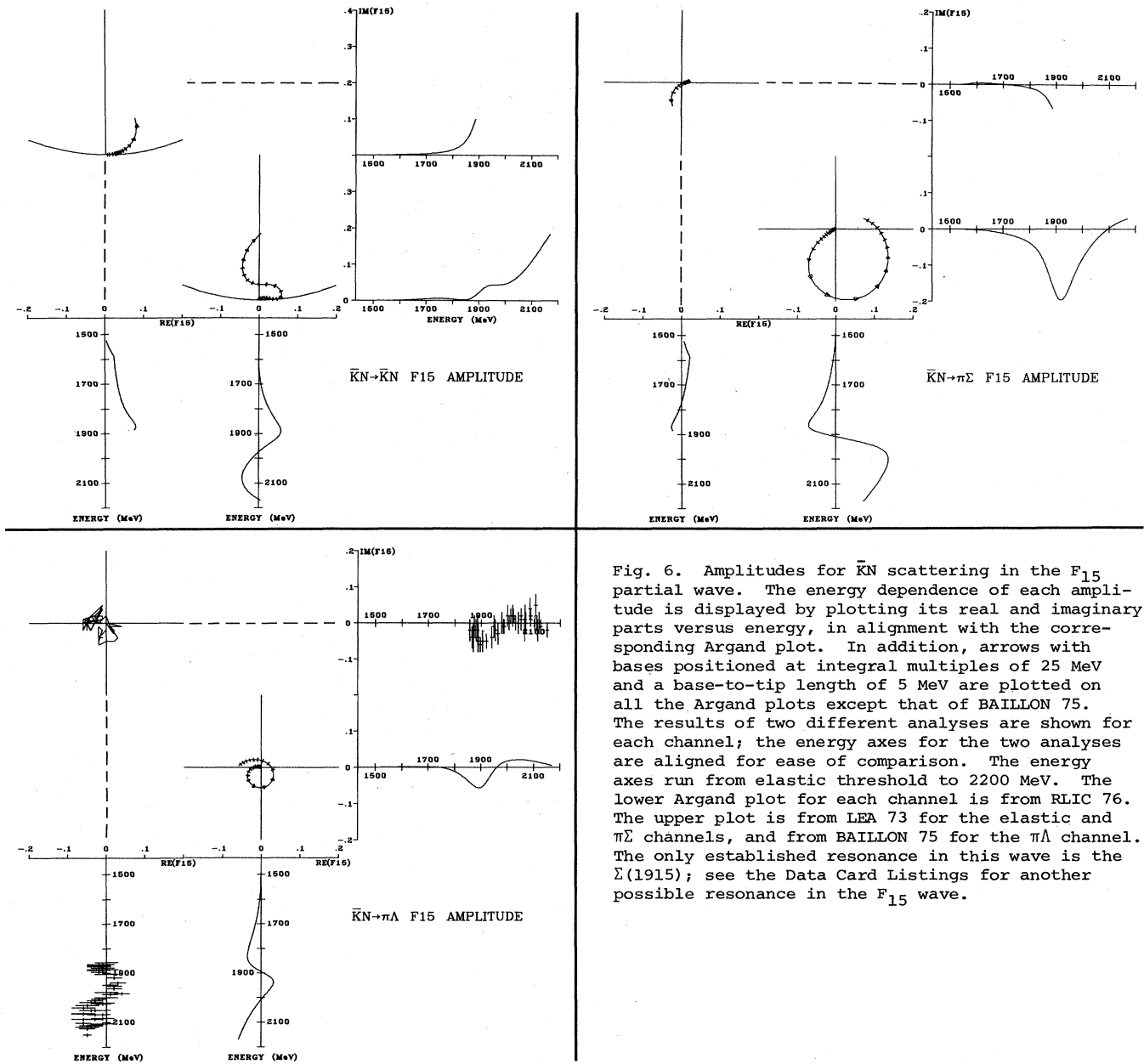


Fig. 6. Amplitudes for $\bar{K}N$ scattering in the F_{15} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots except that of BAILLON 75. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each channel is from RLIC 76. The upper plot is from LEA 73 for the elastic and $\pi\Sigma$ channels, and from BAILLON 75 for the $\pi\Lambda$ channel. The only established resonance in this wave is the $\Sigma(1915)$; see the Data Card Listings for another possible resonance in the F_{15} wave.

Baryons
 Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

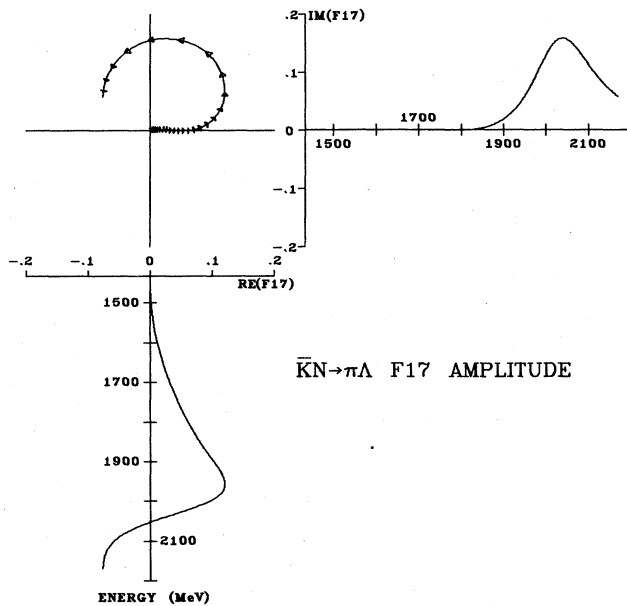
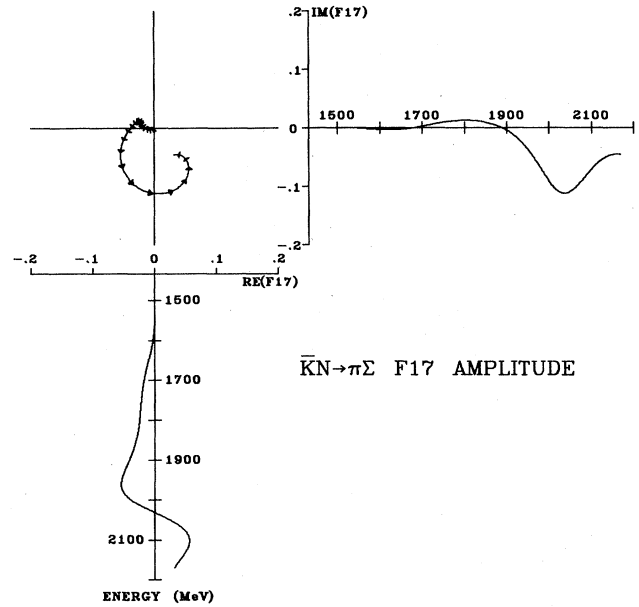
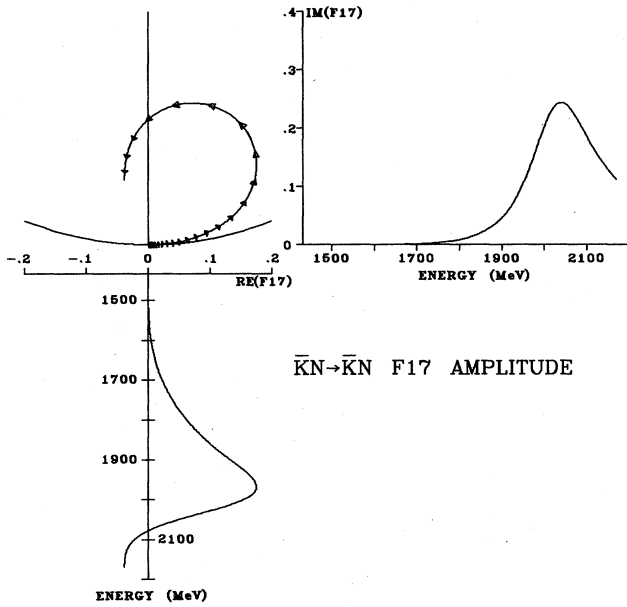


Fig. 7. Amplitudes for $\bar{K}N$ scattering in the F_{17} partial wave. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The energy axes run from elastic threshold to 2200 MeV. All the results shown are from RLIC 76. The only established resonance in the F_{17} wave is the $\Sigma(2030)$.

Data Card Listings

For notation, see key at front of Listings.

Baryons

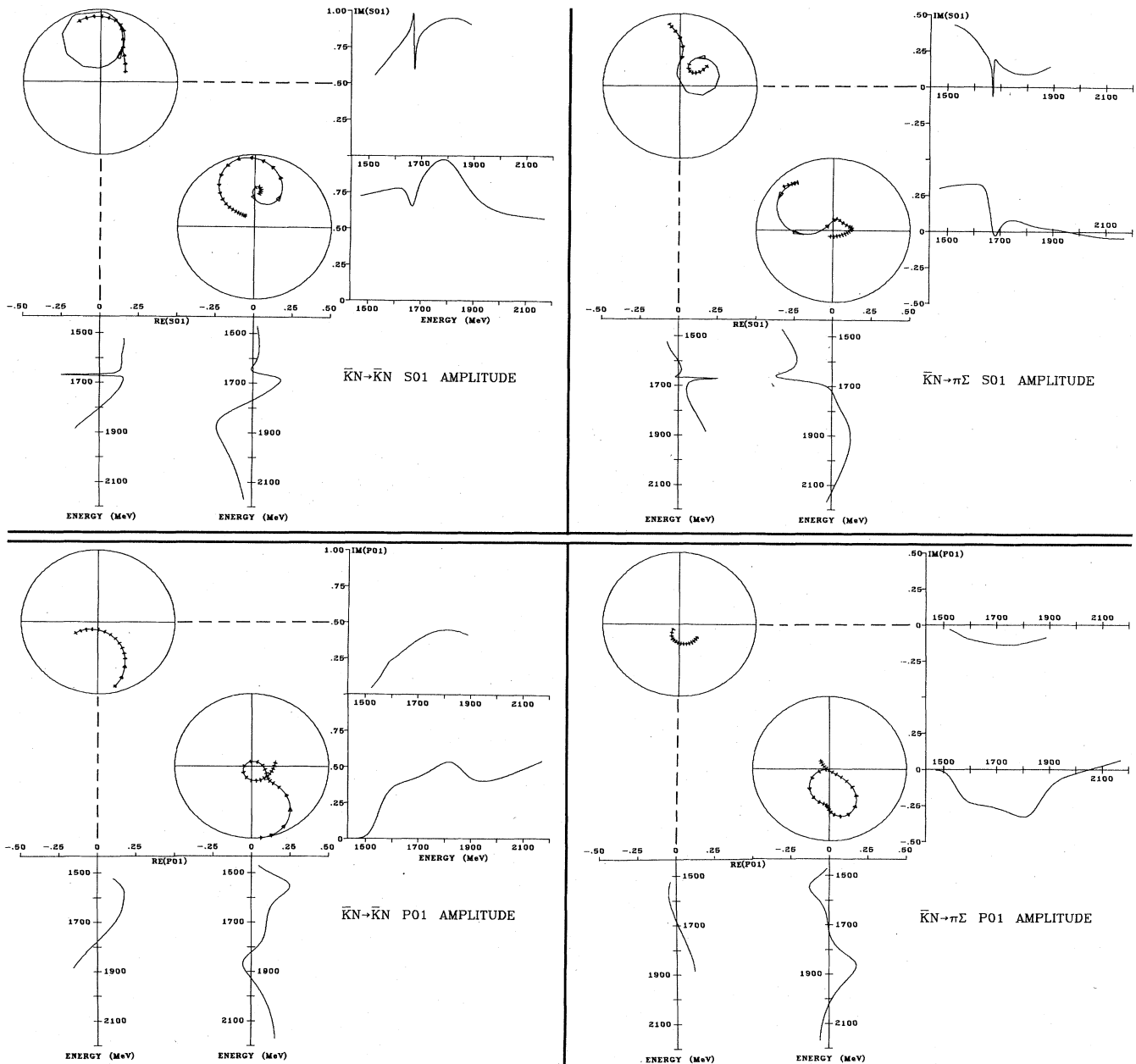
 Λ 's and Σ 's

Fig. 8. Amplitudes for $\bar{K}N$ scattering in the S_{01} and P_{01} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in the S_{01} wave are the $\Lambda(1405)$ (which lies below threshold and does not appear here), the $\Lambda(1670)$, and the $\Lambda(1870)$; the P_{01} wave contains the Λ pole but no established resonances. See the Data Card Listings for possible resonances in the P_{01} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

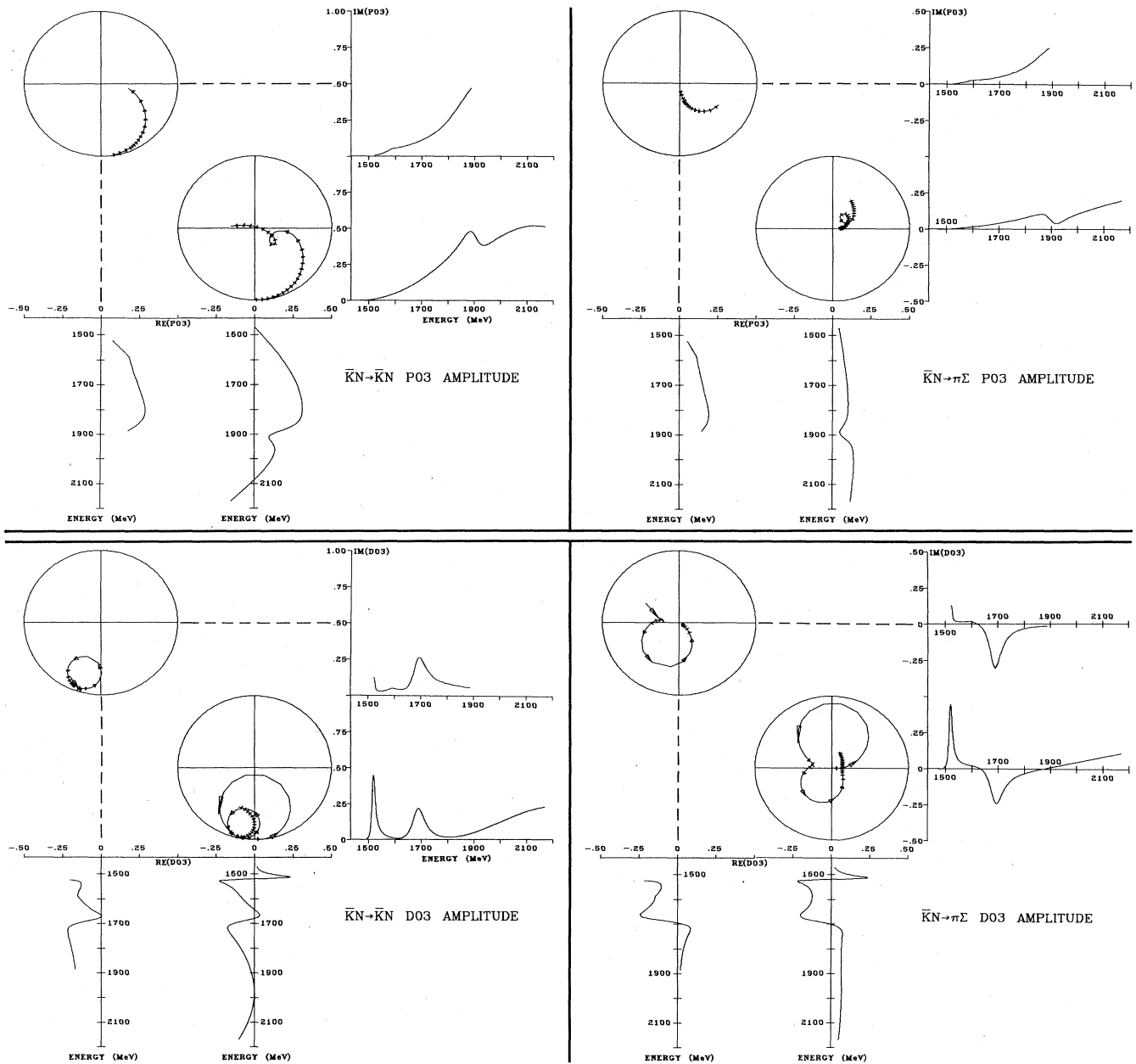


Fig. 9. Amplitudes for $\bar{K}N$ scattering in the P_{03} and D_{03} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in these waves are the $\Lambda(1860)$ in the P_{03} wave and the $\Lambda(1520)$ and $\Lambda(1690)$ in the D_{03} wave. See the Data Card Listings for another possible resonance in the D_{03} wave.

Data Card Listings

For notation, see key at front of Listings.

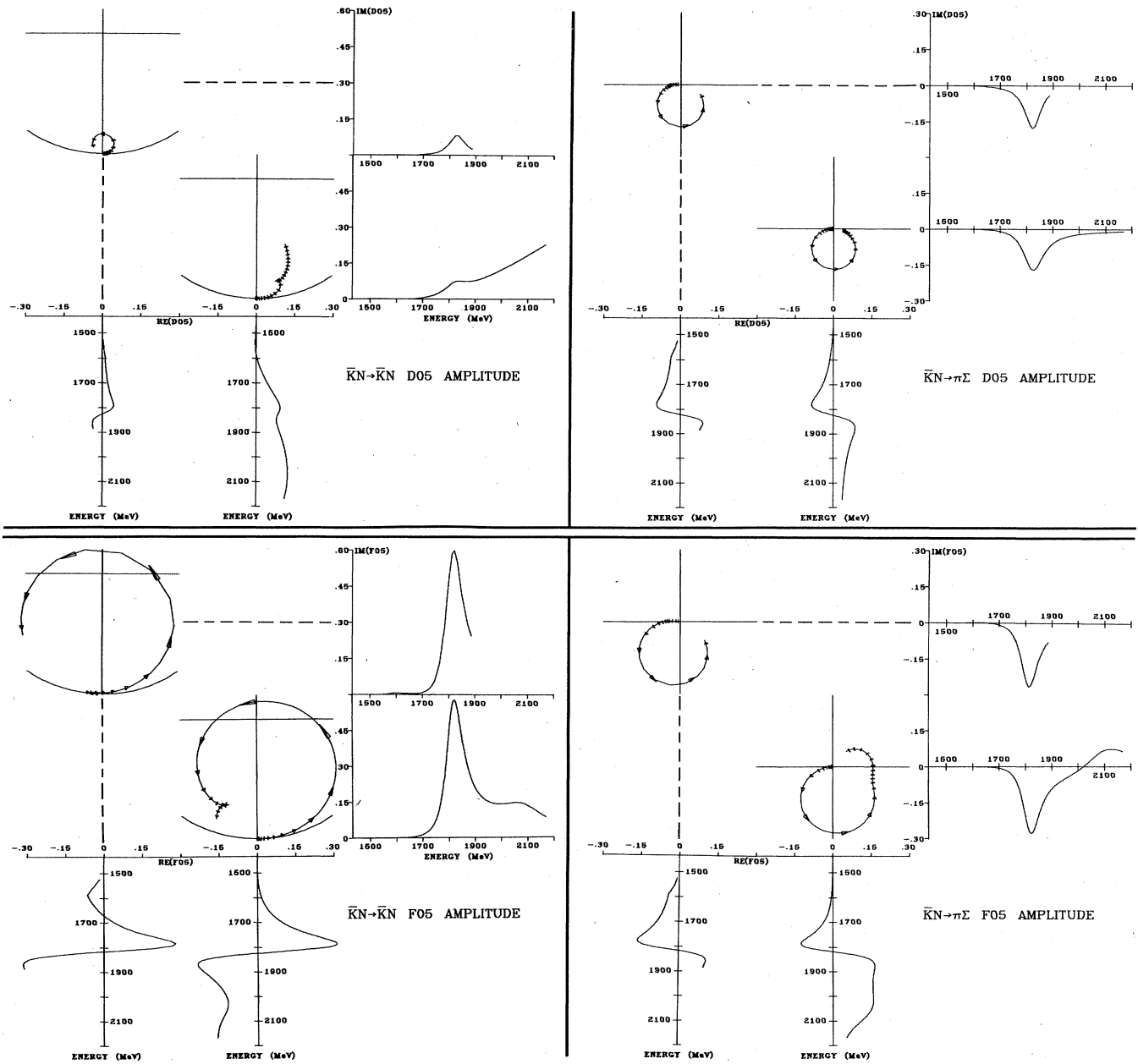
Baryons
 Λ 's and Σ 's

Fig. 10. Amplitudes for $\bar{K}N$ scattering in the D_{05} and F_{05} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The results of two different analyses are shown for each channel; the energy axes for the two analyses are aligned for ease of comparison. The energy axes run from elastic threshold to 2200 MeV. The lower Argand plot for each amplitude is from RLIC 76; the upper plot is from LEA 73. Established resonances in these waves are the $\Lambda(1830)$ in the D_{05} wave and the $\Lambda(1815)$ in the F_{05} wave. See the Data Card Listings for another possible resonance in the F_{05} wave.

Baryons

Λ 's and Σ 's

Data Card Listings

For notation, see key at front of Listings.

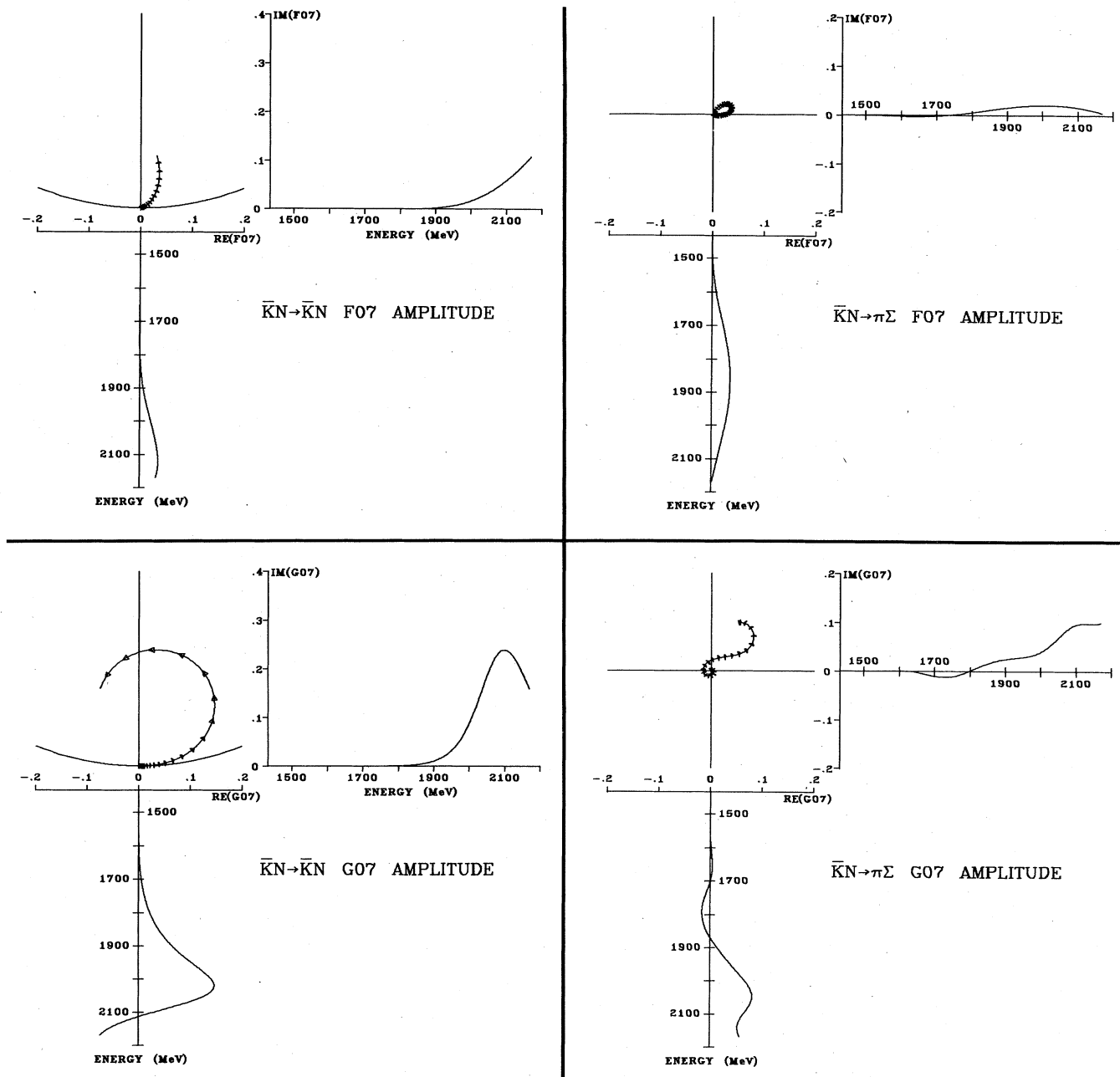


Fig. 11. Amplitudes for $\bar{K}N$ scattering in the F_{07} and G_{07} partial waves. The energy dependence of each amplitude is displayed by plotting its real and imaginary parts versus energy, in alignment with the corresponding Argand plot. In addition, arrows with bases positioned at integral multiples of 25 MeV and a base-to-tip length of 5 MeV are plotted on all the Argand plots. The energy axes run from elastic threshold to 2200 MeV. All the results shown are from RLIC 76. The only established resonance in these waves is the $\Lambda(2100)$ in the G_{07} wave. See the Data Card Listings for a possible resonance in the F_{07} wave.

Data Card Listings

For notation, see key at front of Listings.

Baryons

Λ 's and Σ 's, Λ , $\Lambda(1330)$, $\Lambda(1405)$

Conclusions

Table I is an attempt to evaluate the status of the various Y^* 's. The evaluations are of course partly subjective. A blank indicates that there is no corresponding evidence at all. This may mean either that the relevant couplings are small or that the resonance does not really exist. The Baryon Table includes only the well-established resonances. It seems clear, however, that whereas any particular one of the questionable resonances may disappear with the next analysis, there are probably many new resonances underlying those already established.

References

1. E. Barrelet, Nuovo Cimento **8A**, 331 (1972).
2. P. Litchfield, Proceedings of the IInd Aix-en-Provence Conference (1973).
3. M. Ross and G. Shaw, Ann. Phys. (N.Y.) **13**, 147 (1961).

For other references see the Data Card Listings.

TABLE I. STATUS OF Y^* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --					OTHER CHANNELS
			TOTAL CR. SEC.	KBAR N	LAM PI	SIG PI		
LAM(1115) P01		****						WEAK TO N PI
LAM(1330)		DEAD						
LAM(1405) S01		****		****	F	****		LAM2PI, LAM GAM
LAM(1520) D03		****	****	****	D	****		
LAM(1600) P01		*			R	****		LAM ETA
LAM(1670) S01		****	****	****	B	****		LAM2PI, SIG2PI
LAM(1690) D03		****	****	****	I	****		
LAM(1800) P01		**						
LAM(1800) G09		*						
LAM(1815) F05		****	****	****	D	****		SIG(1385) PI
LAM(1830) D05		****	****	****	E	****		
LAM(1860) P03		***	**	****				
LAM(1870) S01		**		****	N	****		
LAM(2010)		**			F	****		LAM CMG
LAM(2020) F07		*			O	****		
LAM(2100) G07		****	****	****	R	****		XI K, LAM CMG
LAM(2110) F05		**		****	B	****		LAM CMG
LAM(2350)		****	****	****	I	****		
LAM(2355)		***	**	*	D	****		
SIG(1190) P11		****						WEAK TO N PI
SIG(1385) P13		****				****		
SIG(1440) PE		DEAD						
SIG(1480) PE		*		*		*		
SIG(1580) D13		**	**	**		**		
SIG(1620) S11		**	**	**		**		
SIG(1660) P11		**	**	**		**		
SIG(1670) D13		****	**	****	****	****		SEVERAL OTHERS
SIG(1670) PE		**	**	**	**	**		SEVERAL OTHERS
SIG(1690) PE		**	*	**	**	**		LAM 2-PI
SIG(1750) S11		****	****	****	*	****		SIG ETA
SIG(1765) D15		****	****	****	****	****		SEVERAL OTHERS
SIG(1770) P11		*	*	*	*	*		
SIG(1840) P13		*	*	*	*	*		
SIG(1880) P11		**	**	**	**	**		
SIG(1915) F15		****	***	****	****	****		
SIG(1940) D13		****	****	****	****	****		
SIG(2000) S11		*	*	*	*	*		
SIG(2030) F17		****	****	****	****	**		XI K
SIG(2070) F15		*	*	*	*	*		
SIG(2080) P13		**	**	**	**	**		
SIG(2100) G17		*	*	*	*	*		
SIG(2250)		****	****	*	*	*		
SIG(2455)		****	****	*	*	*		
SIG(2620)		****	****	*	*	*		
SIG(3000)		**	**	*	*	*		

**** GOOD, CLEAR, AND UNMISTAKABLE.
 *** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
 ** NEEDS CONFIRMATION.
 * WEAK.
 * ATTRIBUTED TO THE STATE CLOSEST TO WHERE THE CROSS SECTION PEAKS.

S=-1 I=0 HYPERON STATES (Λ)

Λ 18 LAMBDA(1115, JP=1/2+) I=0
 SEE STABLE PARTICLE DATA CARD LISTINGS

$\Lambda(1330)$ BUMPS 87 $Y^*(1330, JP=)$ I=0 PRODUCTION EXPERIMENTS
 SEE THE MINI-REVUE AT THE START OF THE Y^* LISTINGS.
 A PEAK WAS SEEN NEAR 1330 MEV IN THE LAMBDA GAMMA SPECTRUM IN THREE PI-PROPANE EXPERIMENTS (YUNG-CHANG 64, BUBELEV 67, AND BOZOKI 68). ALL MORE RECENT RESULTS INDICATE THAT THERE IS NO RESONANCE NEAR THIS MASS VALUE.

