General Resource Theories in Quantum Mechanics and Beyond: Operational Characterization via Discrimination Tasks

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We establish an operational characterization of general convex resource theories—describing the resource content of not only states but also measurements and channels, both within quantum mechanics and in general probabilistic theories (GPTs)—in the context of state and channel discrimination. We find that discrimination tasks provide a unified operational description for quantification and manipulation of resources by showing that the family of robustness measures can be understood as the maximum advantage provided by any physical resource in several different discrimination tasks, as well as establishing that such discrimination problems can fully characterize the allowed transformations within the given resource theory. Specifically, we introduce a quantifier of resourcefulness of a measurement in any GPT, the generalized robustness of measurement, and show that it admits an operational interpretation as the maximum advantage that a given measurement provides over resourceless measurements in all state discrimination tasks. In the special case of quantum mechanics, we connect discrimination problems with single-shot information theory by showing that the generalized robustness of any measurement can be alternatively understood as the maximal increase in one-shot accessible information when compared to free measurements. We introduce two different approaches to quantifying the resource content of a physical channel based on the generalized robustness measures and show that they quantify the maximum advantage that a resourceful channel can provide in several classes of state and channel discrimination tasks. Furthermore, we endow another measure of resourcefulness of states, the standard robustness, with an operational meaning in general GPTs as the exact quantifier of the maximum advantage that a state can provide in binary channel discrimination tasks. Finally, we establish that several classes of channel and state discrimination tasks form complete families of monotones fully characterizing the transformations of states and measurements, respectively, under general classes of free operations. Our results establish a fundamental connection between the operational tasks of discrimination and core concepts of resource theories—the geometric quantification of resources and resource manipulation—valid for all physical theories beyond quantum mechanics with no additional assumptions about the structure of the GPT required.

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I. INTRODUCTION

The advantages provided by quantum phenomena in the transfer and processing of information allowed for the technological boom currently transforming areas such as communication, computation, cryptography, and sensing [1,2]. The realization that intrinsic physical properties of

Published by the American Physical Society under the terms of the Creative Commons Attribution 4.0 International license. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. quantum mechanics can be regarded precisely as *resources* in information processing tasks sparked an investigation of quantum information in the so-called resource-theoretic setting, aiming to establish the theoretical and practical methods to characterize both the advantages and the limitations associated with different physical properties of quantum systems, measurements, and transformations [3]. Such resource theories are now commonplace in the study of a diverse range of phenomena, such as entanglement [4,5], coherence [6–8], asymmetry [9,10], quantum thermodynamics [11,12], steering [13], non-Markovianity [14–16], magic [17,18], non-Gaussianity [19–21], measurement simulability and incompatibility [22–24], measurement informativeness [25], and quantum memory of channels [26].

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Although applications of the resource-theoretic framework have enhanced systematic studies of many physical settings and significantly contributed to a deeper understanding of our capabilities in manipulating and exploiting such resources, one could wonder whether the generality of the framework allows one to go beyond specific examples and obtain results applicable to a broad class of settings, thus providing a unified picture of resources in general. A complete study of which features are universal among all resources, stemming from only the very foundations of quantum mechanics, therefore remains a major area of investigation, and such an approach of general resource theories has recently gained much attention [3,27-37]. Indeed, although the framework of resource theories began with the characterization of the properties of quantum states, it has recently been adapted to the study of quantum channels [26,37–45] and measurements [24,25,46–48], allowing for the description of dynamic resources on a similar footing to static ones and thus motivating the question of whether all such resources can be described by a unified formalism.

In fact, one can pose an even more fundamental question: can common features of resource theories be understood without relying on quantum mechanics at all? Despite the success of quantum mechanics, the ongoing search for an axiomatic theory of probability and correlations in nature has provided us with insight into physical theories beyond quantum, as well as allowed for a straightforward unification of the methods required to characterize physical theories including classical and quantum probability theory. The formalism of general probabilistic theories (GPTs) [49-52] lends itself perfectly to the investigation of states, measurements, and their transformations at a fully fundamental and general level. It is particularly suited to illuminate which assumptions and which basic features of a theory lead to operational consequences, allowing one to identify the exact axioms one has to accept in order to recover the features of quantum theory [53–60]. This problem leads us to extend the framework of resource theories to general probabilistic theories and investigate a unified characterization of general resources in the extensive formalism of GPTs rather than limiting it to quantum mechanics, as has been previously considered for specific examples of resource theories [61–65].

As the very word "resource" suggests, understanding the operational aspects of resources—how they can be utilized for physical tasks and what limitations a resource theory places on the conversion of physical resources—has central importance both theoretically and practically. However, it frequently requires resource-specific approaches and does not easily generalize to encompass all physically relevant resource theories, and it is thus highly desired to find a fundamental class of operational tasks that would allow for the understanding of the resourcefulness of a given physical property in general settings. A promising candidate for

such a class of operational tasks which, on the one hand, lie at the heart of the nonclassical features of quantum theory [66–74] as well as GPTs [49,50,75–77] and, on the other hand, have found relevance in several existing resource theories, are the tasks of state and channel discrimination. In particular, Piani and Watrous [78] first showed that for every entangled quantum state, there exists a channel discrimination task in which it is more useful than any separable state. Similar results were subsequently found in several different resource theories of states [36,79-82] and measurements [25,48,83], and the work of Takagi et al. [36] finally showed that this property is shared by any convex resource theory of quantum states. It remains to understand how general this property truly is and whether all resources—both static and dynamic, both within quantum mechanics and beyond—can provide explicit advantages in such operational tasks.

A fundamental aspect of any resource theory is its quantification, which aims to measure the amount of inherent resources contained in a given physical object and allows for a quantitative comparison with other objects within the resource theory. This can be approached in many inequivalent ways, and a plethora of possible choices of resource measures exist [3,34]. A natural question which arises in this context is whether the given measure can be understood in an operational sense, assessing exactly the usefulness of a given object in some physical task; however, establishing such an interpretation for a given quantifier is often highly nontrivial. The family of so-called robustness measures [34,84] stands out in this context, as two prominent members of the family have found several applications in operational settings: these are the generalized robustness [33,36,79-82,85-91] and the standard robustness [18,89,92]. They are not only fundamental resource quantifiers faithfully capturing the resourcefulness of given objects with clear geometric interpretations but also significantly relevant to experiments—they are directly observable, that is, can be obtained in an experiment by measuring a single, suitably chosen observable rather than requiring complicated and expensive methods such as state tomography, allowing for the experimental quantification of resources [93,94]. In particular, the investigation of discrimination tasks in resource theories of states [36,79-82] revealed a notable similarity: the advantage that a given resource provides in such discrimination tasks is often quantified precisely by the generalized robustness. The generality of this quantitative relation was unveiled in Ref. [36], which showed that it is true in every convex resource theory of quantum states. This, together with recent progress in the resource theories of measurements which showed a similar interpretation of robustness quantifiers for measurements in specific settings [25,83], suggests that a unified operational interpretation of generalized robustness which could account for dynamic resources in addition to static ones might be possible. However, no such universal interpretation of any of the robustness measures has been obtained thus far, and in fact, despite several known applications of the generalized robustness in the context of discrimination, it has not been known whether the standard robustness plays a role in understanding such problems whatsoever.

The other predominant problem which resource theories are expected to tackle is the manipulation of resources, which asks whether it is possible to transform one resource to another when constrained to employ only the transformations allowed within the given resource theory. It is particularly insightful to understand this question in relation to the operational and quantitative aspects of resource theories: does there exist a family of resource measures or operational tasks which completely characterizes resource manipulation? The work of Skrzypczyk and Linden [25] indicated a potential of discrimination-type problems in this respect by showing that a family of state discrimination tasks fully characterize the simulability of measurements by classical postprocessing. A comprehensive extension of this type of an operational characterization to more general settings which, together with quantification, would complete an operational characterization of general resource theories, has hitherto remained elusive.

A. Summary of results

In this work, we solve the problems raised above under the umbrella of operational tasks of discrimination; specifically, we characterize general convex resource theories of states, measurements, and channels, establishing tools for their quantification, endowing the class of robustness measures with an explicit operational interpretation as the advantage that a physical object can provide in various discrimination tasks, and showing that such discrimination tasks fully characterize the conversion between states or measurements with free operations of the given resource theory. The generality of our methods and results establishes a universal operational description of resources in GPTs, revealing strong connections between several aspects of general resources and showing that the underlying convex structures can provide deep insight into the properties of physical systems also in an operational setting. We stress that all of our results are immediately applicable to broad classes of physically relevant quantum resource theories of states (including entanglement, coherence, magic, asymmetry, athermality, etc.), measurements (informativeness, simulability, separable and positive partial transpose measurements, etc.), and channels (quantum memories, free channels in the resource theories of states, etc.).

We begin by providing an introduction to the main concepts of general probabilistic theories, discrimination problems, and resource quantification in Sec. II. Our investigation starts in Sec. III by extending the results of Ref. [36] beyond quantum mechanics and showing that the generalized robustness of states in any convex resource

theory and any GPT can be understood as the quantifier of the advantage that a given resourceful state can provide in channel discrimination problems. We then introduce the measurement robustness in Sec. IV, which quantifies the resource content of any measurement in convex resource theories of measurements, generalizing the recent approach of Ref. [25] to arbitrary resources and probabilistic theories. In particular, we establish a direct operational interpretation of the measurement robustness in any resource theory by showing that it quantifies exactly the maximum advantage that a measurement can provide in all state discrimination tasks compared to all resourceless measurements. Having established the connection between discrimination tasks and generalized robustness of measurements, in Sec. IVA we further extend this connection to single-shot information theory within the setting of quantum mechanics. We specifically show that the increase in min-accessible information, a single-shot variant of accessible information, of an ensemble of states effected by applying resourceful measurements as compared to free measurements is exactly quantified by the generalized robustness.

The considerations extend beyond the case of states and measurements. In Sec. V, we formalize the quantification of the resource content of channels in two different ways: by measuring the amount of a resource that a channel can generate by acting on a free state, as well as in a more abstract formalism of convex resource theories of channels. In particular, we introduce resource-generating power, as well as generalized robustness of channels defined in general resource theories. We establish operational interpretations of these different robustness measures by showing that the resource-generating power of a channel exactly characterizes the advantage that the channel can enable in resourceless-state discrimination tasks, and by showing that, within quantum theory, the robustness of a given channel or the maximum robustness of a given ensemble of channels quantifies the advantage it provides in a class of state and channel discrimination tasks, respectively.

We additionally consider in Sec. VI another resource measure in the robustness family, the standard robustness, and show that it admits a universal operational interpretation in any convex resource theory of states in the context of quantifying the advantage that a state provides over resourceless states in all balanced binary channel discrimination tasks. This gives a general operational meaning to this quantity for the first time, extending the link between resource quantification and operational advantages in discrimination tasks to the standard robustness of states defined in general settings.

Finally, in Sec. VII we show that different classes of discrimination tasks can form complete sets of monotones in any resource theory, in the sense that a state or a measurement can be transformed into another state or measurement using only free operations of the given resource theory if and only if the former performs better

than the latter in a family of channel or state discrimination tasks. We therefore reveal explicitly that discrimination tasks can be useful not only in the context of quantifying the resource strength but also in fully characterizing the conversions between resource states or measurements. We also extend the results to describe the transformations of state ensembles.

In addition to operationally characterizing two main concepts of resource theories, quantification and manipulation of resources, our results not only endow the robustness measures with direct operational interpretations in general resource theories of states, measurements, and channels but also explicitly demonstrate strong relations between several seemingly unrelated notions—discrimination-type tasks, geometric resource quantification, resource transformations, and one-shot quantum-information-theoretic quantities (see Fig. 1)—as conjectured in Ref. [25] for quantum mechanics. The majority of our results apply to every single physical probabilistic theory in finite dimensions, relying only on the convex structure of the underlying resources and requiring no assumptions about the structure of the GPT beyond the basic axiom referred to as the no-restriction hypothesis [49,53,54]. We make use of methods in convex analysis and in particular conic programming.

II. PRELIMINARIES

A. General probabilistic theories

We briefly outline the basic formalism of general probabilistic theories. We refer the interested reader to Chaps. 1 and 2 of Ref. [52], which provides a modern introduction to the topic and a detailed discussion, derivation, and justification of the concepts.

The physical setting of a GPT can be identified with a convex and closed set of *states* $\Omega(\mathcal{V})$ in a finite-dimensional real complete normed vector space \mathcal{V} and a set of *effects*

contained in the dual vector space V^* , which correspond to the results of physically implementable measurements.

A crucial role is played by the cone generated by $\Omega(\mathcal{V})$, $\mathcal{C} \coloneqq \{\lambda \omega | \lambda \in \mathbb{R}_+, \omega \in \Omega(\mathcal{V})\}$, which we further assume to be pointed [i.e., $\mathcal{C} \cup (-\mathcal{C}) = \{0\}$] and generating (i.e., span $\mathcal{C} = \mathcal{V}$), such that it induces a partial order on the space \mathcal{V} given by $x \leq_{\mathcal{C}} y \Leftrightarrow y - x \in \mathcal{C}$, with $x <_{\mathcal{C}} y \Leftrightarrow y - x \in \text{int}(\mathcal{C})$. The dual cone $\mathcal{C}^* \coloneqq \{E \in \mathcal{V}^* | \langle E, \omega \rangle \geq 0 \ \forall \ \omega \in \mathcal{C}\}$, where we write $\langle E, \omega \rangle$ for $E(\omega)$, then similarly induces a partial order $E \leq_{\mathcal{C}^*} F \Leftrightarrow F - E \in \mathcal{C}^*$ on the dual space.

Associated with each GPT is a fixed *unit effect* $U \succ_{\mathcal{C}^*} 0$ defined as the unique element of the dual space satisfying $\langle U, \omega \rangle = 1$ for all $\omega \in \Omega(\mathcal{V})$; equivalently, this allows one to understand the set of valid states as the set of normalized elements of the cone \mathcal{C} in the sense that $\Omega(\mathcal{V}) = \{\omega \in \mathcal{C} | \langle U, \omega \rangle = 1\}$. Under the so-called no-restriction hypothesis, which we hereafter take as an assumption, the effects are then all functionals E such that $0 \leq_{\mathcal{C}^*} E \leq_{\mathcal{C}^*} U$, which are precisely all linear functionals $E:\Omega(\mathcal{V}) \to [0,1]$. Any finite collection of effects $\{M_i\}_i$ such that $\sum_i M_i = U$ are called a *measurement*, with $\langle M_i, \cdot \rangle$ identifying the probability of measuring the i th outcome. We denote by \mathcal{M} the set of all possible measurements.

Given two GPTs defined by the spaces \mathcal{V} and \mathcal{V}' with the corresponding sets of states $\Omega(\mathcal{V})$ and $\Omega(\mathcal{V}')$, one can then study the transformations between them. The question of physically allowed transformations $\Lambda: \mathcal{V} \to \mathcal{V}'$ between different states, which we refer to as *channels*, is in general heavily dependent on the physical setting of the given theory and additional assumptions placed on it [95,96]. However, since we require the output to be a valid state, two general assumptions can be made: (1) Any valid channel Λ is state-cone preserving, that is, $\Lambda[\mathcal{C}] \subseteq \mathcal{C}'$ where \mathcal{C}' is the cone defined by $\Omega(\mathcal{V}')$ in the output space, and (2) Λ is normalization preserving, that is, $\langle U, x \rangle = \langle U', \Lambda(x) \rangle$ for any $x \in \mathcal{C}$ where U' is the unit effect in the output dual

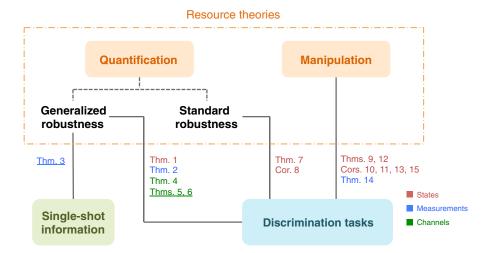


FIG. 1. Schematics showing the connections established by the results in this work. The colors of the theorem labels indicate the types of resources relevant to the results (red, states; blue, measurements; green, channels). The underlined theorems (Theorems 3, 5, and 6) are shown specifically for quantum theory, while the other results are valid in general GPTs.

space. Any valid set of physical transformations in the given GPT, which we denote by $\mathcal{T}(\mathcal{V}, \mathcal{V}')$, will necessarily be a subset of all state-cone- and normalization-preserving operations, although often the inclusion will be strict. To allow for full generality of our results, we place no further restrictions on the set of physically allowed transformations at this point, except for the trivial assumption that the identity map $\mathrm{id}: x \mapsto x$ is physically implementable.

For any $\Lambda: \mathcal{V} \to \mathcal{V}'$, define the dual map $\Lambda^*: \mathcal{V}'^* \to \mathcal{V}^*$ by

$$\langle E, \Lambda(x) \rangle = \langle \Lambda^*(E), x \rangle \quad \forall E \in \mathcal{V}'^*, \qquad x \in \mathcal{V}.$$
 (1)

Without loss of generality, we identify the bidual \mathcal{V}^{**} with \mathcal{V} and say that $\Lambda^{**} = \Lambda$.

Consider now transformations between measurements; that is, maps $\Gamma: \mathcal{V}'^* \to \mathcal{V}^*$ such that for any measurement $\{M_i\}_i, \{\Gamma(M_i)\}_i$ is also a measurement. It is not difficult to see that two conditions need to be satisfied for such a map to always result in valid measurements: One, it needs to preserve the effect cone in the sense that $\Gamma[\mathcal{C}^{\prime*}] \subseteq \mathcal{C}^*$, and two, it needs to preserve the unit effect, $\Gamma(U') = U$, so that $\sum_{i} \Gamma(M_i) = U$. We refer to maps obeying the two conditions as effect cone preserving and unital, respectively. Using the duality relation (1), it is then not difficult to see that the set of all maps dual to the set of state-cone- and normalization-preserving maps are precisely the effect of cone-preserving unital maps. More specifically, a map Λ is state-cone preserving if and only if its dual preserves the effect cone and normalization preserving if and only if its dual is unital.

The concepts of quantum theory can be intuitively understood in this setting. To help translate the notation of the quantum-mechanical setting to more general GPTs, we include a basic comparison in Table I.

B. Discrimination tasks

Consider a finite ensemble of the form $\{p_i, \sigma_i\}_i$, where $\sigma_i \in \Omega(\mathcal{V})$ and p_i are the probabilities corresponding to each state, such that $\sum_i p_i = 1$. The task of *state discrimination* is concerned with the scenario where a state is sampled from the ensemble and one aims to determine which one of the states is in one's possession by measuring it. Specifically, given a measurement $\{M_i\}_i$, we associate the i th measurement outcome with the guess that the sampled state is σ_i . The average probability of successfully obtaining the correct guess is then given by

$$p_{\text{succ}}(\{p_i, \sigma_i\}, \{M_i\}) = \sum_i p_i \langle M_i, \sigma_i \rangle.$$
 (2)

In the context of minimum-error discrimination, one is in particular interested in choosing an optimal measurement which maximizes this quantity. It is a fundamental fact in any GPT that, generalizing the approach in the seminal Holevo-Helstrom theorem [66,67], in the case of

TABLE I. A comparison between the concepts and standard nomenclature used in finite-dimensional quantum mechanics and in more general GPTs.

GPTs	Quantum mechanics
Real vector space \mathcal{V}	Self-adjoint operators acting on a Hilbert space \mathcal{H}
States $\Omega(\mathcal{V})$	Density operators
State cone C	Positive semidefinite operators
Effect cone C^*	Positive semidefinite operators
Unit effect U	Identity operator I
Canonical bilinear	Hilbert-Schmidt inner product
form $\langle E, x \rangle$	Tr(Ex)
Measurement	Positive operator-valued measure (POVM)
Effect	POVM element
State-cone-preserving maps	Positive maps
Effect-cone-preserving maps	Positive maps
Normalization-preserving maps	Trace-preserving maps
Physical transformations	Completely positive
$\mathcal{T}(\mathcal{V},\mathcal{V}')$	trace-preserving maps
Unital maps	Unital maps
Base norm $\ \cdot\ _{\Omega}$	Trace norm $\ \cdot\ _1$
Order unit norm $\ \cdot\ _{\Omega}^{\circ}$	Operator norm $\ \cdot\ _{\infty}$

discriminating between two states the problem can be expressed as [75]

$$\max_{\{M_i\} \in \mathcal{M}} p_{\text{succ}}(\{p_i, \sigma_i\}_{i=0}^1, \{M_i\}_{i=0}^1)
= \frac{1}{2} (\|p_0 \sigma_0 - p_1 \sigma_1\|_{\Omega} + 1),$$
(3)

where $\|\cdot\|_{\Omega}$ is the so-called *base norm* given by

$$||x||_{\Omega} := \min\{\lambda_{+} + \lambda_{-} | x = \lambda_{+} \omega_{+} - \lambda_{-} \omega_{-},$$

$$\lambda_{\pm} \in \mathbb{R}_{+}, \omega_{\pm} \in \Omega(\mathcal{V})\}$$

$$= \max\{\langle E, x \rangle | -U \leq_{\mathcal{C}^{*}} E \leq_{\mathcal{C}^{*}} U\},$$

$$(4)$$

where the equality follows by convex duality [97]. Notice that there exists a measurement which distinguishes an ensemble of two states perfectly if and only if $\|p_0\sigma_0-p_1\sigma_1\|_{\Omega}=1$. In the case that $p_0=p_1=\frac{1}{2}$, we refer to this task as balanced binary discrimination.

We will furthermore make frequent use of the norm dual to the base norm $\|\cdot\|_{\Omega}$ called the order unit norm $\|\cdot\|_{\Omega}^{\circ}$, which can be obtained as

$$||Y||_{\Omega}^{\circ} = \max \{ \langle Y, x \rangle | ||x||_{\Omega} \le 1 \}$$

= \text{max} \{ |\langle Y, \omega \rangle ||\omega \in \Omega \rangle \Omega \rangle \rangl

Notice in particular the set of effects is precisely $\{Y \in \mathcal{C}^* | \|Y\|_{\Omega}^{\circ} \leq 1\}.$

Another setting often encountered in the task of state discrimination is where one is constrained to use only a restricted set of allowed measurements $\{M_i\} \in \mathcal{M}_{\mathcal{F}} \subseteq \mathcal{M}$. This scenario has been found to be of fundamental importance due to the phenomenon of data hiding, that is, the existence of states which can be distinguished perfectly with general measurements but not with local measurements supplemented with classical communication [98,99]. Provided that the set $\mathcal{M}_{\mathcal{F}}$ is informationally complete, that is, the effects contained in $\mathcal{M}_{\mathcal{F}}$ together span the whole space \mathcal{V}^* , the best success probability in this setting can be expressed as [77,100]

$$\sup_{\{M_i\} \in \mathcal{M}_{\mathcal{F}}} p_{\text{succ}}(\{p_i, \sigma_i\}_{i=0}^1, \{M_i\}_{i=0}^1)
= \frac{1}{2} (\|p_0 \sigma_0 - p_1 \sigma_1\|_{\mathcal{M}_{\mathcal{F}}} + 1)$$
(6)

with $\|\cdot\|_{\mathcal{M}_{\mathcal{F}}}$ being the so-called distinguishability norm $\|x\|_{\mathcal{M}_{\mathcal{F}}} \coloneqq \sup_{\{M_i\} \in \mathcal{M}_{\mathcal{F}}} \sum_i |\langle M_i, x \rangle|.$

A related task is concerned with *channel discrimination*, where one of the channels from a given ensemble $\{p_i, \Lambda_i\}$ with each $\Lambda_i \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ occurring with probability p_i is applied to a known state $\omega \in \Omega(\mathcal{V})$. The task then is, by measuring the output state $\Lambda_i(\omega)$, to decide which of the channels was applied. The average probability of guessing correctly with a measurement $\{M_i\}$ is then

$$p_{\text{succ}}(\{p_i, \Lambda_i\}, \{M_i\}, \omega) = \sum_i p_i \langle M_i, \Lambda_i(\omega) \rangle, \quad (7)$$

which is completely equivalent to discriminating the state ensemble $\{p_i, \Lambda_i(\omega)\}$. A more general setting is that of *subchannel discrimination*, in which the object to discriminate is an ensemble $\{\Psi_i\}$ of subchannels, that is, state-cone-preserving maps which are normalization non-increasing in the sense that $\langle U, \omega \rangle \geq \langle U', \Psi_i(\omega) \rangle$, and their sum $\sum_i \Psi_i$ is normalization preserving. We can then write

$$p_{\text{succ}}(\{\Psi_i\}, \{M_i\}, \omega) = \sum_i \langle M_i, \Psi_i(\omega) \rangle.$$
 (8)

Finally, in some cases we allow for the *inconclusive* discrimination; that is, a discrimination task in which an ensemble $\{p_i, \sigma_i\}_{i=0}^{N-1}$ is discriminated with a measurement $\{M_i\}_{i=0}^{N}$, and we associate with the Nth measurement outcome an inconclusive result. Similarly, the average probability of guessing correctly is then

$$p'_{\text{succ}}(\{p_i, \sigma_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^N) = \sum_{i=0}^{N-1} p_i \langle M_i, \sigma_i \rangle, \quad (9)$$

and analogously for the case of channel and subchannel discrimination. Inconclusive discrimination is a broader class of tasks than conclusive discrimination, in the sense that the latter can be considered as a special case of the former where one takes $M_N = 0$.

C. Resources and their quantification

A general resource theory can be identified with a set of objects (here, a set of states, measurements, or channels) together with a set of transformations of said objects that one deems free, in the sense that they are available in the given physical setting at no resource cost. Both definitions will generally depend on the physical setting under consideration; intuitively, an object being free can be understood as it possessing no resource, and a transformation being free means that it is allowed within the given physical constraints. In particular, any such free operation should not generate any resource; that is, any object subjected to a free transformation should have quantitatively "less" of the given resource than it possessed initially. This is formalized precisely by resource monotones, which are functions from the set of states, measurements, or channels to real numbers whose aim is to quantify the resource content of the given object and therefore do not increase under the action of the free transformations. Frequently, further constraints are imposed on functions admissible as valid resource monotones [3,34,101], although we make no additional assumptions at this point. In most cases, the choice of a resource monotone is not unique, and typically, one therefore looks for a choice of monotones which characterize the physical properties of the given resource in addition to merely outputting a number associated with an object.

Having defined the setting of a resource theory in this way, one is then interested in understanding the limitations that it places on one's operational capabilities within the general probabilistic theory. Some of the fundamental questions that can be asked in this context are (i) what physical advantages a resourceful object can provide over a free one, (ii) which transformations are possible with the restricted set of free operations and how to characterize them, (iii) and what exactly can measuring the resourcefulness of an object tell us quantitatively about the usefulness of the object in physical tasks.

The concept which forms a central pillar of this work is convexity. Although ultimately a technical assumption, it stems from deep physical considerations and is a foundation of any GPT [102]. In particular, take the set of states $\Omega(\mathcal{V})$: if one is free to prepare any $\omega_1, \omega_2 \in \Omega(\mathcal{V})$ within the given physical setting of the theory, one should also be allowed to simply forget which one was prepared—such a randomization leads precisely to convex mixtures $p\omega_1 + (1-p)\omega_2$. Dually, this means that allowed measurements should form a convex set, and analogously, the randomization of transformations should be a valid physical procedure. In a similar way, this property can be explicitly required from a given resource theory; if a randomization procedure of free objects is free to be performed, the associated set of objects should be free.

Although it is certainly possible to define resource theories in which convexity does not hold [19,103,104], it is a natural assumption in most physical settings, and the vast majority of established resource theories are indeed convex.

In any such convex resource theory, a very intuitive way to define a quantifier is as follows: given an object A, what is the least amount of mixing $p \in [0,1]$ with another object B such that (1-p)A+pB is free? This is precisely the idea behind robustness measures [84], where in particular the standard (free) robustness asks about the least coefficient p such that B is also a free object, and the generalized (global) robustness asks about the least p when B is any admissible object. In an intuitive sense, this can be understood as the robustness of the resource contained in A to noise in the form of admixing the objects B.

To investigate the properties of such monotones, we now specify to the particular settings of convex resource theories of states, measurements, and channels in GPTs.

III. GENERALIZED ROBUSTNESS OF STATES

The generalized robustness has found a multitude of uses in quantum resource theories. Its operational applications include the tasks of one-shot entanglement dilution [88,89], one-shot coherence distillation [85,86] and dilution [90], phase discrimination with coherent states [81,82], and catalytic resource erasure [33,91].

Notably, it was shown in Ref. [36] that the generalized robustness of states defined in any quantum resource theory serves as an exact quantifier for the maximum advantage that a resource state provides in a class of (sub)channel discrimination tasks. However, the construction considered in the proof relies on the specific structure of quantum theory, and it does not immediately generalize to all GPTs. Here, we introduce a generalization of that result which holds without any assumptions about the underlying GPT, showing that this universal relation applies even beyond the setting of quantum mechanics.

Given a convex and closed set of free states $\mathcal{F} \subseteq \Omega(\mathcal{V})$, the generalized robustness is given for any state $\omega \in \Omega(\mathcal{V})$ as the optimal value of the convex optimization problem

$$R_{\mathcal{F}}(\omega) := \min \left\{ r \in \mathbb{R}_{+} \middle| \frac{\omega + r\delta}{1 + r} \in \mathcal{F}, \delta \in \Omega(\mathcal{V}) \right\}$$
$$= \min \left\{ r \in \mathbb{R}_{+} \middle| \omega \leq_{\mathcal{C}} (1 + r)\sigma, \sigma \in \mathcal{F} \right\}. \tag{10}$$

To ensure that this quantity is well defined for any state, we assume that \mathcal{F} contains at least one interior point of \mathcal{C} . It is straightforward (see, e.g., the Appendix B) to obtain the following equivalent dual problem:

maximize
$$\langle X, \omega \rangle - 1$$

subject to $X \succeq_{\mathcal{C}^*} 0$
 $\langle X, \sigma \rangle \leq 1, \quad \forall \ \sigma \in \mathcal{F}.$ (11)

One can check that the dual problem is strictly feasible by taking the feasible solution X = U/2, and thus strong duality is ensured by Slater's theorem [105], meaning that the solutions to the primal and dual problems coincide.

Recalling that the average probability of success in a channel discrimination task is given by

$$p_{\text{succ}}(\{p_i, \Lambda_i\}, \{M_i\}, \omega) = \sum_i p_i \langle M_i, \Lambda_i(\omega) \rangle, \quad (12)$$

we can now show that the generalized robustness of a state quantifies its maximal advantage over the free states in all such channel discrimination tasks with a fixed choice of measurement.

Theorem 1: For any $\omega \in \Omega(\mathcal{V})$ it holds that

$$\max_{\{M_i\},\{p_i,\Lambda_i\}} \frac{p_{\text{succ}}(\{p_i,\Lambda_i\},\{M_i\},\omega)}{\max_{\sigma \in \mathcal{F}} p_{\text{succ}}(\{p_i,\Lambda_i\},\{M_i\},\sigma)} = 1 + R_{\mathcal{F}}(\omega)$$
(13)

where the maximization is over all finite ensembles of channels $\{p_i, \Lambda_i\}$ with each $\Lambda_i \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ and all measurements $\{M_i\}$ on the output space \mathcal{V}' .

Proof.—To show that the left-hand side is upper bounded by the right-hand side, note that the definition of the robustness implies that, for any state ω , there exists another state δ such that $[(\omega + r\delta)/(1+r)] = \sigma \in \mathcal{F}$ where $r = R_{\mathcal{F}}(\omega)$. This gives for any $\{M_i\}$ and any channel ensemble $\{p_i, \Lambda_i\}$ that

$$\begin{split} p_{\text{succ}}(\{p_{i}, \Lambda_{i}\}, \{M_{i}\}, \omega) \\ &= \sum_{i} p_{i} \langle M_{i}, \Lambda_{i} ((1+r)\sigma - r\delta) \rangle \\ &\leq \sum_{i} p_{i} \langle M_{i}, (1+r)\Lambda_{i}(\sigma) \rangle \\ &\leq (1+r) \underset{\sigma \in \mathcal{F}}{\text{max}} p_{\text{succ}}(\{p_{i}, \Lambda_{i}\}, \{M_{i}\}, \sigma), \end{split} \tag{14}$$

where the first inequality follows since each $M_i \in C^*$ and $\delta \in C$, which concludes the first part of the proof.

To see the opposite inequality, consider an optimal solution X in Eq. (11) for $R_{\mathcal{F}}(\omega)$ and define a measurement by $\{X/\|X\|_{\Omega}^{\circ}, U-X/\|X\|_{\Omega}^{\circ}\}$. This is a valid measurement since each $X \in \mathcal{C}^*$ by the dual form of the robustness in Eq. (11), and so $0 \leq_{\mathcal{C}^*} X/\|X\|_{\Omega}^{\circ} \leq_{\mathcal{C}^*} U$ by definition of the norm $\|\cdot\|_{\Omega}^{\circ}$. Consider now the channel ensemble defined by $p_0=1$, $\Lambda_0=\mathrm{id}$, and $p_1=0$, $\Lambda_1=\Lambda'$ where id denotes the identity map $x\mapsto x$, and Λ' is an arbitrary channel. This gives

$$\max_{\{M_{i}\},\{p_{i},\Lambda_{i}\}} \frac{p_{\text{succ}}(\{p_{i},\Lambda_{i}\},\{M_{i}\},\omega)}{\max_{\sigma\in\mathcal{F}}p_{\text{succ}}(\{p_{i},\Lambda_{i}\},\{M_{i}\},\sigma)}$$

$$\geq \frac{\frac{1}{\|X\|_{\Omega}^{\circ}}\langle X,\omega\rangle}{\frac{1}{\|X\|_{\Omega}^{\circ}}\max_{\sigma\in\mathcal{F}}\langle X,\sigma\rangle}$$

$$\geq 1 + R_{\mathcal{F}}(\omega), \tag{15}$$

where in the last inequality, we use that $\langle X, \sigma \rangle \leq 1$ for any free state σ by the condition in Eq. (11).

We stress that, although the above result optimizes over the set of all ensembles of *physical* transformations $\mathcal{T}(\mathcal{V}, \mathcal{V}')$, the only assumption about the set $\mathcal{T}(\mathcal{V}, \mathcal{V}')$ that we make is that it contains the identity map. This makes the above theorem immediately applicable to every GPT, regardless of how the given physical limitations constrain the implementable set of operations.

We remark that instead of channel discrimination, one can alternatively consider subchannel discrimination—which is generally a broader class of discrimination tasks—and show the same relation, whose proof proceeds analogously. [Equation (14) still holds when one replaces the channel ensemble with a set of subchannels, and one can choose the same channel ensemble to show Eq. (15).]

Importantly, from the dual form in Eq. (11) one can see that the robustness is directly observable, in the sense that it can be obtained by measuring a single effect X (expectation value of observable *X* for the case of quantum mechanics) at the state ω . This ensures that the quantification of the robustness is accessible, allowing in particular to straightforwardly bound the value of $R_{\mathcal{F}}$ based on measurement data obtained in experiment, adapting the approach of socalled quantitative resource witnesses [93,94]. One should also note that the generalized robustness can be computed exactly for certain classes of states in quantum resource theories such as entanglement and Schmidt number k entanglement [34,106,107], coherence and multilevel coherence [34,82,108,109], magic [34,110], and in several cases can be cast as a semidefinite program for any state; this makes the computation of the operationally motivated quantity in Theorem 1 feasible in practice for many relevant cases of resource theories.

IV. GENERALIZED ROBUSTNESS OF MEASUREMENTS

Understanding the discriminative power of restricted sets of measurements is of central importance not only in characterizing the operational consequences precipitated by limitations of physically allowed measurements but often also in studying the very fundamental structure of the underlying GPT [77,100,111]. The phenomenon of data hiding [98,99] has in particular motivated the study of the question: given a measurement, how well can one distinguish physical states with it as compared to some fixed restricted set of measurements? We show that a robustness measure associated with the measurement can provide a precise answer to the question.

First, we formally define the generalized robustness of measurements with respect to some convex and closed set of measurements $\mathcal{M}_{\mathcal{F}}$, which we define as

$$\mathcal{M}_{\mathcal{F}} \coloneqq \{ \{ M_i \}_i \in \mathcal{M} | M_i \in \mathcal{E}_{\mathcal{F}} \ \forall \ i \}, \tag{16}$$

where $\mathcal{E}_{\mathcal{F}} \subseteq \mathcal{C}^*$ is some chosen convex and closed cone of free effects which we are able to access within the constraints of the given resource theory. As examples of such a setting, one can consider local measurements, separable measurements [77], measurements simulable by a given set of measurements [22,24,112], or trivial measurements (proportional to the unit effect); within quantum mechanics, one can furthermore choose, for instance, positive partial transpose (PPT) measurements, incoherent measurements, or (probabilistic mixtures of) Pauli measurements.

We define the generalized robustness of measurement with respect to $\mathcal{E}_{\mathcal{T}}$ for a given measurement $\mathbb{M} = \{M_i\}_i$ as

$$R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}) := \min \{ r \in \mathbb{R}_{+} | M_{i} + rN_{i} \in \mathcal{E}_{\mathcal{F}} \ \forall \ i, \{N_{i}\}_{i} \in \mathcal{M} \}.$$

$$(17)$$

We assume that $\mathcal{M}_{\mathcal{F}}$ contains at least one measurement consisting of effects $M_i \succ_{\mathcal{C}^*} 0$ which are in the interior of \mathcal{C}^* , so that the above quantity is well defined for any measurement. The faithfulness $R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}) = 0$ iff $M_i \in \mathcal{E}_{\mathcal{F}} \ \forall i$, the convexity $R_{\mathcal{E}_{\mathcal{F}}}(p\mathbb{M}+(1-p)\mathbb{M}') \leq pR_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})+(1-p)$ $R_{\mathcal{E}_{\pi}}(\mathbb{M}')$, and the monotonicity of the robustness $R_{\mathcal{E}_{\tau}}(\Gamma(\mathbb{M})) \leq R_{\mathcal{E}_{\tau}}(\mathbb{M})$ for any effect-cone-preserving map Γ s.t. $\Gamma[\mathcal{E}_{\mathcal{F}}] \subseteq \mathcal{E}_{\mathcal{F}}$ follow easily from the definition. It is also straightforward to show the monotonicity under classical postprocessing $R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}') \leq R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})$ holds where $M_b' = \sum_a p(b|a)M_a$ and p(b|a) is any conditional probability distribution. To see this, note that the definition together with the conic structure of $\mathcal{E}_{\mathcal{F}}$ implies that for any i, there exists a free effect $F_i \in \mathcal{E}_{\mathcal{F}}$ such that $[(M_i + rN_i)/(1+r)] = F_i$. Thus, one can write for any b and p(b|a),

$$M_b' = \sum_{a} p(b|a)[(1+r)F_a - rN_a]$$
 (18)

$$= (1+r)F_b' - rN_b', (19)$$

where $F_a \in \mathcal{E}_{\mathcal{F}}$, $N_a \in C^*$ are some effects. Note that $F_b' \coloneqq \sum_a p(b|a) F_a$, $N_b' \coloneqq \sum_a p(b|a) N_a$ are also effects constituting valid measurements, and $F_b' \in \mathcal{E}_{\mathcal{F}}$ due to the convexity of $\mathcal{E}_{\mathcal{F}}$. Since Eq. (19) is one valid decomposition of effects constituting \mathbb{M} , we get $R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}') \leq R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})$ due to the minimization involved in the definition (17).

To observe a close connection between the robustness and discrimination tasks, let us define $N'_i := rN_i$, and rewrite the definition (17) as the following convex optimization problem:

minimize
$$\lambda$$
, (20)

subject to
$$M_i + N_i' \succeq_{\mathcal{E}_{\tau}} 0$$
, (21)

$$N_i' \in \mathcal{C}^* \quad \forall i,$$
 (22)

$$\sum_{i} N_i' = \lambda U. \tag{23}$$

An equivalent dual problem (see, e.g., Appendix B) is written as

maximize
$$-\sum_{i} \langle M_i, \sigma_i \rangle$$
, (24)

subject to
$$\sigma_i \leq_{\mathcal{C}} \eta \quad \forall i,$$
 (25)

$$\eta \in \mathcal{V},$$
(26)

$$\sigma_i \in \mathcal{E}_{\mathcal{T}}^* \quad \forall \quad i,$$
 (27)

$$\langle U, \eta \rangle = 1, \tag{28}$$

and defining $\omega_i = -\sigma_i + \eta$, we can rewrite it as

maximize
$$\sum_{i} \langle M_i, \omega_i \rangle - 1,$$
 (29)

subject to
$$\omega_i \in \mathcal{C} \quad \forall i$$
, (30)

$$\eta \in \mathcal{V},$$
(31)

$$\langle F, \omega_i \rangle \le \langle F, \eta \rangle \quad \forall i, \quad \forall F \in \mathcal{E}_{\mathcal{F}},$$
 (32)

$$\langle U, \eta \rangle = 1, \tag{33}$$

where in Eq. (29) we use that $\sum_i \langle M_i, \eta \rangle = \langle U, \eta \rangle = 1$ by Eq. (33), and in Eq. (32) we write the condition $\eta - \omega_i \in \mathcal{E}_{\mathcal{T}}^*$ explicitly. To see that strong duality holds and thus the optimal value of the dual is equal to the optimal value of the primal problem, choose any $\sigma \succ_{\mathcal{C}} 0$ (and therefore, $\sigma \succ_{\mathcal{E}_{\mathcal{F}}^*} 0$) s.t. $0 < \langle U, \sigma \rangle < 1$, which is guaranteed to exist since the interior of \mathcal{C} is nonempty by assumption, and define $\eta = \sigma/\langle U, \sigma \rangle$. This choice of $\sigma_i = \sigma$ and η can be noticed to strictly satisfy the conditions (25)–(27), and so Slater's theorem ensures that strong duality holds.