REFERENCES FOR $Y^*(1330)$ (PROD. EXP.)
 Y-CHANG 64 DUBNA CONF I 615 YUNG-CHANG, IN, KLADNITSKAYA, + (DUBNA)
 BUBELEV 67 PL 248 246 +CHADRAA, CHUVILO, + (JINR, BUCHAREST, CERN)
 DAHL 67 PR 163 1377 CAHL, HARDY, FESS, KIRZ, MILLER (LRL)
 BOZOKI 68 PL 288 360 +FENYVES, GEMESY, + (BUDAPEST, DUBNA)
 TAN 69 PRL 23 101 T H TAN (SLAC)
 MAYEUR 70 PL 338 441 +VAN BINST, WILQUET++ (BRUX, CERN, TUFT)
 COLAS 75 NP 891 253 COLAS, FARWELL, FERRER, SIX (DRSA)

$\Lambda(1405)$ 37 $Y^*(1405, JP=1/2-)$ I=0 PRODUCTION EXPERIMENTS **S₀₁**
 THIS RESONANCE CAN BE IDENTIFIED WITH THE VIRTUAL BOUND STATE IN THE KBAR-N SYSTEM FOUND IN THE ANALYSIS OF LOW ENERGY K-P INTERACTION. WE LIST SUCH EXPERIMENTS SEPARATELY BELOW. WE USE ONLY PRODUCTION EXPERIMENTS FOR AVERAGING OF MASSES AND WIDTHS.

37 $Y^*(1405)$ MASS (MEV) (PROD. EXP.)

M	(1405.0)	ALSTON 61 HBC	K-P 1.15 BEV/C
M	(1410.0)	ALEXANDER 62 HBC	PI-P 2.1 BEV/C
M	(1405.0)	ALSTON 62 HBC	K-P 1.2-5 BEV/C
M	(1382.0)	(8.0) ENGLER 65 HBC	PI-P, PI+D 1.6B 7/66
M	1400.0	24.0 MUSGRAVE 65 HBC	PBAR P 3-4 BEV/C 7/66
M	67 1400.0	5.0 BIRMINGHAM 66 HBC	K-P 3.5 9/67
M	120 1405.0	5.0 GALTIERI 68 HBC	K-D 2.1-2.7BEV/C 6/68
M	AVG 1402.4	3.5	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
M	STUDENT 1402.4	3.9	AVERAGE USING STUDENT 10(H/1.11) -- SEE TEXT

37 $Y^*(1405)$ WIDTH (MEV) (PROD. EXP.)

W	(20.0)	ALSTON 61 HBC	7/66
W	35.0	5.0 ALEXANDER 62 HBC	
W	(50.0)	ALSTON 62 HBC	
W	(89.0)	(20.0) ENGLER 65 HBC	7/66
W	60.0	20.0 MUSGRAVE 65 HBC	7/66
W	67 50.0	10.0 BIRMINGHAM 66 HBC	K-P 3.5 9/67
W	120 35.0	8.0 GALTIERI 68 HBC	K-C 2.1-2.7BEV/C 6/68
W	AVG 38.1	3.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.0)
W	STUDENT 37.9	4.3	AVERAGE USING STUDENT 10(H/1.11) -- SEE TEXT

37 $Y^*(1405)$ PARTIAL DECAY MODES (PROD. EXP.)

PI	$Y^*(1405)$ INTO SIGMA PI	DECAY MASSES
		1197+139
REFERENCES FOR $Y^*(1405)$ (PROD. EXP.)		
ALSTON 61 PRL 6 698	+ALVAREZ, EBERHARD, GODD, GRAZIANG, + (LFL) I	
ALEXANDE 62 PRL 8 447	ALEXANDER, KALBFLEISCH, MILLER, SMITH (LRL) I	
ALSTON 62 CERN CCNF 311	+ALVAREZ, FERRO-LUZZI, ROSENFELD, + (LRL) I	
ENGLER 65 PRL 15 224	+FISK, KRAEMER, MELTZER, WESTGARD, + (CERN, BNL) IJ	
MUSGRAVE 65 NC 35 735	+PETMEZAS, + (BIRM, CERN, EPOL, LOIC, SACLAY)	
BIRMINGHAM 66 PR 152 1148	BIRMINGHAM, GLASGOW, LOIC, CXFORD, RUTH-ERFORD	
GALTIERI 68 PRL 21 573	BARBARO-GALTIERI, CHADWICK + (LRL, SLAC)	

Baryons

$\Lambda(1405)$, $\Lambda(1520)$

Data Card Listings

For notation, see key at front of Listings.

1405 MEV REGION: EXTRAPOLATIONS BELOW THRESHOLD

24 Y*(1405, JP=1/2-) I=0
EXTRAPOLATION BELOW THRESHOLD

SEE NOTE IN Y*(1405) PRODUCTION EXPERIMENTS -THE DIFFICULTIES IN EXTRAPOLATING FROM THE PHYSICAL REGION TO THE RESONANCE LOCATION ARE DISCUSSED BY DALITZ 67.

THE QUESTION ON WHETHER Y*(1405) IS A KBAR-N BOUND STATE OR A CDD POLE (DALITZ 70, RAJASEKARAN 72 HAS BEEN INVESTIGATED BY CLINE 71, MARTIN 71, GALTIERI 72, AND COBSON 72. THE LAST TWO PAPERS CONCLUDE THAT THE DATA CANNOT TELL THE DIFFERENCE.

24 Y*(1405) MASS (MEV)

Table with columns for mass values, uncertainties, and references. Includes entries for KIM, SAKITT, KITTEL, MARTIN, CHAO, and RUVUE.

24 Y*(1405) WIDTH (MEV)

Table with columns for width values, uncertainties, and references. Includes entries for KIM, SAKITT, KITTEL, MARTIN, CHAO, and RUVUE.

REFERENCES FOR Y*(1405) (FROM EXTRAPOLATIONS)

Table listing references for Y*(1405) with author names and journal information.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards, including authors like ABRAMS, DONALD, KADYK, and others.

$\Lambda(1520)$

D03

38 Y*(1520, JP=3/2-) I=0

PRODUCTION AND FORMATION EXPERIMENTS AGREE QUITE WELL WITH EACH OTHER, THEREFORE, THEY HAVE NOT BEEN SEPARATED FOR THIS PARTICLE.

THE DECAY MODE LAMBDA PI PI IS LARGELY DUE TO Y*(1385) PI. ONLY THE VALUES OF Y*(1385) PI/(LAMBDA PI PI) GIVEN BY MAST 72 AND CORDEN 75 ARE BASED ON REAL 3-BODY PARTIAL WAVE ANALYSES (THE OLDER RESULTS BEING OBTAINED USING CRUDE METHODS). THE DISCREPANCY BETWEEN THE 2 RESULTS IS ESSENTIALLY DUE TO THE DIFFERENT HYPOTHESIS MADE CONCERNING THE SHAPE OF THE EPSILON MESON.

38 Y*(1520) MASS (MEV)

Table with columns for mass values, uncertainties, and references. Includes entries for GALTIERI, WATSON, ALMEIDA, MUSGRAVE, BIRMINGHA, BURKHARDT, KIM, CORDEN, and RLIC.

38 Y*(1520) WIDTH (MEV)

Table with columns for width values, uncertainties, and references. Includes entries for WATSON, MUSGRAVE, BIRMINGHA, DAHL, BURKHARDT, KIM, CORDEN, and RLIC.

38 Y*(1520) PARTIAL DECAY MODES

Table listing partial decay modes for Y*(1520) into various baryon-meson pairs like KBAR N, SIGMA PI, LAMBDA PI PI, etc.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i ± δP_i, where δP_i = √(δP_{ij}δP_{ji}), while the off-diagonal elements are the normalized correlation coefficients (δP_{ij}δP_{ji})/(δP_iδP_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table showing fitted partial decay mode branching fractions and correlation coefficients for various decay channels.

38 Y*(1520) BRANCHING RATIOS

Table listing branching ratios for Y*(1520) into various channels like SIGMA PI/(KBAR N), LAMBDA PI PI/(KBAR N), etc.

Table listing branching ratios for Y*(1520) into various channels like LAMBDA PI PI/(KBAR N), LAMBDA PI PI/(SIGMA PI), etc.

Table listing branching ratios for Y*(1520) into various channels like SIGMA PI/(LAMBDA PI PI), SIGMA PI/(SIGMA PI), etc.

Table listing branching ratios for Y*(1520) into various channels like LAMBDA GAMMA/TOTAL (PERCENT), etc.

Table listing branching ratios for Y*(1520) into various channels like SIGMA GAMMA/TOTAL (PERCENT), etc.

Table listing branching ratios for Y*(1520) into various channels like KBAR N/TOTAL, etc.

Table listing branching ratios for Y*(1520) into various channels like SIGMA PI/TOTAL, etc.

Table listing branching ratios for Y*(1520) into various channels like SIGMA PI/(SIGMA PI), etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1520)$, $\Lambda(1600)$, $\Lambda(1670)$

R9 Y*0(1520) TO Y*1(1385) PI TO LM PI PI/LM PI PI (P8)/(P3) 9/73
R9 MORE THAN 0.10 CLINE 69 DBC K-D TO 2PI LAM N 9/69
R9 B 0.35 0.10 BURKHARDT 71 HBC LAM. 3PI PRD. 3/71
R9 C (1.05) CHAN 72 IPWA K-P TO LAM 2PI 2/73
R9 M 0.82 0.10 MAST 73 IPWA K-P TO 2PI LAM 12/72
R9 .58 .22 CORDEN 75 DBC K-D 1.4-1.8GV/C 4/75*

R10 Y*0(1520) INTO (Y*1(1385) PI)/TOTAL (P7) 3/71
R10 0.041 0.005 CHAN 72 HBC K-P TO LAM 2PI 3/71
R11 Y*0(1520) INTO (LAMBDA PI PI)/TOTAL (P3)
R11 0.10 (0.021) COLLEY 71 DBC K-N 1.5 GEV PRD 10/71
R11 .01 .01 MAST 73 IPWA K-P TO 2PI LAM 9/73
R11 1 BASED ON ASSUMED ELASTICITY OF .467/-02 9/73
R11 3 .091 .006 CORDEN 75 DBC K-D 1.4-1.8GV/C 4/75*

R12 Y*0(1520) INTO (LAMBDA EPSILON)/TOTAL (P9) 4/75*
R12 .20 .08 CORDEN 75 DBC K-D 1.4-1.8GV/C 4/75*

REFERENCES FOR Y*0(1520)

GALTIERI 63 PL 6 296 A BARBARO-GALTIERI, A HUSSAIN, R D TRIPP (LRL)
WATSON 63 PR 131 2248 M B WATSON, M FERRO-LUZZI, R D TRIPP (LRL) IJP
ALMEIDA 64 PL 9 204 S P ALMEIDA, G R LYNCH (CERN)
ARMENTER 65 PL 19 338 ARMENTEROS, F LUZZI, + (CERN, HEID, SACLAY)
MUSGRAVE 65 NC 35 735 +PETEZAS, + (BIRM, CERN, EPOL, LOIC, SACLAY)

BIRMINGHAM 66 PR 152 1148 BIRMINGHAM, GLASSGOW, I. C., OXFORD, RUTHERFORD
DAHL 67 PR 163 1377 DAHL, HADY, HESS, KIRZ, MILLER (LRL)
DAUBER 67 PL 248 525 +MALAMUD, SCHLEIN, SLATER, STOKR (UCLA)
UHLIG 67 PR 155 1448 +CHARLTON, CORDON, GLASSER, YODH, + (UMD, NRL)
MAST 68 PR 21 1715 MAST, ALSTON, BANGERTER, GALTIERI + (LRL)
SCHEUER 68 NP 88 903 SABRE COLLAB. (SACL+AMST+BGNA+REHO+EPOL)

BURKHARD 69 NP 814 106 +FILTHUTH+KLUGE+.. (HEID+EFI+CERN+SACLAY)
CLINE 69 LNC 2 407 +LAUMANN+MAPP (WISC)
GALTIERI 69 LUND 352 BARBARO-GALTIERI, BANGERTER, MAST, TRIPP (LFL)
ALSO 70 DUKE 95 R D TRIPP (LRL)

BURKHARDT 71 NP 827 64 +FILTHUTH, KLUGE, OBERLACK++ (HEID+CERN+SACLAY)
COLLEY 71 NP 831 61 +COX, EASTWOOD, FRY+.. (BIRM+EDIN+GLAS+LOIC)
KIM 71 PRL 27 356 J K KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP

CHAN 72 PRL 28 256 +BUT--SHAFFER, HERTZBACH, KOFLER++ (MASA, YALE)
MAST 73 PR D7 5 +ALSTON-GARNJOST, BANGERTER, +.. (LBL) IJP
MAST2 73 PR D7 3212 +BANGERTER, ALSTON-GARNJOST, + (LBL) IJP
BERTHON 74 NC 21A 146 +BERTHON, TRISTRAM, + (CDFE+RHEL+SACL+STRB) U
CORDEN 75 NP 884 306 CORDEN, COX, CARTNELL, KEVON, ONEALE, + (BIRM)

MAST 76 LBL-4294 (PRPNT) MAST, ALSTON-GARNJOST, BANGERTER+ (LBL)
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP

BERLEY 70 PR D1 1996 +YAMIN, KOFLER, MANN, MEISNER+ (BNL, MASA, YALE) IJP
GOLOWICH 74 PR D1 3861 EUGEN GOLOWICH (SLAC)

$\Lambda(1600)$

P_{01}

101 Y*0(1600, JP=1/2+) I=0 1/76*

101 Y*0(1600) MASS (MEV) 1/76*
M 1 (1570.) KIM 71 DPWA K-MATRIX ANAL. 1/76*
M 1 POSSIBLE EFFECT IN SIGMA PI AND KBAR N CHANNELS.
M 1620.0 10.0 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
M 1573. 25. RLIC 76 DPWA KBAR N MULTICHNL 1/76*

101 Y*0(1600) WIDTH (MEV) 1/76*
W 1 (50.) KIM 71 DPWA K-MATRIX ANAL. 1/76*
W 60.0 10.0 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
W 147. 50. RLIC 76 DPWA KBAR N MULTICHNL 1/76*

101 Y*0(1600) PARTIAL DECAY MODES 1/76*
P1 Y*0(1600) INTO KBAR N DECAY MASSES 497+ 939
P2 Y*0(1600) INTO SIGMA PI 1197+ 139

101 Y*0(1600) BRANCHING RATIOS 1/76*
R1 Y*0(1600) INTO (KBAR NI)/TOTAL (P1) 1/76*
R1 0.25 0.15 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
R1 .24 .04 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

R2 Y*0(1600) FROM KBAR N INTO SIGMA PI SQR(P1*P2) 1/76*
R2 0.28 0.09 LANGBEIN 72 IPWA MULTICHANNEL 1/76*
R2 .16 .04 RLIC 76 DPWA KBAR N MULTICHNL 1/76*

REFERENCES FOR Y*0(1600)
KIM 71 PRL 27 356 J K KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP
LANGBEIN 72 NP 847 477 +WAGNER (MPIM) IJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP

$\Lambda(1670)$

S_{01}

40 Y*0(1670, JP=1/2-) I=0

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.
THIS RESONANCE IS WELL ESTABLISHED.
(SEE THE NOTE FOR THE Y*0(1330)).

40 Y*0(1670) MASS (MEV)
M M (1666.0)OR(1675.0) BERLEY 65 HBC 0 K-P TO LAM ETA 7/66
M M THE FIRST VALUE ASSUMES THE BRANCHING RATIO INTO LAMBDA ETA IS SMALL, THE SECOND THAT IT IS LARGE. BECAUSE THE RESONANCE IS NEAR THE LAMBDA ETA THRESHOLD, THE BRANCHING RATIO AFFECTS THE MOMENTUM DEPENDENCE OF THE TOTAL WIDTH, AND THUS ALSO THE RESONANCE PARAMETERS OBTAINED BY FITTING TO THE DATA.
M N (1663.0) (3.0) ARMENT-1 68 HBC 0 ELASTIC, CH EXCH 11/68
M N (1678.0) (2.0) ARMENT-2 68 HBC 0 K-P TO SIGMA PI 11/68
M A 1674.0 (5.0) ARMENT-3 69 HBC 0 MULTICHANNEL 9/69
M N 1662.0 (3.0) ARMENT-4 69 HBC 0 ELAST, CH, EXC. ED 9/69
M N 1680.0 (1.0) ARMENT-4 69 HBC 0 K-P TO SIG PI. ED 9/69
M 1674.0 BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
M 1683.0 BERLEY 70 HBC 0 SIG PI, EDPA 7/70
M 1670.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1640.0 (40.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 1700. (10.) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
M 1672. (1.) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
M 1665. (5.) PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*
M 1670. (5.) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

40 Y*0(1670) WIDTH (MEV)
W M (22.0)OR(15.0) BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
W N (26.0) (8.0) ARMENT-1 68 HBC 0 SEE NOTE N ABOVE 11/68
W N (26.0) (5.0) ARMENT-2 68 HBC 0 9/69
W A 23.0 (3.0) ARMENT-3 69 HBC 0 9/69
W N 38.0 (15.0) ARMENT-4 69 HBC 0 ELAST, CH EXC. ED 9/69
W A 35.0 (5.0) ARMENT-4 69 HBC 0 K-P TO SIG PI. ED 6/70
W 31.0 BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
W 25.0 (5.0) GALTIERI 70 HBC 0 SIG PI, EDPA 7/70
W 35. KIM 71 DPWA K-MATRIX ANAL. 3/71
W 45.0 (20.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 65. (20.) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
W 19. (2.) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
W 19. (5.) PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*
W 45. (10.) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

40 Y*0(1670) PARTIAL DECAY MODES
P1 Y*0(1670) INTO KBAR N DECAY MASSES 497+ 939
P2 Y*0(1670) INTO LAMBDA ETA 1154+ 548
P3 Y*0(1670) INTO SIGMA PI 1189+ 139
P4 Y*0(1670) INTO SIGMA(1385) PI 139+1384

40 Y*0(1670) BRANCHING RATIOS
R1 Y*0(1670) INTO (KBAR NI)/TOTAL (P1)
R1 P (0.14) (0.04) ARMENT-1 68 HBC 0 OLD DATA 11/68
R1 0.17 ARMENT-3 69 HBC 0 9/69
R1 P 0.14 (0.04) ARMENT-4 69 HBC 0 NEW DATA 9/69
R1 A (0.39) (0.05) CENFORTO 71 HBC 0 K-P, ELAST, CEX 6/70
R1 0.28 KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 0.35 (0.06) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 .36 (.03) HART 73 DPWA EL+CX, .7-.8GEV/C 2/74
R1 .20 (.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*

R2 Y*0(1670) FROM KBAR N TO LAMBDA ETA SQR(P1*P2)
R2 M (0.20) OR 0.23 BERLEY 65 HBC 0 SEE NOTE M ABOVE 7/66
R2 (0.26) ARMENT-3 69 HBC 0 9/69
R2 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 +.20 (.05) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*

R3 Y*0(1670) FROM KBAR N TO SIGMA PI SQR(P1*P3)
R3 1 (-0.25) (0.06) ARMENT-2 68 HBC 0 OLD DATA 9/69
R3 1 (-0.27) ARMENT-3 69 HBC 0 9/69
R3 1 (-0.30) (0.03) ARMENT-4 69 HBC 0 NEW DATA 10/74*
R3 1 PUBLISHED SIGN CHANGED TO AGREE WITH LUND 1969 CONVENTION (SEE TEXT) 10/74*
R3 -0.27 BERLEY 69 HBC 0 K-P TO SIGMA PI 6/70
R3 -0.29 (0.03) LANGBEIN 70 HBC 0 SIG PI, EDPA 7/70
R3 -0.38 KIM 71 DPWA K-MATRIX ANAL. 3/71
R3 -.28 (.05) BAXTER 73 DPWA 0 K- P TO NEUTRALS 10/74*
R3 -.23 (.03) LNDONC 75 HBC 0 K- P TO SIG PI 4/75*

R4 Y*0(1670) FROM KBAR N TO SIGMA(1385) PI SQR(P1*P4)
R4 -.18 .05 PREVOST 74 DPWA 0 K-N TO S(1385)PI 10/74*

Baryons

$\Lambda(1670)$, $\Lambda(1690)$, $\Lambda(1800)$

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR $\Lambda(1670)$

BERLEY 65 PRL 15 641 +CONNOLLY, HART, RAUM, STCNEHILL, + (BNL)IJP
ARMENT-1 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-2 68 NP 88 223 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-3 69 LUND PAPER 229 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
VALUES ARE QUOTED IN LEVI SETTI 69.

BIRMINGHAM 66 PR 152 1148 (BIRMINGHAM, GLASGOW, LOIC, OXFORD, RUTHERFD)
LEVISETTI 69 LUND 339 R LEVI SETTI (RAPPORTEUR) (CHICAGO)

$\Lambda(1690)$

55 $\Lambda(1690)$, JP=3/2- I=0 D_{03}
SEE THE MINI-REVIEW AT THE START OF THE Λ^* LISTINGS.
THIS RESONANCE IS WELL ESTABLISHED.

55 $\Lambda(1690)$ MASS (MEV)

Table with columns for mass (M), width (W), and various parameters for the 55 $\Lambda(1690)$ resonance. Includes text: 'THE $\Lambda(1690)$ IS AT THE EDGE OF THE ENERGY REGION AND THEREBY M ENERGIES ARE INCLUDED IN ARMENTEROS 1.'

55 $\Lambda(1690)$ WIDTH (MEV)

Table with columns for width (W) and various parameters for the 55 $\Lambda(1690)$ resonance. Includes text: 'SEE THE NOTES ACCOMPANYING THE MASSES QUOTED'.

55 $\Lambda(1690)$ PARTIAL DECAY MODES

Table with columns for partial decay modes (P1-P6) and decay masses for the 55 $\Lambda(1690)$ resonance.

55 $\Lambda(1690)$ BRANCHING RATIOS

THE SUM OF ALL THE QUOTED BRANCHING RATIOS IS MORE THAN 1.0. THE TWO-BODY RATIOS ARE FROM PARTIAL WAVE ANALYSES, AND THUS PROBABLY ARE MORE RELIABLE THAN THE THREE-BODY RATIOS, WHICH ARE DETERMINED FROM BUMPS IN CROSS SECTIONS.

Table with columns for branching ratios (R1-R1) and various parameters for the 55 $\Lambda(1690)$ resonance. Includes text: 'R1 N EFFECT IS AT END OF REGION ANALYZED. THIS COULD AFFECT VALUE OF X1.'

Table with columns for $\Lambda(1690)$ resonance parameters (R2-R6) and various parameters.

Table with columns for $\Lambda(1690)$ resonance parameters (R3-R6) and various parameters.

REFERENCES FOR $\Lambda(1690)$

ARMENT-1 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
ARMENT-2 68 NP 88 216 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY) I
ARMENT-3 68 NP 88 223 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
BARTLEY 68 PRL 21 1111 +CHU,DOMO, GREENE, + (TIJFS,FSU, BRANDES) I
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (BIRM,CAVE,RHEL) I
ALSO 67 PRL 18 62 DAVIES, DOWELL, + (BIRM,CAVE,RHEL) I
CONFORTO 68 NP 88 265 +HARSEN, LASINSKI, + (CHICAGO,HEIDEL)IJP
ARMENT-4 69 NP 814 91 ARMENTEROS, BAILLON, + (CERN,HEIDEL,SACLAY)IJP
BERLEY 69 PL 308 430 +HART, RAHM, WILLIS, YAMAMOTO (BNL)IJP
BERTANZA 69 PR 177 2036 +BIGI,CARRARA,CASALI, + (PISA,BNL, YALE)IJP
GALTIERI 70 DUKE 173 A. BARBARO GALTIERI (LRL)IJP
CONFORTO 71 NP 854 41 +LEVI SETTI,LASINSKI, OBERLACK++ (EFI+IED)IJP
KIM 71 PRL 27 356 J. K. KIM (HARV)IJP
ALSO 70 DUKE 161 J. K. KIM (HARV)IJP
LANGBEIN 72 NP 847 477 +WAGNER (NPI)IJP
BAXTER 73 NP 867 125 BAXTER,BUCKINGHAM,CORBETT,DUNN,+ (OXFORD)IJP
HART 73 PURDUE CONF. 311 +RICE,BACASTOW,FUNG,+ (TENN+UCR+MASA+BUFF)IJP
PREVOST 74 NP 869 246 PREVOST,BARLOUTAUD,+ (SACL+CERN+HEID)

PAPERS NOT REFERRED TO IN DATA CARDS

PREVOST 71 AMSTERDAM CONF + CHS COLLABORATION (CERN+HEID+SACL)

$\Lambda(1800)$

77 $\Lambda(1800)$, JP=1/2+ I=0 P_{01}
SEE THE MINI-REVIEW AT THE START OF THE Λ^* LISTINGS.

THE EVIDENCE FOR THIS STATE IS SOMEWHAT CONFUSED. IT WAS FIRST SUGGESTED IN A PARTIAL WAVE ANALYSIS OF KBAR N DATA BY THE BEHAVIOUR OF THE P01 AMPLITUDE WHEN IT WAS PARAMETRIZED AS A TWO-STRAIGHT-LINE BACKGROUND (ARMENTEROS 68).

77 $\Lambda(1800)$ MASS (MEV)

Table with columns for mass (M) and various parameters for the 77 $\Lambda(1800)$ resonance. Includes text: 'A WIDER AND MORE ELASTIC P01 RESONANCE AT ABOUT THE SAME MASS IS SUGGESTED BY THE ANALYSIS OF BAILEY 69.'

77 $\Lambda(1800)$ WIDTH (MEV)

Table with columns for width (W) and various parameters for the 77 $\Lambda(1800)$ resonance. Includes text: 'AVERAGE MEANINGLESS (SCALE FACTOR = 3.1) SEE THE NOTES ACCOMPANYING MASSES QUOTED'.

77 $\Lambda(1800)$ PARTIAL DECAY MODES

Table with columns for partial decay modes (P1-P3) and decay masses for the 77 $\Lambda(1800)$ resonance.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1800)$, $\Lambda(1815)$

77 $\Lambda(1800)$ BRANCHING RATIOS

Table with columns for particle name, mass, width, and branching ratios. Includes entries for R1, R2, R3, and R4.

REFERENCES FOR $\Lambda(1800)$
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
DAVID SAAL BAILEY (LRL LIVERMORE) IJP
ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJP
A BARBARO-GALTIERI (LRL) IJP
J. K. KIM (HARV) IJP
J. K. KIM (HARV) IJP
WAGNER (MPI) IJP

$\Lambda(1800)$

102 $\Lambda(1800)$, JP=9/2- I=0 G_{99}
THIS NARROW RESONANCE SEEMS NECESSARY TO FIT A PEAK IN THE Λ COEFFICIENT OF THE K-P ANGULAR DISTRIBUTION. IT IS NOT REQUIRED IN ANY OTHER CHANNEL.

Table with columns for mass (MEV) and width (MEV). Values: 1808, 5, 27, 5.

Table with columns for partial decay modes. Values: 0.62, 0.02, 0.65, 0.02, 0.58, 0.02, 0.81, 0.01, 0.63, 0.01, 0.52, 0.01, 0.47, 0.02, 0.57, 0.02.

Table with columns for branching ratios. Values: 0.601, 0.023, 0.614, 0.010, 0.601, 0.021.

Table with columns for mass (MEV) and width (MEV). Values: 1813.0, 2.0, 1816.0, 4.0, 1817.0, 2.0, 1819.0, 4.0, 1825.0, 1.0, 1819.0, 1.0, 1830.0, 10.0, 1820.0, 10.0, 1818.0, 2.0, 1810.0, 1.0, 1823.0, 3.0, 1819.0, 13.0, 1822.0, 2.0.

$\Lambda(1815)$

39 $\Lambda(1815)$, JP=5/2+ I=0 F_{05}
SEE THE MINI-REVIEW AT THE START OF THE Λ LISTINGS.

THIS STATE IS WELL ESTABLISHED. MOST OF THE QUOTED ERRORS ARE STATISTICAL ONLY. THE SYSTEMATIC ERRORS DUE TO THE PARTICULAR PARAMETRIZATION USED IN THE P.W.A. ARE NOT INCLUDED. FOR THIS REASON WE DO NOT CALCULATE WEIGHTED AVERAGES FOR MASS AND WIDTH.

Table with columns for mass (MEV) and width (MEV). Values: 1813.0, 2.0, 1816.0, 4.0, 1817.0, 2.0, 1819.0, 4.0, 1825.0, 1.0, 1819.0, 1.0, 1830.0, 10.0, 1820.0, 10.0, 1818.0, 2.0, 1810.0, 1.0, 1823.0, 3.0, 1819.0, 13.0, 1822.0, 2.0.

39 $\Lambda(1815)$ WIDTH (MEV)

Table with columns for particle name, mass, width, and branching ratios. Includes entries for W, N, and R.

39 $\Lambda(1815)$ PARTIAL DECAY MODES

Table with columns for decay mode and mass. Values: 497+939, 1189+139, 139+1384, 1192+139+139, 548+1115.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i , as follows: The diagonal elements are $P_i \pm \delta P_i$, where $\delta P_i = \sqrt{(\delta P_i^2) / (P_i^2 - \delta P_i^2)}$, while the off-diagonal elements are the normalized correlation coefficients $(\delta P_i \delta P_j) / (P_i P_j - \delta P_i \delta P_j)$. For the definitions of the individual P_i , see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with columns for P1, P2, P3, P4, P5. Values: 0.6005+-0.0214, -0.154, 0.1155+-0.0080, -0.1541, 0.0795, 0.1390+-0.0261, -0.4139, 0.0216, -0.7638, 0.1292+-0.0308, -0.0642, 0.0331, 0.0099, -0.2505, 0.0153+-0.0086.

39 $\Lambda(1815)$ BRANCHING RATIOS

ERRORS QUOTED BY EXPERIMENTERS DO NOT INCLUDE UNCERTAINTY DUE TO PARAMETRIZATION USED IN THE P.W.A. THEY SHOULD BE INCREASED.

Table with columns for particle name, mass, width, and branching ratios. Includes entries for R1, R2, R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, R14, R15, R16, R17, R18, R19, R20, R21, R22, R23, R24, R25, R26, R27, R28, R29, R30, R31, R32, R33, R34, R35, R36, R37, R38, R39, R40, R41, R42, R43, R44, R45, R46, R47, R48, R49, R50, R51, R52, R53, R54, R55, R56, R57, R58, R59, R60, R61, R62, R63, R64, R65, R66, R67, R68, R69, R70, R71, R72, R73, R74, R75, R76, R77, R78, R79, R80, R81, R82, R83, R84, R85, R86, R87, R88, R89, R90, R91, R92, R93, R94, R95, R96, R97, R98, R99, R100.

Table with columns for branching ratios. Values: 0.601, 0.023, 0.614, 0.010, 0.601, 0.021, 0.601, 0.023, 0.614, 0.010, 0.601, 0.021.

Table with columns for mass (MEV) and width (MEV). Values: 1813.0, 2.0, 1816.0, 4.0, 1817.0, 2.0, 1819.0, 4.0, 1825.0, 1.0, 1819.0, 1.0, 1830.0, 10.0, 1820.0, 10.0, 1818.0, 2.0, 1810.0, 1.0, 1823.0, 3.0, 1819.0, 13.0, 1822.0, 2.0.

Table with columns for branching ratios. Values: 0.601, 0.023, 0.614, 0.010, 0.601, 0.021.

Table with columns for mass (MEV) and width (MEV). Values: 1813.0, 2.0, 1816.0, 4.0, 1817.0, 2.0, 1819.0, 4.0, 1825.0, 1.0, 1819.0, 1.0, 1830.0, 10.0, 1820.0, 10.0, 1818.0, 2.0, 1810.0, 1.0, 1823.0, 3.0, 1819.0, 13.0, 1822.0, 2.0.

Table with columns for branching ratios. Values: 0.601, 0.023, 0.614, 0.010, 0.601, 0.021.

REFERENCES FOR $\Lambda(1815)$
BIRGE 65 ATHENS CONF 296 +ELY, KALMUS, KERNAN, LOUIE, SAHOURIA, + (LRL) IJP
ARMENT-1 67 PL 248 198 ARMENTEROS, F LUZZI, + (CERN, HEIDEL, SACLAY) IJP
ARMENT-2 67 ZEIT PHYS 202 486 ARMENTEROS, F LUZZI, + (CERN, HEIDEL, SACLAY) IJP
BELL 67 PRL 19 936 R B BELL (LRL) IJP
ARMENT-3 68 NP 88 195 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
ARMENT-4 68 NP 88 216 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHEL+BIRM+CAVE) I

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(1815)$, $\Lambda(1830)$, $\Lambda(1860)$

BRICMAN 70 PL 318 152
BRICMAN1 70 PL 338 511
COOL 70 PR D1 1887
GALTIERI 70 DUKE CONF 173
CONFORTO 71 NP 834 41
KIM 71 PRL 27 356
ALSO 70 DUKE 161
KANE 72 PR 05 1583
LANGBEIN 72 NP 847 477
RADER 73 NC 16A 178
PREVCST 74 NP 869 246
RLIC 76 RL-75-182 (PRPNT)
GOPAL, KALMUS, MCPHERSON, ROSS+
PAPERS NOT REFERRED TO IN DATA CARDS

+FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
+FERRO-LUZZI, LAGNAUX (CERN)
+GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
A BARBARO-GALTIERI (LRL) IJJP
+LEVI SETTI, LASINSKI, OBERLACK++ (EFI+HEID) IJP
J. K. KIM (HARV) IJP
D. F. KANE (LBL) IJP
+WAGNER (MPIM) IJP
+BARLOUTAUD, + (SACL+HEID+CERN+RHEL+CDEF)
PREVOST, BARLOUTAUD, + (SACL+CERN+HEID)

GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP
PAPERS NOT REFERRED TO IN DATA CARDS

THE FOLLOWING PAPERS ARE NEW OF ONLY HISTORICAL INTEREST --

CHAMBERL 62 PR 125 1696
GALTIERI 63 PL 6 296
SODICKSON 64 PR 133 8757
HOLLEY 65 UCLR-16274 THESIS
BIRMINGHAM 66 PR 152 1148
COOL 66 PRL 16 1228
GELFAND 66 PRL 17 1224
ARMENTER 67 NP 83 552
CCNFORTO 68 NP 88 265
LASINSKI 68 PR 163 1792
PREVOST 71 AMSTERDAM CONF
CHAMBERLAIN, CROWE, KEEFE, KERTH, + (LRL) I
A BARBARO-GALTIERI, A HUSSAIN, RD TRIPP (LRL) IJJP
SODICKSON, MANNELLI, FRISCH, WAHLIG (MIT(BNL)) J
W R HOLLEY (LRL) J
BIRMINGHAM, GLASGOW, I.C., OXFORD, RUTHERFORD
+GIACOMELLI, KYCIA, LEONTIC, LUNDBY + (BNL) I
+HARMSEN, LEVI-SETTI, PREDAZZI+ (EFI+ANL) I
ARMENTEROS, FERRO-LUZZI+ (CERN, HEID, SACLAY) IJP
+HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
LASINSKI, LEVI SETTI, PREDAZZI (CHICAGO) JIP
PREVOST 71 AMSTERDAM CONF + CHS COLLABORATION (CERN+HEID+SACL)

$\Lambda(1830)$

56 Y*(1830, JP=5/2-) I=0
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THE BEST EVIDENCE FOR THIS RESONANCE COMES FROM THE SIGMA PI CHANNEL. IT IS WELL ESTABLISHED.

D05

56 Y*(1830) MASS (MEV)
M N 1827.0 (3.0) ARMENTERO 67 HBC 0 K-P TO SIGMA PI 8/67
M N 1837.0 (11.0) BELL 67 HBC 0 K-P TO SIGMA PI 11/67
M N 1807.0 (10.0) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
M N 1840.0 (15.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
M N 1831.0 (5.0) CONFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
M N 1830.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
M K (1172.0) KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1832.0 (5.0) KANE 72 DPWA 0 K-P TO PI SIG 10/71
M 1810.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 1825.0 (10.0) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
M K POSSIBLE EFFECT MAINLY IN SIGMA PI. NOT CLEAR IF UNCORRELATED
M WITH THE 1830 EFFECT
M N ERROR STATIST. ONLY- NO ERROR DUE TO PARTICULAR P.W. ANAL. INCLUDED 1/71

56 Y*(1830) WIDTH (MEV)
W 75.0 (9.0) ARMENTERO 67 HBC 0 K-P TO SIGMA PI 8/67
W 74.0 (18.0) BELL 67 HBC 0 K-P TO SIGMA PI 8/67
W 123.0 (32.0) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
W 150.0 (30.0) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
W 104.0 (35.0) CONFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
W 80.0 KIM 71 DPWA K-MATRIX ANAL. 3/71
W K (20.0) KIM 71 DPWA K-MATRIX ANAL. 3/71
W 88.0 (10.0) KANE 72 DPWA 0 K-P TO PI SIG 10/71
W 60.0 (20.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 94.0 (10.0) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
SEE THE NOTES ACCOMPANYING MASSES QUOTED

56 Y*(1830) PARTIAL DECAY MODES
P1 Y*(1830) INTO KBAR N 497+ 939
P2 Y*(1830) INTO SIGMA PI 1189+ 139
P3 Y*(1830) INTO Y(11385) PI D-WAVE 139+1384
P4 Y*(1830) INTO ETA LAMBDA 548+1115

56 Y*(1830) BRANCHING RATIOS
R1 Y*(1830) INTO (KBAR N)/TOTAL (P1)
R1 0.09 (0.01) ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
R1 0.03 (0.02) BRICMAN1 70 DPWA SIGTOT, ELAS, CHEX 1/71
R1 0.05 (0.02) CCNFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
R1 (0.24) KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 0.10 (0.03) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 .04 (0.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
R2 Y*(1830) FROM KBAR N INTO SIGMA PI SQRTP1*P2)
R2 A (-0.15) (0.02) ARMENTERO 67 DPWA 0 K-P TO SIGMA PI 10/74*
R2 A PUBLISHED SIGN CHANGED TO AGREE WITH LUND 1969 CONVENTION (SEE TEXT) 10/74*
R2 0.19 (0.01) BELL 67 DPWA 0 K-P TO SIGMA PI 11/67
R2 -0.16 (0.03) GALTIERI 70 DPWA 0 K-P TO SIGMA PI 7/70
R2 0.15 (0.01) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 -0.138 (0.0318) KANE 72 DPWA 0 K-P TO PI SIG 10/71
R2 0.27 (0.07) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R2 -0.17 (0.03) RLIC 76 DPWA KBAR N MULTICHNL 1/76*
R3 Y*(1830) FROM KBAR N TO ETA LAMBDA SQRTP1*P4)
R3 -0.044 .020 RADER 73 MPWA 9/73
R4 Y*(1830) FROM KBAR N TO Y*(1385) PI D-WAVE SQRTP1*P3)
R4 +.13 .03 PREVOST 74 DPWA 0 K-N TO S(1385) IPI 10/74*

REFERENCES FOR Y*(1830)
ARMENTEROS, F.-LUZZI, + (CERN, HEIDEL, SACLAY) IJP
R B BELL (LRL) IJP
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
+HARMSEN, LASINSKI, + (CHICAGO, HEIDEL) IJP
IS SUPERSEDED BY CONFORTO 71.
BRICMAN1 70 PL 338 511 +FERRO-LUZZI, LAGNAUX (CERN)
GALTIERI 70 DUKE CONF 173 +BARBARO-GALTIERI (LRL) IJJP

CCNFORTO 71 NP 834 41
KIM 71 PRL 27 356
ALSO 70 DUKE 161
KANE 72 PR 05 1583
LANGBEIN 72 NP 847 477
RADER 73 NC 16A 178
PREVOST 74 NP 869 246
RLIC 76 RL-75-182 (PRPNT)
+LEVI SETTI, LASINSKI, OBERLACK++ (EFI+HEID) IJP
J. K. KIM (HARV) IJP
J. K. KIM (HARV) IJP
D. F. KANE (LBL) IJP
+WAGNER (MPIM) IJP
+BARLOUTAUD, + (SACL+HEID+CERN+RHEL+CDEF)
GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJP
PAPERS NOT REFERRED TO IN DATA CARDS

PREVOST 71 AMSTERDAM CONF + CHS COLLABORATION (CERN+HEID+SACL)

$\Lambda(1860)$ 60 Y*(1860, JP=3/2+) I=0
AVAILABLE DATA (INCLUDING POLARIZATION) AND RECENT PARTIAL WAVE ANALYSES. THE DOMINANT INELASTIC MODES REMAIN UNKNOWN. SEE ALSO Y*(2010) MINI-REVIEW.

P03

60 Y*(1860) MASS (MEV)
M A F07 1864.0 2.0 ARMENTERO 68 DPWA 0 ELASTIC, CH EXCH 11/68
M N 1870.0 5.0 BELL 68 CNTR 0 K-P TOTAL 7/68
M A F07 1871.0 6.0 BRICMAN1 70 CNTR 0 TOTAL AND CH EX 6/70
M 1870.0 6.0 BRICMAN1 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
M N 1883.0 10.0 CONFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
M 1 (1710.) KIM 71 DPWA K-MATRIX ANAL. 3/71
M 1850.0 20.0 LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 2 (1868.) LEA 73 DPWA MULTICHNL K-MTRX 9/73
M 1894. 10. HEMINGWAY 75 DPWA 0 K-P TO KBAR N 11/75*
M 3 (1900.) NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*
M 1900. 5. RLIC 76 DPWA KBAR N MULTICHNL 1/76*
M A THESE TWO ANALYSES GAVE THE F07 ASSIGNMENT, THEY HAVE TO BE 1/71
M A DISCARDED IN VIEW OF CCNFORTO 70 AND BRICMAN1 70 1/71
M N DUE TO PARTICULAR PARAMETERIZATION USED, ERROR CAN BE LARGE 1/71
M 1 POSSIBLE EFFECT MAINLY IN SIGMA PI. WE TENTATIVELY LIST IT HERE. 9/73
M 2 ONLY UNCONSTRAINED STATES FROM TABLE 1 OF LEAYS ARE IN LISTINGS. 1/76*
M 3 FOUND IN ONE OF TWO BEST SOLUTIONS.
M AVERAGE MEANINGLESS (SCALE FACTOR = 2.3)

60 Y*(1860) WIDTH (MEV)
W A F07 39.0 7.0 ARMENTERO 68 DPWA 0 ELASTIC, CH EXCH 11/68
W N 40.0 10.0 BUGG 68 CNTR 0 K-P TOTAL 7/68
W A F07 24.0 15.0 BRICMAN1 70 CNTR 0 TOTAL AND CH EX 6/70
W N 37.0 10.0 BRICMAN1 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
W N 80.0 20.0 CONFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
W 1 (20.) KIM 71 DPWA K-MATRIX ANAL. 3/71
W 125.0 20.0 LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 2 (323.8) KANE 72 DPWA MULTICHANNEL 12/72
W 107. 10. HEMINGWAY 75 DPWA 0 K-P TO KBAR N 11/75*
W 3 (100.) NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*
W 72. 10. RLIC 76 DPWA KBAR N MULTICHNL 1/76*
SEE THE NOTES ACCOMPANYING MASSES QUOTED
AVERAGE MEANINGLESS (SCALE FACTOR = 2.5)

60 Y*(1860) PARTIAL DECAY MODES
P1 Y*(1860) INTO KBAR N 497+ 939
P2 Y*(1860) INTO SIGMA PI 1189+ 139
P3 Y*(1860) INTO LAMBDA OMEGA 1115+ 783

60 Y*(1860) BRANCHING RATIOS
R1 Y*(1860) INTO (KBAR N)/TOTAL (P1)
R1 A F07 0.12 0.02 ARMENTERO 68 HBC 0 ELASTIC, CH EXCH 11/68
R1 (J+1/2)P1= 0.40 BUGG 68 CNTR 0 7/68
R1 A F07 0.07 0.02 BRICMAN1 70 CNTR 0 TOTAL AND CH EX 6/70
R1 0.14 0.02 BRICMAN1 70 DPWA 0 SIGTOT, ELAS, CHEX 1/71
R1 N 0.25 0.03 CONFORTO 71 DPWA 0 ELASTIC, CH EXCH 6/70
R1 0.37 0.05 LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 2 (.32) LEA 73 DPWA MULTICHNL K-MTRX 9/73
R1 .24 .04 HEMINGWAY 75 DPWA 0 K-P TO KBAR N 11/75*
R1 .18 .02 RLIC 76 DPWA KBAR N MULTICHNL 1/76*
R1 SEE THE NOTES ACCOMPANYING MASSES QUOTED
R1 AVERAGE MEANINGLESS (SCALE FACTOR = 2.5)
R2 Y*(1860) INTO SIGMA PI (P2)
R2 P PROBABLY SEEN GALTIERI 68 CBC 0 K-N TO SIG PI PI 11/68
R2 (0.03) OR LESS LANGBEIN 72 IPWA MULTICHANNEL 12/72
R2 P POSSIBLY THIS BUMP SEEN AT 1840+-10 MEV WITH A WIDTH OF 35+-10 MEV
R2 IS THE Y*(1830), WHICH DECAYS STRONGLY TO SIGMA PI. HOWEVER THE
R2 NARROW WIDTH HERE ARGUES FOR ITS BEING THE Y*(1860).
R3 Y*(1860) FROM KBAR N TO SIGMA PI SQRTP1*P2) 9/73
R3 2 (-.15) LEA 73 DPWA MULTICHNL K-MTRX 9/73
R3 -.09 .03 LANGBEIN 72 IPWA KBAR N MULTICHNL 1/76*
R4 Y*(1860) FROM KBAR N INTO LAMBDA OMEGA SQRTP1*P3) 1/76*
R4 3 (.032) NAKKASYA 75 DPWA 0 K-P TO LAM. OMC. 1/76*

REFERENCES FOR Y*(1860)
ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IJP
+GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I
BARBARO-GALTIERI, MATTISON, + (LRL, SACL)

BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
BRICMAN1 70 PL 338 511 +FERRO-LUZZI, LAGNAUX (CERN)
CONFORTO 71 NP 834 41 +LEVI SETTI, LASINSKI, OBERLACK++ (EFI+HEID) IJP
KIM 71 PRL 27 356 J. K. KIM (HARV) IJP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJP
LANGBEIN 72 NP 847 477 +WAGNER (MPIM) IJP

Data Card Listings

Baryons

For notation, see key at front of Listings. $\Lambda(1860), \Lambda(1870), \Lambda(2010), \Lambda(2020), \Lambda(2100)$

LEA 73 NP 856 77 +MARTIN, MOORHOUSE+ (RHEL+LOUC+GLAS+AARHUS)IJP
HEMINGWA 75 NP 891 12 FEMINGWAY, EADES, HARMSEN+ (CERN, HEID, MPIM)IJP
NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)IJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC)IJP
PAPERS NOT REFERRED TO IN DATA CARDS

$\Lambda(1870)$

S_{01}

36 Y*0(1870, JP=1/2-) I=0
THE S01 AMPLITUDE SHOWS A SECOND RESONANCE BEHAVIOR AT ABOUT 1800 MEV IN 4 ANALYSES. THE ELASTICITY OF KIM 71 IS SURPRISINGLY LARGE.

Table with 4 columns: M, (mass), (width), and various decay channels (BRICMAN, KIM, LANGBEIN, RLIC) with their respective DPWA and branching ratios.

36 Y*0(1870) WIDTH (MEV)

Table with 4 columns: W, (mass), (width), and various decay channels (BRICMAN, KIM, LANGBEIN, RLIC) with their respective DPWA and branching ratios.

36 Y*0(1870) PARTIAL DECAY MODES

Table with 2 columns: P1, P2 and DECAY MASSES (497+939, 1197+139).

36 Y*0(1870) BRANCHING RATIOS

Table with 4 columns: R1, R2, (branching ratios), and various decay channels (BRICMAN, KIM, LANGBEIN, RLIC) with their respective DPWA and branching ratios.

REFERENCES FOR Y*0(1870)

BRICMAN 70 PL 338 511 C BRICMAN, M FERRO-LUZZI, J P LAGNALX(CERN)IJP
KIM 71 PRL 27 356 J K KIM (HARV)IJP
ALSO 70 DUKE 161 J. K. KIM (HARV)IJP
LANGBEIN 72 NP 847 477 WAGNER (MPI)IJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC)IJP

$\Lambda(2010)$

89 Y*0(2010,) I=0
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
SEVERAL AMBIGUOUS RESONANCE POSSIBILITIES ARISE FROM THE ANALYSES LISTED HERE. POSSIBLE QUANTUM NUMBERS ARE D3 (GALTIERI 70, SIGMA PI), D3+F5 OR P3+D5 OR P1+D3 (BRANDSTETTER 72, LAMBDA OMEGA). A MORE RECENT LAMBDA OMEGA ANALYSIS (NAKKASYAN 75) FINDS TWO BEST SOLUTIONS, EACH WITH THE Y*0(2100) AND ONE ADDITIONAL RESONANCE WHICH CAN BE EITHER THE Y*0(1860, JP=3/2+) OR THE Y*0(2110, JP=5/2+).

89 Y*0(2010) MASS (MEV)

Table with 4 columns: M, (mass), (width), and various decay channels (GALTIERI, BRANDSTE, KIM) with their respective DPWA and branching ratios.

89 Y*0(2010) WIDTH (MEV)

Table with 4 columns: W, (mass), (width), and various decay channels (GALTIERI, BRANDSTE, KIM) with their respective DPWA and branching ratios.

89 Y*0(2010) PARTIAL DECAY MODES

Table with 2 columns: P1, P2, P3 and DECAY MASSES (497+939, 1197+139, 1115+783).

89 Y*0(2010) BRANCHING RATIOS

Table with 4 columns: R1, R2, (branching ratios), and various decay channels (GALTIERI, BRANDSTE) with their respective DPWA and branching ratios.

REFERENCES FOR Y*0(2010)

GALTIERI 70 DUKE CNF 173 A BARBARC-GALTIERI (LRL)IJP
BRANDSTE 72 NP 839 13 BRANDSTETTER, BUTTERWORTH, + (RHEL+CDEF+SACL)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)IJP

$\Lambda(2020)$

F_{07}

27 Y*0(2020, JP=7/2+) I=0
EFFECTS IN THIS PARTIAL WAVE HAVE BEEN OBSERVED AT SOMEWHAT DIFFERENT ENERGIES IN TWO CHANNELS. HOWEVER, LITCHFIELD 71 NOTE THAT THE NEED FOR THIS STATE IN THEIR ANALYSIS RESTS SOLELY ON A POSSIBLY INCONSISTENT POLARIZATION MEASUREMENT AT 1.784 GEV/C. THE STATE WAS NOT REQUIRED IN THE KBAR N TO KBAR N ANALYSIS OF FEMINGWAY 75, BUT COULD NOT BE CONCLUSIVELY RULED OUT.

27 Y*0(2020) MASS (MEV)

Table with 4 columns: M, (mass), (width), and various decay channels (GALTIERI, LITCHFIE) with their respective DPWA and branching ratios.

27 Y*0(2020) WIDTH (MEV)

Table with 4 columns: W, (mass), (width), and various decay channels (GALTIERI, LITCHFIE) with their respective DPWA and branching ratios.

27 Y*0(2020) PARTIAL DECAY MODES

Table with 2 columns: P1, P2 and DECAY MASSES (497+939, 1197+139).

27 Y*0(2020) BRANCHING RATIOS

Table with 4 columns: R1, R2, (branching ratios), and various decay channels (GALTIERI, BRANDSTE) with their respective DPWA and branching ratios.

REFERENCES FOR Y*0(2020)

GALTIERI 70 DUKE CNF 173 A BARBARC-GALTIERI (LRL)IJP
LITCHFIE 71 NP 830 125 LITCHFIELD, ...+LESQUOY, ... (RHEL+CDEF+SACL)IJP
PAPERS NOT REFERRED TO IN DATA CARDS.
HEMINGWA 75 NP 891 12 FEMINGWAY, EADES, HARMSEN+ (CERN, HEID, MPIM)IJP

$\Lambda(2100)$

G_{07}

41 Y*0(2100, JP=7/2-) I=0
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES. PARAMETERS OF PEAKS SEEN IN CROSS-SECTIONS AND INVARIANT-MASS DISTRIBUTIONS AROUND 2100 MEV ARE GIVEN IN THE NEXT ENTRY. EVENTUALLY THE PARTIAL-WAVE ANALYSES SHOULD GIVE THE BEST RESULTS, AS THEY ISOLATE THE G07 WAVE. THIS SUPERIORITY IS, HOWEVER, PROBABLY NOT YET ATTAINED, AND WE RELY ON BOTH ENTRIES FOR PARAMETERS GIVEN IN THE MAIN BARYON TABLE.

41 Y*0(2100) MASS (MEV)

Table with 4 columns: M, A, L, (mass), (width), and various decay channels (WOHL, BURGUN, BERTHON, GALTIERI, LITCHFIE, KANE, HEMINGWA, NAKKASYA, RLIC) with their respective DPWA and branching ratios.

Baryons

$\Lambda(2100)$, $\Lambda(2110)$

Data Card Listings

For notation, see key at front of Listings.

41 $\Lambda(2100)$ WIDTH (MEV)

W	(145.0)	WCHL	66 HBC	7/66
W	(80.0)	BURGUN	68 DPWA	0 K-P TO XI K 10/69
W	140.0 (15.0)	BERTHONI	70 DPWA	0 K-P TO SIGMA PI 10/70
W	60.0 (25.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA PI 7/70
W	(170.) TO 300.	LITCHFIE	71 DPWA	K-P TO KBAR N 10/71
W	B LARGER VALUE CORRESPONDS TO PURE B.W. LOWER VALUE TO B.W. + BCKGRD			
W	140.0 (50.0) (30.0)	LITCHFIE	71 DPWA	K-P TO SIG PI 10/71
W	208. TO 229.	BRANDSTE	72 DPWA	0 K-P TO LAM. OMG. 1/74
W	144.0 (26.0)	KANE	72 DPWA	0 K-P TO PI SIG 10/71
W	241. (30.)	HEMINGWA	75 DPWA	0 K-P TO KBAR N 11/75*
W	244. OR 302.	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG. 11/75*
W	250. (30.1)	RLIC	76 DPWA	KBAR N MULTICHNL 1/76*

SEE THE NOTES ACCOMPANYING MASSES QUOTED

41 $\Lambda(2100)$ PARTIAL DECAY MODES

P1	$\Lambda(2100)$ INTO KBAR N	497+ 939	DECAY MASSES
P2	$\Lambda(2100)$ INTO SIGMA PI	1197+ 139	
P3	$\Lambda(2100)$ INTO XI K	1321+ 497	
P4	$\Lambda(2100)$ INTO LAMBDA OMEGA	1115+ 783	
P5	$\Lambda(2100)$ INTO ETA LAMBDA	548+1115	

41 $\Lambda(2100)$ BRANCHING RATIOS

R1	$\Lambda(2100)$ INTO (KBAR N)/TOTAL	(P1)	7/66
R1	(0.25)	WCHL 66 HBC	7/66
R1 D	(0.33)	DAUM 68 CNTR	K-P ELA,POL,SIGT 7/70
R1	0.30 (0.03)	LITCHFIE 71 DPWA	K-P TO KBAR N 10/71
R1	.31 (0.03)	HEMINGWA 75 DPWA	0 K-P TO KBAR N 11/75*
R1	.30 (0.03)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

R1 D DAUM 68 ASSUMES (J=1/2)*X VALUE SEEN IN TOTAL CROSS SECTION.

R2 $\Lambda(2100)$ FROM KBAR N INTO SIGMA PI

R2 L	(+0.16) (0.02)	BERTHONI 70 DPWA	0 K-P TO SIGMA PI 10/70
R2	+0.06 (0.03)	GALTIERI 70 DPWA	0 K-P TO SIGMA PI 7/70
R2 L	0.16 (0.03)	LITCHFIE 71 DPWA	K-P TO SIG PI 10/71
R2	+0.096 (0.037)	KANE 72 DPWA	0 K-P TO PI SIG 10/71
R2	+12 (0.04)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

R3 $\Lambda(2100)$ FROM KBAR N TO XI K

R3	(0.05)	TRIPP 67 RVUE	0 K-P TO XI K 8/67
R3 B	(0.09) (0.01)	BURGUN 68 DPWA	0 K-P TO XI K 10/69
R3	(0.003)	MULLER 69 DPWA	0 7/70
R3	0.035 0.018	LITCHFIE 71 DPWA	K-P TO XI K 3/72
R3 B	BURGUN 68 UPDATED BY LITCHFIE 71, WHO TAKES SOLUTION C OF BURGUN		3/72

R4 $\Lambda(2100)$ FROM KBAR N INTO LAMBDA OMEGA

R4 1	(.05) TO .11	BRANDSTE 72 DPWA	0 K-P TO LAM. OMG. 1/74
R4 2	(.122)OR .154	NAKKASYA 75 DPWA	0 K-P TO LAM. OMG. 11/75*

R5 $\Lambda(2100)$ FROM KBAR N TO ETA LAMBDA

R5	-.050 .020	RADER 73 MPWA	9/73
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SEE THE NOTES ACCOMPANYING MASSES QUOTED

REFERENCES FOR $\Lambda(2100)$

WCHL 66 PRL 17 107 C G WOHL, F T SOLMITZ, M L STEVENSON (LRL)JJP
 TRIPP 67 NP 83 10 + LRL,SLC,CERN,HEIDEL,SACLAY
 BURGUN 68 NP 88 447 +MEYER,PAULI, + (SACLAY,COLF,FRANCE,RHEL)
 DAUM 68 NP 87 19 +ERNE,LAGNAUX,SENS,STEUER, UDD (CERN)JP
 CONFIRMS THE SPIN-PARITY ASSIGNMENT.
 MULLER 69 THESIS,UCRL 19372 R A MULLER (LRL)

BERTHONI 70 NP 824 417 +VRANA, BUTTERNORTH, + (CCEF, RHEL, SACLAY)JJP
 GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LPL)JJP
 LITCHFIE 71 NP 830 125 LITCHFIE,....+LESQUOY,+. (RHEL+CDEF+SACL)JJP
 BRANDSTE 72 NP 839 13 BRANDSTETTER,....+TALLINI (RHEL,CDEF,SACL) JJP
 KANE 72 PR D5 1583 D F KANE (LBL)JJP
 RADER 73 NC 16A 178 +BARLOUTAUD, + (SACL+HEID+CERN+RHEL+CDEF)
 HEMINGWA 75 NP 891 12 +HEMINGWAY,EADES,HARMSEN+ (CERN,HEID,MPIM)JJP
 NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)JJP
 RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCPHERSON,ROSS+ (RHEL+LOIC)JJP

$\Lambda(2110)$ **F⁰⁵**

BERTHONI 70 FIND EITHER F⁰⁵ OR D⁰⁵ POSSIBLE IN THE SIG PI CHANNEL, WITH F⁰⁵ SLIGHTLY PREFERRED. IN THE KBAR N CHANNEL, LITCHFIE 71 (SAME GROUP) FIND ONLY D⁰⁵. AS USUAL, THE STATISTICS ARE MUCH BETTER IN THE ELASTIC CHANNEL. ALTHOUGH KANE 72 FINDS AN F⁰⁵ EFFECT, THE UNUSUALLY BROAD WIDTH MAY INVALIDATE A RESONANT INTERPRETATION. HOWEVER RLIC 76 AND BELLEFCN 76 ALSO FIND AN F⁰⁵. THE WEIGHT OF THE EVIDENCE IS THUS IN FAVOR OF F⁰⁵. SEE ALSO THE $\Lambda(2100)$ MINI-REVIEW.

35 $\Lambda(2110)$ MASS (MEV)

M	(2110.)	(10.)	BERTHONI 70 DPWA	- K-P TO SIG PI 1/71
M	D05 2140.	40.	LITCHFIE 71 DPWA	K-P TO KBAR N 10/71
M	A (2141.0)	(6.0)	KANE 72 DPWA	0 K-P TO PI SIG 10/71
M	1 (2103.)	(20.)	NAKKASYA 75 DPWA	0 K-P TO LAM. OMG. 1/76*
M	(2140.)	(20.)	BELLEFCN 76 DPWA	0 K-P TO SIG PI 1/76*
M	(2100.)	(50.)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

M A RESONANCE OUTSIDE RANGE OF DATA.
 M 1 FOUND IN ONE OF TWO BEST SOLUTIONS.