This form of the dual problem allows one to identify the generalized robustness of measurement as an exact quantifier for the advantage in some state discrimination task. Let $\mathcal{A} = \{p_i, \sigma_i\}$ denote a state ensemble to be discriminated. Then, we obtain the following result.

Theorem 2: Let $\mathcal{M}_{\mathcal{F}}$ be the set of measurements whose effects are elements of $\mathcal{E}_{\mathcal{F}}$. Then, for any measurement $\mathbb{M} = \{M_i\}_i$, it holds that

$$\max_{\mathcal{A}} \frac{p_{\text{succ}}(\mathcal{A}, \mathbb{M})}{\max_{\mathbb{F} \in \mathcal{M}_{\mathcal{F}}} p_{\text{succ}}(\mathcal{A}, \mathbb{F})} = 1 + R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}), \quad (34)$$

where the maximization is over all finite ensembles of states $A = \{p_i, \sigma_i\}$.

Proof.—We first show that the left-hand side is smaller than or equal to the right-hand side. Following the

definition of the generalized robustness, one can write $M_i = (1+r)F_i - rN_i$ for some $F_i \in \mathcal{E}_{\mathcal{F}}$, $N_i \in \mathcal{C}^*$ for all i where $r = R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})$. Thus, we get

$$p_{\text{succ}}(\mathcal{A}, \mathbb{M}) = \sum_{i} p_{i} \langle M_{i}, \sigma_{i} \rangle$$

$$= (1+r) \sum_{i} p_{i} \langle F_{i}, \sigma_{i} \rangle - r \sum_{i} p_{i} \langle N_{i}, \sigma_{i} \rangle$$

$$\leq (1+r) \max_{\mathbb{F} \in \mathcal{M}_{r}} p_{\text{succ}}(\mathcal{A}, \mathbb{F}), \tag{35}$$

where the inequality follows since $\sigma_i \in \mathcal{C}$ and $N_i \in \mathcal{C}^*$ for all i, which completes the first part of the proof. To show the converse, consider the set of optimal $\{\omega_i\} \subseteq \mathcal{C}$ that appear in Eq. (29). We choose the probability distribution and states as

$$p_{i} = \frac{\langle U, \omega_{i} \rangle}{\sum_{a} \langle U, \omega_{a} \rangle}, \sigma_{i} = \omega_{i} / \langle U, \omega_{i} \rangle \quad \text{when } \langle U, \omega_{i} \rangle > 0,$$

$$p_{i} = 0 \quad \text{when } \langle U, \omega_{i} \rangle = 0.$$
 (36)

Considering an ensemble defined by the above probability and states $\mathcal{A} = \{p_i, \sigma_i\}$, we obtain for any $\mathbb{F} \in \mathcal{M}_{\mathcal{F}}$,

$$\frac{p_{\text{succ}}(\mathcal{A}, \mathbb{M})}{p_{\text{succ}}(\mathcal{A}, \mathbb{F})} = \frac{\frac{1}{\sum_{a} \langle U, \omega_{a} \rangle} \sum_{i} \langle M_{i}, \omega_{i} \rangle}{\frac{1}{\sum_{a} \langle U, \omega_{a} \rangle} \sum_{i} \langle F_{i}, \omega_{i} \rangle}$$
(37)

$$\geq \frac{1 + R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})}{\sum_{i} \langle F_{i}, \eta \rangle} \tag{38}$$

$$= \frac{1 + R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})}{\langle U, n \rangle} = 1 + R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}), \tag{39}$$

where the inequality is due to Eq. (32), and the second last equality is due to Eq. (33), which concludes the proof.

The above theorem establishes an explicit connection in any GPT between the inherent resourcefulness of a given measurement and the advantage realized in state discrimination tasks with respect to a general set of free measurements, which ensures an operational interpretation of the generalized robustness of measurements, extending the previously considered case of states. Furthermore, it allows for a connection with the data-hiding phenomenon. In the investigation of data hiding in general probabilistic theories, it is common to encounter the so-called data-hiding ratio [77,100], which in our notation can be expressed as

$$\max_{\mathcal{A}'} \frac{\max_{\mathbb{M} \in \mathcal{M}} p_{\text{succ}}(\mathcal{A}', \mathbb{M}) - \frac{1}{2}}{\max_{\mathbb{F} \in \mathcal{M}_{\mathcal{F}}} p_{\text{succ}}(\mathcal{A}', \mathbb{F}) - \frac{1}{2}},$$
(40)

where the ensembles \mathcal{A}' are limited to contain at most two different states. Theorem 2 then shows that maximizing the robustness $R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})$ over all measurements \mathbb{M} provides an alternative ratio of this form, generalized to include state ensembles of arbitrary length.

A. Connections with single-shot information theory

Here, we extend the connection between generalized robustness of measurements and state discrimination tasks to another seemingly different information-theoretic quantity, one-shot accessible information, which generalizes the specific case considered in Ref. [25]. Since entropic quantities are particularly relevant to quantum information theory, in this section we restrict our attention to quantum mechanics (and not general GPTs).

Consider the situation where Alice encodes her classical information into an ensemble $A = \{p_x, \sigma_x\}$, and Bob tries to decode it by making a POVM measurement on a state sampled from A. When this process is independently repeated for asymptotically many times, the amount of information he may learn is quantified by the accessible information $I^{\text{acc}}(A) := \max_{\{N_v\}} I(X:Y)$ where I(X:Y) =H(X) - H(X|Y) is the classical mutual information, X is the random variable associated with the ensemble, and $\{N_{\nu}\}\$ refers to a POVM measurement whose associated random variable is Y [113]. However, quantities based on the von Neumann or Shannon entropy cease to be suitable for more practical nonasymptotic cases, and several kinds of alternative quantities playing major roles in *single-shot* scenarios have been proposed and studied [114–116]. As in Ref. [25], we consider a variant of the single-shot version of accessible information, min-accessible information, for a state ensemble \mathcal{A} defined by

$$I_{\min}^{\text{acc}}(\mathcal{A}) := \max_{\{N_y\} \in \mathcal{M}} I_{\min}(X:Y)$$

$$= -\log \max_{x} p_x + \max_{\{N_y\} \in \mathcal{M}} \log \sum_{y} \max_{x} p_x \text{Tr}[\sigma_x N_y],$$
(41)

where $I_{\min}(X:Y) := H_{\min}(X) - H_{\min}(X|Y)$ is a single-shot variant of mutual information [116], and $H_{\min}(X) =$ $-\log \max_{x} p_{x}$, $H_{\min}(X|Y) = -\log \sum_{y} \max_{x} p(x, y)$ are the min-entropy and min-conditional entropy [114,115,117]. We shall see that the accessible information for state ensembles enables the information-theoretic characterization of the generalized robustness of measurements. To see the relation, it is insightful to see the measurements as channels; in particular, consider the measure-and-prepare channel associated with the measurement M with POVMs $\{M_i\}$ defined by $\Lambda_{\mathbb{M}}(\cdot) := \sum_{j} \operatorname{Tr}[\cdot M_{j}]|j\rangle\langle j|$. Denoting the ensemble one would possess by applying the channel Λ to the ensemble $\mathcal{A} = \{p_x, \sigma_x\}$ by $\mathcal{A}_{\Lambda_{M}} := \{p_x, \Lambda(\sigma_x)\}$, we obtain the following relation between the maximal increase in minaccessible information due to the given resource measurement and the generalized robustness of that measurement.

Theorem 3: For any measurement $\mathbb{M} \in \mathcal{M}$, it holds that

$$\max_{\mathcal{A}} [I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}}}) - \max_{\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}} I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}'}})] = \log [1 + R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M})],$$
(42)

where $\mathcal{M}_{\mathcal{F}}$ is the set of free measurements consisting of POVM elements in $\mathcal{E}_{\mathcal{F}}$.

Proof.—Consider the ensemble $A = \{p_x, \sigma_x\}$. Using the definition (41), we get

$$I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\text{M}}}) = -\log \max_{x} p_{x} + \max_{x} \log \sum_{y} \max_{x} \left[p_{x} \sum_{a} \text{Tr}[\sigma_{x} M_{a}] \langle a | N_{y} | a \rangle \right]$$
(43)

$$= -\log \max_{x} p_{x} + \log \sum_{y} \max_{x} \left[p_{x} \sum_{a} \text{Tr}[\sigma_{x} M_{a}] \delta_{ay} \right]$$
(44)

$$= -\log \max_{x} p_{x} + \log \max_{\{q(x|y)\}} \sum_{y} \sum_{x} q(x|y) p_{x} \operatorname{Tr}[\sigma_{x} M_{y}]. \tag{45}$$

The second equality is because $N_y = |y\rangle\langle y|$ can be chosen as optimal POVM elements, and the third equality is because $\max_x f(x) = \max_{\{q(x)\}} \sum_x q(x) f(x)$ for any function f(x) and probability distribution q(x). Note that we can interchange the summation over y and the maximization over probability distributions $\{q(x|y)\}$ because the summation over y is maximized when q(x|y) maximizes the summant for each y. Then, we get for any ensemble \mathcal{A} ,

$$I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}}}) - \max_{\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}} I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}'}}) = \log \frac{\max_{\{q(x|y)\}} \sum_{y} \sum_{x} q(x|y) p_{x} \text{Tr}[\sigma_{x} M_{y}]}{\max_{\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}} \max_{\{q(x|y)\}} \sum_{y} \sum_{x} q(x|y) p_{x} \text{Tr}[\sigma_{x} M'_{y}]}$$
(46)

$$\leq \log\left[1 + R_{\mathcal{E}_{\tau}}(\mathbb{M})\right],\tag{47}$$

where the inequality can be proved in a way similar to Eq. (35). To see the converse, note that for any $\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}$ and conditional probability distribution q(x|y), another measurement \mathbb{M}'' defined by the POVM $M_x'' = \sum_y q(x|y)M_y'$ is also a member of $\mathcal{M}_{\mathcal{F}}$ due to the convexity of $\mathcal{E}_{\mathcal{F}}$ that defines $\mathcal{M}_{\mathcal{F}}$. Thus, Eq. (46) is equivalently written as

$$I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}}}) - \max_{\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}} I_{\min}^{\text{acc}}(\mathcal{A}_{\Lambda_{\mathbb{M}'}})$$

$$= \log \frac{\max_{\{q(x|y)\}} \sum_{y} \sum_{x} q(x|y) p_{x} \text{Tr}[\sigma_{x} M_{y}]}{\max_{\mathbb{M}' \in \mathcal{M}_{\mathcal{F}}} \sum_{x} p_{x} \text{Tr}[\sigma_{x} M'_{x}]}. \quad (48)$$

The equality in Eq. (47) can be achieved by setting $q(x|y) = \delta_{xy}$ and taking the ensemble defined by Eq. (36), which concludes the proof.

Theorem 3 generalizes the result in Ref. [25] for the case of theory of informativeness of measurements to general quantum resource theories of measurements. As a result, Theorem 3, together with Theorem 2, establishes the general connections between discrimination tasks, robustness measures, and single-shot information theory conjectured in Ref. [25] for the case of quantum measurements.

V. ROBUSTNESS MEASURES FOR CHANNELS

In addition to states and measurements, the state transformations themselves can also be regarded as resourceful, and their operational characterization in this way has recently become an active area of investigation [26,37,39–43,45].

Here, in particular, we discuss two general ways to approach the resource content of channels. One is to relate them to an underlying resource theory for states: if a channel $\Lambda \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ is a free operation in a resource theory characterized by the set of free states \mathcal{F} , then we know that $\Lambda[\mathcal{F}] \subseteq \mathcal{F}'$; the resourcefulness of a nonfree operation can then be measured exactly by understanding how much resource it can create, which is known as the resource-generating power. The second approach is to define an arbitrary set of transformation $\mathcal{O}_{\mathcal{F}}$, which we deem as free, and quantify the resource content of a channel by defining a robustness measures in a similar way as we have done for states and measurements.

A. Robustness-generating power of channels

Let $\mathcal{F} \subseteq \Omega(\mathcal{V})$ be a set of free states in a given resource theory, with \mathcal{F}' the corresponding free states in another space \mathcal{V}' . If a transformation $\Lambda \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ is a free operation in this resource theory, we know for certain that $\sigma \in \mathcal{F} \Rightarrow \Lambda(\sigma) \in \mathcal{F}'$. As we have seen before, the resource content of a state is naturally and operationally quantified by the robustness $R_{\mathcal{F}}(\omega)$. Motivated by this, we like to understand the best robustness achievable after the application of a channel on a resourceless state. The robustness-generating power of a map Λ is then defined as [42,118]

$$\begin{split} P_{\mathcal{F}}(\Lambda) &\coloneqq \max_{\sigma \in \mathcal{F}} R_{\mathcal{F}'}[\Lambda(\sigma)] \\ &= \max\{ \langle W, \Lambda(\sigma) \rangle - 1 | \sigma \in \mathcal{F}, W \in \mathcal{C}'^*, \\ &\langle W, \pi \rangle \leq 1 \ \forall \ \pi \in \mathcal{F}' \}, \end{split} \tag{49}$$

where the second line follows from the dual form of the robustness.

We now show that this quantity admits an operational interpretation in any GPT as the best advantage that Λ can enable in state discrimination tasks of free-state ensembles with a given measurement, where the transformation Λ is applied to the ensembles prior to measurement. Specifically, consider a state discrimination task with the possibility of an inconclusive outcome, where the average success probability is

$$p_{\text{succ}}'(\{p_i, \sigma_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^{N}, \Lambda) := \sum_{i=0}^{N-1} p_i \langle M_i, \Lambda(\sigma_i) \rangle. \quad (50)$$

We then have the following.

Theorem 4: Given a map $\Lambda \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$, its robustness-generating power is equivalently given by

$$\max_{\{M_i\}} \frac{\max_{\{p_i,\sigma_i\},\sigma_i \in \mathcal{F}} p'_{\text{succ}}(\{p_i,\sigma_i\},\{M_i\},\Lambda)}{\max_{\{q_i,\pi_i\},\pi_i \in \mathcal{F}'} p'_{\text{succ}}(\{q_i,\pi_i\},\{M_i\},\text{id})} = 1 + P_{\mathcal{F}}(\Lambda),$$
(51)

where id denotes the identity map on \mathcal{V}' .

Proof.—One direction is as usual. Notice that the definition of the robustness-generating power implies that, for any free state $\sigma_i \in \mathcal{F}$, there exists another state $\delta_i \in \Omega(\mathcal{V}')$ such that $[\Lambda(\sigma_i) + r_i\delta_i]/(1 + r_i) = \rho_i \in \mathcal{F}'$ where $r_i \leq r = P_{\mathcal{F}}(\Lambda)$. This gives for any $\{M_i\}$ and any free-state ensemble $\{p_i, \sigma_i\}$ that

$$\begin{aligned} p_{\text{succ}}'(\{p_{i}, \sigma_{i}\}, \{M_{i}\}, \Lambda) \\ &= \sum_{i} p_{i} \langle M_{i}, (1+r_{i})\rho_{i} - r_{i}\delta_{i} \rangle \\ &\leq \sum_{i} p_{i} \langle M_{i}, (1+r)\rho_{i} \rangle \\ &\leq (1+r) \underset{x \in \mathcal{F}}{\max} p_{\text{succ}}'(\{q_{i}, \pi_{i}\}, \{M_{i}\}, \text{id}). \end{aligned} \tag{52}$$

On the other hand, consider the optimal solution W in Eq. (49) and define a measurement by $\{W/\|W\|_{\Omega'}^{\circ},\ U'-W/\|W\|_{\Omega'}^{\circ}\}$. We can then choose the single-element ensemble $\{1,\sigma\}$ where σ is an optimal state in Eq. (49). This gives

$$\max_{\{M_{i}\}} \frac{\max_{\{p_{i},\sigma_{i}\},\sigma_{i}\in\mathcal{F}} p'_{\text{succ}}(\{p_{i},\sigma_{i}\},\{M_{i}\},\Lambda)}{\max_{\{q_{i},\pi_{i}\},\pi_{i}\in\mathcal{F}'} p'_{\text{succ}}(\{q_{i},\pi_{i}\},\{M_{i}\},\text{id})}$$

$$\geq \frac{\frac{1}{\|W\|_{\Omega'}^{\circ}} \langle W,\Lambda(\sigma)\rangle}{\frac{1}{\|W\|_{\Omega'}^{\circ}} \max_{\pi\in\mathcal{F}'} \langle W,\pi\rangle}$$

$$\geq 1 + P_{\mathcal{F}}(\Lambda), \tag{53}$$

where in the last inequality we use that $\pi \in \mathcal{F}'$, and so $\langle W, \pi \rangle \leq 1$ by the conditions on W in Eq. (49).

The above result establishes a universal operational meaning of the robustness-generating power for any resource theory and any choice of free operations in the given theory.

We remark that Ref. [37] showed an operational interpretation in the context of binary channel discrimination of another measure of channel resourcefulness within quantum mechanics, the resource-generating power as quantified by the trace norm distance to the set of free states: $\max_{\sigma \in \mathcal{F}} \min_{\sigma' \in \mathcal{F}'} \|\Lambda(\sigma) - \sigma'\|_1$. It can be noticed that the proof in Ref. [37] uses only the fundamental properties of the trace norm as the base norm and thus holds in the same way in any general probabilistic theory.

B. Generalized robustness of channels

Here, we take another approach to quantify the resource-fulness of channels, in which we aim to directly evaluate the intrinsic resourcefulness of a given channel without the aid of an underlying resource theory of states. Having seen the discussions for states and measurements, a natural approach to take is to measure the resourcefulness with respect to a given set of free channels $\mathcal{O}_{\mathcal{F}}$. Analogous to the cases of states and measurements, we propose the generalized robustness of channels and discuss an operational meaning of it via channel discrimination (see also Ref. [42] for a discussion of this quantity in the resource theory of quantum coherence).

The operational characterization of this measure depends on several properties which general GPTs need not satisfy, and in particular makes heavy use of the Choi-Jamiołkowski isomorphism. Although equivalent forms of this isomorphism can be obtained in GPTs beyond quantum mechanics under suitable assumptions [54], for the simplicity of the discussion, we limit ourselves to quantum theory. Thus, in this subsection, operator inequalities should be understood in terms of positive semidefiniteness (we do not distinguish between the state and effect cones as they are isomorphic to each other).