35 $\Lambda(2110)$ WIDTH (MEV)

W	(185.)	(30.)	BERTHONI 70 DPWA	- K-P TO SIG PI 1/71
W	D05 120.	40.	LITCHFIE 71 DPWA	K-P TO KBAR N 10/71
W	A (504.0)	(10.0)	KANE 72 DPWA	0 K-P TO PI SIG 10/71
W	1 (391.)	(20.)	NAKKASYA 75 DPWA	0 K-P TO LAM. OMG. 1/76*
W	(140.)	(20.)	BELLEFCN 76 DPWA	0 K-P TO SIG PI 1/76*
W	(200.)	(50.)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

35 $\Lambda(2110)$ PARTIAL DECAY MODES

P1	$\Lambda(2110)$ INTO KBAR N	497+ 939	DECAY MASSES
P2	$\Lambda(2110)$ INTO SIGMA PI	1197+ 139	
P3	$\Lambda(2110)$ INTO LAMBDA OMEGA	1115+ 783	

35 $\Lambda(2110)$ BRANCHING RATIOS

R1	$\Lambda(2110)$ FROM KBAR N TO SIGMA PI	BERTHONI 70 DPWA	SQRT(P1*P2) K-P TO SIG PI 1/71
R1	(+0.156) (0.013)	KANE 72 DPWA	0 K-P TO PI SIG 10/71
R1	(+14) (0.01)	BELLEFCN 76 DPWA	0 K-P TO SIG PI 1/76*
R1	(+10) (0.03)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

R2 $\Lambda(2110)$ INTO (KBAR N)/TOTAL

R2	D05 0.14	0.04	LITCHFIE 71 DPWA	K-P TO KBAR N 10/71
R2	(0.7)	(0.03)	RLIC 76 DPWA	KBAR N MULTICHNL 1/76*

R3 $\Lambda(2110)$ FROM KBAR N INTO LAMBDA OMEGA

R3	(.112)	NAKKASYA 75 DPWA	0 K-P TO LAM. OMG. 1/76*
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REFERENCES FOR $\Lambda(2110)$

BERTHONI 70 NP 824 417 +VRANA,BUTTERNORTH,+ (CDEF,RHEL,SACLAY)JJP
 LITCHFIE 71 NP 830 125 LITCHFIE,....+LESQUOY,+. (RHEL+CDEF+SACL)JJP
 KANE 72 PR D5 1583 D F KANE (LBL)JJP
 NAKKASYA 75 NP 893 85 A. NAKKASYAN (CERN)JJP
 BELLEFCN 76 SBMTD. TO NP DE BELLEFCN,BILLOIR,BERTHON+ (CDEF+SACL)JJP
 RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCPHERSON,ROSS+ (RHEL+LOIC)JJP

2100 MEV REGION - PRODUCTION AND σ_{TOTAL} EXP'TS

25 $\Lambda(2100)$, JP= 1 I=0 PRODUCTION EXPERIMENTS
 SEE THE MINI-REVIEW AT THE START OF THE Λ LISTINGS.

SEE THE NOTE TO THE G07 $\Lambda(2100)$, WHICH PRECEDES THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE $\Lambda(2100)$, BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

25 $\Lambda(2100)$ MASS (MEV) (PROD. EXP.)

M	(2097.0)	(6.0)	BOCK 65 HBC	PBAR P 5.7 BEV/C 7/66
M	2100.0	(7.0)	BUGG 68 CNTR	K-P, D TCTAL 6/68
M	2121.0	(5.0)	BRICMAN 70 CNTR	0 TOTAL AND CH EX 6/70
M	2107.0	(10.0)	COOL 70 CNTR	K-P, D TOTAL 10/70
M	(2135.0)	(20.0)	LU 70 CNTR	0 GAMMA P TO K+ Λ 1/71

25 $\Lambda(2100)$ WIDTH (MEV) (PROD. EXP.)

W	(24.0)	(14.0)	(24.0)	BOCK 65 HBC	INTO KBAR N (P1) 7/66
W	140.0	(15.0)		BUGG 68 CNTR	6/68
W	147.0	(15.0)		BRICMAN 70 CNTR	0 TOTAL AND CH EX 6/70
W	185.0			COOL 70 CNTR	K-P, D TOTAL 10/70
W	(40.0)			LU 70 CNTR	0 GAMMA P TO K+ Λ 1/71

25 $\Lambda(2100)$ PARTIAL DECAY MODES (PROD. EXP.)

P1	$\Lambda(2100)$ INTO KBAR N	497+ 939	DECAY MASSES
P2	$\Lambda(2100)$ INTO KBAR N PI	497+ 939+ 139	
P3	$\Lambda(2100)$ INTO LAMBDA ETA	1115+ 548	
P4	$\Lambda(2100)$ INTO LAMBDA OMEGA	1115+ 783	

25 $\Lambda(2100)$ BRANCHING RATIOS (PROD. EXP.)

R1	$\Lambda(2100)$ INTO (KBAR N)/TOTAL	(P1)	7/66
R1	THESE VALUES OF ELASTICITIES ASSUME J=7/2 --		
R1	0.305	BUGG 68 CNTR	0 TOTAL AND CH EX 6/68
R1	0.24 (0.02)	BRICMAN 70 CNTR	0 TOTAL AND CH EX 6/70
R1	0.4	COOL 70 CNTR	K-P, D TOTAL 10/70

R2 $\Lambda(2100)$ INTO KBAR N PI

R2	SEEN	BOCK 65 HBC	(P2) 6/68
----	------	-------------	-----------

R3 $\Lambda(2100)$ FROM KBAR N INTO LAMBDA ETA

R3	(0.09) OR LESS	FLATTE 2 67 HBC	SQRT(P1*P3) 0 K-P TO LAM ETA 6/68
----	----------------	-----------------	-----------------------------------

R4 $\Lambda(2100)$ INTO (LAMBDA OMEGA)/TOTAL

R4	(0.1) OR LESS	FLATTE 1 67 HBC	(P4) 0 K-P TO LAM OMEGA 8/67
----	---------------	-----------------	------------------------------

REFERENCES FOR $\Lambda(2100)$ (PROD. EXP.)

BOCK 65 PL 17 166 +COOPER,FRENCH,KINSON, + (CERN,SACLAY)
 FLATTE 1 67 PR 155 1517 S M FLATTE (LRL)
 FLATTE 2 67 PR 163 1441 S M FLATTE, C G WOHL (LRL)
 BUGG 68 PR 168 1466 +GILMORE,KNIGHT, + (RHEL,BIRM,CAVE) I

BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN,CAEN,SACLAY)
 COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
 LU 70 PR D2 1846 +GREENBERG, HUGHES, MINEHART, MORI, + (YALE)

PAPERS NOT REFERRED TO IN DATA CARDS

COOL 66 PRL 16 1228 +GIACOMELLI,KYCIA,LENTIC,LI,LUNDBY,+ (BNL) I
 SUPERSEDED BY COOL 70.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Lambda(2350)$, $\Lambda(2585)$, Σ^+ , Σ^- , Σ^0 , $\Sigma(1385)$

$\Lambda(2350)$ BUMPS

42 Y*(2350, JP=) I=0 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
DAUM 68 FAVORS JP=7/2- OR 9/2+. BRICHAN 70 FAVORS 9/2+.
LASINSKI 71 SUGGESTS THREE STATES IN THIS REGION
USING A PGMERK + RESONANCES MODEL. THERE IS NOW ALSO
ONE FORMATION EXPERIMENT WHICH WE INCLUDE HERE, DE BELLEFCN 76, WHICH
FINDS 9/2- FRGM A DPWA OF KBAR N TO SIGMA P1.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Bugg, Brichan, Cool, and Bellefcn data.

Table with 4 columns: W, mass (MeV), width (MeV), and production experiments. Rows include Bugg, Brichan, Cool, and Bellefcn data.

Table with 4 columns: P1, P2, mass (MeV), and production experiments. Rows include decay masses for Kbar N and Sigma P1.

Table with 4 columns: R1, R2, mass (MeV), and production experiments. Rows include branching ratios for Kbar N into Sigma P1.

REFERENCES FOR Y*(2350) (PROD. EXP.)
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHEL, BIRM, CAVE) I
DAUM 68 NP 87 19 +ERNE, LAGNAUX, SENS, STEUER, UDO (CERN)JP
BRICHAN 70 PL 318 152 +FERRO LUZZI, PERRERAU, + (CERN, CAEN, SACLAY) I
COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
LU 70 PR D2 1846 +GREENBERG, HUGES, MINEHART, MORI, + (YALE) I
BELLEFCN 76 SBMTD. TO NP DE BELLEFCN, BILLOIR, BERTHON+ (CDF+SACL) IJP
PAPERS NOT REFERRED IN DATA CARDS
+GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, + (BNL) I
COOL 66 PRL 16 1228 T A LASINSKI (EFI) IJP
SUPERSEDED BY COOL 70. DE BELLEFCN, BERTHON, BILLOIR+ (CDF+SACL) I
LASINSKI 71 NP 829 125 PRESENTLY LISTED UNDER Y*(2250), BUT ISOSPIN UNDETERMINED.

$\Lambda(2585)$ BUMPS

7 Y*(2585, JP=) I=0 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Abrams and Lu data.

Table with 4 columns: W, mass (MeV), width (MeV), and production experiments. Rows include Abrams and Lu data.

Table with 4 columns: P1, P2, mass (MeV), and production experiments. Rows include decay masses for Kbar N.

Table with 4 columns: R1, R2, mass (MeV), and production experiments. Rows include branching ratios for Kbar N into Tctal.

REFERENCES FOR Y*(2585) (PROD. EXP.)
COOL 66 PRL 16 1228 +GIACOMELLI, KYCIA, LEONTIC, LUNDBY + (BNL) I
ABRAMS 70 PR 10 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I
BRICHAN 70 PL 318 152 +FERRO LUZZI, PERRERAU, + (CERN, CAEN, SACLAY) I
LU 70 PR D2 1846 +GREENBERG, HUGES, MINEHART, MORI, + (YALE) I
PAPERS NOT REFERRED TO IN DATA CARDS

S=-1 I=1 HYPERON STATES (Σ)

Σ^+ 19 SIGMA+(1189, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

Σ^- 20 SIGMA-(1198, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

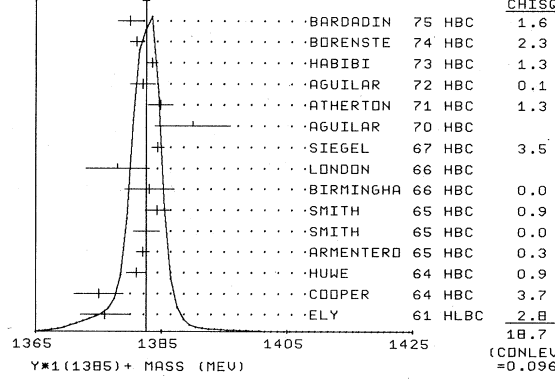
Σ^0 21 SIGMA0(1193, JP=1/2+) I=1
SEE STABLE PARTICLE DATA CARD LISTINGS

$\Sigma(1385)$ 43 Y*(1385, JP=3/2+) I=1

FOR DISCUSSION OF INCONSISTENCY OF ERRORS AND OUR
MODIFICATIONS, SEE NUTE CN K*(892)
FOR THE TABLES WE USE ONLY THE UNSTARRED DATA, WHICH
ATTEMPTS TO OBTAIN THE SEPARATE CHARGE-STATE MASSES AN
WIDTHS. SEE HOWEVER THE IDEOGRAMS INSERTED IN LISTING
THESE INDICATE SERVICUS SYSTEMATICS, PERHAPS ARISING FROM INTERFERENCE E
FFECTS THAT CHANGE WITH PRODUCTION MECHANISM AND BEAM MOMENTUM.

Table with 4 columns: M, mass (MeV), width (MeV), and production experiments. Rows include Alston, Borge, Martin, Colley, Baltay, Musgrave, Ammann, Atherton, Curtis, Thmas, Borenste, and others.

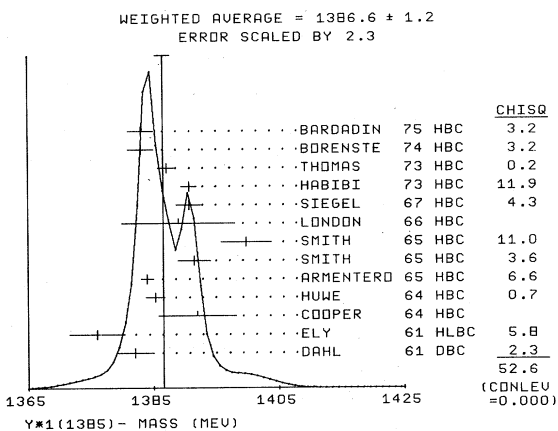
WEIGHTED AVERAGE = 1382.53 ± 0.50
ERROR SCALED BY 1.2



Baryons

$\Sigma(1385)$

M-	93	1382.0	3.0	DAHL	61 DBC	-	K-C 0.45 BEV/C		
M-	E 224	1376.0	4.4	ELY	61 HLBC	-			
M-	E			COOPER	64 HBC	-	BY US, BECAUSE < STATIST. ERR.	10/69	
M-	200	1392.0	6.2	COOPER	64 HBC	-			
M-	1086	1385.3	1.5	HUWE	64 HBC	-			
M-	1380	1384.0	1.0	ARMENTERO	65 HBC	-			
M-	S 120	1391.9	2.6	SMITH	65 HBC	-	K-P 1.8 BEV/C	9/66	
M-	S 58	1399.8	4.0	SMITH	65 HBC	-	K-P 1.95 BEV/C	9/66	
M-	S			SMITH	65 HBC	-	BY US, BECAUSE < STATIST. ER.	10/69	
M-	S			SMITH	65 HBC	-	BY US, BECAUSE < STATIST. ER.	10/69	
M-	1389.0		9.0	LNDCN	66 HBC	-		7/66	
M-	370	1390.7	2.0	SIEGEL	67 HBC	-	K-P AT 2.1 GEV/C	10/69	
M-	1900	1390.7	1.2	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73	
M-	722	1387.1	1.3	THOMAS	73 HBC	-	PI-P TO PI-K+L M	9/73	
M-	3060(1389.1)	(1.1)		BERTHON	74 HBC	-	0 QUASI 2 BODY CS	10/74*	
M-	2303	1385.3	2.1	BORENSTE	74 HBC	-	+OK-P TO(1385)+PIS	10/74*	
M-	1383.		2.1	BARDADIN	75 HBC	-	K-P 14.3 GEV/C	1/76*	
M-	AVG	1386.6	1.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.3)					
M-	STUDENT	1386.15	0.80	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
(SEE IDEOGRAM BELOW)									



43 (Y*-) - (Y**+) MASS DIFFERENCE (MEV)

D R	(0.0)	(4.2)	ELY	61 HLBC	+ -	K-P 1.11 BEV/C	8/66
D R	(17.1)	(7.1)	COOPER	64 HBC	-		10/69
D R	(4.3)	(2.2)	HUWE	64 HBC	+ -	K-P 1.22 BEV/C	8/66
D R	(2.0)	(1.5)	ARMENTERO	65 HBC	+ -	K-P 1.9-1.2 BEV/C	8/66
D R	(7.2)	(2.1)	SMITH	65 HBC	+ -	K-P 1.8 BEV/C	9/66
D R	(17.2)	(2.0)	SMITH	65 HBC	+ -	K-P 1.95 BEV/C	9/66
D R	(11.0)	(9.0)	LONDON	66 HBC	+ -	K-P 2.24 BEV/C	8/66
D R	9.0	6.0	LNDCN	66 HBC	-	LAMBDA 3 PI EVTS	7/66
D R	(6.3)	(2.0)	SIEGEL	67 HBC	-	K-P AT 2.1 GEV/C	10/69
D R	(7.2)	(1.4)	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73

REDUNDANT WITH DATA IN MASS LISTING.

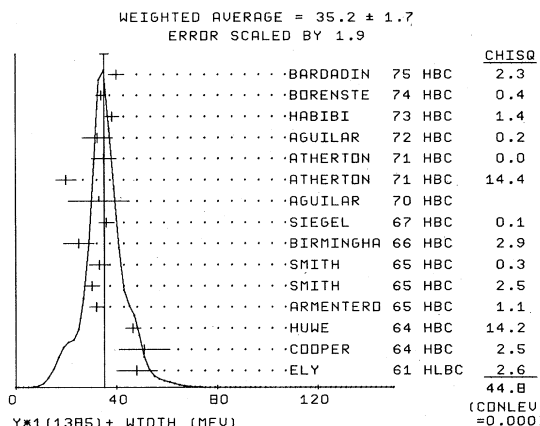
43 Y*1(1385) WIDTH (MEV)

W	(64.0)		ALSTEN	60 HBC	+ -				
W	(40.0)		BERGE	61 HBC	+ -				
W	(20.0)	OR LESS	MARTIN	61 HBC	+0				
W	(80.0)	(10.0)	COLLEY	62 HLBC	-0				
W	(30.0)	(9.0)	CURTIS	63 DSPK	0				
W	(26.0)	(5.0)	BALTAY	65 HBC	+ -			7/66	
W	(38.0)	(9.0)	MUSGRAVE	65 HBC	+0			7/66	
W	61.	10.	AMMANN	73 DBC	-	K-N 4.5 GEV/C	9/73		
W	89.	23.	ATHERTON	75 HBC	-	PBAR P 5.7 GEV/C	1/76*		
W	AVG	65.5	10.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)					
W	STUDENT	65.2	10.2	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					

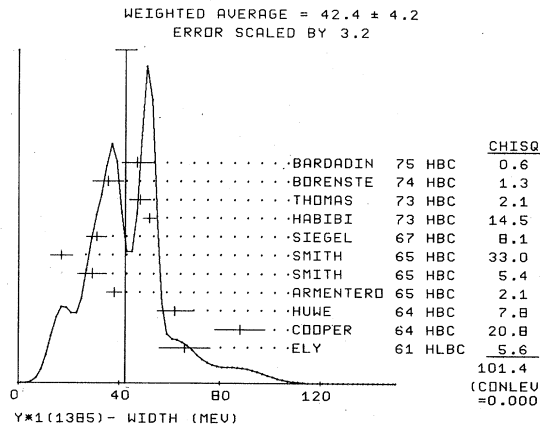
W0	330	39.3	6.1	THOMAS	73 HBC	-	PI-P TO PI0KOLM	9/73	
W0	3100	CONSISTENT WITH	W0=W**+*	BORENSTE	74 HBC	-	+OK-P TO(1385)+PIS	10/74*	
W+	48.0	8.0	ELY	61 HLBC	+				
W+	51.0	10.0	COOPER	64 HBC	+				
W+	46.5	3.0	HUWE	64 HBC	+				
W+	32.0	3.0	ARMENTERO	65 HBC	+				
W+	30.3	3.1	SMITH	65 HBC	+	K-P 1.8 BEV/C	9/66		
W+	33.1	3.8	SMITH	65 HBC	+	K-P 1.95 BEV/C	9/66		
W+	40	25.0	6.0	BIRMINGHA	66 HBC	+	K-P 3.5	9/67	
W+	1260	36.0	3.0	SIEGEL	67 HBC	-	K-P AT 2.1 GEV/C	10/69	
W+	33.	12.0	AGUILAR	70 HBC	+	K-P 4 GEV/SIG.PI	5/70		
W+	40	20.	4.	ATHERTON	71 HBC	-	LAM PI* + C.C.	10/71	
W+	AVG	35.2	3.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)					
W+	STUDENT	35.0	1.2	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
(SEE IDEOGRAM BELOW)									
W+ 1	RESULTS FROM LAM PI* + PI- AND LAM PI* PI- PIO COMBINED BY US.								
W+ R	FIT B.W. AND NO BCKGRD								
W+ R	40	35.2	3.7	ATHERTON	71 HBC	-	LAM PI* + C.C.	10/71	
W+ R	40	32.5	6.0	AGUILAR	72 HBC	-	K-P TO LAM+PIS	10/74*	
W+ R	2300	38.3	2.6	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73	
W+ 1	6846	34.	2.	BORENSTE	74 HBC	-	+OK-P TO(1385)+PIS	10/74*	
W+ 1				BARDADIN	75 HBC	-	K-P 14.3 GEV/C	1/76*	

Data Card Listings

For notation, see key at front of Listings.



W-	(40.0)		DAHL	61 DBC	-				
W-	66.0	10.0	ELY	61 HLBC	-				
W-	88.0	10.0	COOPER	64 HBC	-				
W-	82.0	7.0	HUWE	64 HBC	-				
W-	38.0	3.0	ARMENTERO	65 HBC	-				
W-	29.2	5.7	SMITH	65 HBC	-	K-P 1.80 BEV/C	9/66		
W-	17.1	4.4	SMITH	65 HBC	-	K-P 1.95 BEV/C	9/66		
W-	370	31.0	4.0	SIEGEL	67 HBC	-	K-P AT 2.1 GEV/C	10/69	
W-	1900	51.9	2.5	HABIBI	73 HBC	-	K-P TO 2PI LAM	9/73	
W-	722	48.2	4.0	THOMAS	73 HBC	-	PI-P TO PI-KOLM	9/73	
W- 1	2303	35.5	6.	BORENSTE	74 HBC	-	+OK-P TO(1385)+PIS	10/74*	
W-		47.	6.	BARDADIN	75 HBC	-	K-P 14.3 GEV/C	1/76*	
W-	AVG	42.4	4.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.2)					
W-	STUDENT	42.7	2.7	AVERAGE USING STUDENT10(H/1.11) -- SEE TEXT					
(SEE IDEOGRAM BELOW)									



43 Y*1(1385) REAL PART OF POLE POSITION

RE+	1379.0	1.0	LICHTENB 74	+	EXTRAP HABIBI73	4/75*
RE-	1383.0	1.0	LICHTENB 74	-	EXTRAP HABIBI73	4/75*

43 Y*1(1385) IMAGINARY PART OF POLE POSITION

IM+	17.5	1.5	LICHTENB 74	+	EXTRAP HABIBI73	4/75*
IM-	22.5	1.5	LICHTENB 74	-	EXTRAP HABIBI73	4/75*

43 Y*1(1385) PARTIAL DECAY MODES

PI	Decay Masses
P1	Y*1(1385) INTO LAMBDA PI 1115+ 139
P2	Y*1(1385) INTO SIGMA PI 1197+ 139
P3	Y*1(1385) INTO LAMBDA GAMMA 1115+ 0
P4	Y*1(1385) INTO KBAR N 493+ 938
P5	Y*1(1385) INTO SIGMA GAMMA 1197+ 0

Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(1385), Σ(1440), Σ(1480), Σ(1580)

43 Y*(1385) BRANCHING RATIOS

Table with columns for particle name, branching ratio, and reference. Includes entries for Y*(1385) INTO (SIGMA PI)/(LAMBDA PI) and Y*(1385) INTO LAMBDA GAMMA.

WEIGHTED AVERAGE = 0.140 ± 0.015
ERROR SCALED BY 1.0

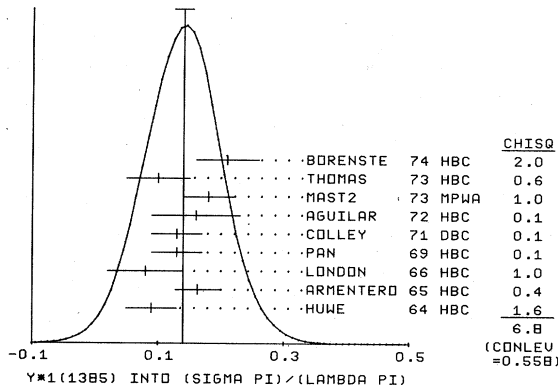


Table with columns for particle name, branching ratio, and reference. Includes entries for Y*(1385) FROM KBAR N TO LAMBDA PI and Y*(1385) INTO (LAMBDA GAMMA)/(LAMBDA PI).

REFERENCES FOR Y*(1385)

Large table of references for Y*(1385) from various researchers and institutions, including ALSTON, BASTIEN, BERGE, DAHL, ELY, MARTIN, ALSTON, COLLEY, CURTIS, COOPER, HUWE, ARMENTEROS, BIRNINGHAM, LONDON, SIEGEL, PAN, AGUILAR, ATHERTON, COLLEY, AGUILAR, MEISNER, AMMANN, MAST2, HABIBI, THOMAS, BERTHON, BARDADIN, COLAS, MALAMUD, SHAFER, HUNGERBU.

Σ(1440) BUMPS

80 Y*(1440, JP= 1/2-) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

CLINE 68 FINC A NARROW PEAK AT 1440 MEV (JUST ABOVE THE KBAR N THRESHOLD) IN THE LAMBDA PI INVARIANT MASS FOR K- D TO LAMBDA PI- P EVENTS. THEY DISCUSS ALTERNATE INTERPRETATIONS -- THAT IT IS A RESONANCE OR A KINEMATIC EFFECT. IN CLINE 68 THE K- BEAM MOMENTUM IS 0.4 GEV/C. IN A STUDY OF THE SAME REACTION WITH A MOMENTUM OF 1.1 GEV/C, ALEXANDER 69 FIND NO PEAK. IN ADDITION, THEY ARE ABLE TO EXPLAIN THE RESULTS OF BOTH EXPERIMENTS WITHOUT INVOKING A NEW RESONANCE. A REANALYSIS OF THE CLINE 68 DATA MADE BY BUNNELL 70 SHOW AGREEMENT OF THE DATA WITH THE ALEXANDER 69 INTERPRETATION.

REFERENCES FOR Y*(1440) (PROD. EXP.)

Table of references for Y*(1440) production experiments, listing authors like CLINE, ALEXANDER, BUNNELL and their respective publications.

Σ(1480) BUMPS

23 Y*(1480, JP= 1/2-) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

PEAKS ARE SEEN IN LAMBDA PI AND SIGMA PI SPECTRA IN THE REACTION PI+P TO K+ PI Y AT 1.7 GEV/C. ALSO THE Y POLARIZATION OSCILLATES IN THE SAME REGION.

SEE MILLER 70 FOR A DISCUSSION OF THIS STATE. HE SUGGESTS A POSSIBLE ALTERNATE EXPLANATION IN TERMS OF A REFLECTION OF N(1216) DECAY TO LAMBDA K. HOWEVER, SUCH AN EXPLANATION FOR THE K+ SIGMA+ PI0 CHANNEL SEEMS UNLIKELY (SEE PAN 70). IN TERMS OF KN(1670) DECAY INTO SIGMA K, IN ADDITION SUCH REFLECTIONS WOULD ALSO HAVE TO ACCOUNT FOR THE OSCILLATION OF THE Y POLARIZATION IN THE 1480 MASS REGION. HANSON 71, WITH FEWER DATA THAN PAN 70, CAN NEITHER CONFIRM NOR DENY THE EXISTENCE OF THIS STATE. MAST 75 SEES NO STRUCTURE IN THIS MASS REGION IN K- P TO LAMBDA PI0.

23 Y*(1480) MASS (MEV) (PROD. EXP.)

Table of Y*(1480) mass measurements from various experiments, listing authors like HANSON, PAN, CLINE and their results.

23 Y*(1480) WIDTH (MEV) (PROD. EXP.)

Table of Y*(1480) width measurements from various experiments, listing authors like HANSON, PAN, CLINE and their results.

23 Y*(1480) PARTIAL DECAY MODES (PROD. EXP.)

Table of partial decay modes for Y*(1480), listing decay channels like INTO KEAR N, INTO LAMBDA PI, INTO SIGMA PI.

23 Y*(1480) BRANCHING RATIOS (PROD. EXP.)

Table of branching ratios for Y*(1480) from various experiments, listing authors like HANSON, PAN, CLINE and their results.

REFERENCES FOR Y*(1480) (PROD. EXP.)

Table of references for Y*(1480) production experiments, listing authors like PAN, CLINE, YU-LI, MILLER, HANSON, MAST and their respective publications.

Σ(1580)

00 Y*(1580, JP=3/2-) I=1 D13 4/75*

OBSERVED IN K- N I=1 TOTAL CS AT BNL (LI 73, CARROLL 73), AND IN PWA OF K- P --> LAMBDA PI FOR CM ENERGIES=1560-1600 MEV BY LITCHFIELD 74. LITCHFIELD 74 FINDS JP=3/2-.

00 Y*(1580) MASS (MEV)

Table of Y*(1580) mass measurements, listing authors like LITCHFIELD and their results.

00 Y*(1580) WIDTH (MEV)

Table of Y*(1580) width measurements, listing authors like LITCHFIELD and their results.

Baryons

$\Sigma(1580)$, $\Sigma(1620)$

Data Card Listings

For notation, see key at front of Listings.

```

00 Y*1(1580) PARTIAL DECAY MODES
-----
P1 Y*1(1580) INTO KBAR N          497+ 939
P2 Y*1(1580) INTO LAMBDA PI      1115+ 139
P3 Y*1(1580) INTO SIGMA PI       1197+ 139
-----
00 Y*1(1580) BRANCHING RATIOS
-----
R1 Y*1(1580) INTO KBAR N/TOTAL      (P1)
R1 L +.03 .01 LITCHFIELD 74 DPWA  KBAR N MULTICHNL 4/75*
R1 L MAIN EFFECT OBSERVED BY LITCHFIELD 74 IS IN PI LAMBDA FINAL STATE,
R1 L KBAR N AND SIGMA PI COUPLINGS ALSO ESTIMATED FROM MULTICHANNEL FIT 4/75*
R1 L INCLUDING TOTAL CROSS SECTION DATA (LI 73). 4/75*
R2 Y*1(1580) FROM KBAR N TO LAMBDA PI SQRT(P1*P2) 4/75*
R2 L +.10 .02 LITCHFIELD 74 DPWA  O K- P TO LAM PI 4/75*
R3 Y*1(1580) FROM KBAR N TO SIGMA PI SQRT(P1*P3) 4/75*
R3 L +.03 .04 LITCHFIELD 74 DPWA  KBAR N MULTICHNL 4/75*
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*****
REFERENCES FOR Y*1(1580)
LITCHFIE 74 PL 51E 509 LITCHFIELD (CERN)IJP
PAPERS NOT REFERRED TO IN CATA CARDS.
CARROLL 73 APS BRKLY MTG 208 CARRCLL,CHIANG,KYCIA,LI,MAZUR,MICHAEL+(BNL)I
LI 73 PURDUE CONF. 283 LI (BNL)I
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Note on $\Sigma(1620)$

This state was first suggested by the BNL-CCNY collaboration (CRENNELL 68) who presented evidence for it in the reaction $K^- n \rightarrow \Sigma(1620)^+ \pi^- \pi^-$ with $\Sigma(1620)^+$ decaying into $\Lambda\pi^+$. Since then there have been conflicting reports about this state.

Total Cross-Section Experiment

A measurement of the $K^- p$ and $K^- d$ total cross sections in the 0.4 to 1.1 GeV/c range has been reported by the BNL group (CARROLL 73, LI 73). A clear bump about 40 MeV wide and 3-4 mb high [corresponding to $(J+1/2)x \sim 0.1$] is seen at a c.m. energy of 1590 MeV in the $I=1 K^- N$ cross section.

Formation Experiments

There is evidence from several partial-wave analyses for one or two fairly narrow states within ~ 50 MeV of the effect seen in production; see the entries for $\Sigma(1580, 3/2^-)$, $\Sigma(1620, 1/2^-)$, and $\Sigma(1660, 1/2^+)$. Note however that the various analyses do not agree on the widths and branching ratios of these states.

Production Experiments

A good review of the production experiments has been given by MILLER 70. Here the evidence is only in the $\Lambda\pi$ channel. The BNL-CCNY collaboration, with increased data, CRENNELL 69, still claim the effect in the $\Lambda\pi$ channel (no evidence seen in $\bar{K}N$ or $\bar{K}N\pi$). SABRE 70 studied the same reaction at 3.0

GeV/c with comparable statistics and do not see any evidence for it in the $\Lambda\pi$ channel; on the contrary, they believe it to be a spurious peak resulting from misidentified Σ^0 from the production of $\Sigma(1670)$ decaying into $\Sigma^0 \pi^+$. CRENNELL 69 give counter arguments to show that this is not the case in their data and the controversy goes on. AMMANN 70 studied the same reaction at 4.5 GeV/c and report a state at 1640 MeV, again decaying only into $\Lambda\pi$ (no evidence seen in $\Sigma\pi$ or $\bar{K}N$ channels). Upper limits on production cross sections for a 25 GeV/c Σ^- beam are reported by HUNGERBUHLER 74.

In conclusion, for understanding of the $\Sigma(1620)$ we probably have to wait for more data and for a more complete understanding of the entire mass region from 1600 to 1700 MeV. The closeness of the $\Sigma(1620)$ mass to 1670 MeV is suggestive that this effect may be related to what goes on in that region (see discussion below).

$\Sigma(1620)$ S11
 32 Y*1(1620, JP=1/2- I=1 THE S11 STATE AT 1647 MEV REPORTED BY VANHORN 75 IS INTERMEDIATE IN MASS BETWEEN THE SIGMA(1620) AND SIGMA(1750). WE TENTATIVELY LIST IT UNDER SIGMA(1750).

32 Y*1(1620) MASS (MEV)					
M	(1620.)				
M	1630.0	(10.0)	KIM LANGBEIN	71 DPWA 72 IPWA	K-MATRIX ANAL. MULTICHANNEL 3/71 12/72

32 Y*1(1620) WIDTH (MEV)					
W	(40.1				
W	65.0	(20.0)	KIM LANGBEIN	71 DPWA 72 IPWA	K-MATRIX ANAL. MULTICHANNEL 3/71 12/72

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32 Y*1(1620) PARTIAL DECAY MODES
-----
P1 Y*1(1620) INTO KBAR N          497+ 939
P2 Y*1(1620) INTO SIGMA PI      1197+ 139
P3 Y*1(1620) INTO LAMBDA PI     1115+ 134
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32 Y*1(1620) BRANCHING RATIOS
-----
R1 Y*1(1620) INTO KBAR N          (P1)
R1 (0.05) KIM 71 DPWA K-MATRIX ANAL. 3/71
R1 A 0.05 OR LESS WONG 71 DPWA K-MATRIX ANAL. 10/71
R1 0.22 (0.02) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R1 A K-MATRIX FIT(NEGLECTS 3-BODY CHANNELS) REQUIRES NO RESONANCE 10/71
R2 Y*1(1620) FROM KEAR N TO SIGMA PI SQRT(P1*P2)
R2 (0.08) KIM 71 DPWA K-MATRIX ANAL. 3/71
R2 0.40 (0.06) LANGBEIN 72 IPWA MULTICHANNEL 12/72
R3 Y*1(1620) FROM KBAR N TO LAMBDA PI SQRT(P1*P3)
R3 (0.15) KIM 71 DPWA K-MATRIX ANAL. 3/71
R3 NOT SEEN BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
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*****
REFERENCES FOR Y*1(1620)
KIM 71 PRL 27 356 J K KIM (HARV)IJP
ALSO 70 DUKE 161 J. K. KIM (HARV)IJP
WONG 71 NC 2A 353 N S WONG (YALE)IJP
LANGBEIN 72 NP B47 477 +WAGNER (PIMI)IJP
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN,RHEL)IJP
PAPERS NOT REFERRED TO IN DATA CARDS
VANHORN 75 NP B87 145 A. J. VAN HORN (LBL)IJP
ALSO 75 NP B87 157 A. J. VAN HORN (LBL)IJP
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Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(1620), Σ(1660), Σ(1670)

1620 MEV REGION - PRODUCTION EXPERIMENTS

78 Y*(1620, JP=) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

THIS RESONANCE NEEDS CONFIRMATION. THE RESULTS OF CRENNELL 69 AT 3.9 GEV/C ARE NOT CONFIRMED BY THE SABRE COLLABORATION AT 3.0 GEV/C (SABRE 70). HOWEVER IN AN EXPERIMENT AT 4.5 GEV/C, AMMANN 70 SEE A PEAK AT 1642 MEV WHICH ON THE BASIS OF BRANCHING RATIOS THEY DO NOT ASSOCIATE WITH THE Y*(1670). SEE MILLER 70 FOR A REVIEW OF THESE CONFLICTS.

78 Y*(1620) MASS (MEV) (PROD. EXP.)

Table with columns M, N, mass values, and experiment names like CRENNELL, BLUMENFEL, AMMANN.

78 Y*(1620) WIDTH (MEV) (PROD. EXP.)

Table with columns M, N, width values, and experiment names like CRENNELL, BLUMENFEL, AMMANN.

78 Y*(1620) PARTIAL DECAY MODES (PROD. EXP.)

Table listing decay modes like Y*(1620) INTO KBAR N, LAMBDA PI, etc.

78 Y*(1620) BRANCHING RATIOS (PROD. EXP.)

Table listing branching ratios for various decay channels like (P4)/(P3), (P1)/(P2), etc.

REFERENCES FOR Y*(1620) (PROD. EXP.)

CRENNELL 68 PRL 21 648 + DELANEY, FLAMINGO, KARSHON, + (BNL,CUNY) I
BLUMENFEL 69 PL 298 58 + BLUMENFELD, KALBFLEISCH (BNL) I
CRENNELL 69 LUND PAPER 183 + KARSHON, LAI, ONEIL, SCARR, + (BNL,CUNY) I
RESULTS ARE QUOTED IN LEVI SETTI 69.

Σ(1660)

79 Y*(1660, JP=1/2+) I=1 P11
SEE THE MINI-REVUE AT THE START OF THE Y* LISTINGS.

THE PARTIAL-WAVE ANALYSIS OF K- N TO SIGMA PI BY ARMENTEROS 70 SUGGEST SUCH A RESONANCE. NOW FOUND ALSO IN SOME, BUT NOT ALL, MORE RECENT ANALYSES.

79 Y*(1660) MASS (MEV)

Table with columns M, N, mass values, and experiment names like ARMENTEROS, LEA, HART, BAILLON, PONTE, VANHORN, RLIC.

79 Y*(1660) WIDTH (MEV)

Table with columns M, N, width values, and experiment names like ARMENTEROS, KIM, LEA, HART, BAILLON, PONTE, VANHORN, RLIC.

79 Y*(1660) PARTIAL DECAY MODES

Table listing decay modes like Y*(1660) INTO KBAR N, SIGMA PI, LAMBDA PI.

79 Y*(1660) BRANCHING RATIOS

Table listing branching ratios for various decay channels like (P1)/(P2), (P1)/(P3), etc.

REFERENCES FOR Y*(1660)

ARMENTEROS 70 DUKE 123 ARMENTEROS, BAILLON, + (CERN,HEIDEL IJIP
KIM 71 PRL 27 356 J. K. KIM (HARV) IJIP
ALSO 70 DUKE 161 J. K. KIM (HARV) IJIP
HART 73 PURDUE CGNF. 311 +RICE, BACASTON, FUNG, + (TENN+UCR+MASA+BUFF) IJIP
LEA 73 NP 856 77 +MARTIN, MOORHOUSE + (RHEL+LOUC+GLAS+AARHUS) IJIP
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IJIP
PONTE 75 PRD 12 2597 +HERTZBACH, BUTTON-SHAFER* (MASA+TENN+UCR) IJIP
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL) IJIP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL) IJIP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS + (RHEL+LOUC) IJIP

Note on Σ(1670)

Production Experiments

The measured Σπ/Σππ branching ratio for produced Σ(1670)'s is strongly dependent on momentum transfer. This was first discovered by EBERHARD 69 who suggested the existence of two Y1* with the same mass and quantum numbers; one object with a large Σππ [mainly Λ(1405)π] decay mode produced peripherally, and another one with a large Σπ decay mode produced at larger angles. This observation has been confirmed by AGUILAR-BENITEZ 70, APSELL 74, ESTES 74, and BLOKZIJL 75. When determined, the most likely quantum numbers are 3/2- [for both Σπ and Λ(1405)π]. There is also the possibility of a third Y1* state, referred to as Σ(1690) in the Data Card Listings, with a large Λπ/Σπ branching ratio and somewhat larger mass. The large branching ratio is the main justification for this hypothesis and needs confirmation. These problems have been reviewed by EBERHARD 73 and MILLER 70.

Baryons
Sigma(1670)

Data Card Listings
For notation, see key at front of Listings.

Formation Experiments

Two states are also observed near this mass in formation. One of these, the Sigma(1670, 3/2-), has the same quantum numbers as those observed in production and a large Sigma pi/Sigma pi pi branching ratio. It may well correspond to the produced Sigma(1670) seen at larger angles. The other, the Sigma(1660, 1/2+) has different quantum numbers from those seen in production, and its Sigma pi/Sigma pi pi branching ratio is unknown. Thus its relation to the produced Sigma(1670) remains obscure. [See also the mini-review on Sigma(1620)].

Sigma(1670) 44 Y*(1670, JP=3/2-) I=1 D13

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE NOTE ABOVE

WELL ESTABLISHED RESONANCE. IT HAS BEEN SEEN IN BOTH FORMATION AND PRODUCTION EXPERIMENTS. HOWEVER THE BRANCHING RATIOS OBTAINED BY THESE TWO METHODS SHOW LARGE INCONSISTENCIES.

SEE LISTING OF PRODUCTION EXPERIMENTS BELOW

AS FOR THE QUANTUM NUMBERS, THE ANALYSES OF LAMBDA PI CHANNEL (IN FORMATION EXP.) AS WELL AS THE SIGMA PI CHANNEL AGREE ON JP=3/2-.

Table with 4 columns: Particle ID, Mass (MeV), Production Experiment, and Branching Ratio. Includes entries for Y*(1670) MASS (MEV) with various experiments like BERLEY, ARMENTER, etc.

Table with 4 columns: Particle ID, Width (MeV), Production Experiment, and Branching Ratio. Includes entries for Y*(1670) WIDTH (MEV) with various experiments like BERLEY, ARMENTER, etc.

44 Y*(1670) PARTIAL DECAY MODES

Table with 3 columns: Decay Mode, Decay Masses, and Branching Ratio. Includes entries like Y*(1670) INTO KEAR N, Y*(1670) INTO LAMBDA PI, etc.

44 Y*(1670) BRANCHING RATIOS

Table with 4 columns: Particle ID, Branching Ratio, Production Experiment, and Branching Ratio. Includes entries for Y*(1670) INTO (KEAR N)/TOTAL and Y*(1670) INTO (LAMBDA PI PI)/TOTAL.

Table with 4 columns: Particle ID, Production Experiment, Branching Ratio, and Branching Ratio. Includes entries for Y*(1670) INTO (SIGMA PI PI)/TOTAL, Y*(1670) INTO (Y*(1405) PI)/TOTAL, etc.

REFERENCES FOR Y*(1670)

Table with 4 columns: Reference Name, Reference Title, Reference Author, and Reference Journal. Includes entries like BERLEY 64 DUBNA CONF I 565, ARMENTEROS, BAILLON, etc.

Sigma(1670) BUMPS

51 Y*(1670, JP=) I=1 PRODUCTION EXPERIMENTS

SEE NOTE PRECEDING Y*(1670) PROBABLY THERE ARE TWO STATES AT SAME MASS WITH SAME QUANTUM NUMBERS, ONE DECAYING INTO SIGMA PI AND LAMBDA PI, THE OTHER INTO Y*(1405) PI. BRANCHING RATIOS NOT DISENTANGLED YET, WE LIST THEM TOGETHER FOR NOW.

51 Y*(1670) MASS (MEV) (PROD. EXP.)

Table with 4 columns: Particle ID, Mass (MeV), Production Experiment, and Branching Ratio. Includes entries like ALEXANDER 62 HBC, ALVAREZ 63 HBC, etc.

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Sigma(1670)$, $\Sigma(1690)$

M 1200 1688+-2. OR 1683+-5. BERTHON 74 HBC 0 QUASI 2 BODY CS 4/75*
M 3 1665. 3. BLOKZIJL 75 HBC K-P AT 4.2 GEV 1/76*
M 3 PARAMETERS DETERMINED FROM (SIGMA PI+) PRODUCTION. 1/76*

51 Y*(1670) MASS (MEV) (PROD. EXP.)
W (45.0) ALEXANDER 62 HBC -0
W 40.0 10.0 ALVAREZ 63 HBC +
W (30.0) (15.0) BUGG 68 CNTR 11/66

51 Y*(1670) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*(1670) INTO KBAR N 497+ 939
P2 Y*(1670) INTO LAMBDA PI 1115+ 139
P3 Y*(1670) INTO SIGMA PI 1197+ 139

51 Y*(1670) BRANCHING RATIOS (PROD. EXP.)
R1 Y*(1670) INTO (KBAR N)/(SIGMA PI) (P1)/(P3)
R1 0 (0.19) OR LESS ALVAREZ 63 HBC + K-P 1.15 BEV/C
R1 (0.51)+ .25 OR MORE SMITH 63 HBC -0

R2 Y*(1670) INTO (LAMB. PI)/(SIG PI) (P2)/(P3)
R2 130 (1.20) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R2 (1.2) SMITH 63 HBC -0
R2 0.15 0.07 MUSEL 64 HBC +

R3 Y*(1670) INTO (LAMB. PI PI)/(SIG PI) (P4)/(P3)
R3 90 (0.56) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R3 (0.17) SMITH 63 HBC -0
R3 (0.6) OR LESS LONDON 66 HBC + K-P AT 2.25 BEV/C 7/66

R4 Y*(1670) INTO (SIGMA PI PI)/(SIG PI) (P5)/(P3)
R4 180 (0.56) ALVAREZ 63 HBC + K-P 1.15 BEV/C
R4 LARGEST AT SMALL ANGLES ESTES 74 HBC 0 K-P 2.1+2.66GEV/C 11/75*

R5 Y*(1670) INTO (Y*(1405) PI)/(SIG PI) (P7)/(P3)
R5 50 3. 1.6 LONDON 66 HBC + K-P 2.25 BEV/C 7/66
R5 P 17 (0.58) (0.20) PRIMER 68 HBC + K-P 4.6-5. GEV/C 7/68

R6 Y*(1670) INTO (SIGMA PI)/(SIGMA PI PI) (P3)/(P5)
R6 .4 OR LESS BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67
R6 0.30 0.15 LONDON 66 HBC + K-P 2.25 BEV/C 7/66

R7 Y*(1670) INTO (Y*(1405) PI)/(SIGMA PI PI) (P7)/(P5)
R7 0.90 0.10 0.16 EBERHARD 65 HBC + K-P 2.45 BEV/C 7/66
R7 B 1.00 -0.02 APSELL 74 HBC K-P 2.87 GEV/C 4/75*

R8 Y*(1670) INTO (Y*(1405) PI)/(Y*(1385) PI) (P7)/(P6)
R8 (0.8) OR LESS EBERHARD 65 HBC + K-P 2.45 BEV/C 7/66
R9 Y*(1670) INTO (LAMBDA PI PI)/(SIGMA PI PI) (P4)/(P5)
R9 0.35 0.2 BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67

R10 Y*(1670) INTO (LAMBDA PI)/(SIGMA PI PI) (P2)/(P5)
R10 (1.2) OR LESS BIRMINGHAM 66 HBC + K-P AT 3.5 GEV/C 11/67
R11 Y*(1670) INTO (LAMBDA PI)/(LAMBDA PI + SIG PI) (P2)/(P2+P3)
R11 (0.6) OR LESS AGUILAR 70 HBC 5/70

R12 Y*(1670) INTO (Y*(1385) PI)/(SIGMA PI) (P6)/(P3)
R12 (.251)+/--.05) OR LESS BLOKZIJL 75 HBC K-P AT 4.2 GEV 1/76*

51 Y*(1670) QUANTUM NUMBER DETERMINATION (PROD. EXP.)
Q1 JP=3/2+ LEVEQUE 65 HBC INTO Y*(1405)+PI 11/68
Q3 JP=3/2- EBERHARD 67 HBC + INTO Y*(1405) PI 11/68
Q4 400 JP=3/2- BUTTCH-SH 68 HBC + INTO SIGZERO+PI 11/68

REFERENCES FOR Y*(1670) (PROD. EXP.)
ALEXANDE 62 CERN CONF 320 ALEXANDER, JACOBS, KALBFLEISCH, MILLER, + (LRL) I
ALVAREZ 63 PRL 10 184 +ALSTON, FERRO-LUZZI, HUWE, + (LRL) I
SMITH 63 ATHENS CONF 67 G A SMITH (LRL)

APSELL 74 PRD 10 1419 APSELL, FCRO, GUREVITCH, BRAN, UME, SYRA, TUFT II
BERTHON 74 NC 21A 146 BERTHON, TRISTRAM, + (CDF+RHEL+SACL+STRB) L
ESTES 74 LBL-3827 (THESIS) R. D. ESTES (LRL)

LEVEQUE 65 PL 18 69 + (SACLAY, EPOL, GLASGOW, LCLC, CFX, RHEL) JP
LEE 66 PRL 17 45 Y Y LEE, D D REEDER, R W HARTUNG (WISC) JP
EBERHARD 67 PR 163 1446 +PRIPSTEIN, SHIVELY, KRUSE, SWANSON (LRL, ILL) JP

58 Y*(1690) MASS (MEV) (PROD. EXP.)
M 30(1715.0) (12.0) COLLEY 67 HBC + K-P 6 GEV/C 8/67
M P 60(1694.0) (24.0) PRIMER 68 HBC + K-P 4.6-5 GEV/C 7/68
M N (1700.0) (6.0) SIMS 68 HBC - K-N TO LAMB PI 11/68

58 Y*(1690) WIDTH (MEV) (PROD. EXP.)
W 30 (100.0) (35.0) COLLEY 67 HBC + 8/67
W 60 (105.0) (35.0) PRIMER 68 HBC + 7/68
W N 46 (62.0) (14.0) SIMS 68 HBC - SEE NOTE N ABOVE 11/68
W 46 (25.0) (10.0) BLUMENFEL 69 HBC + 9/69
W (130.0) (25.0) MOTT 69 HBC + 9/69

58 Y*(1690) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*(1690) INTO KBAR N 497+ 939
P2 Y*(1690) INTO LAMBDA PI 1115+ 139
P3 Y*(1690) INTO SIGMA PI 1197+ 139
P4 Y*(1690) INTO Y*(1385) PI 1384+ 139
P5 Y*(1690) INTO LAMBDA PI PI (INCLUDING P4) 1115+ 139+ 139

58 Y*(1690) BRANCHING RATIOS (PROD. EXP.)
R1 Y*(1690) INTO (KBAR N)/(LAMBDA PI) (P1)/(P2)
R1 18 0.4 0.25 COLLEY 67 HBC + 6/30 EVENTS 8/67
R1 (0.2) OR LESS MOTT 69 HBC + 9/69

R2 Y*(1690) INTO (SIGMA PI)/(LAMBDA PI) (P3)/(P2)
R2 0.3 0.3 COLLEY 67 HBC + 4/30 EVENTS 8/67
R2 (0.4) OR LESS CL=.90 MOTT 69 HBC + 9/69

R3 Y*(1690) INTO (Y*(1385) PI)/(LAMBDA PI) (P4)/(P2)
R3 (0.5) OR LESS MOTT 69 HBC + 9/69
R4 Y*(1690) INTO (LAMBDA PI PI)/(LAMBDA PI) (P5)/(P2)
R4 0.5 0.25 COLLEY 67 HBC + 15/30 EVENTS 8/67
R4 2.0 0.6 BLUMENFEL 69 HBC + 31/15 EVENTS 9/69

R4 AVG 0.72 0.53 AVERAGE (ERROR INCLUDES SCALE FACTOR CF 2.5)
R4 STUDENT 0.67 0.28 AVERAGE USING STUDENT10(H/1.1) -- SEE TEXT

R5 Y*(1690) INTO (Y*(1385) PI)/(LAMBDA PI PI) (P4)/(P5)
R5 SMALL COLLEY 67 HBC + 8/67
R5 LARGE SIMS 68 HBC - K-N TO L2PI 11/68

Baryons

Σ(1690), Σ(1750), Σ(1765)

Data Card Listings

For notation, see key at front of Listings.

REFERENCES FOR Y*(1690) (PROD. EXP.)

COLLEY 67 PL 248 489
DERRICK 67 PRL 18 266
REPLACED BY HOIT 65.
PRIMER 68 PRL 20 610
SIMS 68 PRL 21 1413
BLUMENFE 69 PL 298 58
HOIT 69 PR 177 1966

AGUILAR 70 PRL 25 58

Σ(1750)

S11

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THIS STATE CORRESPONDS TO THE SIGMA ETA THRESHOLD BUMP
BUT ITS INTERPRETATION IN TERMS OF A RESONANCE IS NOT
CONCLUSIVE (JONES 74) -- MORE DATA ARE NEEDED. BY
ANALOGY WITH THE SIMILAR N ETA AND LAMBDA ETA THRESHOLD
EFFECTS, WHICH ARE ALMOST CERTAINLY RESONANCES, IT SEEMS VERY LIKELY
THAT THIS TOO IS A RESONANCE.

THERE IS ALSO EVIDENCE FOR THIS STATE IN MANY PARTIAL-WAVE ANALYSES,
MOSTLY IN THE LAMBDA PI CHANNEL. THIS EVIDENCE SHOULD BE CONSIDERED
WITH CAUTION AS THE MASSES AND/OR BRANCHING RATIOS OF THE REPORTED
STATES ARE OFTEN INCONSISTENT.

57 Y*(1750) MASS (MEV)

Table with columns for mass (MEV), decay modes, and references. Includes entries for NEAR SIGMA ETA THRESHOLD, ABOUT 1750.0, (1757.0) (10.0), etc.

57 Y*(1750) WIDTH (MEV)

Table with columns for width (MEV), decay modes, and references. Includes entries for ABOUT 50.0, ABOUT 80.0, (155.0) (10.0), etc.

57 Y*(1750) PARTIAL DECAY MODES

Table with columns for decay masses and partial decay modes. Includes entries for Y*(1750) INTO KBAR N, Y*(1750) INTO SIGMA ETA, etc.

57 Y*(1750) BRANCHING RATIOS

Table with columns for branching ratios and decay modes. Includes entries for Y*(1750) INTO (KBAR N)/TOTAL, Y*(1750) FROM KBAR N INTO SIGMA ETA, etc.

REFERENCES FOR Y*(1750)

CLINE, OLSSON (WISCONSIN) IIP
MEYER (RAPPOURTEUR) (SACLAY) IIP
ARMENTERO, BAILLON, + (CERN, HEIDEL) IIP
+LEVI SETTI, LASINSKI, OBERLACK+ (EFI+HEID) IIP
KIM (HARV) IIP
J. K. KIM (HARV) IIP
+WAGNER (MFM) IIP

BAXTER, BUCKINGHAM, CORBETT, CUNN, + (OXFORD) IIP
CHU, BARTLEY, + (SUNY PLATTSBURGH+TUFTS+BRAN) IIP
JONES (U. CHICAGO) IIP
DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
PREVOST, BARLOUTAUD, + (SACL+ CERN+HEID)

BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IIP
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL) IIP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL) IIP

RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IIP

PAPERS NOT REFERRED TO IN CATA CARDS

FERRO-LU 66 BERKELEY CONF 183 M FERRO LUZZI (RAPPOURTEUR) (CERN)
ARMENTERO 68 NP 88 183 ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IIP
ARMENTERO 69 LUND CNFN PAPER ARMENTEROS, BAILLON, + (CERN, HEIDEL, SACLAY) IIP
HARRISON 70 FSU-HEP TO 3 I W.C. HARRISON (THEIST) (FSU)

Σ(1765)

D15

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

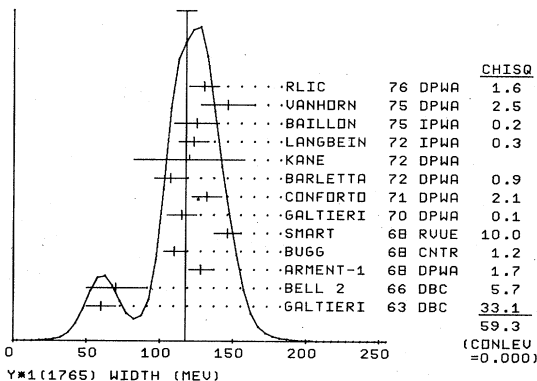
45 Y*(1765) MASS (MEV)

Table with columns for mass (MEV), decay modes, and references. Includes entries for 1765.0 10.0, 1755.0 10.0, 1760.0 10.0, etc.

45 Y*(1765) WIDTH (MEV)

Table with columns for width (MEV), decay modes, and references. Includes entries for 60.0 10.0, 70.0 20.0, 128.0 8.0, etc.

WEIGHTED AVERAGE = 117.5 ± 6.8
ERROR SCALED BY 2.3



Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(1765), Σ(1770)

45 Y*(1765) PARTIAL DECAY MODES

Table with columns for decay mode (P1-P7), mass (DECAY MASSES), and branching fractions.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P_i, as follows: The diagonal elements are P_i + 6P_i, where 6P_i = sqrt(6P_i * 6P_i), while the off-diagonal elements are the normalized correlation coefficients (6P_i * 6P_j) / (6P_i * 6P_j). For the definitions of the individual P_i, see the listings above; only those P_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Matrix of correlation coefficients for P1 through P6.

45 Y*(1765) BRANCHING RATIOS

ERRORS QUOTED BY EXPERIMENTERS DO NOT INCLUDE UNCERTAINTY DUE TO PARAMETRIZATION USED IN THE P.W.A. THEY SHOULD BE INCREASED.

Table of branching ratios for various decay channels (R1-R11) with associated uncertainties and references.

WEIGHTED AVERAGE = 0.409 ± 0.018
ERROR SCALED BY 2.8

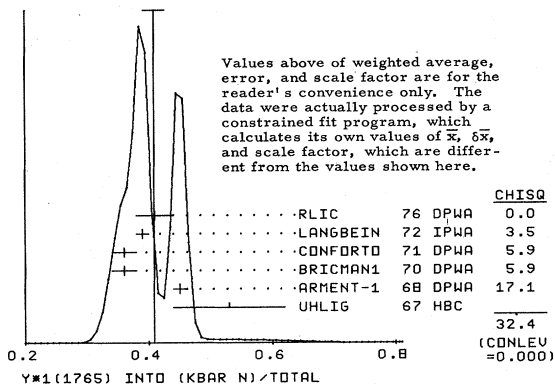


Table of partial decay modes for Y*(1765) from KBAR N into Lambda PI, including decay masses and branching fractions.

Table of branching ratios for Y*(1765) from KBAR N into Y*(1385) PI D-WAVE, including decay masses and branching fractions.

Table of partial decay modes for Y*(1765) from KBAR N into Sigma PI, including decay masses and branching fractions.

Table of references for Y*(1765) from various experiments and publications.

Σ(1770) 100 Y*(1770, JP=1/2+) I=1 P11

Table of mass (MEV) for Y*(1770) decay modes, including decay masses and branching fractions.

Table of width (MEV) for Y*(1770) decay modes, including decay masses and branching fractions.

Table of partial decay modes for Y*(1770) from KBAR N into Lambda PI, including decay masses and branching fractions.