Given a convex and closed set of channels $\mathcal{O}_{\mathcal{F}} \subseteq \mathcal{T}(\mathcal{V}, \mathcal{V}')$, we define the generalized robustness measure for channel $\Lambda \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ with respect to $\mathcal{O}_{\mathcal{F}}$ as

$$R_{\mathcal{O}_{\mathcal{F}}}(\Lambda) := \min \left\{ r \in \mathbb{R}_{+} \middle| \frac{\Lambda + r\Theta}{1 + r} \in \mathcal{O}_{\mathcal{F}}, \Theta \in \mathcal{T}(\mathcal{V}, \mathcal{V}') \right\}. \tag{54}$$

For further analysis, we utilize the Choi representation of channels. Let $J_{\Lambda} \coloneqq \mathbb{I} \otimes \Lambda(|\tilde{\Phi}^{+}\rangle\langle \tilde{\Phi}^{+}|)$ where $|\tilde{\Phi}^{+}\rangle \coloneqq \sum_{j} |jj\rangle$ is the unnormalized maximally entangled state. It is well known that Λ is completely positive if and only if $J_{\Lambda} \succeq 0$, trace preserving if and only if $\mathrm{Tr}_{\mathcal{V}}[J_{\Lambda}] = \mathbb{I}_{\mathcal{V}}$. Let $\mathcal{O}_{\mathcal{F}}^{J} \subseteq \mathcal{V} \otimes \mathcal{V}'$ be the set of Choi matrices corresponding to channels in $\mathcal{O}_{\mathcal{F}}$. We assume that $\mathcal{O}_{\mathcal{F}}^{J}$ contains at least one interior point of $\mathcal{C} \otimes \mathcal{C}'$, i.e., a channel Γ s.t. $J_{\Gamma} \succ 0$, so that

Eq. (54) is well defined for any channel. Then, introducing the variable $\Xi = \Lambda + r\Theta$, Eq. (54) can be rewritten as the following optimization problem:

minimize
$$r$$
, (55)

subject to
$$J_{\Lambda} \leq J_{\Xi}$$
, (56)

$$J_{\Xi} \in \operatorname{cone}(\mathcal{O}_{\mathcal{F}}^{J}),$$
 (57)

$$\operatorname{Tr}_{\mathcal{V}}[J_{\Xi}] = (1+r)\mathbb{I}_{\mathcal{V}},\tag{58}$$

where $cone(\mathcal{O}_{\mathcal{F}}^{J})$ is the cone generated by $\mathcal{O}_{\mathcal{F}}^{J}$. The corresponding dual problem is written as (see, e.g., Appendix B)

maximize
$$Tr[YJ_{\Lambda}] - 1$$
, (59)

subject to
$$Y = -Z + X \otimes \mathbb{I} \succeq 0$$
, (60)

$$X \in \mathcal{V}, \qquad \operatorname{Tr}[X] = 1,$$
 (61)

$$Z \in \mathcal{V} \otimes \mathcal{V}'$$
, $\operatorname{Tr}[ZJ_{\Xi}] \geq 0$, $\forall J_{\Xi} \in \mathcal{O}_{\mathcal{F}}^{J}$. (62)

It can be confirmed that the strong duality holds by taking $X = \mathbb{I}_{\mathcal{V}}/d$ and $Z = \mathbb{I}_{\mathcal{V}\mathcal{V}}/(2d)$ where $d = \dim \mathcal{V}$. Using Eqs. (60) and (61), one can write Eq. (62) as

$$Tr[(-Y + X \otimes \mathbb{I})J_{\Xi}] = -Tr[YJ_{\Xi}] + Tr[\Xi(X)]$$
(63)

$$= -\text{Tr}[YJ_{\Xi}] + 1 \ge 0. \tag{64}$$

Since the objective function does not include X or Z, we reach the following equivalent formulation;

maximize
$$Tr[YJ_{\Lambda}] - 1$$
, (65)

subject to
$$Y \succeq 0$$
, (66)

$$\operatorname{Tr}[YJ_{\Xi}] \le 1, \quad \forall \ J_{\Xi} \in \mathcal{O}_{\mathcal{F}}^{J}.$$
 (67)

This quantity finds a connection with two different types of discrimination problems: on the one hand, when considered for a single channel, it can characterize the advantage that the channel provides over all free channels in state discrimination tasks with given input states and measurement. On the other hand, when considered for an ensemble of channels, it can be used to express the advantage provided by the ensemble in a class of channel discrimination problems over ensembles composed of free channels.

Specifically, let us first consider the problem of discriminating an ensemble of quantum states by an

application of the channel $\mathbb{I} \otimes \Lambda$. The average success probability for this task can be expressed as

$$p_{\text{succ}}(\{p_j,\sigma_j\},\{M_j\},\mathbb{I}\otimes\Lambda) = \sum_j p_j \text{Tr}[\mathbb{I}\otimes\Lambda(\sigma_j)M_j]. \tag{68}$$

We then have the following result.

Theorem 5: Let $\mathcal{O}_{\mathcal{F}}$ be a convex and closed set of free channels. Then

$$\max_{\substack{\{p_{j},\sigma_{j}\},\{M_{j}\}\\}} \frac{p_{\text{succ}}(\{p_{j},\sigma_{j}\},\{M_{j}\},\mathbb{I}\otimes\Lambda)}{\max_{\Xi\in\mathcal{O}_{\mathcal{F}}}p_{\text{succ}}(\{p_{j},\sigma_{j}\},\{M_{j}\},\mathbb{I}\otimes\Xi)}$$

$$= 1 + R_{\mathcal{O}_{\mathcal{F}}}(\Lambda), \tag{69}$$

where the maximization is over all state ensembles $\{p_j, \sigma_j\}$ with $\sigma_j \in \Omega(\mathcal{V} \otimes \mathcal{V}')$ and all measurements $\{M_j\}$ in the space $\mathcal{V} \otimes \mathcal{V}'$.

Proof.—The proof again proceeds analogously to the proof of Theorem 1. Specifically, for any given measurement and state ensemble, we have

$$p_{\text{succ}}(\{p_j,\sigma_j\},\{M_j\},\mathbb{I}\otimes\Lambda) = \underset{j}{\sum} p_j \text{Tr}[\mathbb{I}\otimes\Lambda(\sigma_j)M_j] \eqno(70)$$

$$\leq \sum_{j} p_{j} (1+r) \operatorname{Tr}[\mathbb{I} \otimes \tilde{\Xi}(\sigma_{j}) M_{j}]$$
 (71)

$$\leq (1+r) \max_{\Xi \in \mathcal{O}_{\mathcal{F}}} p_{\text{succ}}(\{p_j, \sigma_j\}, \{M_j\}, \mathbb{I} \otimes \Xi), \quad (72)$$

where $r=R_{\mathcal{O}_{\mathcal{F}}}(\Lambda)$ and $\tilde{\Xi}\in\mathcal{O}_{\mathcal{F}}$ is a free channel for an optimal decomposition of Λ . The converse follows by choosing the measurement $\{Y/\|Y\|_{\infty}, \mathbb{I}-Y/\|Y\|_{\infty}\}$ where Y is an optimal solution in Eq. (65) and the ensemble $\{p_i,\sigma_i\}_{i=0}^1$ defined as $p_0=1,\sigma_0=|\Phi^+\rangle\langle\Phi^+|$, and $p_1=0$ where $|\Phi^+\rangle$ is the maximally entangled state and σ_1 is an arbitrary state. For this choice, it holds that

$$\frac{p_{\text{succ}}(\{p_{j}, \sigma_{j}\}, \{M_{j}\}, \mathbb{I} \otimes \Lambda)}{\max_{\Xi \in \mathcal{O}_{\mathcal{F}}} p_{\text{succ}}(\{p_{j}, \sigma_{j}\}, \{M_{j}\}, \mathbb{I} \otimes \Xi)}$$

$$= \frac{\text{Tr}[\mathbb{I} \otimes \Lambda(|\Phi^{+}\rangle\langle\Phi^{+}|)Y]}{\max_{\Xi \in \mathcal{O}_{\mathcal{F}}} \text{Tr}[\mathbb{I} \otimes \Xi(|\Phi^{+}\rangle\langle\Phi^{+}|)Y]}$$

$$= \frac{\text{Tr}[J_{\Lambda}Y]}{\max_{\Xi \in \mathcal{O}_{\mathcal{F}}} \text{Tr}[J_{\Xi}Y]}$$

$$\geq 1 + R_{\mathcal{O}_{\mathcal{F}}}(\Lambda), \tag{73}$$

where the inequality follows from Eq. (67).

Another approach is to consider, instead of a single channel Λ , an ensemble of channels, and characterize

its performance in a class of channel discrimination problems.

Consider a scenario where a channel sampled from some prior distribution is applied to one part of the input bipartite state and a collective measurement is made on the output system, where we allow for an inconclusive measurement outcome. The success probability of this channel discrimination problem is written as

$$p'_{\text{succ}}(\{p_i, \mathbb{I} \otimes \Lambda_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^{N}, \omega)$$

$$= \sum_{i=0}^{N-1} p_j \text{Tr}[\{\mathbb{I} \otimes \Lambda_j(\omega)\} M_j]. \tag{74}$$

For a given ensemble of N-1 channels $\{p_i, \Lambda_i\}_{i=0}^{N-1}$, where we assume that each $p_i > 0$ for simplicity, let us consider the maximum robustness for the ensemble:

$$\tilde{R}_{\mathcal{O}_{\mathcal{F}}}(\{p_i, \Lambda_i\}) := \max_{i} R_{\mathcal{O}_{\mathcal{F}}}(\Lambda_j). \tag{75}$$

Then, we obtain the following result, which connects the maximal advantage for this class of channel discrimination and the maximum robustness of channel ensembles.

Theorem 6: Let $\mathcal{O}_{\mathcal{F}}$ be a convex and closed set of free channels. Then,

$$\begin{split} \max_{\substack{\boldsymbol{\omega} \in \Omega(\mathbb{V}\otimes \mathbb{V}) \\ \{M_{j}\}_{j=0}^{N}}} & \frac{p_{\text{succ}}'(\{p_{j}, \mathbb{I} \otimes \Lambda_{j}\}_{j=0}^{N-1}, \{M_{j}\}_{j=0}^{N}, \boldsymbol{\omega})}{\max_{\Xi_{j} \in \mathcal{O}_{\mathcal{F}}} p_{\text{succ}}'(\{p_{j}, \mathbb{I} \otimes \Xi_{j}\}_{j=0}^{N-1}, \{M_{j}\}_{j=0}^{N}, \boldsymbol{\omega})} \\ &= 1 + \tilde{R}_{\mathcal{O}_{\mathcal{F}}}(\{p_{j}, \Lambda_{j}\}). \end{split} \tag{76}$$

Proof.—One direction of the inequality is shown as usual,

$$p'_{\text{succ}}(\{p_j, \mathbb{I} \otimes \Lambda_j\}_{j=0}^{N-1}, \{M_i\}_{j=0}^N, \omega)$$

$$= \sum_{i=0}^{N-1} p_j \text{Tr}[\{\mathbb{I} \otimes \Lambda_j(\omega)\} M_j]$$
(77)

$$\leq \sum_{j=0}^{N-1} p_j (1+r_j) \text{Tr}[\{\mathbb{I} \otimes \Xi_j(\omega)\} M_j]$$
 (78)

$$\leq (1+r_{j^{\star}}) \max_{\Xi_{j} \in \mathcal{O}_{\mathcal{F}}} p'_{\text{succ}}(\{p_{j}, \mathbb{I} \otimes \Xi_{j}\}_{j=0}^{N-1}, \{M_{j}\}_{j=0}^{N}, \omega), \quad (79)$$

where we set $r_j\coloneqq R_{\mathcal{O}_{\mathcal{F}}}(\Lambda_j)$ and $j^\star=\operatorname{argmax}_j r_j$. To show the left-hand side is greater than or equal to the right-hand side, consider the state $\omega=|\Phi^+\rangle\langle\Phi^+|$ and the POVM defined by $M_{j^\star}=Y_{j^\star}/\|Y_{j^\star}\|_\infty$ where Y_{j^\star} is an optimal solution for Λ_{j^\star} in Eq. (65), $M_j=0$ for $j\neq j^\star$, and $M_N=\mathbb{I}-M_{j^\star}$. For this choice, we then have

$$\frac{p'_{\text{succ}}(\{p_{j}, \mathbb{I} \otimes \Lambda_{j}\}_{j=0}^{N-1}, \{M_{i}\}_{j=0}^{N}, \omega)}{\max_{\Xi_{j} \in \mathcal{O}_{\mathcal{F}}} p'_{\text{succ}}(\{p_{j}, \mathbb{I} \otimes \Xi_{j}\}_{j=0}^{N-1}, \{M_{i}\}_{j=0}^{N}, \omega)}$$

$$= \frac{\sum_{j=0}^{N-1} p_{j} \text{Tr}[\{\mathbb{I} \otimes \Lambda_{j}(|\Phi^{+}\rangle\langle\Phi^{+}|)\}M_{j}]}{\max_{\Xi_{i} \in \mathcal{O}_{\mathcal{F}}} \sum_{j=0}^{N-1} p_{j} \text{Tr}[\{\mathbb{I} \otimes \Xi_{j}(|\Phi^{+}\rangle\langle\Phi^{+}|)\}M_{j}]}$$
(80)

$$= \frac{p_{j^{\star}} \operatorname{Tr}[J_{\Lambda_{j^{\star}}} Y_{j^{\star}}]}{\max_{\Xi \in \mathcal{O}_{\mathcal{F}}} p_{j^{\star}} \operatorname{Tr}[J_{\Xi} Y_{j^{\star}}]}$$
(81)

$$\geq 1 + \tilde{R}_{\mathcal{O}_{\pi}}(\{p_i, \Lambda_i\}), \tag{82}$$

where the last inequality is due to Eq. (67).

We remark that, in the specific case of the resource theory of quantum coherence and the choice of operations $\mathcal{O}_{\mathcal{F}}$ as the so-called maximally incoherent operations (the largest set of operations preserving the set of free states) [6], it was shown that the two approaches to the quantification of channel resources coincide, and we have $P_{\mathcal{F}}(\Lambda) = R_{\mathcal{O}_{\mathcal{F}}}(\Lambda)$ for any channel Λ [42]. An investigation of more general cases under which this happens is an interesting open question in the operational characterization of resources of transformations.

VI. STANDARD ROBUSTNESS OF STATES

We have seen that the generalized robustness is a fundamental measure capturing the resourcefulness contained in general types of resources with operational significance. In addition to the generalized robustness, another important member of the class of robustness measures is the standard (free) robustness. This is also a valid resource monotone in any convex resource theory of states [34] with some known operational interpretations; the standard robustness of entanglement plays a role in activation of quantum teleportation [92] and one-shot cost for entanglement dilution under the nonentangling operations [89], and the standard robustness of magic is related to classical simulation overhead for quantum Clifford circuits with magic-state inputs [18]. Furthermore, this quantifier in several quantum resource theories can be evaluated analytically for some states [18,84,109,119] and admits computable forms as semidefinite or linear programs [18,108,109]. However, the standard robustness has not found use in discrimination-type problems thus far, and a universal operational meaning of the standard robustness in general resource theories, whether in quantum mechanics or beyond, has not been established. We address this question for the standard robustness of states and indeed give such a general operational meaning in terms of the advantage for the most fundamental type of discrimination task: balanced binary channel discrimination.

An interesting difference between the generalized robustness and standard robustness is that, unlike generalized robustness, standard robustness can diverge for any resource state in some important theories such as the resource theory of coherence and asymmetry. It may thus appear that interpreting such a divergent quantity in an operational setting would be implausible. We address this issue by considering a natural figure of merit for the advantage in balanced binary channel discrimination in a way that it encompasses this seemingly problematic circumstances as well at the limit of diverging standard robustness. Note that the following argument is valid in any GPT with an additional reasonable assumption introduced below.

The standard robustness of a state $\omega \in \Omega(\mathcal{V})$ with respect to a convex and closed set of free states $\mathcal{F} \subseteq \Omega(\mathcal{V})$ is defined as

$$R_{\mathcal{F}}^{\mathcal{F}}(\omega) := \inf \left\{ r \in \mathbb{R}_{+} \middle| \frac{\omega + r\pi}{1 + r} \in \mathcal{F}, \pi \in \mathcal{F} \right\}.$$
 (83)

It is straightforward to verify that the standard robustness is faithful $(R_{\mathcal{F}}^{\mathcal{F}}(\omega) = 0 \Leftrightarrow \omega \in \mathcal{F})$, convex, and monotonic under free operations in the sense that $R_{\mathcal{F}}^{\mathcal{F}}[\Lambda(\omega)] \leq R_{\mathcal{F}}^{\mathcal{F}}(\omega)$ for any linear map such that $\Lambda[\mathcal{F}] \subseteq \mathcal{F}$ [34]. The dual optimization problem can be obtained as

maximize
$$\langle X, \omega \rangle - 1$$
,
subject to $0 \le \langle X, \sigma \rangle \le 1 \quad \forall \ \sigma \in \mathcal{F}$. (84)

We stress that this quantity is finite for any ω only when the set $\mathcal F$ spans the whole space $\mathcal V$; we therefore make this assumption for the discussed quantities to be well defined, although we later see that it is not necessary for the operational interpretation of $R_{\mathcal F}^{\mathcal F}$. Under this assumption, the infimum in Eq. (83) is achieved, and strong duality always holds. Indeed, the cone generated by the set $\mathcal F$ is pointed by the pointedness of $\mathcal C$, and the dual of any pointed cone is generating; hence, it contains an interior point, and so we can always choose a strictly feasible solution.

We show that the standard robustness characterizes the maximum advantage that a resource state provides over free states in channel discrimination. For this task to be well defined, one needs to specify what constitutes the set of physically allowed state transformations $\mathcal{T}(\mathcal{V},\mathcal{V}')$ or channels in the given GPT. As discussed before, the very basic assumption one can make about the channels is that they preserve the state cone as well as the normalization of states. Throughout this section, we make an additional assumption: Measure-and-prepare channels of the form $\Lambda(\omega) = \sum_i \langle M_i, \omega \rangle \omega_i'$ with $\{M_i\}_i$ being a measurement, and $\{\omega_i'\}_i$ a collection of states in the output space are allowed physical transformations.

The justification for allowing measure-and-prepare channels follows from the fact that measurement and state preparation can be considered as the two building blocks of any GPT and are necessarily physically implementable [51,52,102]; one then needs only to allow for

classical information about the measurement outcome to be transferred in order to implement any such measure-and-prepare channel. In particular, this assumption is clearly satisfied in any theory which extends quantum mechanics or even classical probability theory. We denote by $\mathcal{T}(\mathcal{V},\mathcal{V}')$ the set of all physically allowed operations within the given GPT and assume throughout this section that it satisfies the assumption above. In fact, for the purpose of our proof, it suffices to assume that only binary measure-and-prepare channels based on two-outcome measurements are allowed.

Consider now a binary channel discrimination problem where one of two channels is applied according to the ensemble $\{1/2,\Lambda_i\}_{i=0}^1$, with each $\Lambda_i \in \mathcal{T}(\mathcal{V},\mathcal{V}')$. The only assumption about the output space \mathcal{V}' we make is that it contains at least two distinct states, as otherwise the task becomes trivial. Since the two channels Λ_i are equiprobable, a random guess gives the success probability 1/2, and thus, a meaningful figure of merit of this task is how much one can increase the success probability by suitably choosing a measurement and inferring the applied channel from the measurement outcome. Motivated by this observation, we consider the following quantity:

$$p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega)$$

$$:= p_{\text{succ}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega) - \frac{1}{2}, \quad (85)$$

and the one being maximized with respect to measurement,

$$p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \omega) = \max_{\{M_i\}_{i=0}^1 \in \mathcal{M}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega).$$
(86)

Then, we obtain the following result, stating that the maximum advantage in terms of success probability gain that a resource state provides over free states is characterized exactly by the standard robustness measure.

Theorem 7: It holds that

$$\max_{\Lambda_0, \Lambda_1 \in \mathcal{T}(\mathcal{V}, \mathcal{V})} \frac{p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \omega)}{\max_{\sigma \in \mathcal{F}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \sigma)} = 1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega),$$
(87)

where the maximization is over all possible ensembles $\{1/2, \Lambda_i\}_{i=0}^1$.