Table of branching ratios for Y*(1770) from KBAR N into Lambda PI, including decay masses and branching fractions.

Baryons

$\Sigma(1770)$, $\Sigma(1840)$, $\Sigma(1880)$, $\Sigma(1915)$

R2 Y*1(1770) INTO (KBAR N)/TOTAL 1/76*
R2 .14 .04 RLIC 76 DPWA (P1) KBAR N MULTICHNL 1/76*

REFERENCES FOR Y*1(1770)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IJJP
RLIC 76 RL-75-182 (PRPNT) GOPAL, KALMUS, MCPHERSON, ROSS+ (RHEL+LOIC) IJJP

$\Sigma(1840)$

01 Y*1(1840, JP=3/2+) I=1 P13

SEE THE MINI-REVIEWS PRECEDING THE Y*0'S.
FOR THE TIME BEING, WE LIST ALL RESONANCE CLAIMS IN THE
P13 WAVE IN THE 1700-1900 MEV MASS REGION TOGETHER UNDER THIS HEADING.

01 Y*1(1840) MASS (MEV)
M 1840.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
M 1 (1720.0) (30.) BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

01 Y*1(1840) WIDTH (MEV)
W 120.0 (10.0) LANGBEIN 72 IPWA MULTICHANNEL 12/72
W 1 (120.) (30.) BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

01 Y*1(1840) PARTIAL DECAY MODES
P1 Y*1(1840) INTO KBAR N 497+ 939
P2 Y*1(1840) INTO SIGMA PI 1197+ 139
P3 Y*1(1840) INTO LAMBDA PI 1115+ 134

01 Y*1(1840) BRANCHING RATIOS
R1 Y*1(1840) INTO (KBAR N)/TOTAL (P1) 12/72
R1 0.37 (0.13) LANGBEIN 72 IPWA MULTICHANNEL
R2 Y*1(1840) FROM KBAR N INTO SIGMA PI SQR(P1*P2) 12/72
R2 0.15 (0.04) LANGBEIN 72 IPWA MULTICHANNEL

REFERENCES FOR Y*1(1840)
LANGBEIN 72 NP 847 477 *WAGNER (MPIM) IJJP
DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DESY, NORDITA, LOUC)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL) IJJP
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL) IJJP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL) IJJP

$\Sigma(1880)$

67 Y*1(1880, JP=1/2+) I=1 P11

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL
PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE
IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN
BARYON TABLE.

67 Y*1(1880) MASS (MEV)
M 1882.0 40.0 SMART 68 DPWA -0 K- N TO LAM PI 7/68
M (1850.0) BAILLEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
M ABOUT 1850.0 ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70

67 Y*1(1880) WIDTH (MEV)
W 222.0 150.0 SMART 68 DPWA -0 K- N TO LAM PI 7/68
W (200.0) BAILLEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
W ABOUT 30.0 ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70

Data Card Listings

For notation, see key at front of Listings.

67 Y*1(1880) PARTIAL DECAY MODES

P1 Y*1(1880) INTO KBAR N 497+ 939
P2 Y*1(1880) INTO SIGMA PI 1115+ 134
P3 Y*1(1880) INTO LAMBDA PI 1197+ 139

67 Y*1(1880) BRANCHING RATIOS

R1 Y*1(1880) INTO (KBAR N)/TOTAL (P1)
R1 (0.22) BAILLEY 69 DPWA 0 ELASTIC, CH EXCH 10/70
R1 (0.20) ARMENTERO 70 IPWA -0 ELASTIC, CH EXCH 6/70
R1 2 (.51) LEA 73 DPWA MULTICHNL K-MTRX 9/73

REFERENCES FOR Y*1(1880)
SMART 68 PR 169 1330 W M SMART (LRL) IJJP
BAILLEY 69 THESIS UCLRL-50617 DAVID SAAL BAILLEY (LRL LIVERMORE) IJJP
ARMENTERO 70 DUKE CONF 123 ARMENTEROS, BAILLON, + (CERN, HEIDEL) IJJP
GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL) IJJP
LITCHFIELD 70 NP 822 269 P J LITCHFIELD (RUTHERFORD) IJJP
KANE 72 PR 05 1583 D F KANE (LBL)

$\Sigma(1915)$

46 Y*1(1915, JP=5/2+) I=1 F15

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

THIS RESONANCE WAS FIRST SEEN IN THE TOTAL-CROSS-SECTION
MEASUREMENTS OF COOL 66, IN THIS ENTRY, HOWEVER,
WE LIST ONLY THE RESULTS FROM PARTIAL-WAVE ANALYSES.
SEE THE NEXT ENTRY FOR THE PARAMETERS OF PEAKS SEEN AROUND 1900-1950
MEV IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. WE MAKE THIS
SEPARATION BECAUSE ONLY THE PARTIAL-WAVE ANALYSES ISOLATE THE F15 WAVE
(OR AT LEAST ATTEMPT TO -- THE SIGNAL IS WEAK). THIS MASS REGION IS
COMPLICATED AND POORLY UNDERSTOOD AND THE PEAKS MAY CONTAIN MORE THAN
JUST THE Y*1(1915). SEE ALSO THE NOTE TO THE NEXT ENTRY.

46 Y*1(1915) MASS (MEV)

M 1902.0 11.0 SMART 68 DPWA -0 K- N TO LAMBDA PI 7/68
M 1910.0 20.0 BERTHON 70 DPWA 0 K-P TO LAMBDA PI 7/70
M 1900.0 15.0 BERTHON 70 DPWA 0 K-P TO SIGMA PI 10/70
M N 1936.0 (3.0) BRICMANI 70 DPWA SIGTOT, ELAS, CHEX 1/71
M 1903.0 10.0 COX 70 DPWA - K- N TO LAMBDA PI 6/70
M 1905.0 30.0 GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
M 1895.0 10.0 LITCHFIELD 70 DPWA -0 K- N TO LAMBDA PI 6/70
M B (1985.0) (21.0) ISLAM 71 DPWA KN--PI-SIG 12/72
M B DISCREPANCY DUE POSSIBLY TO INSUFFICIENT STATISTICS
M 1910. 15. LITCHFIELD 71 DPWA K-P TO KBAR N 10/71
M 1925.0 8.0 KANE 72 DPWA 0 K-P TO PI SIG 10/71
M 1920. 30. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

46 Y*1(1915) WIDTH (MEV)

W A (50.0) (20.0) ARMENTERO 67 DPWA 0 ELASTIC, CH EXCH 11/67
W 52.0 25.0 SMART 68 DPWA -0 K- N TO LAMBDA PI 7/68
W 60.0 20.0 BERTHON 70 DPWA 0 K-P TO LAMBDA PI 7/70
W 75.0 20.0 BERTHON 70 DPWA 0 K-P TO SIGMA PI 10/70
W 135.0 12.0 BRICMANI 70 DPWA SIGTOT, ELAS, CHEX 1/71
W 77.0 27.0 COX 70 DPWA - K- N TO LAMBDA PI 6/70
W 70.0 20.0 GALTIERI 70 DPWA 0 K-P TO LAMBDA PI 7/70
W 70.0 15.0 LITCHFIELD 70 DPWA -0 K- N TO LAMBDA PI 6/70
W B (159.0) (80.0) ISLAM 71 DPWA KN--PI-SIG 12/72
W 70. 15. LITCHFIELD 71 DPWA K-P TO KBAR N 10/71
W 146.0 22.0 KANE 72 DPWA 0 K-P TO PI SIG 10/71
W 70. 20. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

46 Y*1(1915) PARTIAL DECAY MODES

P1 Y*1(1915) INTO KBAR N 497+ 939
P2 Y*1(1915) INTO LAMBDA PI 1115+ 139
P3 Y*1(1915) INTO SIGMA PI 1197+ 139

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Sigma(1915)$, $\Sigma(1940)$

46 Y*(1915) BRANCHING RATIOS

Table with columns for particle name, mass, width, and branching ratios. Includes entries for Y*(1915) INTO (KBAR N)/TOTAL and FROM KBAR N INTO LAMBDA PI.

REFERENCES FOR Y*(1915)

Table listing references for Y*(1915) from various sources like ARMENTER, BERTHON, BRICMAN, etc.

1915 MEV REGION - PRODUCTION AND σ_{TOTAL} EXPTS

29 Y*(1915, J_p=) I=1 PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS. SEE THE NOTES TO THE Y*(1915) AND Y*(1940), WHICH IMMEDIATELY PRECEDE AND FOLLOW THIS ENTRY.

29 Y*(1915) MASS (MEV) (PRCD. EXP.)

Table showing mass measurements for Y*(1915) from various experiments like BUGG, BRICMAN, COOL, etc.

29 Y*(1915) WIDTH (MEV) (PRCD. EXP.)

Table showing width measurements for Y*(1915) from various experiments like BUGG, BRICMAN, COOL, etc.

29 Y*(1915) PARTIAL DECAY MODES (PRCD. EXP.)

Table listing partial decay modes for Y*(1915) such as INTO KBAR N, INTO LAMBDA PI, INTO SIGMA PI.

29 Y*(1915) BRANCHING RATIOS (PROD. EXP.)

Table with columns for particle name, mass, width, and branching ratios. Includes entries for Y*(1915) INTO (KBAR N)/TOTAL and FROM KBAR N INTO LAMBDA PI.

REFERENCES FOR Y*(1915) (PROD. EXP.)

Table listing references for Y*(1915) from various sources like BOCK, COOL, BUGG, etc.

PAPERS NOT REFERRED TO IN DATA CARDS

Table listing papers not referred to in data cards, such as PRIMER, AGUILAR, etc.

$\Sigma(1940)$

98 Y*(1940, J_p=3/2-) I=1 D₁₃

SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES IN THIS REGION.

98 Y*(1940) MASS (MEV)

Table showing mass measurements for Y*(1940) from various experiments like M, M 2, etc.

98 Y*(1940) WIDTH (MEV)

Table showing width measurements for Y*(1940) from various experiments like W, W 2, etc.

98 Y*(1940) PARTIAL DECAY MODES

Table listing partial decay modes for Y*(1940) such as INTO KBAR N, INTO LAMBDA PI, INTO SIGMA PI.

98 Y*(1940) BRANCHING RATIOS

Table with columns for particle name, mass, width, and branching ratios. Includes entries for Y*(1940) FROM KBAR N INTO LAMBDA PI.

Baryons

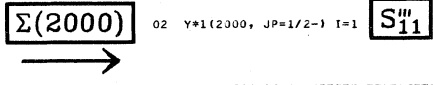
$\Sigma(1940)$, $\Sigma(2000)$, $\Sigma(2030)$

Data Card Listings

For notation, see key at front of Listings.

R3 2 Y*1(1940) INTO KBAR N (P1) 9/73
R3 2 (-21) LEA 73 DPWA MULTICHLN K-MTRX 9/73
R3 LESS THAN .04 RLIC 76 DPWA KBAR N MULTICHLN 1/76*

REFERENCES FOR Y*1(1940)
GALTIERI 70 DUKE CCNF 173 A BARBARO-GALTIERI (LRL)IJP
LITCHFIE 70 NP B22 269 P J LITCHFIELD (RUTHERFORD)IJP



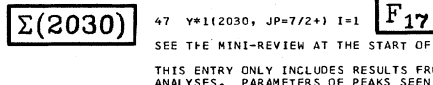
02 Y*1(2000) MASS (MEV)
M 2004. 40. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
M 1955. 15. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

02 Y*1(2000) WIDTH (MEV)
W 116. 40. VANHORN 75 DPWA O K-P TO LAM P10 11/75*
W 170. 40. RLIC 76 DPWA KBAR N MULTICHLN 1/76*

02 Y*1(2000) PARTIAL DECAY MODES
P1 Y*1(2000) INTO KBAR N 497+ 939
P2 Y*1(2000) INTO LAMBDA PI 1115+ 134
P3 Y*1(2000) INTO SIGMA PI 1197+ 139

02 Y*1(2000) BRANCHING RATIOS
R1 Y*1(2000) FROM KBAR N INTO LAMBDA PI SQRT(P1*P2)
R1 NOT SEEN BAILLON 75 IPWA KBAR N TO LAM PI 11/75*

REFERENCES FOR Y*1(2000)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL)IJP
VANHORN 75 NP 887 145 A. J. VAN HORN (LRL)IJP



SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES.

02 Y*1(2030) MASS (MEV)
M 2030.0 10.0 WOHLL 66 HBC O K-P TO LAM P10 7/66

47 Y*1(2030) MASS (MEV)
M 2030.0 (20.0) WOHLL 66 HBC O K-P TO LAM P10 7/66
M 2032.0 6.0 SMART 68 DPWA - K-N TO LAMBDA PI 6/68

47 Y*1(2030) WIDTH (MEV)
W 170.0 (17.0) WOHLL 66 HBC O 7/66
W 160.0 16.0 SMART 68 DPWA - K-N TO LAMBDA PI 6/68

47 Y*1(2030) PARTIAL DECAY MODES
P1 Y*1(2030) INTO KBAR N 497+ 939
P2 Y*1(2030) INTO LAMBDA PI 1115+ 134

47 Y*1(2030) BRANCHING RATIOS
R1 Y*1(2030) INTO (KBAR N)/TOTAL (P1)
R1 (0.25) WOHLL 66 HBC O K-P CH EX 7/66

REFERENCES FOR Y*1(2030)
BAILLON 75 NP 894 39 P. BAILLON, P. J. LITCHFIELD (CERN, RHEL)IJP
VANHORN 75 NP 887 145 A. J. VAN HORN (LRL)IJP

47 Y*1(2030) PARTIAL DECAY MODES
P1 Y*1(2030) INTO KBAR N 497+ 939
P2 Y*1(2030) INTO LAMBDA PI 1115+ 134

Data Card Listings

For notation, see key at front of Listings.

Baryons

Σ(2030), Σ(2070), Σ(2080)

R8 Y*1(2030) FRCH KEAR N TO KBAR DELTA(1232) F-WAVE SQRTP1*P8
R8 3 .16 .03 LITCHF13 74 DPWA 0 K-P TO KBAR DEL 10/74*
R8 3 (-17) (.03) CORDEN2 75 DBC - KBAR PI-NUCLEON 11/75*

R9 Y*1(2030) FROM KEAR N TO KBAR DELTA(1232) H-WAVE SQRTP1*P9
R9 .00 .02 LITCHF13 74 DPWA 0 K-P TO KBAR DEL 10/74*

REFERENCES FOR Y*1(2030)
WOHL 66 PRL 17 107 C G WOHL, F T SOLMITZ, M L STEVENSON (LRL)IJP
+ LEITH, + (LRL,SLAC,CERN,HEIDEL,SACLAY)
TRIPP 67 NP 83 10 +MEYER,PAULI,TALLINI + (SACL+CDEF+RHEL)
BURGUN 68 NP 88 447 +ERNE,LAGNAUX,SENS,STEUER,UDO (CERN)IJP
DAUM 68 NP 87 19 +RANGE,VRANA,+ (COL FRANCE, RHEL, SACLAY)IJP

CCNFIRMS THE SPIN-PARITY ASSIGNMENT.
SMART 68 PR 169 1336 W M SMART (LRL)IJP
MULLER 69 THESIS,UCLR 19372 R A MULLER (LRL)
+VRANA,VRANA,+ (COL FRANCE, RHEL, SACLAY)IJP
+RANGAN,VRANA,+ (COL FRANCE, RHEL, SACLAY)IJP
+VRANA, BUTTERTHORTH,+ (CDEF, RHEL, SACLAY)IJP
+ISLAM, COLLEY, + (BIRM,EDIN,GLAS,LOIC)IJP
A BARBARO-GALTIERI (LRL)IJP
P J LITCHFIELD (RUTHERFORD)IJP

+MORTON, NEGUS, GOYAL, MILLER (GLAS, LOIC)IJP
LITCHFIELD,....+LESQUOY,+. (RHEL+CDEF+SACL)IJP
D F KANE (LBL)IJP
DEVENISH 74 NP 881 330 DEVENISH,FROGGATT,MARTIN(DESY,NORDITA,LOUC)
LITCHF11 74 NP 874 12 +COX,DARTNELL,KENYON,ONEALE,SUMOROK+ (BIRM)IJP
LITCHF12 74 NP 874 19 +HEMINGWAY,EADES,HARMSEN+ (CERN,HEID,MPH)IJP
LITCHF13 74 NP 874 39 A. J. VAN HORN (LBL)IJP
A. J. VAN HORN (LBL)IJP

BAILLON 75 NP 894 39 P. BAILLON,P. J. LITCHFIELD (CERN,RHEL)IJP
CORDEN2 75 NP 892 365 +COX,DARTNELL,KENYON,ONEALE,SUMOROK+ (BIRM)IJP
HEMINGWA 75 NP 891 12 +HEMINGWAY,EADES,HARMSEN+ (CERN,HEID,MPH)IJP
VANHORN 75 NP 887 145 A. J. VAN HORN (LBL)IJP
ALSO 75 NP 887 157 A. J. VAN HORN (LBL)IJP

RLIC 76 RL-75-182 (PRPNT) GOPAL,KALMUS,MCPHERSON,ROSS+ (RHEL+LOIC)IJP

2030 MEV REGION - PRODUCTION AND σTOTAL EXP'TS

28 Y*1(2030, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

SEE THE NOTE TO THE F17 Y*1(2030), WHICH PRECEDES THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS-SECTION PEAKS ARE AT LEAST DOMINANTLY ASSOCIATED WITH THE Y*1(2030), BUT MAY CONTAIN A SMALL CONTRIBUTION FROM THE SUGGESTED BUT NOT ESTABLISHED OTHER RESONANCES IN THIS REGION.

28 Y*1(2030) MASS (MEV) (PROD. EXP.)
M (2022.0) (20.0) BLANPIED 65 CNTR 0 GAMMA P TO K+ Y*
M 2020.0 7.0 BUGG 68 CNTR K-P, D TOTAL 6/68
M 2049.0 4.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
M 2025.0 10.0 COOL 70 CNTR K-P, D TOTAL 10/70
M (2025.0) (20.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71
M AVERAGE MEANINGLESS (SCALE FACTOR = 2.8)

28 Y*1(2030) WIDTH (MEV) (PROD. EXP.)
W (120.0) (20.0) BLANPIED 65 CNTR 0
W 130.0 10.0 BUGG 68 CNTR 6/68
W 126.0 11.0 BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
W 165.0 0.0 COOL 70 CNTR K-P, D TOTAL 10/70
W (180.0) LU 70 CNTR 0 GAMMA P TO K+ Y* 1/71
W AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

28 Y*1(2030) PARTIAL DECAY MODES (PROD. EXP.)
P1 Y*1(2030) INTO KBAR N 497+ 929
P2 Y*1(2030) INTO KBAR N PI 497+ 935+ 139

28 Y*1(2030) BRANCHING RATIOS (PROD. EXP.)
R1 Y*1(2030) INTO (KEAR N)/TOTAL
R1 THESE VALUES OF ELASTICITIES ASSUME J=7/2 - (P1)
R1 0.131 BUGG 68 CNTR 6/68
R1 0.27 (0.02) BRICMAN 70 CNTR 0 TOTAL AND CH EX 6/70
R1 0.12 .00 COOL 70 CNTR K-P, D TOTAL 10/70
R2 Y*1(2030) INTO KBAR N PI (P2)
R2 SEEN BOCK HBC

REFERENCES FOR Y*1(2030) (PROD. EXP.)
BLANPIED 65 PRL 14 741 +GREENBERG,HUGHES,KITCHING,LU,+ (YALE(CEA))
COOL 66 PRL 16 1228 +GIACOMELLI,KYCIA,LEONTIC,LI,LUNDBY,+ (BNL I
BUGG 68 PR 168 1466 +GILMORE,KNIGHT, + (RHEL,BIRM,CAVE) I
BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU,+ (CERN,CAEN,SACLAY)
COOL 70 PR D1 1867 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
LU 70 PR D2 1846 +GREENBERG, HUGHES, WINEHART, MORI,+ (YALE)

Σ(2070)

F15

THIS STATE HAS BEEN SUGGESTED BY ONLY ONE PARTIAL WAVE ANALYSIS ACROSS THIS REGION. IT NEEDS CONFIRMATION THE RESONANCE PROPOSED BY KANE IS TOO BROAD TO BE USED AS EVIDENCE.

34 Y*1(2070) MASS (MEV)
M (2070.) (10.) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
M (2057.0) KANE 72 DPWA K-P TO SIGMA PI 1/73

34 Y*1(2070) WIDTH (MEV)
W (140.) (20.) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
W (906.0) KANE 72 DPWA K-P TO SIGMA PI 1/73

34 Y*1(2070) PARTIAL DECAY MODES
P1 Y*1(2070) INTO KBAR N 497+ 939
P2 Y*1(2070) INTO SIGMA PI 1197+ 139

34 Y*1(2070) BRANCHING RATIOS
R1 Y*1(2070) FROM KEAR N TO SIGMA SQRTP1*P2
R1 (+.12) (.02) BERTHON1 70 DPWA - K- P TO SIG PI 1/71
R1 (+0.104) KANE 72 DPWA K-P TO SIGMA PI 1/73

REFERENCES FOR Y*1(2070)
BERTHON1 70 NP 824 417 +VRANA,BUTTERTHORTH,+ (CDEF,RHEL,SACLAY)IJP
KANE 72 PR D5 1583 D F KANE (LBL)

Σ(2080)

P13

SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

88 Y*1(2080) MASS (MEV)
M (2082.0) (4.0) COX 70 DPWA - K- N TO LAM PI 6/70
M (2070.0) (30.0) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
M 1 2120. 40. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
M 1 FROM SOLUTION 1 OF BAILLON 75.
M 2 2140. 40. BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
M 2 FROM SOLUTION 2 OF BAILLON 75.
M 2140. 30. BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*
M AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

88 Y*1(2080) WIDTH (MEV)
W (87.0) (20.0) COX 70 DPWA - K- N TO LAM PI 6/70
W (250.0) (40.0) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
W 1 240. 50. BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
W 2 200. 50. BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
W 180. 20. BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*
W AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

88 Y*1(2080) PARTIAL DECAY MODES
P1 Y*1(2080) INTO KBAR N 497+ 939
P2 Y*1(2080) INTO LAMBDA PI 1115+ 139

88 Y*1(2080) BRANCHING RATIOS
R1 Y*1(2080) FROM KEAR N TO LAMBDA PI SQRTP1*P2
R1 (-0.16) (0.03) COX 70 DPWA - K- N TO LAM PI 6/70
R1 (-0.09) (0.03) LITCHFIELD 70 DPWA -0 K- N TO LAM PI 6/70
R1 1 -.13 .04 BAILLON 75 IPWA KBAR N TO LAM PI 11/75*
R1 2 -.13 .04 BAILLON 75 IPWA KBAR N TO LAM PI 1/76*
R1 .19 .03 BELLEFOL 75 DPWA 0 K- P TO LAM P IO 11/75*
R1 AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

REFERENCES FOR Y*1(2080)
COX 70 NP 819 61 +ISLAM, COLLEY, + (BIRM,EDIN,GLAS,LOIC)IJP
LITCHFIE 70 NP 822 269 P J LITCHFIELD (RUTHERFORD)IJP
BAILLON 75 NP 894 39 P. BAILLON,P. J. LITCHFIELD (CERN,RHEL)IJP
BELLEFOL 75 NP 890 1 DE BELLEFON,BERTHON,BRUNET+ (CDEF,SACL)IJP

Baryons

$\Sigma(2100)$, $\Sigma(2250)$, $\Sigma(2455)$

Data Card Listings

For notation, see key at front of Listings.

$\Sigma(2100)$

26 $Y^*(2100, JP=7/2^-) I=1$ **G17**
SEE THE MINI-REVIEW AT THE START OF THE Y^* LISTINGS.
SUCH A RESONANCE IS SUGGESTED BY SOME BUT NOT ALL PARTIAL-WAVE ANALYSES ACROSS THIS REGION. UNTIL THERE IS MORE EVIDENCE, WE OMIT THIS STATE FROM THE MAIN BARYON TABLE.

Table with 2 columns: Mass (MEV) and branching ratios for $\Sigma(2100)$.

Table with 2 columns: Width (MEV) and branching ratios for $\Sigma(2100)$.

Table with 2 columns: Partial decay modes and decay masses for $\Sigma(2100)$.

Table with 2 columns: Branching ratios for $\Sigma(2100)$ to various states.

REFERENCES FOR $Y^*(2100)$
GALTIERI 70 DUKE CNF 173 A BARBARO-GALTIERI (LRL)JJP

$\Sigma(2250)$ BUMPS

48 $Y^*(2250, JP=) I=1$ PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y^* LISTINGS.
THE PARTIAL-WAVE ANALYSIS RESULTS ARE TOO WEAK TO WARRANT SEPARATING THEM FROM THE PRODUCTION AND CROSS-SECTION EXPERIMENTS.
LASINSKI 71 IN KBAR N, USING A POMERON-RESONANCES MODEL, AND BELLEFON1 75,76 IN DPWA'S OF LAMBDA PI AND SIGMA PI SUGGEST THE PRESENCE OF TWO RESONANCES AROUND THIS MASS VALUE.

Table with 2 columns: Mass (MEV) and branching ratios for $\Sigma(2250)$.

Table with 2 columns: Width (MEV) and branching ratios for $\Sigma(2250)$.

Table with 2 columns: Partial decay modes and decay masses for $\Sigma(2250)$.

Table with 2 columns: Branching ratios for $\Sigma(2250)$ to various states.

48 $Y^*(2250)$ BRANCHING RATIOS (PROD. EXP.)

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

Table with 2 columns: Branching ratios for $Y^*(2250)$ into various states.

$\Sigma(2455)$ BUMPS

53 $Y^*(2455, JP=) I=1$ PRODUCTION EXPERIMENTS

SEE THE MINI-REVIEW AT THE START OF THE Y^* LISTINGS.
THERE IS ALSO SOME SLIGHT EVIDENCE FOR Y^* STATES IN THIS MASS REGION FROM THE REACTION $\gamma + p \rightarrow K^+ +$ MISSING MASS -- SEE GREENBERG 68.

Table with 2 columns: Mass (MEV) and branching ratios for $\Sigma(2455)$.

Table with 2 columns: Width (MEV) and branching ratios for $\Sigma(2455)$.

Table with 2 columns: Partial decay modes and decay masses for $\Sigma(2455)$.

Table with 2 columns: Branching ratios for $\Sigma(2455)$ to various states.

Table with 2 columns: Branching ratios for $\Sigma(2455)$ to various states.

Table with 2 columns: Branching ratios for $\Sigma(2455)$ to various states.

Data Card Listings

Baryons

For notation, see key at front of Listings. $\Sigma(2620), \Sigma(3000), \text{EXOTIC HYPERONS}, \Xi^{\prime}s, \Xi^{-}, \Xi^0$

**$\Sigma(2620)$
BUMPS**

54 Y*1(2620, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.

54 Y*1(2620) MASS (MEV) (PROD. EXP.)

M	2620.0	15.0	ABRAMS	70 CNTR	K-P, D TOTAL	10/70
M	2542.	22.	DIBIANCA	75 DBC	XI K PI	1/76*
M	AVERAGE MEANINGLESS (SCALE FACTOR = 2.9)					

54 Y*1(2620) WIDTH (MEV) (PROD. EXP.)

W	(175.0)		ABRAMS	70 CNTR	K-P, D TOTAL	10/70
W	221.	81.	DIBIANCA	75 DBC	XI K PI	1/76*

54 Y*1(2620) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y*1(2620) INTO KBAR N	DECAY MASSES
		497+ 939

54 Y*1(2620) BRANCHING RATIOS (PROD. EXP.)

R1	Y*1(2620) INTO (KBAR N)/TOTAL	(P1)	
R1	J IS NOT KNOWN. THE FOLLOWING IS (J+1/2)*P1.		
R1	(0.32)	ABRAMS	70 CNTR K-P, D TOTAL
R1	0.36	BRICMAN	70 CNTR 0 TOTAL AND CH EX

REFERENCES FOR Y*1(2620) (PROD. EXP.)

ABRAMS 67 PRL 19 678 +CGOL, GIACOMELLI, KYCIA, LEONIC, LI, + (BNL)
SUPERSEDED BY ABRAMS 70.
ABRAMS 70 PR 10 1917 +CGOL, GIACOMELLI, KYCIA, LEONIC, + (BNL) I
BRICMAN 70 PL 31B 152 +FERRO LUZZI, PERREAU, + (CERN, CAEN, SACLAY)
DIBIANCA 75 NP 898 137 DIBIANCA, ENDORFR (CERN)

**$\Sigma(3000)$
BUMPS**

59 Y*1(3000, JP=) I=1 PRODUCTION EXPERIMENTS
SEE THE MINI-REVIEW AT THE START OF THE Y* LISTINGS.
ENHANCEMENT IN LAMBDA PI AND KBAR N INVARIANT MASS SPECTRA AND IN MISSING MASS OF NEUTRALS RECOLLING AGAINST K0. EVIDENCE NOT CONCLUSIVE. OMITTED FROM TABLE.

59 Y*1(3000) MASS (MEV) (PROD. EXP.)

M	(3000.0)	EHRlich	66 HBC	0 P1-P	7.91 BEV/C	9/66
---	----------	---------	--------	--------	------------	------

59 Y*1(3000) PARTIAL DECAY MODES (PROD. EXP.)

P1	Y*1(3000) INTO KBAR N	DECAY MASSES
P2	Y*1(3000) INTO LAMBDA PI	497+ 939 1115+ 139

REFERENCES FOR Y*1(3000) (PROD. EXP.)

EHRlich 66 PR 152 1194 R EHRlich, W SELOVE, H YUTA (PENN(BNL)) I

EXOTIC HYPERON CROSS SECTION LIMITS

THIS IS NOT A COMPLETE LIST. WE TABULATE ONLY FROM 1970 ON.

CS	UNITS MICROBARN			
CS	G (20.) OR LESS	GALTIERI 68 DBC	K-N TO SG-PI-PI0	7/70
CS	G ABOVE LIMIT FOR MASS < 2.15 GEV AND GAMMA < 60 MEV- (2.1 GEV/C K-)			7/70
CS	A (40.) OR LESS	GALTIERI 68 DBC	K-N TO SG-PI-PI0	7/70
CS	A ABOVE LIMIT FOR MASS < 2.3 GEV AND GAMMA < 120 MEV- (2.7 GEV/C K-)			7/70

REFERENCES FOR EXOTIC HYPERONS

GALTIERI 68 PRL 21 573 A. BARBARO-GALTIERI, CHADWICK + (LRL+SLAC)

Note on Ξ Resonances

The Ξ resonance situation has long been unsettled. This is because 1) they can only be produced as part of a final state, $K^+p \rightarrow \Xi^* + \text{others}$, and 2) they are so produced with very small cross sections ($< 50 \mu\text{b}$). Thus the numbers of events available are small, and the analysis is more complicated than if direct formation were possible. Only the $\Xi(1530)$ is really well established. There are apparently at least two Ξ states in the 1800-2000 MeV region, and there are indications of several more above 2000 MeV, but the situation is unclear. Two large bubble chamber experiments reported at the 1975 ANL Symposium on New Directions in Hadron Spectroscopy (BLOKZIJL 75 and MORRIS 75) may be able to clarify the situation. In the meantime, we are forced to group together rather disparate observations and await new results. Figures in the listings point out disagreements among various experiments. The table following this note gives our evaluation of the status of the Ξ resonances, based on the data available at this time.

STATUS OF XI* RESONANCES
THOSE WITH AN OVERALL STATUS OF *** OR **** ARE INCLUDED IN THE MAIN BARYON TABLE. THE OTHERS AWAIT CONFIRMATION.

PARTICLE	LIJ	OVERALL STATUS	STATUS AS SEEN IN --				
			XI PI	LAM K	SIG K	XI* PI	OTHER CHANNELS
XI(1320) P11		****					
XI(1530) P13		****					WEAK TO LAM PI
XI(1630)		**					
XI(1820)		***			**	**	
XI(1940)		***			**	**	
XI(2030)		**			**	**	3-BODY DECAYS
XI(2250)		*			**	**	3-BODY DECAYS
XI(2500)		**			**	**	3-BODY DECAYS

**** GOOD, CLEAR, AND UNMISTAKABLE.
*** GOOD, BUT IN NEED OF CLARIFICATION OR NOT ABSOLUTELY CERTAIN.
** NEEDS CONFIRMATION.
* WEAK.

S=-2 I=1/2 HYPERON STATES (Ξ)

Ξ^-

22 XI-(1321, JP=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

Ξ^0

23 XI(1314, JP=1/2) I=1/2
SEE STABLE PARTICLE DATA CARD LISTINGS

Baryons

$\Xi(1530)$, $\Xi(1630)$

$\Xi(1530)$

49 $\Xi^{*1/2}(1530)$, JP=3/2+ I=1/2 P13
THIS IS THE ONLY REALLY WELL-ESTABLISHED Ξ^{*} . THE QUANTUM NUMBERS 3/2+ ARE FAVORED BY THE DATA.

WE DO NOT USE DETERMINATIONS OF THE MASS AND THE WIDTH OF THIS STATE UNLESS THEY ARE ACCOMPANIED BY SOME DISCUSSION OF SYSTEMATICS AND RESOLUTION.

49 $\Xi^{*1/2}(1530)$ MASS (MEV)

Table with columns for mass (M), width (W), and fit (FIT) values for various experiments and theoretical models. Includes average values and error scales.

49 ($\Xi^{*1/2}$) - (Ξ) MASS DIFFERENCE (MEV)

Table showing mass differences between $\Xi^{*1/2}$ and Ξ for different experiments. Includes average values and error scales.

49 $\Xi^{*1/2}(1530)$ WIDTH (MEV)

Table showing widths for various experiments and theoretical models. Includes average values and error scales.

WEIGHTED AVERAGE = 9.14 ± 0.48
ERROR SCALED BY 1.0

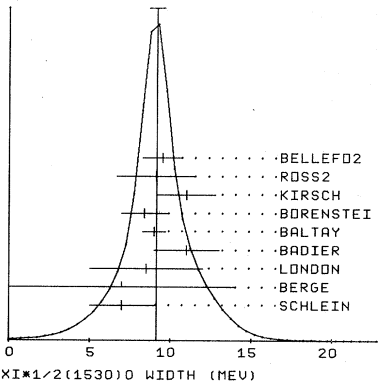


Table with columns for width (W), average (AVG), and student (STUDENT) values. Includes average values and error scales.

49 $\Xi^{*1/2}(1530)$ REAL PART OF POLE POSITION

Table showing real part of pole position for different experiments. Includes average values and error scales.

Data Card Listings

For notation, see key at front of Listings.

49 $\Xi^{*1/2}(1530)$ IMAGINARY PART OF POLE POSITION 4/75*

Table with columns for imaginary part of pole position, including values for Lichtenberg and other experiments.

49 $\Xi^{*1/2}(1530)$ PARTIAL DECAY MODES

Table showing decay masses for $\Xi^{*1/2}(1530)$ into Ξ pi and Ξ gamma.

49 $\Xi^{*1/2}(1530)$ BRANCHING RATIOS (MEV)

Table showing branching ratios for $\Xi^{*1/2}(1530)$ into Ξ pi and Ξ gamma.

REFERENCES FOR $\Xi^{*1/2}(1530)$

- List of references for $\Xi^{*1/2}(1530)$ including authors like PJERRCU, SCHLEIN, BADIER, etc., and their respective publications.

SHAFER 66 PR 142 883 BLTTCN-SHAHER, LINDSEY, MURRAY, SMITH (LRL) JP
A SPIN-PARITY DETERMINATION.
HABIBI 73 NEVIS 199(THESIS) HABIBI (COLU)
HUNGERBU 74 PRD 10 2051 HUNGERBUHLER, MAJKA, + (YALE, FNAL, BNL, PITT)
BRIEFEL 75 PRD 12 1859 +GUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)

$\Xi(1630)$

21 $\Xi^{*1/2}(1630)$, JP= 1 I=1/2

THIS EFFECT NEEDS CONFIRMATION.
BARTSCH 69 SEE A SMALL, BROAD ENHANCEMENT NEAR 1650 MEV - IT IS NOT CLEAR THAT IT IS THE SAME PHENOMENON AS BMST 70, WHO FIND CS=3.6+1.6 MICROBARN AT 2.87 GEV/C INCIDENT K- MOMENTUM.
BORENSTEIN 72 SEE NO EFFECT IN THIS REGION. THEY FIND CS<2 MICROBARN AT 2.18.
ROSS 72 ARGUE THAT THE EFFECT THEY SEE IS NOT THE SAME AS THAT SEEN BY BMST 70, AND FIND CS=2+1 MICROBARN AT 3.3 GEV/C.
BELLEFON 75B FIND A CS OF AROUND 10 MICROBARN NEAR 2 GEV/C, BUT LESS THAN 3 MICROBARN AROUND 2.3 GEV/C.

21 $\Xi^{*1/2}(1630)$ MASS (MEV)

Table with columns for mass (M), width (W), and fit (FIT) values for various experiments and theoretical models.

21 $\Xi^{*1/2}(1630)$ WIDTH (MEV)

Table showing widths for various experiments and theoretical models. Includes average values and error scales.

21 $\Xi^{*1/2}(1630)$ PARTIAL DECAY MODES

Table showing decay masses for $\Xi^{*1/2}(1630)$ into Ξ pi and Ξ gamma.

REFERENCES FOR $\Xi^{*1/2}(1630)$

- List of references for $\Xi^{*1/2}(1630)$ including authors like APSELL, BARTSCH, BORENSTEIN, etc., and their respective publications.

Data Card Listings

For notation, see key at front of Listings.

Baryons
Xi(1820), Xi(1940)

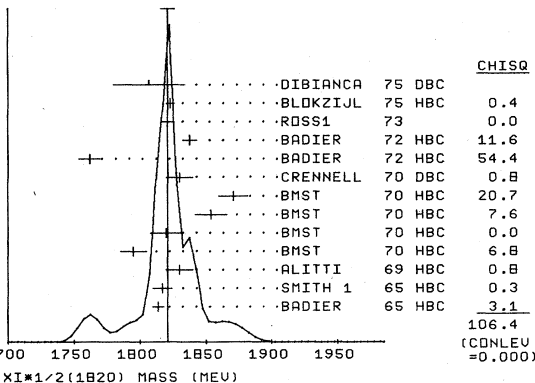
Xi(1820)

50 XI*1/2(1820, JP=) I=1/2

AS THE ACCOMPANYING IDEOGRAMS ILLUSTRATE, THE SITUATION IS CONFUSED, UNTIL SOME FUTURE CLARIFICATION, WE LIST UNDER XI(1820) EVERYTHING REPORTED IN THE MASS RANGE 1750-1875 MEV. WHEN BRANCHING RATIOS ARE REPORTED, WE QUOTE THEM, BUT ONLY THE MOST QUALITATIVE CONCLUSIONS ARE JUSTIFIED.

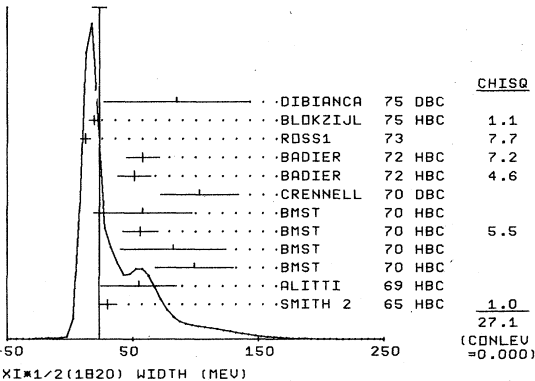
50 XI*1/2(1820) MASS (MEV)

Table with columns for mass (MEV), source, and branching ratios for Xi(1820). Includes entries like HALSTEINS 63 FBC, BADIER 65 HBC, SMITH 1 65 HBC, etc.



50 XI*1/2(1820) WIDTH (MEV)

Table with columns for width (MEV), source, and branching ratios for Xi(1820). Includes entries like HALSTEINS 63 FBC, BADIER 65 HBC, SMITH 2 65 HBC, etc.



50 XI*1/2(1820) PARTIAL DECAY MODES

Table of partial decay modes for Xi(1820) including decay masses for various channels like XI*1/2(1820) INTO LAMBDA KBAR, etc.

50 XI*1/2(1820) BRANCHING RATIOS

Table of branching ratios for Xi(1820) for various decay channels, including R1, R2, R3, R4, R5, R6, R7, R8, R9, R10.

REFERENCES FOR XI*1/2(1820)

List of references for Xi(1820) including HALSTEIN 63 SIENA CONF 173, BADIER 65 PL 16 171, SMITH 1 65 PRL 16 25, etc.

Xi(1940)

52 XI*1/2(1940, JP=) I=1/2

WE LIST UNDER XI(1940) EVERYTHING REPORTED IN THE MASS RANGE 1875-2300 MEV.

52 XI*1/2(1940) MASS (MEV)

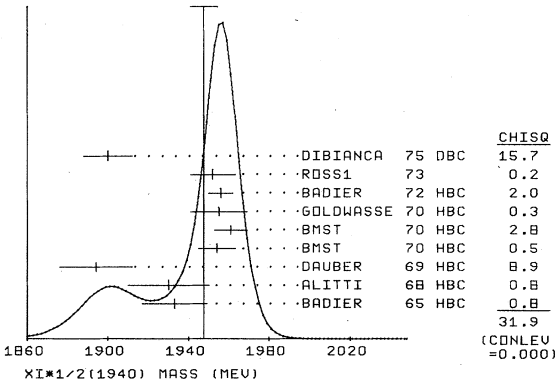
Table with columns for mass (MEV), source, and branching ratios for Xi(1940). Includes entries like BADIER 65 HBC, ALITTI 68 HBC, DAUBER 69 HBC, etc.

Baryons

$\Xi(1940)$, $\Xi(2030)$, $\Xi(2250)$

Data Card Listings

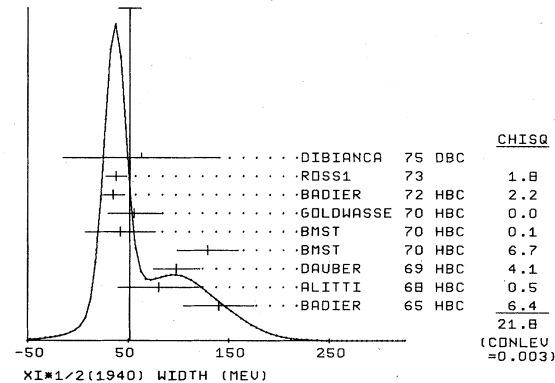
For notation, see key at front of Listings.



52 $\Xi^{*1/2}(1940)$ WIDTH (MEV)

W	Mass (MEV)	CHI SQ
35	140.0	35.0
27	80.0	40.0
66	98.0	23.0
110	129.0	30.0
40	42.0	35.0
21	56.0	26.0
29	35.0	11.0
38	10.0	10.0
63	78.0	78.0

AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)
(SEE IDECGRAM BELOW)



52 $\Xi^{*1/2}(1940)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi$	1321+ 139
P2	$\Xi^{*1/2}(1940) \rightarrow \Xi^{*}(1530) \pi$	153+ 135
P3	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi \pi$ (EXCLUDING P2)	1321+ 139+ 139
P4	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi \pi$	1314+ 139
P5	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi \pi$	1321+ 134

52 $\Xi^{*1/2}(1940)$ BRANCHING RATIOS

THE $\Xi(1940)$ IS SEEN MAINLY IN $\Xi \pi$ AND SOME IN $\Xi^{*}(1530) \pi$. IT HAS BEEN LOOKED FOR IN OTHER CHANNELS BUT NOT SEEN.

R	Decay Mode	Branching Ratio
R1	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi$	2.8
R2	$\Xi^{*1/2}(1940) \rightarrow \Xi^{*}(1530) \pi$	0.3
R3	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi \pi$	0.6
R4	$\Xi^{*1/2}(1940) \rightarrow \Xi \pi \pi$	0.7

REFERENCES FOR $\Xi^{*1/2}(1940)$

BADIER 45 PL 16 171
ALITTI 68 PRL 21 1119
DAUBER 69 PR 179 1262
APSELL 70 PRL 24 777
BMST 70 DUKE 317
GOLDWASSER 70 PR 1D 196
BADIER 72 NP 837,429
ROSSI 73 PURDUE CONF. 345
DIBIANCA 75 NP 898 137

+DEMUNLIN, GELDBERG, + (EPOL, SACLAY, AMST) I
+FLAMINIO, METZGER, RADOJICIC, + (BNL, SYRACUSE) I
+BERGE, HUBBARD, MERRILL, MULLER (LRL) I
+ (BRANDEIS, MARYLAND, SYRACUSE, TUFTS) I
BRANDEIS+MARYLAND+SYRACUSE+TUFTS COLLABOR.
E L GOLDWASSER, F SCHULTZ (ILLINOIS)
+BARRELET, CHARLTON, VIDEAU (EPOL)
RCSS, LLOYD, RADOJICIC (OXFORD)

DIBIANCA, ENDORFR (CARN)

PAPERS NOT REFERRED TO IN DATA CARDS

APSELL 69 PRL 23 884 + (BRANDEIS, MARYLAND, SYRACUSE, TUFTS)
SUPERSEDED BY BMST 70.
SCHMIDT 73 PURDUE CONF. 363 + SCHMIDT (BRANDEIS)
BRIEFEL 75 PRD 12 1859 +GOUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)

$\Xi(2030)$

68 $\Xi^{*1/2}(2030, JP=)$ I=1/2
THE STATE AT 2129 MEV REPORTED BY BLOKZIJL 75 IS REALLY INTERMEDIATE BETWEEN $\Xi^{*}(2130)$ AND $\Xi^{*}(2250)$. WE LIST IT HERE TEMPORARILY.

68 $\Xi^{*1/2}(2030)$ MASS (MEV)

M	Mass (MEV)	CHI SQ
42	2030.0	10.0
40	2058.0	17.0
15	2019.0	7.0
25	2129.0	5.0
2044.0	8.0	

AVERAGE MEANINGLESS (SCALE FACTOR = 7.4)

68 $\Xi^{*1/2}(2030)$ WIDTH (MEV)

W	Mass (MEV)	CHI SQ
45.0	40.0	20.0
57.0	30.0	
15	35.0	17.0
25	21.0	8.0
60.0	24.0	

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

68 $\Xi^{*1/2}(2030)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^{*1/2}(2030) \rightarrow \Xi \pi$	1321+ 139
P2	$\Xi^{*1/2}(2030) \rightarrow \Lambda \text{B} \text{K} \text{B} \text{A} \text{R}$	1115+ 497
P3	$\Xi^{*1/2}(2030) \rightarrow \Sigma \text{B} \text{K} \text{A} \text{R}$	1197+ 497
P4	$\Xi^{*1/2}(2030) \rightarrow \Xi^{*}(1530) \pi$	153+ 139
P5	$\Xi^{*1/2}(2030) \rightarrow \Lambda \text{B} \text{C} \text{A}$ (OR $\Sigma \text{B} \text{K} \text{A} \text{R}$)	1115+ 497+ 139

68 $\Xi^{*1/2}(2030)$ BRANCHING RATIOS

R	Decay Mode	Branching Ratio
R1	$\Xi^{*1/2}(2030) \rightarrow \Xi \pi$	0.30
R2	$\Xi^{*1/2}(2030) \rightarrow \Lambda \text{B} \text{K} \text{B} \text{A} \text{R}$	0.25
R3	$\Xi^{*1/2}(2030) \rightarrow \Sigma \text{B} \text{K} \text{A} \text{R}$	0.75
R4	$\Xi^{*1/2}(2030) \rightarrow \Xi^{*}(1530) \pi$	0.15
R5	$\Xi^{*1/2}(2030) \rightarrow \Lambda \text{B} \text{C} \text{A}$ (OR $\Sigma \text{B} \text{K} \text{A} \text{R}$)	0.20

REFERENCES FOR $\Xi^{*1/2}(2030)$

ALITTI 69 PRL 22 79 +BARNES, FLAMINIO, METZGER, + (BNL, SYRACUSE) I
BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LOIC, VIENNA)
ROSSI 73 PURDUE CONF. 345 RCSS, LLOYD, RADOJICIC (OXFORD)

BLOKZIJL 75 ANL-HEP-CP-75-58 +DEGROOT, HOOGLAND+ (AMST+CERN+NIJH+OXF)
75 PALERMO CONF. +DEGROOT, HOOGLAND+ (AMST+CERN+NIJH+OXF)
DIBIANCA 75 NP 898 137 DIBIANCA, ENDORFR (CARN)

$\Xi(2250)$

22 $\Xi^{*1/2}(2250, JP=)$
THE EVIDENCE FOR THIS STATE IS WEAK. BARTSCH 69 SEE A BUMP OF NOT MUCH STATISTICAL SIGNIFICANCE IN $\Lambda \text{B} \text{C} \text{A}$ - $\text{K} \text{B} \text{A} \text{R}$ - π , $\Sigma \text{B} \text{K} \text{A} \text{R}$ - π , AND Ξ - π - π MASS SPECTRA. GOLDWASSER 70 SEE A NARROWER BUMP IN Ξ - π - π AT A HIGHER MASS. PERHAPS THEY ARE THE SAME STATE, PERHAPS THEY ARE NOT, BUT SEE ALSO MORRIS 75.

22 $\Xi^{*1/2}(2250)$ MASS (MEV)

M	Mass (MEV)	CHI SQ
35	2244.0	52.0
18	2295.0	15.0

AVERAGE MEANINGLESS (SCALE FACTOR = 1.0)

22 $\Xi^{*1/2}(2250)$ WIDTH (MEV)

W	Mass (MEV)	CHI SQ
130.0	80.0	2/74
LESS THAN	30.0	2/74

22 $\Xi^{*1/2}(2250)$ PARTIAL DECAY MODES

P	Decay Mode	Decay Masses
P1	$\Xi^{*1/2}(2250) \rightarrow \Xi \pi$	1321+ 139+ 139
P2	$\Xi^{*1/2}(2250) \rightarrow \Lambda \text{B} \text{C} \text{A}$	1115+ 497+ 139
P3	$\Xi^{*1/2}(2250) \rightarrow \Sigma \text{B} \text{K} \text{A} \text{R}$	1197+ 497+ 139

Data Card Listings

For notation, see key at front of Listings.

Baryons

$\Xi(2250)$, $\Xi(2500)$, Ω^-

REFERENCES FOR $\Xi^{*1/2}(2250)$

BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LOIC, VIENNA)
 GOLDWASS 70 PR 10 1960 E L GOLDWASSER, P F SCHULTZ (ILLINOIS)

PAPERS NOT REFERRED TO IN DATA CARDS.

MORRIS 75 ANL-HEP-CP-75-58 MORRIS, CH, PARKER, SMITH, WHITMORE (MSU)

$\Xi(2500)$

99 $\Xi^{*1/2}(2500)$, JP=) I=1/2

IT IS QUITE POSSIBLE THAT THE REASON THE EXPERIMENTS
 DISAGREE ABOUT THE MASS AND WIDTH IS THAT THEY ARE
 SEEING DIFFERENT Ξ^{*} 'S. FOR NOW, HOWEVER, WE GROUP
 THEM TOGETHER.

99 $\Xi^{*1/2}(2500)$ MASS (MEV)

M	30	2430.0	20.0	ALITTI	69 HBC	-	K-P 4.6-5 GEV/C	9/69
M	45	2500.0	10.0	BARTSCH	69 HBC	-0	K-P 10 GEV/C	9/69
M		2392.	27.	DIBIANCA	75 DBC		XI 2PI	1/76*
M	AVERAGE MEANINGLESS (SCALE FACTOR = 3.2)							

99 $\Xi^{*1/2}(2500)$ WIDTH (MEV)

W	150.0	60.0	40.0	ALITTI	69 HBC	-		9/69
W	59.0	27.0		BARTSCH	69 HBC	-0		9/69
W	75.	69.		DIBIANCA	75 DBC		XI 2PI	1/76*
W	AVERAGE MEANINGLESS (SCALE FACTOR = 1.1)							

99 $\Xi^{*1/2}(2500)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\Xi^{*1/2}(2500)$ INTO XI PI	1321+ 139
P2	$\Xi^{*1/2}(2500)$ INTO LAMBDA KBAR	1115+ 497
P3	$\Xi^{*1/2}(2500)$ INTO SIGMA KBAR	1197+ 457
P4	$\Xi^{*1/2}(2500)$ INTO $\Xi^{*1/2}(1530)$ PI	1533+ 139
P5	$\Xi^{*1/2}(2500)$ INTO LAMBDA (OR SIGMA) KBAR PI	1115+ 497+ 139
P6	$\Xi^{*1/2}(2500)$ INTO XI PI PI	1321+ 139+ 139

99 $\Xi^{*1/2}(2500)$ BRANCHING RATIOS

R1	$\Xi^{*1/2}(2500)$ INTO (XI PI)/(MODES P1 THRU P4)	(P1)/(P1+P2+P3+P4)	
R1	(0.5) OR LESS	ALITTI 69 HBC	1 STD DEV LIMIT 9/69
R2	$\Xi^{*1/2}(2500)$ INTO (LAM KBAR)/(MODES P1 THRU P4)	(P2)/(P1+P2+P3+P4)	
R2	0.5 0.2	ALITTI 69 HBC	- 9/69
R3	$\Xi^{*1/2}(2500)$ INTO (SIG KBAR)/(MODES P1 THRU P4)	(P3)/(P1+P2+P3+P4)	
R3	0.5 0.2	ALITTI 69 HBC	- 9/69
R4	$\Xi^{*1/2}(2500)$ INTO (XI* PI)/(MODES P1 THRU P4)	(P4)/(P1+P2+P3+P4)	
R4	(0.2) OR LESS	ALITTI 69 HBC	1 STD DEV LIMIT 9/69
R5	$\Xi^{*1/2}(2500)$ INTO (LAMBDA (OR SIGMA) KBAR PI)/TOTAL	(P5)	
R5	SEEN	BARTSCH 69 HBC	-0 9/69
R6	$\Xi^{*1/2}(2500)$ INTO (XI PI PI)/TOTAL	(P6)	
R6	SEEN	BARTSCH 69 HBC	-0 9/69

 REFERENCES FOR $\Xi^{*1/2}(2500)$
 ALITTI 69 PRL 22 79 +BARNES, FLAMINIO, METZGER, + (BNL, SYRACUSE) I
 BARTSCH 69 PL 288 439 + (AACHEN, BERLIN, CERN, LOIC, VIENNA)
 DIBIANCA 75 NP 898 137 DIBIANCA, ENDORFR (CERN)

S=-3 I=0 HYPERON STATE (Ω)

Ω^-

24 OMEGA-(1675, JP=3/2+) I=0

SEE STABLE PARTICLE DATA CARD LISTINGS

Appendix I

TEST OF $\Delta I=1/2$ RULE FOR K DECAYS

The quantities of interest for making tests of theoretical predictions regarding the $\Delta I=1/2$ rule for K decay are usually partial decay rates for single channels or special sums of channels. It is not possible to compute the errors on sums, differences, and ratios of partial decay rates from the information given in the Table of Stable Particles because of the presence of off-diagonal terms in the error matrix. For this reason we give some of these quantities in Table I. Throughout this Appendix, italics are used to indicate that a quantity has changed by more than one (old) standard deviation since our previous edition, and S gives the scale factor included in the quoted error because of inconsistencies in the data (see footnote at end of Stable Particle Table for definition of S).

Table I. (000) and (+-0) refer to the sign of the pions into which the K_L decays.

$\Gamma_{K_{L3}^+} = \Gamma_{K_{e3}^+} + \Gamma_{K_{\mu 3}^+}$	$= (6.483 \pm 0.090) 10^6 \text{ sec}^{-1}$
$\Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}$	$= 0.663 \pm 0.018 \quad S=1.7^*$
$\Gamma_{K_{\mu 3}^+} / \Gamma_{K_{\mu 3}^+}$	$= 3.227 \pm 0.083$
$\Gamma_{K_{L3}^0} = \Gamma_{K_{e3}^0} + \Gamma_{K_{\mu 3}^0}$	$= (12.75 \pm 0.15) 10^6 \text{ sec}^{-1} \quad S=1.1^*$
$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0}$	$= 0.696 \pm 0.017$
$\Gamma_{K^0(000)} / \Gamma_{K^0(+ - 0)}$	$= 1.747 \pm 0.070 \quad S=1.1^*$

1. Leptonic decay rates

The $\Gamma_{K_{L3}}$ rates are useful in testing the leptonic $\Delta I=1/2$ rule in the way suggested by Trilling.¹ The predictions are

$$\Gamma_{K_{L3}^0} / 2\Gamma_{K_{L3}^+} = 1.012, \text{ a phase-space factor,}^2$$

and

$$\Gamma_{K_{\mu 3}^0} / \Gamma_{K_{e3}^0} = \Gamma_{K_{\mu 3}^+} / \Gamma_{K_{e3}^+}.$$

From Table 1,

$$\Gamma_{K_{L3}^0} / 2\Gamma_{K_{L3}^+} = 0.983 \pm 0.018$$

and

$$\frac{\Gamma_{K_{\mu 3}^0}}{\Gamma_{K_{e3}^0}} \left[\frac{\Gamma_{K_{\mu 3}^+}}{\Gamma_{K_{e3}^+}} \right]^{-1} = 1.050 \pm 0.038.$$

These results seem to show a less than 2σ disagreement with the predictions, but the errors should be regarded with caution in view of the internal disagreements in the data. (Note the ideograms in the Data Listings for the charged K meson.)

2. Three-pion decays

We follow here the tests done by Mast et al.,³ based on the general analysis of K decays suggested by Zemach.⁴ Both decay rates (Γ) and slopes (g , the energy dependence of the Dalitz plot distributions) are used. The $\Delta I=1/2$ rule predicts that the following test quantities are all equal to zero:

$$\text{Test 1} = \frac{2}{3} \frac{\Gamma_{K^0(000)}}{\phi_1} \left[\frac{\Gamma_{K^0(+ - 0)}}{\phi_2} \right]^{-1} - 1,$$

$$\text{Test 2} = \frac{1}{4} \frac{\Gamma_{K_{\mu 3}^+}}{\phi_3} \left[\frac{\Gamma_{K_{\mu 3}^+}}{\phi_4} \right]^{-1} - 1,$$

$$\text{Test 3} = \frac{1}{2} \frac{\Gamma_{K_{\mu 3}^+}}{\phi_3} \left[\frac{\Gamma_{K^0(+ - 0)}}{\phi_2} \right]^{-1} - 1,$$

$$\text{Test 4} = \frac{1}{2} g_{K_{\mu 3}^+} + g_{K_{\mu 3}^+},$$

$$\text{Test 5} = g_{K^0(+ - 0)} + g_{K_{\mu 3}^+} - \frac{1}{2} g_{K_{\mu 3}^+}.$$

The ϕ_i are phase-space factors which have been calculated as described in Mast et al.³ by use of a relativistic formulation and the masses and slopes from this edition. The factors labeled UDP are the relative areas of the Dalitz plots, assuming a uniform distribution. The NUDD include the observed slopes (see below). The CNUDD have been calculated by including the final-state Coulomb interaction.

The values are:

	Method		
	UDP	NUDD	CNUDD
$\phi_1(000) =$	1.489	1.489	1.444
$\phi_2(+ - 0) =$	1.221	1.300	1.284
$\phi_3(++ -) =$	1.000	1.000	1.000
$\phi_4(+00) =$	1.247	1.180	1.144

For convenience, we repeat the slope parameters tabulated in the Stable Particle Table. They are as follows:

$g_{K_{\mu 3}^+}$	$= -0.214 \pm 0.005$	$S=1.7^*$
$g_{K_{\mu 3}^-}$	$= -0.214 \pm 0.007$	$S=2.7^*$
$\bar{g}_{K_{\mu 3}^+}$	$= -0.214 \pm 0.004$	
$g_{K_{\mu 3}^0}$	$= 0.550 \pm 0.020$	$S=1.6^*$
$g_{K^0(+ - 0)}$	$= 0.646 \pm 0.014$	$S=2.5^*$

A difference in the τ^+ and τ^- slopes would be an indication of CP violation in this decay. Since no difference is observed at this time, we average the two and use this value in Test 4 and Test 5.

We use the CNUDP factors and the rates and slopes reported in this edition to compute the five test quantities which the $\Delta I=1/2$ rule predicts to be zero. The results are:

$$\begin{aligned} \text{Test 1} &= 0.036 \pm 0.042 \\ \text{Test 2} &= -0.077 \pm 0.024 \\ \text{Test 3} &= 0.227 \pm 0.021 \\ \text{Test 4} &= 0.061 \pm 0.011 \\ \text{Test 5} &= 0.157 \pm 0.018 \end{aligned}$$

The three-pion final state can be in isospin states $I = 1, 2, 3$. Tests 1 and 2 test the existence of isospin $I = 3$ in the final state. Since the rate tests (Tests 1, 2, and 3) could differ from zero by as much as 0.1 owing to the mass differences and the occurrence of big slopes⁵, no evidence for $I=3$ is found. Test 4 is related to the $I=2$ amplitude in the final state and indicates the presence of $I=2$. Tests 3 and 5 give information on the $\Delta I=3/2$ part of the $I=1$ amplitude relative to the $\Delta I=1/2$ part. Both tests indicate the presence of $\Delta I=3/2$.

References

1. G. Trilling, K-Meson Decays, UCRL-16473, (updated from Argonne Conference Proceedings, 1965, p. 115).
2. N. Brene (CERN), private communication. In our Jan. 1968 edition we had erroneously used 1.04.
3. T. S. Mast, L. K. Gershwin, M. Alston-Garnjost, R. O. Bangerter, A. Barbaro-Galtieri, J. J. Murray, F. T. Solmitz, and R. D. Tripp, Phys. Rev. **183**, 1200 (1969).
4. C. Zemach, Phys. Rev. **133**, B1201 (1964).
5. C. Bouchiat and M. Veltman, Topical Conference on Weak Interactions, CERN 69-7 (1969), p. 225.

Appendix II

A. SU(3) CLASSIFICATION OF BARYON RESONANCES

It is established that a symmetry higher than SU(3) is necessary to classify the known baryon resonances. However, many higher-symmetry schemes have been proposed, and even for SU(6) various versions exist (for a review see Dalitz¹). Since it is not clear which one of these schemes best fits the data, we do not review them here, but we report once again fits of baryon states into SU(3) multiplets.

For the reader's convenience, we collect here the relevant formulae.

Exact SU(3) symmetry predicts that all the members of a multiplet should have the same mass and the same couplings for decays into other multiplets. It has been found, however, that the members of the octet of stable baryons lie within 20% of their mean mass; therefore a symmetry-breaking interaction has been introduced by Gell-Mann and Okubo independently.² In addition, for the isospin-0 vector mesons (ω and ϕ), an additional symmetry-breaking interaction has been introduced by Sakurai³ to take care of octet-singlet mixing. The relevant formulae for masses and decay rates are given below.

Mass Formulae

Broken SU(3) gives:

$$\text{Decuplet} \quad \Delta - \Sigma = \Sigma - \Xi^* = \Xi^* - \Omega \quad \text{GMO} \quad (1)$$

$$\text{Octet} \quad 2(N + \Xi) = 3\Lambda + \Sigma \quad \text{GMO} \quad (2)$$

$$\left. \begin{array}{l} \text{Octet-} \\ \text{Singlet} \end{array} \right\} \sin^2 \theta = \frac{\Lambda - M_B}{\Lambda - \Lambda'} \quad \text{Mixing} \\ \text{mixing} \quad \left\{ \begin{array}{l} \Lambda - \Lambda' \\ \Lambda - \Lambda' \end{array} \right. \text{angle}^2 \quad (3)$$

$$M_B = \frac{2(N + \Xi) - \Sigma}{3} \quad \text{GMO} \quad (4)$$

Here GMO stands for the Gell-Mann-Okubo formula; the particle symbol indicates its mass. The formulae would be the same if squared masses were used. For the nonet case, Λ is the "mostly-octet" particle, Λ' is the "mostly-singlet" particle.