Proof.—By the definition of the standard robustness, there exist $\tilde{\tau}, \tilde{\sigma} \in \mathcal{F}$ such that $\omega = [1 + R_{\mathcal{F}}^{\mathcal{F}}(\omega)]\tilde{\tau} - R_{\mathcal{F}}^{\mathcal{F}}(\omega)\tilde{\sigma}$. For any ensemble of channels $\{1/2, \Lambda_i\}_{i=0}^1$, let $\{\tilde{M}_i\}_{i=0}^1$ be a measurement maximizing the quantity $p_{\text{succ}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega)$. Then,

$$p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \omega) = \frac{1}{2} \sum_{i=0}^1 \langle \tilde{M}_i, \Lambda_i(\omega) \rangle - \frac{1}{2}$$
 (88)

$$= \frac{1}{2} [1 + R_{\mathcal{F}}^{\mathcal{F}}(\omega)] \sum_{i=0}^{1} \langle \tilde{M}_i, \Lambda_i(\tilde{\tau}) \rangle - \frac{1}{2} R_{\mathcal{F}}^{\mathcal{F}}(\omega) \sum_{i=0}^{1} \langle \tilde{M}_i, \Lambda_i(\tilde{\sigma}) \rangle - \frac{1}{2}$$

$$(89)$$

$$= \left[1 + R_{\mathcal{F}}^{\mathcal{F}}(\omega)\right] \left(\frac{1}{2} \sum_{i=0}^{1} \langle \tilde{M}_i, \Lambda_i(\tilde{\tau}) \rangle - \frac{1}{2}\right) - R_{\mathcal{F}}^{\mathcal{F}}(\omega) \left(\frac{1}{2} \sum_{i=0}^{1} \langle \tilde{M}_i, \Lambda_i(\tilde{\sigma}) \rangle - \frac{1}{2}\right)$$
(90)

$$\leq [1 + R_{\mathcal{F}}^{\mathcal{F}}(\omega)] \max_{\sigma \in \mathcal{F}} \max_{\{M_i\}} \left(\frac{1}{2} \sum_{i=0}^{1} \langle M_i, \Lambda_i(\sigma) \rangle - \frac{1}{2} \right) \tag{91}$$

$$+ R_{\mathcal{F}}^{\mathcal{F}}(\omega) \max_{\sigma \in \mathcal{F}} \max_{\{M_i\}} \left(\frac{1}{2} - \frac{1}{2} \sum_{i=0}^{1} \langle M_i, \Lambda_i(\sigma) \rangle \right). \tag{92}$$

Let us change the variable in the second term by $M_0 = U - M_0'$ and $M_1 = U - M_1'$ where $0 \leq_{\mathcal{C}^*} M_0', M_1' \leq_{\mathcal{C}^*} U$, $M_0' + M_1' = U$. Then, the second term becomes

$$\frac{1}{2} - \frac{1}{2} \sum_{i=0}^{1} \langle M_i, \Lambda_i(\sigma) \rangle = \frac{1}{2} - \frac{1}{2} \sum_{i=0}^{1} \langle U - M_i', \Lambda_i(\sigma) \rangle = \frac{1}{2} \sum_{i=0}^{1} \langle M_i', \Lambda_i(\sigma) \rangle - \frac{1}{2}, \tag{93}$$

and the maximum in Eq. (92) is taken with respect to the new variables $\{M'_i\}$. Hence, we get for any Λ_0 , Λ_1 ,

$$\begin{aligned} & p_{\text{gain}}(\left\{1/2, \Lambda_{i}\right\}_{i=0}^{1}, \omega) \\ & \leq \left[1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega)\right] \max_{\sigma \in \mathcal{F}} \max_{\left\{M_{i}\right\}} \left(\frac{1}{2} \sum_{i=0}^{1} \left\langle M_{i}, \Lambda_{i}(\sigma) \right\rangle - \frac{1}{2}\right) \end{aligned} \tag{94}$$

$$= [1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega)] \max_{\sigma \in \mathcal{F}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \sigma).$$
 (95)

To show the opposite inequality, with the change of variables X' := 2X - U we rewrite the optimization problem (84) as

$$1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega) = \max\{\langle X', \omega \rangle | -1 \le \langle X', \sigma \rangle \le 1 \ \forall \ \sigma \in \mathcal{F}\}$$

$$(96)$$

for any state ω . Let X' be an optimal solution in the above and $\eta_0, \eta_1 \in \Omega(\mathcal{V}')$ be any two distinct states, and consider the linear maps defined by

$$\Lambda_0(\omega) \coloneqq \frac{1}{2} \langle U + \frac{X'}{\|X'\|_{\Omega}^{\circ}}, \omega \rangle \eta_0 + \frac{1}{2} \langle U - \frac{X'}{\|X'\|_{\Omega}^{\circ}}, \omega \rangle \eta_1,$$

$$\tag{97}$$

$$\Lambda_{1}(\omega) := \frac{1}{2} \langle U - \frac{X'}{\|X'\|_{\Omega}^{\circ}}, \omega \rangle \eta_{0} + \frac{1}{2} \langle U + \frac{X'}{\|X'\|_{\Omega}^{\circ}}, \omega \rangle \eta_{1}.$$

$$(98)$$

Now, since

$$\langle U \pm \frac{X'}{\|X'\|_{\Omega}^{\circ}}, \rho \rangle = 1 \pm \frac{\langle X', \rho \rangle}{\max_{\pi \in \Omega(\mathcal{V})} |\langle X', \pi \rangle|} \in [0, 2],$$
(99)

for any $\rho \in \Omega(\mathcal{V})$, the set $\{\frac{1}{2}[U+X'/\|X'\|_{\Omega}^{\circ}], \frac{1}{2}[U-X'/\|X'\|_{\Omega}^{\circ}]\}$ constitutes a valid measurement, and so Λ_0 and Λ_1 are measure-and-prepare channels. Hence,

 $\Lambda_0, \Lambda_1 \in \mathcal{T}(\mathcal{V}, \mathcal{V}')$ by our assumption about the set of allowed transformations. We then get

$$\max_{\{M_i\}} p_{\text{succ}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega) - \frac{1}{2}$$

$$= \frac{1}{4} \|\Lambda_0(\omega) - \Lambda_1(\omega)\|_{\Omega}$$

$$= \frac{1}{4} \left\|\frac{\langle X', \omega \rangle}{\|X'\|_{\Omega}^{\circ}} (\eta_0 - \eta_1)\right\|_{\Omega}$$

$$= \frac{1}{4\|X'\|_{\Omega}^{\circ}} |\langle X', \omega \rangle| \|\eta_0 - \eta_1\|_{\Omega}$$
(100)

using the absolute homogeneity of the norm $\|\cdot\|_{\Omega}$. Hence, for this choice of Λ_0 , Λ_1 , we get

$$\frac{p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \omega)}{\max_{\sigma \in \mathcal{F}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \sigma)}$$

$$= \frac{4\|X'\|_{\Omega}^{\circ} |\langle X', \omega \rangle| \|\eta_0 - \eta_1\|_{\Omega}}{4\|X'\|_{\Omega}^{\circ} \max_{\sigma \in \mathcal{F}} |\langle X', \sigma \rangle| \|\eta_0 - \eta_1\|_{\Omega}} \qquad (101)$$

$$= \frac{|\langle X', \omega \rangle|}{\max_{\sigma \in \mathcal{F}} |\langle X', \sigma \rangle|} \qquad (102)$$

$$\geq 1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\rho),\tag{103}$$

where the inequality is due to the fact that $|\langle X', \sigma \rangle| \le 1$ by Eq. (96).

Although the use of general measurements can in general provide a significant advantage over discrimination with restricted sets of measurements [77,100], one can notice by following the proof of Theorem 7 that, in the context of quantifying the advantage provided by a resource state over all free states, the restriction of measurements is inconsequential; the standard robustness still acts as the exact quantifier of the advantage in binary channel discrimination, regardless of how the allowed measurements are restricted. Specifically, we have the following.

Corollary 8: For any informationally complete closed set of measurements $\mathcal{M}_{\mathcal{F}}$, it holds that

$$\max_{\Lambda_0, \Lambda_1 \in \mathcal{T}(\mathcal{V}, \mathcal{V}')} \frac{\max_{\{M_i\} \in \mathcal{M}_{\mathcal{F}}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \omega)}{\max_{\{M_i\} \in \mathcal{M}_{\mathcal{F}}} \max_{\sigma \in \mathcal{F}} p_{\text{gain}}(\{1/2, \Lambda_i\}_{i=0}^1, \{M_i\}_{i=0}^1, \sigma)} = 1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega),$$
(104)

where the maximization is over all possible ensembles $\{1/2, \Lambda_i\}_{i=0}^1$.

Proof.—This follows exactly in the same way as the proof of Theorem 7 by replacing the norm $\|\cdot\|_{\Omega}$ with the distinguishability norm $\|\cdot\|_{\mathcal{M}_{\mathcal{F}}}$.

Another fact which can be noticed from the proof of Theorem 7 in Eq. (96) is that, for any state ω , the quantity $1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega)$ in fact corresponds to the base norm induced by the set \mathcal{F} [34,120]; specifically,

$$\begin{split} 1 + 2R_{\mathcal{F}}^{\mathcal{F}}(\omega) \\ &= \|\omega\|_{\mathcal{F}} \\ &= \min\{\lambda_{+} + \lambda_{-} | \omega = \lambda_{+}\sigma_{+} - \lambda_{-}\sigma_{-}, \lambda_{\pm} \in \mathbb{R}_{+}, \sigma_{\pm} \in \mathcal{F}\}. \end{split} \tag{105}$$

Although we assume that the standard robustness takes finite value in the course of the above argument, Theorem 7

and Corollary 8 successfully capture the cases where the standard robustness diverges for resource theories of states such as the theory of coherence and asymmetry. The results then imply that one can always find two channels for which no free state can enable one to perform the task better than the random guess in such theories.

We further remark an interesting connection between our results and the result in Ref. [18] where it was found that for the case of theory of magic, the same quantity $1+2R_{\mathcal{F}}^{\mathcal{F}}$ quantifies an upper bound for the overhead of classical simulation of quantum Clifford circuits with magic-state inputs. (Note that the name "robustness of magic" was used in Ref. [18] to refer to the quantity $1+2R_{\mathcal{F}}^{\mathcal{F}}$ instead of $R_{\mathcal{F}}^{\mathcal{F}}$.) It would be an interesting problem to investigate whether this is merely a coincidence or whether the two seemingly very different tasks, channel discrimination and classical simulation of quantum circuits, are related at a deeper level through the standard robustness measure.

VII. COMPLETE SETS OF MONOTONES

It is not difficult to see that resource monotones in any resource theory provide nontrivial necessary conditions on the manipulation of resources with free operations because of their monotonicity properties: namely, if one object contains a larger amount of resources than another with respect to any monotone, then the transformation from the less resourceful to the more resourceful one with any free operation is prohibited. However, a single monotone fails in most cases to completely characterize the resource transformations; it is necessary that the amount of resource measured by the monotone does not increase under free operations, but it is usually not sufficient to ensure the existence of a free operation realizing that transformation. Finding the necessary and sufficient conditions for the existence of a transformation by means of free operations for a given input and output object is one of the most important questions to address in resource theories, as it underlies the operational capabilities of a given resource theory.

We call a (possibly infinite) family of monotones a complete set of monotones if it fully characterizes the necessary and sufficient conditions for the existence of a free transformation. Such sets of monotones were discussed in several specific settings [25,26,32,121–131], but no set of general conditions with a clear operational meaning was previously known to provide a comprehensive characterization of transformations in general resource theories.

We now show that performance of a state or measurement in a class of channel or state discrimination tasks precisely serves as a complete set of monotones for general resource theories defined in any GPT. This, together with the results obtained above, completes the operational characterization of quantification and exact state transformations in general resource theories in terms of discrimination tasks.

A. Complete sets of monotones for states

Let $\mathcal{O} \subseteq \mathcal{T}(\mathcal{V},\mathcal{V})$ be a convex and closed set of transformations which contains the identity map $x \mapsto x$ and is closed under concatenation in the sense that $\Lambda_1, \Lambda_2 \in \mathcal{O}$ means that $\Lambda_2 \circ \Lambda_1 \in \mathcal{O}$. This set of assumptions is particularly natural if \mathcal{O} is identified with a set of free operations in a convex resource theory, where the action of the identity channel trivially cannot generate any resource and neither should the application of two free channels; however, our results are completely general, and we do not need to explicitly assume any relation with resource theories.

Consider now the problem of channel discrimination with a possible inconclusive measurement outcome of channels from the set \mathcal{O} , that is, ensembles of the form $\{p_i, \Lambda_i\}$ with each $\Lambda_i \in \mathcal{O}$ and p_i being the corresponding probabilities. We first show that the performance of two given states ω, ω' optimized over *all* such discrimination problems leads to a complete set of monotones under the operations \mathcal{O} , although as we see later, this condition can be significantly simplified.

To begin, let us consider a fixed (N+1)-outcome measurement $\mathbb{M} = \{M_i\}_{i=0}^N$ and state ω , for which the best achievable probability of success in all free channel discrimination problems is

$$\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega) \coloneqq \max_{\substack{\{p_i, \Lambda_i\}_{i=0}^{N-1} \\ \Lambda_i \in \mathcal{O}}} \sum_{i=0}^{N-1} p_i \langle M_i, \Lambda_i(\omega) \rangle. \quad (106)$$

We additionally consider the case where the probability distribution $\{p_i\}$ is fixed *a priori*, and the optimization is only over the channels $\{\Lambda_i\}$ themselves; specifically, in this case we have

$$\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega, \{p_i\}_{i=0}^{N-1}) \coloneqq \max_{\substack{\{\Lambda_i\}_{i=0}^{N-1}\\\Lambda_i \in \mathcal{O}}} \sum_{i=0}^{N-1} p_i \langle M_i, \Lambda_i(\omega) \rangle. \tag{107}$$

We now have the following result. The proof is based on an approach similar to the one taken to characterize measurement informativeness in Ref. [25] (see also Ref. [125]).

Theorem 9: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if either of the following conditions is satisfied:

- (i) it holds that $\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega) \geq \tilde{p}'_{\text{succ}}(\mathbb{M}, \omega')$ for any measurement $\mathbb{M} \in \mathcal{M}$,
- (ii) for a fixed probability distribution $\{p_i\}_{i=0}^{N-1}$, it holds that $\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega, \{p_i\}) \geq \tilde{p}'_{\text{succ}}(\mathbb{M}, \omega', \{p_i\})$ for any measurement $\mathbb{M} = \{M_i\}_{i=0}^N \in \mathcal{M}$.

Proof.—The proof proceeds in exactly the same way for both of the conditions. In the following, we consider the case where the probability distribution $\{p_i\}$ can vary [case (i)], but it can equivalently be taken to be fixed [case (ii)] for the remainder of the proof.

Suppose first that $\omega' = \tilde{\Lambda}(\omega)$ with $\tilde{\Lambda} \in \mathcal{O}$. Then, for any $\mathbb{M} = \{M_i\}$, we have

$$\tilde{p}'_{\text{succ}}(\{M_i\}, \omega') = \max_{\{p_i, \Lambda_i\}, \Lambda_i \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i(\omega') \rangle \qquad (108)$$

$$= \max_{\{p_i, \Lambda_i\}, \Lambda_i \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i \circ \tilde{\Lambda}(\omega) \rangle \quad (109)$$

$$\leq \max_{\{p_i, \Lambda_i'\}, \Lambda_i' \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i'(\omega) \rangle$$

$$= \tilde{p}'_{\text{succ}}(\{M_i\}, \omega), \tag{110}$$

where the inequality is due to the closedness of \mathcal{O} under concatenation.

To show the converse, suppose that $\forall \{M_i\}$, we have $\tilde{p}'_{\text{succ}}(\{M_i\}, \omega) \geq \tilde{p}'_{\text{succ}}(\{M_i\}, \omega')$. This implies that

$$0 \leq \inf_{\{M_{i}\}\in\mathcal{M}} \left[\max_{\{p_{i},\Lambda_{i}\},\Lambda_{i}\in\mathcal{O}} \sum_{i} p_{i} \langle M_{i},\Lambda_{i}(\omega) \rangle - \max_{\{q_{i},\Theta_{i}\},\Theta_{i}\in\mathcal{O}} \sum_{i} q_{i} \langle M_{i},\Theta_{i}(\omega') \rangle \right]$$

$$(111)$$

$$\leq \inf_{\{M_i\} \in \mathcal{M}} \max_{\{p_i, \Lambda_i\}, \Lambda_i \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i(\omega) - \omega' \rangle \qquad (112)$$

$$\leq \min_{\substack{\{M_i\}_{i=0}^{N} \in \mathcal{M} \\ \Lambda_i \in \mathcal{O}}} \max_{\substack{\{p_i, \Lambda_i\}_{i=0}^{N-1}, \\ \Lambda_i \in \mathcal{O}}} \sum_{i} p_i \langle M_i, \Lambda_i(\omega) - \omega' \rangle \tag{113}$$

$$= \max_{\substack{\{p_i, \Lambda_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^{N} \in \mathcal{M}}} \min_{\{M_i\}_{i=0}^{N}} \sum_{i} p_i \langle M_i, \Lambda_i(\omega) - \omega' \rangle, \qquad (114)$$

where the second inequality is obtained by setting each $q_i = p_i$ and each Θ_i to be the identity map in the second term, the third inequality is because we restricted the minimization over N+1-outcome measurements where $N \ge 2$ is an arbitrary integer, and the equality is due to Sion's minimax theorem [132] by the convexity and compactness of \mathcal{M} , \mathcal{O} , and the set of ensembles $\{p_i, \Lambda_i\}_{i=0}^{N-1}$ (since N is a finite integer), as well as the linearity of the objective function with respect to Λ_i and M_i .

To show that there exists $\Lambda \in \mathcal{O}$ such that $\Lambda(\omega) = \omega'$, suppose to the contrary that it does not hold true. Since any Λ_i is a physical channel and thus normalization preserving, we get that

$$\langle U, \Lambda_i(\omega) - \omega' \rangle = 0 \quad \forall i.$$
 (115)

In particular, since $\Lambda_i(\omega) - \omega' \neq 0$ by assumption, we cannot have that $\Lambda_i(\omega) - \omega' \in \mathcal{C}$ as this necessarily implies that $\langle U, \Lambda_i(\omega) - \omega' \rangle > 0$. Therefore, by the hyperplane separation theorem [97], for every Λ_i there exists an effect $E_i \in \mathcal{C}^*$ such that $\langle E_i, \Lambda_i(\omega) - \omega' \rangle < 0$. We now construct an incomplete measurement $\{M_i\}_{i=0}^{N-1}$ by

$$M_i \coloneqq \frac{E_i}{\|\sum_j E_j\|_{\Omega}^{\circ}},\tag{116}$$

such that $\sum_i M_i \leq_{\mathcal{C}^*} U$, and so $\{M_i\}_{i=0}^{N-1}$ can be completed to a measurement $\{M_i\}_{i=0}^N$ by appending another effect $M_N \coloneqq U - \sum_i M_i$ which does not affect the measurement outcomes corresponding to the ensemble $\{p_i, \Lambda_i\}_{i=0}^{N-1}$. We then have for this choice of $\{M_i\}$ that

$$\sum_{i} p_{i} \langle M_{i}, \Lambda_{i}(\omega) - \omega' \rangle$$

$$= \sum_{i} \frac{p_{i}}{\|\sum_{i} E_{i}\|_{O}^{\circ}} \langle E_{i}, \Lambda_{i}(\omega) - \omega' \rangle < 0.$$
(117)

Since such a measurement can be constructed for any ensemble $\{p_i, \Lambda_i\}$, this is in contradiction with Eq. (114), which states that there exists a choice of an ensemble $\{p_i, \Lambda_i\}$ such that for every measurement $\{M_i\}$ we have $\sum_i p_i \langle M_i, \Lambda_i(\omega) - \omega' \rangle \ge 0$. We conclude that our original assumption must be wrong, and there exists $\Lambda \in \mathcal{O}$ such that $\Lambda(\omega) = \omega'$.

The result immediately establishes a general relation between channel discrimination problems and state transformations under any fixed set of operations \mathcal{O} . Remarkably, the freedom in choosing the probability distribution $\{p_i\}$ allows us to significantly simplify this characterization and show that much *smaller* classes of channel discrimination problems already constitute complete sets of monotones. In particular, by taking $\{p_i\}_{i=0}^1$ with $p_0 = p_1 = \frac{1}{2}$, we reduce the problem to the much more straightforward task of binary channel discrimination.

Corollary 10: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if for any three-outcome measurement $\mathbb{M} = \{M_i\}_{i=0}^2$ it holds that $\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega, \{\frac{1}{2}, \frac{1}{2}\}) \geq \tilde{p}'_{\text{succ}}(\mathbb{M}, \omega', \{\frac{1}{2}, \frac{1}{2}\})$.

This greatly reduces the difficulty of determining whether a free transformation between two given states exists.

Another particularly interesting case is obtained by considering the binary probability distribution defined as $p_0 = 1$, $p_1 = 0$. Although going beyond what one could consider a physically motivated "discrimination" task, this allows us to obtain the following characterization of a complete set of monotones under the operations \mathcal{O} .

Corollary 11: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if for any $E \succeq_{\mathcal{C}^*} 0$ it holds that

$$\max_{\Lambda \in \mathcal{O}} \langle E, \Lambda(\omega) \rangle \ge \max_{\Lambda \in \mathcal{O}} \langle E, \Lambda(\omega') \rangle. \tag{118}$$

Proof.—It follows from Theorem 9 by noting that for any measurement \mathbb{M} , we have

$$\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega, \{1, 0\}) = \max_{\{\Lambda_i\}_{i=0}^1 \subset \mathcal{O}} \langle M_0, \Lambda_0(\omega) \rangle, \qquad (119)$$

so the conditions of the theorem reduce to verifying whether the inequality $\tilde{p}'_{\text{succ}}(\mathbb{M}, \omega, \{1, 0\}) \geq \tilde{p}'_{\text{succ}}(\mathbb{M}, \omega', \{1, 0\})$ is satisfied for any operator $M_0 \in \mathcal{C}^*$ such that there exists a valid measurement $\{M_i\}_{i=0}^2$. This can be easily verified to be precisely the set $0 \leq_{\mathcal{C}^*} M_0 \leq_{\mathcal{C}^*} U$. Without loss of generality, we can then relax the constraint $M_0 \leq_{\mathcal{C}^*} U$ as any $E \in \mathcal{C}^* \setminus \{0\}$ can be renormalized as $M_0 = E/\|E\|_{\Omega}^{\circ}$.

We can further make an observation that in Corollary 11 it suffices to optimize over effects E which are normalized in a suitable manner. Specifically, it suffices to verify whether Eq. (118) holds for any operator E in a chosen base of the cone C^* . Within quantum mechanics, or indeed in any GPT where $C \cong C^*$, this means in particular that there exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if

$$\max_{\Lambda \in \mathcal{O}} \langle \sigma, \Lambda(\omega) \rangle \ge \max_{\Lambda \in \mathcal{O}} \langle \sigma, \Lambda(\omega') \rangle \tag{120}$$

holds for any *state* σ . This recovers a result of Refs. [133,134] obtained with different methods in the context of quantum resource theories. We note that the class of monotones in Eq. (118) has previously been considered in the resource theory of quantum coherence [135], albeit without an operational application to state transformations under the free operations.

We further note that in Theorem 9 and Corollary 10, one could instead consider the tasks of subchannel discrimination from a chosen set \mathcal{O} of normalization nonincreasing maps. The proofs proceed analogously.

Alternatively, we can establish a complete set of monotones by considering a modification of the task: we now consider channel discrimination (without inconclusive outcomes) over all valid choices of channel ensembles but allow for the application of a single chosen prior transformation from the set $\mathcal O$ to the ensemble before applying the channels to be discriminated. The success probability for this task for a choice of channel ensemble $\{p_i, \Lambda_i\}$ and measurement $\{M_i\}$ is given by

$$\tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega) \coloneqq \max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i \circ \Xi(\omega) \rangle. \tag{121}$$

Note that we separate the probability distribution $\{p_i\}$ from the corresponding channels $\{\Lambda_i\}$ for reasons which become clear shortly. We now show that this success probability serves as a complete set of monotones for state transformations under the operations \mathcal{O} in two different ways.

Theorem 12: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if either of the following conditions is satisfied:

- (i) it holds that $\tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega) \ge \tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega')$ for all channel ensembles $\{p_i, \Lambda_i\}$ and measurements $\{M_i\}$,
- (ii) for a fixed set of channels $\{\Lambda_i\}_{i=0}^{N-1}$ containing the identity channel id, it holds that $\tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega) \geq \tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega')$ for all

probability distributions $\{p_i\}_{i=0}^{N-1}$ and measurements $\{M_i\}_{i=0}^{N-1}$.

Note that one could also consider subchannel discrimination in (i) instead of channel discrimination. The proof proceeds in a way similar to Theorem 9, and we include it in Appendix A for completeness.

An interesting application of Theorem 12 is obtained by choosing the set of channels $\{\Lambda_i\}_{i=0}^1$ where $\Lambda_0 = \mathrm{id}$ and $\Lambda_1 = \Theta$ is a fixed transformation. The theorem then gives the following.

Corollary 13: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if for all two-element probability distributions $\{p_i\}$ and two-outcome measurements $\{M_i\}$ it holds that

$$\tilde{p}_{\text{succ}}(\{p_i\}, \{\text{id}, \Theta\}, \{M_i\}, \omega)$$

$$\geq \tilde{p}_{\text{succ}}(\{p_i\}, \{\text{id}, \Theta\}, \{M_i\}, \omega')$$
(122)

for some a priori fixed choice of $\Theta \in \mathcal{T}(\mathcal{V}, \mathcal{V})$.

One can interpret this scenario as detecting the noise introduced by Θ with the help of prior processing with the operations \mathcal{O} . Remarkably, Corollary 13 then shows that a single noise model completely determines the capability of the input state ω as a noise-detecting resource aided by the operations \mathcal{O} —if ω is better than ω' at detecting *some* type of noise Θ for any noise strength and detection strategy, ω is more capable than ω' in detecting any other noise introduced by a physical transformation. This tells us for instance that, using a standard quantum-mechanical example [136], considering only the family of depolarizing noise models is sufficient to assess the usefulness of a given state for all possible preprocessing assisted noise-detection tasks considered here. It is notable that this nontrivial fact can be shown via the seemingly unrelated problem of resource manipulation thanks to Theorem 12.

We return to the problem of state transformations in the next section, where we consider transformations of state ensembles instead of single states.

B. Complete set of monotones for measurements

A familiar picture of transforming resources involves channels applied to states, transforming one state to another. However, any meaningful information processing task includes a measurement at the end, so it is reasonable to consider states, channels, and measurements as parts of a single, consolidated family. It is therefore insightful to understand channels from an alternative dual perspective: namely, not as operations transforming states but as operations transforming measurements. Motivated by this observation, we extend the above consideration on the transformation of states to the transformation of measurements. We show that a similar reasoning allows for an operational characterization of measurement transformations in the context of state discrimination.

Let $\mathcal{O}_{\mathcal{E}}$ be a convex and closed set of effect-conepreserving unital maps $\mathcal{V}^* \to \mathcal{V}^*$ which furthermore includes the identity map $E \mapsto E$ and is closed under concatenation; i.e., if $\Gamma_1, \Gamma_2 \in \mathcal{O}_{\mathcal{E}}$ then $\Gamma_2 \circ \Gamma_1 \in \mathcal{O}_{\mathcal{E}}$. These assumptions are again particularly natural in the context of a resource theory, but this is not assumed.

We now consider a variant of state discrimination where instead of immediately making a measurement to discriminate the state ensemble, we apply a prior transformation to the measurement effects. We see that by restricting a set of allowed prior operations to the chosen set $\mathcal{O}_{\mathcal{E}}$, the success probability of this task serves as a complete set of monotones for the measurements.

Recalling that any effect-cone-preserving and unital operation has a corresponding dual operation which is normalization- and state-cone-preserving, this can be equivalently understood as a task of state discrimination where the transformations from the set $\{\Lambda | \Lambda^* \in \mathcal{O}_{\mathcal{E}}\}$ are applied to the states. The success probability of the measurement $\mathbb{M} = \{M_i\}_i$ in distinguishing the ensemble $\mathcal{A} = \{p_i, \sigma_i\}$ in this setting is then

$$\tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}) = \max_{\Lambda^* \in \mathcal{O}_{\mathcal{E}}} \sum_{i} p_i \langle M_i, \Lambda(\sigma_i) \rangle.$$
 (123)

We now show that this success probability serves as a set of complete monotones for measurements under the free operations $\mathcal{O}_{\mathcal{E}}$.

Theorem 14: Given measurements $\mathbb{M} = \{M_i\}_{i=0}^N$ and $\mathbb{M}' = \{M_i'\}_{i=0}^N$, there exists $\Gamma \in \mathcal{O}_{\mathcal{E}}$ such that $\mathbb{M}' = \Gamma(\mathbb{M})$ if and only if for all ensembles $\mathcal{A} = \{p_i, \sigma_i\}_{i=0}^N$ it holds that $\tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}) \geq \tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}')$.

Proof.—For one direction, suppose $\mathbb{M}' = \Gamma(\mathbb{M})$. Then for any $\mathcal{A} = \{p_i, \sigma_i\}$,

$$\tilde{p}_{\mathrm{succ}}(\mathcal{A}, \mathbb{M}') = \max_{\Lambda^* \in \mathcal{O}_{\mathcal{E}}} \sum_a p_a \langle \Lambda^* \circ \Gamma(M_a), \sigma_a \rangle \qquad (124)$$

$$\leq \max_{\tilde{\Lambda}^* \in \mathcal{O}_{\mathcal{E}}} \sum_{a} p_a \langle M_a, \tilde{\Lambda}(\sigma_a) \rangle = \tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}), \qquad (125)$$

where the inequality is due to the closedness of $\mathcal{O}_{\mathcal{E}}$ under concatenation. To show the converse, suppose $\forall \, \mathcal{A}, \, \tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}) \geq \tilde{p}_{\text{succ}}(\mathcal{A}, \mathbb{M}')$. It implies that

$$\begin{split} 0 & \leq \min_{\mathcal{A}} \left[\max_{\Lambda^* \in \mathcal{O}_{\mathcal{E}}} \sum_{a} p_a \langle M_a, \Lambda(\sigma_a) \rangle \right. \\ & \left. - \max_{\Delta^* \in \mathcal{O}_{\mathcal{E}}} \sum_{a} p_a \langle M_a', \Delta(\sigma_a) \rangle \right] \end{split} \tag{126}$$

$$\leq \min_{\mathcal{A}} \max_{\Lambda^* \in \mathcal{O}_{\mathcal{E}}} \sum_{a} p_a \langle \Lambda^*(M_a) - M_a', \sigma_a \rangle \qquad (127)$$

$$= \max_{\Lambda^* \in \mathcal{O}_{\mathcal{E}}} \min_{\mathcal{A}} \sum_{a} p_a \langle \Lambda^*(M_a) - M_a', \sigma_a \rangle, \quad (128)$$

where the second inequality is obtained by setting Δ^* being the identity for the second term, and the equality is due to Sion's minimax theorem [132] because of the compactness and convexity of $\mathcal{A}, \mathcal{O}_{\mathcal{E}}$, and the linearity of the objective function with respect to Λ and $p_a\sigma_a$.

We now show that there exists $\Lambda^* \in \mathcal{O}_{\mathcal{E}}$ such that $\Lambda^*(M_a) - M_a' = 0 \ \forall \ a$, thus concluding the proof. To this end, suppose to the contrary that such an operation does not exist. Since Λ^* is unital by assumption, we get for any Λ ,

$$\sum_{a} \Lambda^*(M_a) - M'_a = \Lambda^*(U) - U = 0.$$
 (129)

In particular, it holds that

$$\langle \sum_{a} \Lambda^*(M_a) - M'_a, \omega \rangle = 0 \quad \forall \ \omega \in \mathcal{C}, \quad (130)$$

which implies that we cannot have $\Lambda^*(M_a) - M_a' \in \mathcal{C}^*$ for all a, as this necessarily means that $\Lambda^*(M_a) - M_a'$ are identically zero. Therefore, for any choice of Λ^* there exists an index a^* and a state $\sigma \in \Omega(\mathcal{V})$ such that

$$\langle \Lambda^*(M_{a^*}) - M'_{a^*}, \sigma \rangle < 0. \tag{131}$$

Choosing the ensemble $\{p_a, \sigma_a\}$ such that $p_a = 0$ if $a \neq a^*$ and $p_{a^*} = 1, \sigma_{a^*} = \sigma$ then gives

$$\sum_{a} p_a \langle \Lambda^*(M_a) - M_a', \sigma_a \rangle < 0.$$
 (132)

We have therefore reached a contradiction, as Eq. (128) says that there exists a choice of Λ^* such that any ensemble $\{p_a,\sigma_a\}$ gives $\sum_a p_a \langle \Lambda^*(M_a) - M_a',\sigma_a \rangle \geq 0$. We conclude that our original assumption must be wrong, and therefore there exists $\Gamma = \Lambda^* \in \mathcal{O}_{\mathcal{E}}$ such that $M_a' = \Gamma(M_a) \ \forall \ a$.

By using the aforementioned dual interpretation of this task as operations applied to states rather than measurements, we can additionally obtain a complete set of monotones for the transformations between state *ensembles*. Such tasks have been considered in quantum information theory in different contexts [46,120,123,126,137,138] and indeed find use in several resource theories which employ generalizations of majorization [121,124,128]. To this end, we consider two different types of tasks: one, the conclusive state discrimination just as above, with probability of success given by

$$\tilde{p}_{\text{succ}}(\{p_a, \sigma_a\}, \{M_a\}) = \max_{\Lambda \in \mathcal{O}} \sum_{a} p_a \langle M_a, \Lambda(\sigma_a) \rangle \quad (133)$$

with \mathcal{O} being a set of operations defined as in Sec. VII A, and two, the inconclusive state discrimination task characterized by

$$\tilde{p}'_{\text{succ}}(\{p_a, \sigma_a\}_{a=0}^{N-1}, \{M_a\}_{a=0}^{N}) = \max_{\Lambda \in \mathcal{O}} \sum_{a=0}^{N-1} p_a \langle M_a, \Lambda(\sigma_a) \rangle.$$
(134)

The following is then a straightforward adaptation of the concepts of Theorem 14.

Corollary 15: Let $\{\sigma_i\}_{i=0}^{N-1}$ and $\{\sigma_i'\}_{i=0}^{N-1}$ be two collections of states. Then, there exists $\Lambda \in \mathcal{O}$ such that $\sigma_i' = \Lambda(\sigma_i) \ \forall \ i$ if and only if either of the following conditions is satisfied:

- (i) $N \ge 2$ and for all N-outcome measurements \mathbb{M} and all probability distributions $\{p_i\}_{i=0}^{N-1}$ it holds that $\tilde{p}_{\text{succ}}(\{p_i, \sigma_i\}, \mathbb{M}) \ge \tilde{p}_{\text{succ}}(\{p_i, \sigma_i'\}, \mathbb{M})$,
- (ii) $N \ge 1$ and for all (N+1)-outcome measurements $\mathbb{M} = \{M_i\}_{i=0}^N$ it holds that $\tilde{p}'_{\text{succ}}(\{p_i, \sigma_i\}, \mathbb{M}) \ge \tilde{p}'_{\text{succ}}(\{p_i, \sigma_i'\}, \mathbb{M})$, where $\{p_i\}_{i=0}^{N-1}$ is any fixed probability distribution such that $p_i > 0 \ \forall i$, which in particular can be taken to be $p_i = 1/N$.

We include the full proof in Appendix A for completeness. We remark that for N = 1, we recover Corollary 11.

VIII. CONCLUSIONS

We provided a general operational characterization of quantification and manipulation of resources, the two core concepts of resource theories, in terms of the performance of state and channel discrimination tasks. The generality of our work is threefold: our formulations encompass general convex resource theories associated with general types of resource objects (states, measurements, and channels), and major parts of the results are valid for general probabilistic theories beyond quantum mechanics. In particular, we found that robustness measures play central roles in bridging the quantification of resources and the success probability in discrimination tasks, specifically establishing that the maximum advantages in classes of discrimination tasks realized by resource objects are exactly quantified by the corresponding robustness measures. We also characterized the manipulation of resources associated with states and measurements by considering families of discrimination tasks where their success probabilities serve as complete sets of monotones that fully characterize the transformation of resources under free operations. In the case of quantum mechanics, we further extended the above connections between discrimination tasks and resource-theoretic concepts to single-shot information theory.

In addition to providing fundamental insights spanning a broad class of physical theories, the results are immediately applicable to a wide range of physical resources in quantum-information theory. The resource theories of coherence, entanglement, magic, athermality, asymmetry, as well as their generalizations in the form of multipartite entanglement or multilevel entanglement and coherence all fit the

framework introduced herein, and therefore, all of our results apply to them immediately. In the case of measurements, many significant insights can be gained from studying classes of measurements such as separable, PPT, incoherent, or Pauli measurements, all of which are again special cases of the resource theories considered in this work. In the characterization of channels, we obtain results applicable on the one hand to sets of free operations in the aforementioned state-based resource theories, and on the other hand obtain an operational characterization of quantum channels which can be applied in the study of arbitrary channel-based resource theories such as the resource theory of quantum memories (non-entanglement-breaking channels).

Our results furthermore reveal interesting connections between the quantification of resources and the phenomenon of data hiding. Although data hiding has been mostly discussed in the context of the theory of entanglement, where one compares the capability of local operations and classical communication (LOCC) measurements (or other restricted sets such as separable or PPT measurements) to that of arbitrary measurements for state discrimination, one could consider more general data-hiding procedures depending on the given physical setting. For instance, in the scenario where only one party has the ability to produce magic (so-called "magic factory"), it is sensible to encode information in a way that Pauli measurements have less capability of decoding it than arbitrary measurements, with the data-hiding ratio characterized by the difference between the capabilities of these two sets of measurements in state discrimination. One can further think of various other scenarios such as (i) one party having access to a restricted but larger set of measurements than the other (not necessarily arbitrary measurements), (ii) one party having more access to resource states or channels (not measurements) than the other, in which the data could be encoded in the form of channel discrimination tasks, (iii) or the encoded data requiring discrimination tasks more intricate than the standard binary discrimination. Our general formulations encompass such variants, allowing for considerable flexibility of the encoding strategies.

The generality of the results also provides insights into the foundation of quantum mechanics. As we show in this work, the operational advantage realized by any type of resource object is not unique to quantum mechanics but rather a universal phenomenon stemming solely from the convexity of the underlying cones and thus shared by general GPTs. Our results in particular imply that there is no separation between the generalized robustness (*a priori* a geometric concept) and the advantage provided in the considered classes of discrimination tasks (explicitly operational tasks) in any GPT; therefore, one cannot hope to separate a given theory from quantum mechanics by finding a gap between these quantities. Our results additionally provide an experimentally accessible way of bounding the geometric resource measures as well as characterizing

resource transformations in any GPT by relating them with discrimination tasks.

An interesting problem we leave for future work is to give general operational meaning to standard robustness of measurements and channels, which would solidify the fundamental operational significance of the standard robustness measure alongside that of the generalized robustness as established in this work. In light of the series of results obtained in this work, it can be expected that discrimination tasks are also suitable for characterizing these measures at a high level of generality. Additionally, it remains to understand whether general quantitative relations between the generalized and standard robustness measures can be found and whether there exists a way to generalize our results to more members of the robustness family besides the two we considered. We also remark that this work raises an interesting question about a unified understanding of different operational tasks via robustness measures: in addition to the discrimination tasks studied in this work, robustness measures have appeared in very different operational contexts, albeit mostly in a resourcespecific fashion. One could then speculate that these operational settings may be deeply connected by a more fundamental class of tasks whose performance is somehow characterized by the robustness measures.

Finally, in addition to the transformation of states and measurements, one could consider the transformation of channels realized by *superchannels* [26,40,73,133, 139–141]. The full understanding of superchannels is still on the way, and we hope that our results will help propel this journey forward from the perspective of general resource theories.

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Note added.— Recently, we became aware of the independent related works by Uola et al. [142] as well as by Oszmaniec and Biswas [143], which obtained results similar to Theorem 2 on the relation between state discrimination tasks and quantification of resources associated with measurements within quantum mechanics.

APPENDIX A: PROOFS OF RESULTS IN SEC. VII

Theorem 12: There exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$ if and only if either of the following conditions is satisfied:

- (i) it holds that $\tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega) \ge \tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega')$ for all channel ensembles $\{p_i, \Lambda_i\}$ and measurements $\{M_i\}$,
- (ii) for a fixed set of channels $\{\Lambda_i\}_{i=0}^{N-1}$ containing the identity channel id, it holds that $\tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega) \geq \tilde{p}_{\text{succ}}(\{p_i\}, \{\Lambda_i\}, \{M_i\}, \omega')$ for all probability distributions $\{p_i\}_{i=0}^{N-1}$ and measurements $\{M_i\}_{i=0}^{N-1}$.

Proof.—As before, in the proof we optimize over all sets of channels $\{\Lambda_i\}$ [case (i)], but one could equivalently consider a fixed choice of $\{\Lambda_i\}$ setting $\Lambda_0 = \text{id}$ [case (ii)] and the proof proceeds the same.

For the "only if" direction, suppose $\omega' = \Lambda(\omega)$. Then for any $\{p_i, \Lambda_i\}$ and $\{M_i\}$,

$$\begin{split} \tilde{p}_{\text{succ}}(\{p_i, \Lambda_i\}, \{M_i\}, \omega') \\ &= \max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i \circ \Xi(\omega') \rangle \end{split} \tag{A1}$$

$$= \max_{\Xi \in \mathcal{O}} \sum_{i} p_{i} \langle M_{i}, \Lambda_{i} \circ \Xi \circ \Lambda(\omega) \rangle \tag{A2}$$

$$\leq \max_{\Lambda' \in \mathcal{O}} \sum_{i} p_{i} \langle M_{i}, \Lambda_{i} \circ \Lambda'(\omega) \rangle$$

$$= \tilde{p}_{\text{succ}}(\{p_{i}, \Lambda_{i}\}, \{M_{i}\}, \omega), \tag{A3}$$

where the inequality is due to the closedness of \mathcal{O} under concatenation.