Decay Rates

In terms of a relativistically invariant matrix element T , the decay rate for two-body decay of a resonance of mass M_R is

$$\Gamma \propto \frac{|T|^2 R_2}{M_R}, \quad (5)$$

where $R_2 = k/M_R$ is the two-body phase space factor. Since the numerator is an invariant, and since Γ must transform as $1/E$, we introduce the denominator M_R .⁴

For meson decays (see below) the rates are calculated according to Eq. (5); for baryon resonance decays into $1/2^+$ baryons and 0^- mesons, one next takes into account the fact that spin sums in $|T|^2$ introduce another factor M_R , cancelling the $1/M_R$. We are then left with

$$\Gamma = \frac{|T|^2 k}{M_R} M_N, \text{ for baryons} \quad (5')$$

$$= \frac{|T|^2 k}{M_R^2} M_N^2, \text{ for mesons.} \quad (5'')$$

The powers of the nucleon mass M_N or M_N^2 have been introduced so that we can treat $|T|$ as dimensionless.

$|T|^2$ contains centrifugal barrier factors, which we call B_l . We then have

$$\left. \begin{array}{l} \text{Decuplet} \\ \text{Singlet} \end{array} \right\} \Gamma = (c_g)^2 B_l(k) \frac{M_N}{M_R} k \quad (6)$$

$$\text{Octet} \quad \Gamma = (c_D g_D + c_F g_F)^2 B_l(k) \frac{M_N}{M_R} k \quad (7)$$

$$\left. \begin{array}{l} \text{Octet-Singlet} \\ \text{mixing} \end{array} \right\} \begin{array}{l} \Lambda = G_8 \cos \theta + G_1 \sin \theta \\ \Lambda' = -G_8 \sin \theta + G_1 \cos \theta \end{array} \quad (8)$$

$$\text{with} \quad \begin{array}{l} G_8 = c_D g_D + c_F g_F \\ G_1 = c_1 g_1. \end{array} \quad (9)$$

Here c_i are the SU(3) coefficients with the sign convention adopted in this article [see note in the Table of SU(3) Isoscalar Factors and Fig. 2 in the text]. M_N is the nucleon mass, M_R is the resonance mass for which Γ is calculated, k is the center-of-mass momentum for the channel being considered, and g_i are the relevant couplings. For the case of singlet-octet mixing, formula (8) has to be used in conjunction with (6) and (7). G_8 and G_1 represent the couplings for the multiplet, and Λ and Λ' represent the couplings for the physical states.

The relation between g_D , g_F , and the parameter α is

$$\alpha = \left[1 + \frac{\sqrt{5}}{3} \frac{g_F}{g_D} \right]^{-1}. \quad (10)$$

Exact SU(3) predicts that the couplings g_i for all the members of a multiplet are the same; however, since the symmetry is broken for the masses, it is probably broken for the widths. In the case of the $3/2^+$ decuplet, for broken SU(3) a sum rule has been derived by Becchi⁵ and by Gupta⁶ independently. It relates the g_i for the members of the decuplet by the relation

$$2(\Delta + \Xi) = 3\Sigma^*(\Lambda\pi) + \Sigma^*(\Sigma\pi), \quad (11)$$

where $\Sigma^*(\Lambda\pi)$ is the coupling for the $\Sigma(1385) \rightarrow \Lambda\pi$ decay and $\Sigma^*(\Sigma\pi)$ is the coupling for the decay $\Sigma(1385) \rightarrow \Sigma\pi$.

As mentioned in the text (Sec. IV B) the determination of the relative signs of resonant amplitudes can be useful in making an SU(3) assignment of resonances. In fact the resonant amplitude $T \propto \sqrt{\kappa_e \kappa_i} \propto G_e G_i$, where the subscript e refers to the elastic channel and the G_e , G_i are the couplings of Eqs. (6) through (9). Assuming that all g_i are positive, the sign of the G_i are dependent upon the sign of the Clebsh-Gordon coefficients c_i . Once a sign convention is adopted (we use the Levi-Setti⁷ convention, see Fig. 2 in the text) and the signs for a Σ state ($I=1$) and a Λ state ($I=0$) of known SU(3) assignment have been chosen for reference, the signs of all the other amplitudes can be useful in determining multiplet assignments. For exact SU(3) all the decays of members of a decuplet have the same sign. For octets the relative sign depends upon the value of g_D/g_F and the mixing angle, as seen from Eqs.(7) through (9).

Fits to the Data

Fits of baryon decay rates within SU(3) can be found in, among others, papers by Tripp,^{8,9} Levi-Setti,⁷ Samios,¹⁰ and Plane.¹¹ The most recent fits were made by Barbaro-Galtieri¹² and Samios.¹³

In fitting the data a choice for B_l has to be made. Plane¹¹ tried two forms for B_l :

(a) The form $B_l = (kr)^{2l} D_l(kr)$, r being the radius of interaction and D_l the polynomials in kr given by Blatt and Weisskopf.¹⁴ Usually r is taken to be 1 fermi.⁸

(b) The form $B_l = k^{2l}$.

However, for final results form (b) was chosen. A discussion of the differences among these two forms has been given by Barbaro-Galtieri.¹⁵ It turns out that not only the values of the couplings, g_i , depend upon the form used for B_L , but also the value obtained for the mixing angle. For the $3/2^-$ singlet, $\Lambda(1520)$, and the isospin-0 member of the octet, $\Lambda(1690)$, the mixing angles obtained in the two cases are

$$\theta_a = (-16.1^{+1.4}_{-1.3})^\circ, \theta_b = (-27.5^{+3.6}_{-3.4})^\circ,$$

in disagreement by a few standard deviations. It turns out that if a radius of interaction of $r = 0.15$ fermi is used for form (a), the two values of θ agree. This value of r does not fit resonance shapes when used in the Breit-Wigner resonant form.

Samios¹³ used form (b) for B_L .

Table I is a summary of the fits made by us (update of Barbaro-Galtieri¹²) using the barrier factor form (a) and exact SU(3). A few comments follow.

1/2⁻ Nonet (Baryon-Eta Resonances)

For this nonet Eq. (7) was multiplied by the factor

$$\left[\frac{M_R - M_B}{\bar{M}_R - \bar{M}_B} \right]^2,$$

where M_B is the decay baryon and $\bar{M}_R - \bar{M}_B = 564$ MeV is the difference of the mean $1/2^-$ and $1/2^+$ baryon octet masses. This kinematic factor comes from PCAC

arguments (i.e., the assumption that the axial vector current remains an octet in the presence of symmetry breaking) and it was advocated by Graham.¹⁶ For the $1/2^-$ nonet it was used in this form first by Gell-Mann.¹⁷

3/2⁺ Decuplet

The agreement among the coupling constants obtained for the four rates in this decuplet is very bad. The fit made using form (b) for B_L has $\chi^2=50$ for 3 degrees of freedom; the one made with form (a) for B_L has $\chi^2/DF=24/3$. The broken SU(3) relation (11), however, is very well satisfied.

B. SU(3) CLASSIFICATION OF MESON RESONANCES

All of the discussion above applies, except that for bosons the GMO formula is usually applied to the square of the masses, as opposed to the first power for fermions. Thus for example, Eq. (2) becomes

$$4\hat{K} = 3\hat{\eta} + \hat{\pi}. \tag{2'}$$

The symbol \hat{K} was introduced by Glashow and Socolow¹⁸ for the square of the K mass, etc.

Because of the difference between Eqs. (5') and (5''), there is also an extra factor of (M_N/M_R) in Eqs. (6) and (7). The three established nonets ($0^-, 1^-, 2^+$) and their mixing angles are listed at the bottom of the Meson Table.

Table I. SU(3) baryon multiplets with two or more known members. Values of θ and α [defined by Eqs. (8) and (10)] are the result of fits made to all the measured two-body decay rates of each multiplet.

J^P	Octet members ^a				Singlet	$\theta(\text{deg})^b$	α
$1/2^-$	N(1535)	$\Lambda(1670)$	$\Sigma(1750)$	$[\Xi(1825)]$	$\Lambda(1405)$	3 ± 5	$0.93 \pm .11$
$3/2^-$	N(1520)	$\Lambda(1690)$	$\Sigma(1670)$	$[\Xi(1815)]$	$\Lambda(1520)$	-23 ± 4	$0.31 \pm .05$
$5/2^-$	N(1670)	$\Lambda(1830)$	$\Sigma(1765)$				$1.17 \pm .04$
$5/2^+$	N(1688)	$\Lambda(1815)$	$\Sigma(1915)$				$0.65 \pm .03$
Decuplet members ^c					g_{10}		
$3/2^+$	$\Delta(1232)$	$\Sigma(1385)$	$\Xi(1530)$	Ω^-	$1.78-2.29$	$\chi^2/DF=50/3$	
$7/2^+$	$\Delta(1950)$	$\Sigma(2030)$					

^aMasses in parentheses are the nominal masses used in the Baryon Table. The Ξ members have masses as calculated by using formulae (1) and (2) with the mixing angle θ derived from the decay widths.

^bSee text for a discussion of the $3/2^-$ mixing angle.

^cCoupling constants from Ref. 12.

References

1. R. H. Dalitz, *Baryonic Spectroscopy and Its Immediate Future*, paper presented at the Summer Symposium, New Directions in Hadron Spectroscopy, Argonne National Laboratory (August 1975).
2. M. Gell-Mann, *Phys. Rev.* **125**, 1067 (1962); S. Okubo, *Prog. Theor. Phys. (Kyoto)* **27**, 949 (1962).
3. J. J. Sakurai, *Phys. Rev. Letters* **9**, 472 (1962).
4. See R. P. Feynman, *Theory of Fundamental Processes*, W. A. Benjamin, Inc., New York, 1962.
5. C. Becchi, E. Eberle, and G. Morpurgo, *Phys. Rev.* **136B**, 808 (1964).
6. V. Gupta and V. Singh, *Phys. Rev.* **135B**, 1442 (1964).
7. R. Levi-Setti, in *Proceedings of the Lund International Conference on Elementary Particles, Lund, 1969*.
8. R. D. Tripp, in *Proceedings of the 14th International Conference on High Energy Physics, Vienna, 1968*, p. 173.
9. R. D. Tripp, in *Proceedings of the 3rd Hawaiian Topical Conference on Particle Physics; UCRL-19361* (1969).
10. N. P. Samios, in *Proceedings of the 15th International Conference on High Energy Physics, Kiev, 1970*, p. 187.
11. D. E. Plane et al., *Nuclear Physics* **B22**, 93 (1970). Also J. Meyer and D. E. Plane, *Nuclear Physics* **B25**, 428 (1971).
12. A. Barbaro-Galtieri, LBL-1366 and in *Proceedings of the 16th International Conference on High Energy Physics*, National Accelerator Laboratory, Vol. 1, page 159 (1972).
13. N. P. Samios, M. Goldberg, and B. T. Meadows in *Hadrons and SU(3): A Critical Review*, BNL Report BNL-17851 (1973).
14. J. M. Blatt and V. F. Weisskopf, *Theoretical Nuclear Physics*, Wiley, New York, 1952.
15. A. Barbaro-Galtieri, in *Properties of Fundamental Interactions*, Erice, July 8-26, 1971, edited by A. Zichichi, Editrice Compositori, page 533 (1973).
16. R. Graham, S. Pakvasa, and K. Raman, *Phys. Rev.* **163**, 1774 (1957).
17. M. Gell-Mann, R. Oakes, and B. Renner, *Phys. Rev.* **175**, 2195 (1968).
18. The formula has been calculated from analogy with the formula for mixing of meson states, first put in this form by S. L. Glashow and R. H. Socolow, *Phys. Rev. Letters* **15**, 329 (1966). For the baryon formula see A. Barbaro-Galtieri, *Phenomenology of Resonances and Particle Supermultiplets*, UCRL-17054 (1966).

Appendix III

TEST OF $\Delta I=1/2$ RULE FOR HYPERON DECAYS

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1. Nonleptonic decay Amplitudes

In this edition we again use the new convention for the amplitudes A and B adopted in 1973. Some theorists have suggested that dimensionless amplitudes are more useful to them than the ones appearing in the literature. Berge¹ used a convention with A and B in units of $\text{sec}^{-1/2}$. Samios² used a convention which gave A and B in units of $(\text{MeV}\text{-sec})^{-1/2}$. Following is the convention suggested by Jackson³, which gives dimensionless A and B.

The effective Lagrangian density for nonleptonic hyperon decays ($B_1 \rightarrow B_2 + \pi$) can be written

$$L_{\text{eff}} = G\mu_c^2 [\bar{\psi}_2(A+B\gamma_5)\psi_1]\phi_\pi,$$

where $G=10^{-5}m_p^{-2}$ is a coupling constant characteristic of first-order weak decays, μ_c is the charged pion mass, and A and B are dimensionless complex numbers giving the relative amplitudes of the parity-violating and parity-conserving decays, respectively. The matrix γ_5 is to be taken in the Pauli form, $\gamma_5 = \begin{pmatrix} 0 & -I \\ -I & 0 \end{pmatrix}$. The invariant amplitude for the decay is

$$M = G\mu_c^2 [\bar{u}(p)(A+B\gamma_5)u(P)],$$

where P is the 4-momentum of the decaying hyperon of mass M, and p is the 4-momentum of the baryon decay product of mass m. With the normalization convention, $\bar{u}_i u_i = 2m_i$, the Pauli form of the matrix element in the rest frame of the decaying hyperon is

$$M = G\mu_c^2 (\chi_2 | \sqrt{2M(E+m)}A + \sqrt{2M(E-m)}B\vec{\sigma}\cdot\hat{q} | \chi_1),$$

where E is the total energy of the final baryon and \hat{q} is a unit vector in the direction of motion of the final baryon. Comparison with Sec. VI D of the text shows that the amplitudes s and p defined there are proportional to A and B:

$$\frac{p}{s} = \left(\frac{E-m}{E+m} \right)^{1/2} \frac{B}{A} = \left[\frac{(M-m)^2 - \mu^2}{(M+m)^2 - \mu^2} \right]^{1/2} \frac{B}{A}.$$

Here μ is the mass of the pion entering the decay. The parameters α , β , and γ can therefore be expressed in terms of A and B, rather than s and p, if desired.

The decay rate for $B_1 \rightarrow B_2 + \pi$ is

$$\Gamma = \frac{G^2 \mu_c^4}{8\pi q} \left\{ \left[\frac{(M+m)^2 - \mu^2}{M^2} \right] |A|^2 + \left[\frac{(M-m)^2 - \mu^2}{M^2} \right] |B|^2 \right\}.$$

where q is the c.m. momentum of the decay products. For reference, the dimensionless constant in this expression has the value $(G^2 \mu_c^4 / 8\pi) = 1.9488 \times 10^{-15}$.

Table I summarizes the amplitudes A and B for the nonleptonic decays of the Λ , Σ , and Ξ hyperons. These amplitudes have been calculated by using the experimental data for mean lives, branching ratios, and the decay asymmetry α given in the Stable Particle Table of this Review. Time-reversal invariance is assumed and final-state interactions are neglected, so A and B are taken to be relatively real. The subscript on the hyperon refers to the sign of the decaying pion. The statistical correlation coefficient

$$C_{AB} = \frac{\langle \Delta A \Delta B \rangle}{\sqrt{\langle \Delta A^2 \rangle \langle \Delta B^2 \rangle}}$$

is also given. The absolute signs of A and B have been assigned, using the following convention. Taking $A(\Lambda^0)$ as positive, the other S -wave decay amplitudes are chosen to give an approximate fit to the triangular relationships

$$\sqrt{2}A(\Sigma_0^+) + A(\Sigma_+^+) = A(\Sigma^-) \text{ and } \sqrt{3}A(\Sigma_0^+) + A(\Lambda^0) = 2A(\Xi^-).$$

The signs of the B amplitudes relative to those of the corresponding A amplitudes are determined by the sign of the appropriate α decay parameter.

Table I

$M \rightarrow m + \mu$	A	B	C_{AB}
$\Lambda^0 \rightarrow p + \pi^-$	1.48 ± 0.01	10.17 ± 0.24	-0.272
$\Lambda_0^0 \rightarrow n + \pi^0$	-1.08 ± 0.02	-7.28 ± 0.59	-0.747
$\Sigma_+^+ \rightarrow n + \pi^+$	0.06 ± 0.02	19.06 ± 0.16	0.003
$\Sigma_0^+ \rightarrow p + \pi^0$	1.48 ± 0.05	-12.04 ± 0.59	0.918
$\Sigma^- \rightarrow n + \pi^-$	1.93 ± 0.01	-0.65 ± 0.08	-0.024
$\Xi_0^0 \rightarrow \Lambda + \pi^0$	1.53 ± 0.03	-5.90 ± 1.11	0.347
$\Xi^- \rightarrow \Lambda + \pi^-$	2.04 ± 0.02	-6.71 ± 0.38	0.198

2. Tests of the $\Delta I=1/2$ Rule

(a) Λ Decay

For Λ decay the $\Delta I=1/2$ rule predicts that $\Gamma_0/\Gamma_- = 0.50$ and $\alpha_0 = \alpha_-$. In order to determine the magnitude of possible $\Delta I=3/2$ amplitudes present we write the linear expressions⁴ for the $\Delta I=3/2$ A- and B-wave amplitudes in terms of $\Delta\alpha$, where $\Delta\alpha$ is the measured value of α_0/α_- minus the predicted value, and in terms of $\Delta\Gamma$ similarly defined. Evaluating these we find

$$\Delta\alpha = -1.53 (A_3/A_1) + 1.60 (B_3/B_1),$$

$$\Delta\Gamma = 1.83 (A_3/A_1) + 0.26 (B_3/B_1).$$

Here the $\Delta I=3/2$ amplitudes are expressed relative to the $\Delta I=1/2$ amplitudes. The numerical values of the coefficients depend on the ratio B/A . The uncertainties in the coefficients are small compared to the uncertainties in $\Delta\alpha$ and $\Delta\Gamma$. Final-state πN interactions have been included in these relations but have a very small effect. From the Stable Particle Table,

$$\Delta\alpha = 0.006 \pm 0.066, \quad \Delta\Gamma = 0.058 \pm 0.012,$$

and hence

$$(A_3/A_1) = 0.027 \pm 0.008$$

and

$$(B_3/B_1) = 0.030 \pm 0.037.$$

The possible 3% $\Delta I=3/2$ A-wave amplitude is due to the disagreement of decay rates with prediction. At this level the results are sensitive to electromagnetic corrections. However, in Λ decay the phase space correction and the other radiative corrections appear to be about equal in magnitude and have opposite signs,^{5,6} and hence cancel each other in the correction to the decay rates.

(b) Ξ Decay

The analysis for Ξ decay is very similar to that for Λ decay. If the $\Delta I=1/2$ rule is valid, $\Gamma_0(\Xi^0)/\Gamma_-(\Xi^-) = 0.50$ and $\alpha_0 = \alpha_-$. For this case the expressions linear in $\Delta I=3/2$ A- and B-wave amplitudes are⁴

$$\Delta\alpha = 1.38 (A_3/A_1) - 1.38 (B_3/B_1),$$

$$\Delta\Gamma = -1.44 (A_3/A_1) - 0.06 (B_3/B_1).$$

From the Stable Particle Table,

$$\Delta\alpha = 0.12 \pm 0.21, \quad \Delta\Gamma = 0.058 \pm 0.024,$$

and we find

$$(A_3/A_1) = -0.035 \pm 0.017$$

and

$$(B_3/B_1) = -0.13 \pm 0.15.$$

(c) Σ Decay

The traditional test of the $\Delta I=1/2$ rule in Σ decay is that the amplitudes satisfy the relationship

$$\sqrt{2} \Sigma_0^+ + \Sigma_+^+ - \Sigma^- = 0.$$

Graphically this is equivalent to closing the Σ triangle when the amplitudes are plotted on A, B axes. Including $\Delta I \geq 3/2$ amplitudes in Σ decay analysis, the " Σ triangle" relationship becomes

$$\sqrt{2} A_0 + A_+ - A_- = -3\sqrt{2/5} A_3 + \frac{2}{\sqrt{15}} A_5,$$

where A_3 and A_5 are $\Delta I=3/2$ and $\Delta I=5/2$ amplitudes, respectively. There is a similar equation for the B amplitudes. From Table I,

$$\sqrt{2} A_0 + A_+ - A_- = 0.22 \pm 0.09$$

and

$$\sqrt{2}B_0 + B_+ - B_- = 2.7 \pm 1.1 .$$

If we neglect the $\Delta I=5/2$ amplitudes and assume all amplitudes to be real we can solve for possible $\Delta I=3/2$ amplitudes. The result is

$$\frac{A_3}{A_-} = -0.060 \pm 0.026$$

and

$$\frac{B_3}{B_+} = -0.074 \pm 0.030 .$$

Thus for hyperon decay, present experimental data limit $\Delta I=3/2$ amplitudes to less than about 5%.

3. The Lee-Sugawara Relation

From Table I the Lee-Sugawara relation,^{7,8} $\sqrt{3}\Sigma_0^+ + \Lambda^0 - 2\Sigma^- = 0$, is satisfied to -0.05 ± 0.12 for the A amplitudes, and to 2.7 ± 2.0 for the B amplitudes.

References

1. J. P. Berge, in Proceedings of the 13th International Conference on High-Energy Physics, Berkeley, (1966) (University of California Press, Berkeley, 1967), p. 46.
2. N. P. Samios, International Conference on Weak Interactions, Argonne, (1965), p. 189.
3. J. D. Jackson, private communication (1973).
4. See O. E. Overseth and S. Pakvasa, Phys. Rev. **184**, 1663 (1969). The expression for Γ_0/Γ_- for Λ decay should read

$$\frac{\Gamma_0}{\Gamma_-} \approx \frac{1}{2} \left\{ 1 + 3\sqrt{2} \times \left[\frac{S_{11}S_{33}\cos(\delta_1 - \delta_3) + P_{11}P_{33}\cos(\delta_{11} - \delta_{31})}{S_{11}^2 + P_{11}^2} \right] \right\} .$$

5. See A. A. Belavin and I. M. Narodetsky, Yadern. Fiz. **8**, 978 (1968) [Soviet J. Nucl. Phys. **8**, 568 (1969)].
6. G. W. Intemann, private communication (1973).
7. See B. W. Lee, Phys. Rev. Lett. **12**, 83 (1964).
8. See H. Sugawara, Prog. Theor. Phys. **31**, 213 (1964).

Appendix IV

GROWTH OF INFORMATION

From time to time we have presented figures demonstrating the amount of experimental work which has gone into spectroscopy, and the amount of new information available as a result. The 1976 versions of these figures are shown as Figs. 1, 2, and 3.

Figure 1 is a simple count of the number of meson resonances listed in the Tables, categorized as those "understood" -- i.e., all quantum numbers are believed known -- and those simply "listed". For the 1976 edition, there is an increase in both of these categories because of the discovery of the J/ψ and related particles.

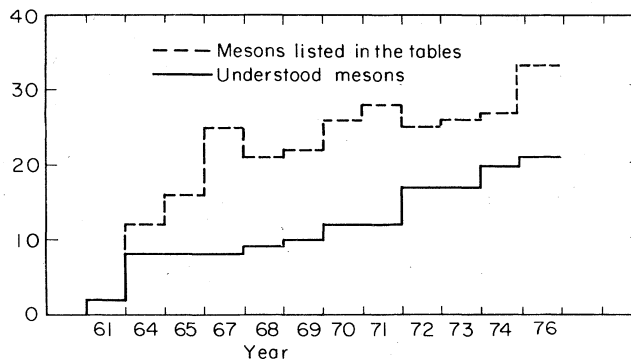


Fig. 1. Number of meson resonances listed in the Tables (dashed line) and those for which all quantum numbers are known (solid line), as a function of year of publication of the Review of Particle Properties. Note abscissa omits years in which no publication occurred.

In Figure 2 we present similar information for the baryon resonances, but concentrating here on the "growth of understanding". That is, the number of known baryons (we include for this figure only those with known J^P) has grown only very slowly with time (dashed line); the real progress has been in the measurement of the properties of those baryons. Therefore we show as the solid curve a count of the number of baryonic properties -- mass, width, and branching ratios. Most of these results are from partial-wave analyses.

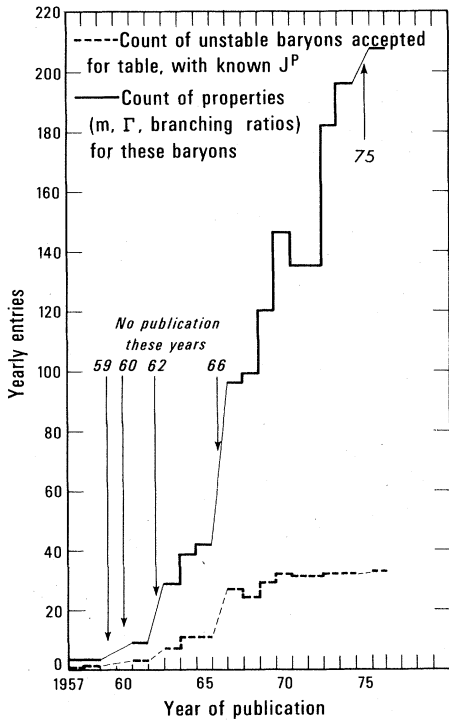


Fig. 2. Total amount of information (mass + width + branching ratios) on baryon resonances listed in the Tables, restricted to those with well-established J^P (solid line). Dashed line shows numbers of such resonances listed. Abscissa shows year of publication of Review of Particle Properties.

Finally, in Figure 3 we show a count of the number of new results put in the Listings each year, shown according to the type of detector that was used in obtaining the result, and according to the type of particle. N^* and Z^* particles have been omitted from this figure, because for these we use mainly the results of partial-wave analyses, rather than the primary data.

The fall in productivity in recent years doubtless reflects the tight budget situation and the declining emphasis on spectroscopy. Intense activity on the J/ψ and related particles is evident in the curve for meson resonances.

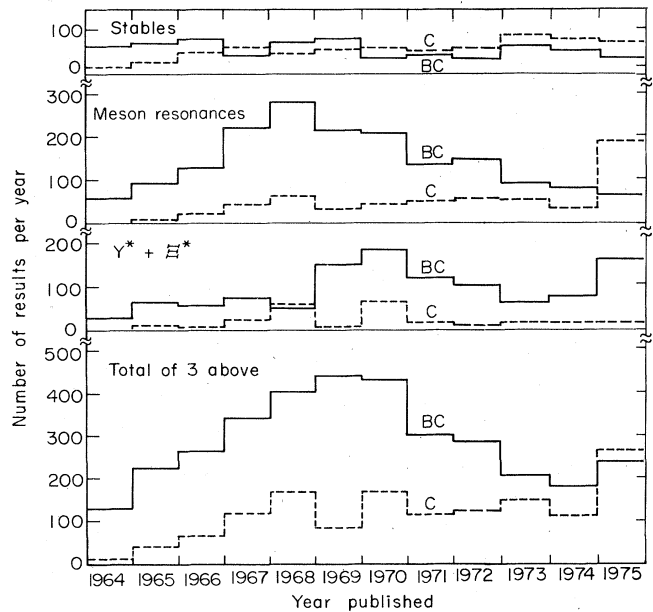


Fig. 3. The rate of production of data on particle properties, as a function of year of publication of the original result, based on data cards added to the Listings. Solid line shows bubble chamber results, dashed line counter or other electronic systems results.