On the other hand, assuming $\tilde{p}_{\text{succ}}(\{p_i, \Lambda_i\}, \{M_i\}, \omega) \ge \tilde{p}_{\text{succ}}(\{p_i, \Lambda_i\}, \{M_i\}, \omega')$ holds for all $\{p_i, \Lambda_i\}$ and $\{M_i\}$ implies

$$0 \le \inf_{\{p_i, \Lambda_i\}, \{M_i\}} \left[\max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i \circ \Xi(\omega) \rangle - \max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i \circ \Xi(\omega') \rangle \right]$$
(A4)

$$\leq \inf_{\{p_i, \Lambda_i\}, \{M_i\}} \max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i(\Xi(\omega) - \omega') \rangle \tag{A5}$$

$$\leq \min_{\{p_i, \Lambda_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^{N-1}} \max_{\Xi \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda_i(\Xi(\omega) - \omega') \rangle \qquad (A6)$$

$$= \max_{\Xi \in \mathcal{O}} \min_{\{p_i, \Lambda_i\}_{i=0}^{N-1}, \{M_i\}_{i=0}^{N-1}} \sum_i p_i \langle M_i, \Lambda_i(\Xi(\omega) - \omega') \rangle, \quad (A7)$$

where the second inequality is obtained by setting Ξ as the identity for the second term, the third inequality is because we restricted the minimization over the *N*-element sets where $N \ge 2$ is an arbitrary integer, and the equality is due

to the minimax theorem because of the convexity and compactness of the sets of channels, measurements, and \mathcal{O} and the linearity of the objective function with respect to Ξ , $p_i\Lambda_i$, and M_i .

Suppose now that there does not exist $\Lambda \in \mathcal{O}$ such that $\Lambda(\omega) = \omega'$. Since each $\Xi \in \mathcal{O}$ preserves the normalization of states, we have $\langle U, \Xi(\omega) - \omega' \rangle = 0$ which implies that $\Xi(\omega) - \omega' \notin \mathcal{C}$, and so there exists an effect E such that $\langle E, \Xi(\omega) - \omega' \rangle < 0$. Take N = 2 in Eq. (A6) and define the measurement $\{M_0, M_1\} = \{E, U - E\}$ and the channel ensemble defined by $p_0 = 1$, $\Lambda_0 = \operatorname{id}$, and $p_1 = 0$, $\Lambda_1 = \Lambda'$ where id denotes the identity map and Λ' is an arbitrary channel. We then get

$$\sum_{i} p_{i} \langle M_{i}, \Lambda_{i}(\Xi(\omega) - \omega') \rangle = \langle E, \Xi(\omega) - \omega' \rangle < 0$$
 (A8)

for any Ξ , which contradicts Eq. (A7), and so there exists $\Lambda \in \mathcal{O}$ such that $\omega' = \Lambda(\omega)$.

Corollary 15: Let $\{\sigma_i\}_{i=0}^{N-1}$ and $\{\sigma_i'\}_{i=0}^{N-1}$ be two collections of states. Then, there exists $\Lambda \in \mathcal{O}$ such that $\sigma_i' = \Lambda(\sigma_i) \ \forall \ i$ if and only if either of the following conditions is satisfied:

- (i) N ≥ 2 and for all N-outcome measurements M and all probability distributions {p_i}_{i=0}^{N-1} it holds that p̃_{succ}({p_i, σ_i}, M) ≥ p̃_{succ}({p_i, σ_i'}, M),
 (ii) N ≥ 1 and for all (N + 1)-outcome measurements
- (ii) $N \ge 1$ and for all (N+1)-outcome measurements $\mathbb{M} = \{M_i\}_{i=0}^N$ it holds that $\tilde{p}'_{\text{succ}}(\{p_i, \sigma_i\}, \mathbb{M}) \ge \tilde{p}'_{\text{succ}}(\{p_i, \sigma_i'\}, \mathbb{M})$, where $\{p_i\}_{i=0}^{N-1}$ is any fixed probability distribution such that $p_i > 0 \ \forall i$, which in particular can be taken to be $p_i = 1/N$.

Proof.—The "only if" part follows analogously to Theorem 9. For the other implication, assume the desired operation $\Lambda \in \mathcal{O}$ does not exist, and consider the case (i) first. Notice that if $\tilde{p}_{\text{succ}}(\{p_i,\sigma_i\},\mathbb{M}) \geq \tilde{p}_{\text{succ}}(\{p_i,\sigma_i'\},\mathbb{M})$ for all N-outcome measurements and all probability distributions, then

$$0 \leq \min_{\{M_i\}, \{p_i\}} \left[\max_{\Lambda \in \mathcal{O}} \sum_{i} p_i \langle M_i, \Lambda(\sigma_i) \rangle - \max_{\Theta \in \mathcal{O}} \sum_{i} p_i \langle M_i, \Theta(\sigma'_i) \rangle \right]$$
(A9)

$$\leq \min_{\{M_i\},\{p_i\}} \max_{\Lambda \in \mathcal{O}} \sum_i p_i \langle M_i, \Lambda(\sigma_i) - \sigma_i' \rangle \qquad (A10)$$

$$= \max_{\Lambda \in \mathcal{O}} \min_{\{M_i\}, \{p_i\}} \sum_i p_i \langle M_i, \Lambda(\sigma_i) - \sigma_i' \rangle \qquad (\text{A11})$$

by Sion's minimax theorem. But since for any $\Lambda\in\mathcal{O}$ we have

$$\langle U, \Lambda(\sigma_i) - \sigma'_i \rangle = 0 \quad \forall i,$$
 (A12)

and $\Lambda(\sigma_i) - \sigma_i'$ are not uniformly 0, this means that there must exist an index i^* such that $\Lambda(\sigma_{i^*}) - \sigma_{i^*}' \notin \mathcal{C}$, and so

there exists an effect E such that $\langle E, \Lambda(\sigma_{i^*}) - \sigma'_{i^*} \rangle < 0$. Choosing an arbitrary index $i^{\diamond} \in \{0, ..., N-1\} \setminus \{i^*\}$, where such a choice is guaranteed to exist because $N \geq 2$ by assumption, we see that the choice of measurement $\mathbb{M} = \{M_i\}_{i=0}^{N-1}$ as $M_{i^*} = E/\|E\|_{\Omega}^{\circ}$, $M_{i^{\diamond}} = U - E/\|E\|_{\Omega}^{\circ}$, $M_i = 0 \ \forall \ i \in \{0, ..., N-1\} \setminus \{i^*, i^{\diamond}\}$ together with the probability distribution $\{p_i\}$ defined as $p_{i^*} = 1$, $p_i = 0 \ \forall \ i \neq i^*$ contradicts Eq. (A11).

Similarly, in case (ii) we get that

$$0 \le \max_{\Lambda \in \mathcal{O}} \min_{\{M_i\}_{i=0}^N} \sum_i p_i \langle M_i, \Lambda(\sigma_i) - \sigma_i' \rangle$$
 (A13)

for a fixed probability distribution, and choosing the measurement $\mathbb{M}=\{M_i\}_{i=0}^N$ as $M_{i^\star}=E/\|E\|_\Omega^\circ$, $M_i=0 \ \forall \ i\in \{0,...,N-1\}\backslash\{i^\star\}$, with $M_N=U-E/\|E\|_\Omega^\circ$ completes the proof.

APPENDIX B: DUALITY IN CONIC OPTIMIZATION

For completeness, we include a derivation of the dual form of the optimization problems which we employ throughout the manuscript. This section is based on standard arguments found, e.g., in Refs. [105,144].

Consider first some real complete normed vector spaces W, W' and the optimization problem whose optimal value is given by

$$\inf \{ \langle A, x \rangle | \Lambda(x) = y, x \in \mathcal{K} \}, \tag{B1}$$

where $A \in \mathcal{W}^*$, $y \in \mathcal{W}'$ are given, $\Lambda: \mathcal{W} \to \mathcal{W}'$ is some linear function, and $\mathcal{K} \subseteq \mathcal{W}$ is a closed and convex cone. All of the optimization problems considered in this work can be expressed in this form (and indeed, so can any convex optimization problem over a closed and convex set), which we see explicitly.

We refer to the above as the *primal* problem and to the set $\{x \in \mathcal{W} | \Lambda(x) = y, x \in \mathcal{K}\}$ as the feasible set. Define the Lagrangian

$$L(x; Q, Z) \coloneqq \langle A, x \rangle - \langle Z, \Lambda(x) - y \rangle - \langle Q, x \rangle, \quad (B2)$$

where $Q \in \mathcal{W}^*, Z \in \mathcal{W}'^*$ are the so-called Lagrange multipliers. Notice now that for every x such that $\Lambda(x) - y \neq 0$, there must exist a choice of $Z \in \mathcal{W}'^*$ such that $\langle Z, \Lambda(x) - y \rangle < 0$; similarly, by the strongly separating hyperplane theorem [97], for any $x \notin \mathcal{K}$ there will exist a choice of $Q \in \mathcal{K}^*$ such that $\langle Q, x \rangle < 0$, where $\mathcal{K}^* = \{Y \in \mathcal{W}^* | \langle Y, k \rangle \geq 0 \ \forall \ k \in \mathcal{K} \}$ is the cone dual to \mathcal{K} . This allows us to write

$$\sup_{\substack{Q \in \mathcal{K}^* \\ Z \in \mathcal{W}^*}} L(x; Q, Z) = \begin{cases} \langle A, x \rangle & \text{if } \Lambda(x) = y \text{ and } x \in \mathcal{K}, \\ \infty & \text{otherwise,} \end{cases}$$
(B3)

which in particular means that

$$p = \inf_{x \in \mathcal{W}} \sup_{\substack{Q \in \mathcal{K}^* \\ Z \in \mathcal{W}^{f_e}}} L(x; Q, Z).$$
 (B4)

The *dual* problem is then defined by interchanging the minimization and maximization in the above:

$$d := \sup_{\substack{Q \in \mathcal{K}^* \\ Z \cap \mathcal{M}^*}} \inf_{x \in \mathcal{W}} L(x; Q, Z). \tag{B5}$$

Noticing that

$$\inf_{x \in \mathcal{W}} L(x; Q, Z) = \inf_{x \in \mathcal{W}} \langle A - \Lambda^*(Z) - Q, x \rangle + \langle Z, y \rangle$$

$$= \begin{cases} \langle Z, y \rangle & \text{if } A - \Lambda^*(Z) - Q = 0, \\ -\infty & \text{otherwise,} \end{cases}$$
(B6)

since $\langle A - \Lambda^*(Z) - Q, x \rangle$ is a linear function of x, we can equivalently write

$$d = \sup \{ \langle Z, y \rangle | A - \Lambda^*(Z) - Q = 0, Q \in \mathcal{K}^* \}$$

= \sup \{ \langle Z, y \rangle | A - \Lambda^*(Z) \in \mathcal{K}^* \rangle. (B7)

We often refer to this form as the dual form of the primal optimization problem p.

It is not difficult to see that $p \ge d$ in general, a fact often called weak Lagrange duality. Crucially, Slater's theorem (see, e.g., Refs. [105,145]) states that if there exists a feasible solution x such that x is in the (relative) interior of \mathcal{K} —called a strictly feasible solution—then p=d. We refer to this property as Slater's condition, and the equivalence between the primal and dual problems as strong Lagrange duality.

Consider first the generalized robustness of states, which can be expressed as

$$R_{\mathcal{F}}(\omega) + 1 = \min\{\langle U, \sigma \rangle | \sigma - \omega \in \mathcal{C}, \sigma \in \text{cone}(\mathcal{F})\}.$$
 (B8)

Note that Eq. (B1) is reduced to Eq. (B8) by choosing $W = V \oplus V$, W' = V, $K = \operatorname{cone}(\mathcal{F}) \oplus \mathcal{C}$, $\Lambda(x_1 \oplus x_2) = x_1 - x_2$, $A = U \oplus 0$ and $y = \omega$; analogous forms can be obtained for the other considered measures. Writing the Lagrangian as $L(\omega; X, Z) = \langle U, \sigma \rangle - \langle X, \sigma - \omega \rangle - \langle Z, \sigma \rangle$, following the steps above we straightforwardly obtain the dual as

$$\max\{\langle X, \omega \rangle | X \in \mathcal{C}^*, U - X \in \mathcal{F}^*\}, \tag{B9}$$

as announced in Eq. (11). Analogously, the standard robustness $R_{\mathcal{F}}^{\mathcal{F}}$ can be obtained by changing the constraint $\sigma - \omega \in \mathcal{C}$ to $\sigma - \omega \in \text{cone}(\mathcal{F})$.

In the case of the robustness of measurement, for any (N+1)-outcome measurement $\mathbb{M} = \{M_i\}_{i=0}^N$, we can write

$$\begin{split} R_{\mathcal{E}_{\mathcal{F}}}(\mathbb{M}) &= \min \Big\{ \lambda \mid M_i + N_i \in \mathcal{E}_{\mathcal{F}} \ \forall \ i, N_i \in \mathcal{C}^* \ \forall \ i, \\ \lambda U - \sum_i N_i &= 0_{\mathcal{V}^*} \Big\}, \end{split} \tag{B10}$$

which gives the Lagrangian as

$$L(\{N_i\}, \lambda; \{\sigma_i\}, \{\delta_i\}, \eta)$$

$$= \lambda - \sum_{i} \langle M_i + N_i, \sigma_i \rangle - \sum_{i} \langle N_i, \delta_i \rangle - \langle \lambda U - \sum_{i} N_i, \eta \rangle$$

$$= \lambda (1 - \langle U, \eta \rangle) + \sum_{i} \langle N_i, \eta - \sigma_i - \delta_i \rangle - \sum_{i} \langle M_i, \sigma_i \rangle.$$
(B11)

Optimizing over the Lagrange multipliers $\sigma_i \in \mathcal{E}_{\mathcal{F}}^*$, $\delta_i \in \mathcal{C}$, and $\eta \in \{0_{\mathcal{V}^*}\}^* = \mathcal{V}$, we get the dual as in Eqs. (24)–(28).

In the case of the robustness of channels, we have the problem as

$$R_{\mathcal{O}_{\mathcal{F}}}(\Lambda) = \min\{\lambda \mid J_{\Xi} \in \operatorname{cone}(\mathcal{O}_{\mathcal{F}}^{J}), J_{\Xi} - J_{\Lambda} \succeq 0,$$
$$\operatorname{Tr}_{\mathcal{V}} J_{\Xi} - (1 + \lambda) \mathbb{I}_{\mathcal{V}} = 0_{\mathcal{V}}\}.$$
(B12)

Writing the Lagrangian as

$$\begin{split} L(J_\Xi,\lambda;X,Y,Z) &= \lambda - \langle Z,J_\Xi \rangle - \langle Y,J_\Xi - J_\Lambda \rangle \\ &- \langle X,(1+\lambda)\mathbb{I}_{\mathcal{V}} - \mathrm{Tr}_{\mathcal{V}}J_\Xi \rangle \\ &= \lambda(1-\langle X,\mathbb{I}_{\mathcal{V}} \rangle) - \langle Z+Y-X \otimes \mathbb{I}_{\mathcal{V}},J_\Xi \rangle \\ &+ \langle Y,J_\Lambda \rangle - \langle X,\mathbb{I}_{\mathcal{V}} \rangle, \end{split} \tag{B13}$$

where we used that $\operatorname{Tr}_{\mathcal{V}}(\cdot)^* = \cdot \otimes \mathbb{I}_{\mathcal{V}}$, an optimization over the Lagrange coefficients $X \in \mathcal{V}$, $Y \succeq 0$, and $Z \in \mathcal{O}_{\mathcal{F}}^{J^*}$ gives the desired dual problem (59)–(62).

- J. P. Dowling and G. J. Milburn, *Quantum Technology: The Second Quantum Revolution*, Phil. Trans. R. Soc. A 361, 1655 (2003).
- [2] A. Acín, I. Bloch, H. Buhrman, T. Calarco, C. Eichler, J. Eisert, D. Esteve, N. Gisin, S. J. Glaser, F. Jelezko, S. Kuhr, M. Lewenstein, M. F. Riedel, P. O. Schmidt, R. Thew, A. Wallraff, I. Walmsley, and F. K. Wilhelm, *The Quantum Technologies Roadmap: A European Community View*, New J. Phys. 20, 080201 (2018).
- [3] E. Chitambar and G. Gour, *Quantum Resource Theories*, Rev. Mod. Phys. **91**, 025001 (2019).
- [4] M. B. Plenio and S. Virmani, An Introduction to Entanglement Measures, Quantum Inf. Comput. 7, 001 (2007).
- [5] R. Horodecki, P. Horodecki, M. Horodecki, and K. Horodecki, *Quantum Entanglement*, Rev. Mod. Phys. **81**, 865 (2009).
- [6] J. Åberg, *Quantifying Superposition*, arXiv:quant-ph/0612146.

- [7] T. Baumgratz, M. Cramer, and M. B. Plenio, *Quantifying Coherence*, Phys. Rev. Lett. **113**, 140401 (2014).
- [8] A. Streltsov, G. Adesso, and M. B. Plenio, *Colloquium: Quantum Coherence as a Resource*, Rev. Mod. Phys. 89, 041003 (2017).
- [9] G. Gour and R. W. Spekkens, The Resource Theory of Quantum Reference Frames: Manipulations and Monotones, New J. Phys. 10, 033023 (2008).
- [10] I. Marvian and R. W. Spekkens, How to Quantify Coherence: Distinguishing Speakable and Unspeakable Notions, Phys. Rev. A 94, 052324 (2016).
- [11] F. G. S. L. Brandão, M. Horodecki, J. Oppenheim, J. M. Renes, and R. W. Spekkens, *Resource Theory of Quantum States Out of Thermal Equilibrium*, Phys. Rev. Lett. 111, 250404 (2013).
- [12] F. Brandão, M. Horodecki, N. Ng, J. Oppenheim, and S. Wehner, *The Second Laws of Quantum Thermodynamics*, Proc. Natl. Acad. Sci. U.S.A. 112, 3275 (2015).
- [13] R. Gallego and L. Aolita, Resource Theory of Steering, Phys. Rev. X 5, 041008 (2015).
- [14] A. Rivas, S. F. Huelga, and M. B. Plenio, *Entanglement and Non-Markovianity of Quantum Evolutions*, Phys. Rev. Lett. **105**, 050403 (2010).
- [15] S. Bhattacharya, B. Bhattacharya, and A. Majumdar, Resource Theory of Non-Markovianity: A Thermodynamic Perspective, arXiv:1803.06881.
- [16] E. Wakakuwa, Operational Resource Theory of Non-Markovianity, arXiv:1709.07248.
- [17] V. Veitch, S. A. H. Mousavian, D. Gottesman, and J. Emerson, *The Resource Theory of Stabilizer Quantum Computation*, New J. Phys. **16**, 013009 (2014).
- [18] M. Howard and E. Campbell, *Application of a Resource Theory for Magic States to Fault-Tolerant Quantum Computing*, Phys. Rev. Lett. **118**, 090501 (2017).
- [19] M. G. Genoni, M. G. A. Paris, and K. Banaszek, *Quantifying the Non-Gaussian Character of a Quantum State by Quantum Relative Entropy*, Phys. Rev. A **78**, 060303(R) (2008).
- [20] R. Takagi and Q. Zhuang, Convex Resource Theory of Non-Gaussianity, Phys. Rev. A 97, 062337 (2018).
- [21] F. Albarelli, M. G. Genoni, M. G. A. Paris, and A. Ferraro, Resource Theory of Quantum Non-Gaussianity and Wigner Negativity, Phys. Rev. A 98, 052350 (2018).
- [22] M. Oszmaniec, L. Guerini, P. Wittek, and A. Acín, Simulating Positive-Operator-Valued Measures with Projective Measurements, Phys. Rev. Lett. 119, 190501 (2017).
- [23] M. Oszmaniec, F. B. Maciejewski, and Z. Puchała, All Quantum Measurements Can Be Simulated Using Projective Measurements and Postselection, Phys. Rev. A 100, 012351 (2019).
- [24] L. Guerini, J. Bavaresco, M. T. Cunha, and A. Acín, *Operational Framework for Quantum Measurement Simulability*, J. Math. Phys. (N.Y.) **58**, 092102 (2017).
- [25] P. Skrzypczyk and N. Linden, *Robustness of Measurement, Discrimination Games, and Accessible Information*, Phys. Rev. Lett. **122**, 140403 (2019).
- [26] D. Rosset, F. Buscemi, and Y.-C. Liang, Resource Theory of Quantum Memories and Their Faithful Verification with Minimal Assumptions, Phys. Rev. X 8, 021033 (2018).

- [27] M. Horodecki and J. Oppenheim, (*Quantumness in the Context of*) Resource Theories, Int. J. Mod. Phys. B **27**, 1345019 (2013).
- [28] F. G. S. L. Brandão and G. Gour, Reversible Framework for Quantum Resource Theories, Phys. Rev. Lett. 115, 070503 (2015).
- [29] L. Del Rio, L. Kraemer, and R. Renner, Resource Theories of Knowledge, arXiv:1511.08818.
- [30] B. Coecke, T. Fritz, and R. W. Spekkens, A Mathematical Theory of Resources, Inf. Comput. 250, 59 (2016).
- [31] Z.-W. Liu, X. Hu, and S. Lloyd, *Resource Destroying Maps*, Phys. Rev. Lett. **118**, 060502 (2017).
- [32] G. Gour, Quantum Resource Theories in the Single-Shot Regime, Phys. Rev. A 95, 062314 (2017).
- [33] A. Anshu, M.-H. Hsieh, and R. Jain, *Quantifying Resources in General Resource Theory with Catalysts*, Phys. Rev. Lett. **121**, 190504 (2018).
- [34] B. Regula, Convex Geometry of Quantum Resource Quantification, J. Phys. A 51, 045303 (2018).
- [35] L. Lami, B. Regula, X. Wang, R. Nichols, A. Winter, and G. Adesso, *Gaussian Quantum Resource Theories*, Phys. Rev. A 98, 022335 (2018).
- [36] R. Takagi, B. Regula, K. Bu, Z.-W. Liu, and G. Adesso, *Operational Advantage of Quantum Resources in Sub*channel Discrimination, Phys. Rev. Lett. **122**, 140402 (2019).
- [37] L. Li, K. Bu, and Z.-W. Liu, Quantifying the Resource Content of Quantum Channels: An Operational Approach, arXiv:1812.02572.
- [38] S. Pirandola, R. Laurenza, C. Ottaviani, and L. Banchi, *Fundamental Limits of Repeaterless Quantum Communications*, Nat. Commun. **8**, 15043 (2017).
- [39] K. B. Dana, M. G. Díaz, M. Mejatty, and A. Winter, Resource Theory of Coherence: Beyond States, Phys. Rev. A 95, 062327 (2017).
- [40] G. Gour, Comparison of Quantum Channels by Superchannels, IEEE Trans. Inf. Theory **65**, 5880 (2019).
- [41] M. M. Wilde, Entanglement Cost and Quantum Channel Simulation, Phys. Rev. A 98, 042338 (2018).
- [42] M. G. Díaz, K. Fang, X. Wang, M. Rosati, M. Skotiniotis, J. Calsamiglia, and A. Winter, *Using and Reusing Coherence to Realize Quantum Processes*, Quantum **2**, 100 (2018).
- [43] T. Theurer, D. Egloff, L. Zhang, and M. B. Plenio, *Quantifying Operations with an Application to Coherence*, Phys. Rev. Lett. **122**, 190405 (2019).
- [44] Q. Zhuang, P. W. Shor, and J. H. Shapiro, *Resource Theory of Non-Gaussian Operations*, Phys. Rev. A **97**, 052317 (2018).
- [45] J. R. Seddon and E. Campbell, Quantifying Magic for Multi-Qubit Operations, Proc. R. Soc. A, 475, 20190251 (2019).
- [46] T. Heinosaari, J. Kiukas, and D. Reitzner, *Noise Robustness of the Incompatibility of Quantum Measurements*, Phys. Rev. A **92**, 022115 (2015).
- [47] E. Haapasalo, Robustness of Incompatibility for Quantum Devices, J. Phys. A 48, 255303 (2015).
- [48] C. Carmeli, T. Heinosaari, and A. Toigo, *State Discrimination with Postmeasurement Information and Incompatibility of Quantum Measurements*, Phys. Rev. A **98**, 012126 (2018).

- [49] G. Ludwig, An Axiomatic Basis for Quantum Mechanics: Volume 1 Derivation of Hilbert Space Structure (Springer-Verlag, Berlin, 1985).
- [50] Foundations of Quantum Mechanics and Ordered Linear Spaces, edited by A. Hartkämper and H. Neumann (Springer, New York, 1974).
- [51] E. B. Davies and J. T. Lewis, An Operational Approach to Quantum Probability, Commun. Math. Phys. 17, 239 (1970).
- [52] L. Lami, Non-Classical Correlations in Quantum Mechanics and Beyond, Ph.D. thesis, Universitat Autònoma de Barcelona, 2018.
- [53] J. Barrett, Information Processing in Generalized Probabilistic Theories, Phys. Rev. A 75, 032304 (2007).
- [54] G. Chiribella, G. M. D'Ariano, and P. Perinotti, *Probabilistic Theories with Purification*, Phys. Rev. A **81**, 062348 (2010).
- [55] G. Chiribella, G. M. D'Ariano, and P. Perinotti, *Informational Derivation of Quantum Theory*, Phys. Rev. A 84, 012311 (2011).
- [56] H. Barnum and A. Wilce, *Information Processing in Convex Operational Theories*, Electron. Notes Theor. Comput. Sci. **270**, 3 (2011).
- [57] L. Masanes and M. P. Müller, A Derivation of Quantum Theory from Physical Requirements, New J. Phys. 13, 063001 (2011).
- [58] H. Barnum, M. P. Müller, and C. Ududec, *Higher-Order Interference and Single-System Postulates Characterizing Quantum Theory*, New J. Phys. **16**, 123029 (2014).
- [59] C. M. Lee and J. Barrett, Computation in Generalised Probabilistic Theories, New J. Phys. 17, 083001 (2015).
- [60] C. M. Lee and J. H. Selby, A No-Go Theorem for Theories That Decohere to Quantum Mechanics, Proc. R. Soc. A 474, 20170732 (2018).
- [61] G. Chiribella and C. M. Scandolo, *Entanglement and Thermodynamics in General Probabilistic Theories*, New J. Phys. **17**, 103027 (2015).
- [62] G. Chiribella and C. M. Scandolo, *Operational Axioms for Diagonalizing States*, Electron. Proc. Theor. Comput. Sci. 195, 96 (2015).
- [63] G. Chiribella and C.M. Scandolo, *Entanglement as an Axiomatic Foundation for Statistical Mechanics*, arXiv: 1608.04459.
- [64] G. Chiribella and C.M. Scandolo, *Microcanonical Thermodynamics in General Physical Theories*, New J. Phys. **19**, 123043 (2017).
- [65] C. M. Scandolo, Information-Theoretic Foundations of Thermodynamics in General Probabilistic Theories, Ph. D. thesis, University of Oxford, 2018.
- [66] A. Holevo, Statistical Decision Theory for Quantum Systems, J. Multivariate Anal. 3, 337 (1973).
- [67] C. W. Helstrom, *Quantum Detection and Estimation Theory* (Academic Press, New York, 1976).
- [68] A. Y. Kitaev, Quantum Computations: Algorithms and Error Correction, Russ. Math. Surv. 52, 1191 (1997).
- [69] A. Chefles, *Quantum State Discrimination*, Contemp. Phys. **41**, 401 (2000).
- [70] A. M. Childs, J. Preskill, and J. Renes, *Quantum Infor-mation and Precision Measurement*, J. Mod. Opt. 47, 155 (2000).

- [71] A. Acin, Statistical Distinguishability between Unitary Operations, Phys. Rev. Lett. 87, 177901 (2001).
- [72] J. Bae and L.-C. Kwek, Quantum State Discrimination and Its Applications, J. Phys. A 48, 083001 (2015).
- [73] A. Jenčová, Base Norms and Discrimination of Generalized Quantum Channels, J. Math. Phys. (N.Y.) 55, 022201 (2014).
- [74] J. Watrous, *The Theory of Quantum Information* (Cambridge University Press, Cambridge, England, 2018).
- [75] G. Kimura, K. Nuida, and H. Imai, *Distinguishability Measures and Entropies for General Probabilistic Theories*, Rep. Math. Phys. **66**, 175 (2010).
- [76] J. Bae, D.-G. Kim, and L.-C. Kwek, Structure of Optimal State Discrimination in Generalized Probabilistic Theories, Entropy 18, 39 (2016).
- [77] L. Lami, C. Palazuelos, and A. Winter, *Ultimate Data Hiding in Quantum Mechanics and Beyond*, Commun. Math. Phys. **361**, 661 (2018).
- [78] M. Piani and J. Watrous, All Entangled States Are Useful for Channel Discrimination, Phys. Rev. Lett. 102, 250501 (2009).
- [79] M. Piani and J. Watrous, Necessary and Sufficient Quantum Information Characterization of Einstein-Podolsky-Rosen Steering, Phys. Rev. Lett. 114, 060404 (2015).
- [80] J. Bae, D. Chruściński, and M. Piani, *More Entanglement Implies Higher Performance in Channel Discrimination Tasks*, Phys. Rev. Lett. **122**, 140404 (2019).
- [81] C. Napoli, T. R. Bromley, M. Cianciaruso, M. Piani, N. Johnston, and G. Adesso, *Robustness of Coherence: An Operational and Observable Measure of Quantum Coherence*, Phys. Rev. Lett. 116, 150502 (2016).
- [82] M. Piani, M. Cianciaruso, T. R. Bromley, C. Napoli, N. Johnston, and G. Adesso, *Robustness of Asymmetry and Coherence of Quantum States*, Phys. Rev. A 93, 042107 (2016).
- [83] P. Skrzypczyk, I. Šupić, and D. Cavalcanti, All Sets of Incompatible Measurements Give an Advantage in Quantum State Discrimination, Phys. Rev. Lett. 122, 130403 (2019).
- [84] G. Vidal and R. Tarrach, *Robustness of Entanglement*, Phys. Rev. A **59**, 141 (1999).
- [85] K. Bu, U. Singh, S.-M. Fei, A. K. Pati, and J. Wu, Maximum Relative Entropy of Coherence: An Operational Coherence Measure, Phys. Rev. Lett. 119, 150405 (2017).
- [86] B. Regula, K. Fang, X. Wang, and G. Adesso, *One-Shot Coherence Distillation*, Phys. Rev. Lett. **121**, 010401 (2018).
- [87] N. Datta, Min- and Max-Relative Entropies and a New Entanglement Monotone, IEEE Trans. Inf. Theory 55, 2816 (2009).
- [88] F. G. S. L. Brandão and M. B. Plenio, A Reversible Theory of Entanglement and Its Relation to the Second Law, Commun. Math. Phys. **295**, 829 (2010).
- [89] F. G. S. L. Brandão and N. Datta, One-Shot Rates for Entanglement Manipulation under Non-entangling Maps, IEEE Trans. Inf. Theory 57, 1754 (2011).
- [90] Q. Zhao, Y. Liu, X. Yuan, E. Chitambar, and X. Ma, *One-Shot Coherence Dilution*, Phys. Rev. Lett. **120**, 070403 (2018).

- [91] M. Berta and C. Majenz, Disentanglement Cost of Quantum States, Phys. Rev. Lett. 121, 190503 (2018).
- [92] F. G. S. L. Brandão, Entanglement Activation and the Robustness of Quantum Correlations, Phys. Rev. A 76, 030301(R) (2007).
- [93] F. G. S. L. Brandão, *Quantifying Entanglement with Witness Operators*, Phys. Rev. A **72**, 022310 (2005).
- [94] J. Eisert, F. G. S. L. Brandão, and K. M. R. Audenaert, Quantitative Entanglement Witnesses, New J. Phys. 9, 46 (2007).
- [95] C. M. Edwards, Classes of Operations in Quantum Theory, Commun. Math. Phys. 20, 26 (1971).
- [96] C. M. Edwards, The Theory of Pure Operations, Commun. Math. Phys. 24, 260 (1972).
- [97] R. T. Rockafellar, *Convex Analysis* (Princeton University Press, Princeton, NJ, 2015).
- [98] B. M. Terhal, D. P. DiVincenzo, and D. W. Leung, *Hiding Bits in Bell States*, Phys. Rev. Lett. 86, 5807 (2001).
- [99] D. P. DiVincenzo, D. W. Leung, and B. M. Terhal, Quantum Data Hiding, IEEE Trans. Inf. Theory 48, 580 (2002).
- [100] W. Matthews, S. Wehner, and A. Winter, *Distinguishability* of Quantum States Under Restricted Families of Measurements with an Application to Quantum Data Hiding, Commun. Math. Phys. **291**, 813 (2009).
- [101] G. Vidal, Entanglement Monotones, J. Mod. Opt. 47, 355 (2000).
- [102] G. Ludwig, Foundations of Quantum Mechanics I, Theoretical and Mathematical Physics (Springer-Verlag, Berlin, 1983).
- [103] K. Modi, A. Brodutch, H. Cable, T. Paterek, and V. Vedral, The Classical-Quantum Boundary for Correlations: Discord and Related Measures, Rev. Mod. Phys. 84, 1655 (2012).
- [104] Z.-W. Liu, R. Takagi, and S. Lloyd, *Diagonal Quantum Discord*, J. Phys. A 52, 135301 (2019).
- [105] S. Boyd and L. Vandenberghe, *Convex Optimization* (Cambridge University Press, New York, 2004).
- [106] A. W. Harrow and M. A. Nielsen, Robustness of Quantum Gates in the Presence of Noise, Phys. Rev. A 68, 012308 (2003).
- [107] M. Steiner, Generalized Robustness of Entanglement, Phys. Rev. A 67, 054305 (2003).
- [108] M. Ringbauer, T. R. Bromley, M. Cianciaruso, L. Lami, W. Y. S. Lau, G. Adesso, A. G. White, A. Fedrizzi, and M. Piani, Certification and Quantification of Multilevel Quantum Coherence, Phys. Rev. X 8, 041007 (2018).
- [109] N. Johnston, C.-K. Li, S. Plosker, Y.-T. Poon, and B. Regula, *Evaluating the Robustness of k-Coherence and k-Entanglement*, Phys. Rev. A **98**, 022328 (2018).
- [110] S. Bravyi, D. Browne, P. Calpin, E. Campbell, D. Gosset, and M. Howard, Simulation of Quantum Circuits by Low-Rank Stabilizer Decompositions, Quantum 3, 181 (2019).
- [111] G. Aubrun, L. Lami, C. Palazuelos, S. J. Szarek, and A. Winter, *Universal Gaps for XOR Games from Estimates on Tensor Norm Ratios*, arXiv:1809.10616.
- [112] S. N. Filippov, T. Heinosaari, and L. Leppäjärvi, *Simulability of Observables in General Probabilistic Theories*, Phys. Rev. A **97**, 062102 (2018).
- [113] M. M. Wilde, Quantum Information Theory (Cambridge University Press, Cambridge, England, 2013).

- [114] R. Renner and S. Wolf, in *Proceedings of International Symposium on Information Theory*, 2004, ISIT 2004 (IEEE, Chicago, IL, USA, 2004), p. 233.
- [115] R. Renner, Security of Quantum Key Distribution, Int. J. Quantum. Inform. **06**, 1 (2008).
- [116] N. Ciganović, N. J. Beaudry, and R. Renner, Smooth Max-Information as One-Shot Generalization for Mutual Information, IEEE Trans. Inf. Theory 60, 1573 (2014).
- [117] R. König, R. Renner, and C. Schaffner, *The Operational Meaning of Min- and Max-Entropy*, IEEE Trans. Inf. Theory 55, 4337 (2009).
- [118] P. Zanardi, C. Zalka, and L. Faoro, Entangling Power of Quantum Evolutions, Phys. Rev. A 62, 030301(R) (2000).
- [119] M. A. Jafarizadeh, M. Mirzaee, and M. Rezaee, Exact Calculation of Robustness of Entanglement via Convex Semi-Definite Programming, Int. J. Quantum. Inform. 03, 511 (2005).
- [120] D. Reeb, M. J. Kastoryano, and M. M. Wolf, *Hilbert's Projective Metric in Quantum Information Theory*, J. Math. Phys. (N.Y.) **52**, 082201 (2011).
- [121] M. A. Nielsen, *Conditions for a Class of Entanglement Transformations*, Phys. Rev. Lett. **83**, 436 (1999).
- [122] F. Buscemi, M. Keyl, G. M. D'Ariano, P. Perinotti, and R. F. Werner, *Clean Positive Operator Valued Measures*, J. Math. Phys. (N.Y.) 46, 082109 (2005).
- [123] F. Buscemi, All Entangled Quantum States Are Nonlocal, Phys. Rev. Lett. 108, 200401 (2012).
- [124] M. Horodecki and J. Oppenheim, Fundamental Limitations for Quantum and Nanoscale Thermodynamics, Nat. Commun. 4, 2059 (2013).
- [125] F. Buscemi, Degradable Channels, Less Noisy Channels, and Quantum Statistical Morphisms: An Equivalence Relation, Probl. Inf. Transm. 52, 201 (2016).
- [126] F. Buscemi and G. Gour, Quantum Relative Lorenz Curves, Phys. Rev. A 95, 012110 (2017).
- [127] F. Buscemi, Comparison of Noisy Channels and Reverse Data-Processing Theorems, in Proceedings of the 2017 IEEE Information Theory Workshop (ITW), Kaohsiung, Taiwan, 2017 (IEEE, 2017), pp. 489–493.
- [128] G. Gour, D. Jennings, F. Buscemi, R. Duan, and I. Marvian, Quantum Majorization and a Complete Set of Entropic Conditions for Quantum Thermodynamics, Nat. Commun. 9, 5352 (2018).
- [129] K. Matsumoto, A Quantum Version of Randomization Criterion, arXiv:1012.2650.
- [130] F. Buscemi, Comparison of Quantum Statistical Models: Equivalent Conditions for Sufficiency, Commun. Math. Phys. **310**, 625 (2012).
- [131] A. Jenčová, Comparison of Quantum Channels and Statistical Experiments, Proceedings of the IEEE International Symposium on Information Theory (ISIT), Barcelona, Spain, 2012 (IEEE, 2012), pp. 2249–2253.
- [132] M. Sion, *On General Minimax Theorems*, Pac. J. Math. **8**, 171 (1958).
- [133] G. Gour and C. M. Scandolo, *The Entanglement of a Bipartite Channel*, arXiv:1907.02552.
- [134] M. W. Girard, Convex Analysis in Quantum Information, Ph. D. thesis, University of Calgary, 2017.

- [135] K. C. Tan, S. Choi, and H. Jeong, Optimizing Nontrivial Quantum Observables Using Coherence, New J. Phys. 21, 023013 (2019).
- [136] M. A. Nielsen and I. L. Chuang, *Quantum Computation and Quantum Information*, 10th ed. (Cambridge University Press, New York, 2011).
- [137] P. M. Alberti and A. Uhlmann, A Problem Relating to Positive Linear Maps on Matrix Algebras, Rep. Math. Phys. 18, 163 (1980).
- [138] A. Chefles, R. Jozsa, and A. Winter, *On the Existence of Physical Transformations between Sets of Quantum States*, Int. J. Quantum. Inform. **02**, 11 (2004).
- [139] G. Chiribella, G. M. D'Ariano, and P. Perinotti, *Transforming Quantum Operations: Quantum Supermaps*, Europhys. Lett. **83**, 30004 (2008).

- [140] K. Życzkowski, Quartic Quantum Theory: An Extension of the Standard Quantum Mechanics, J. Phys. A 41, 355302 (2008).
- [141] G. Chiribella, G. M. D'Ariano, and P. Perinotti, *Quantum Circuit Architecture*, Phys. Rev. Lett. 101, 060401 (2008).
- [142] R. Uola, T. Kraft, J. Shang, X.-D. Yu, and O. Gühne, Quantifying Quantum Resources with Conic Programming, Phys. Rev. Lett. 122, 130404 (2019).
- [143] M. Oszmaniec and T. Biswas, Operational Relevance of Resource Theories of Quantum Measurements, Quantum 3, 133 (2019).
- [144] J.-B. Hiriart-Urruty and C. Lemaréchal, *Fundamentals of Convex Analysis* (Springer-Verlag, Berlin, 2001).
- [145] J. P. Ponstein, *Approaches to the Theory of Optimization* (Cambridge University Press, Cambridge, England, 2004).