Direct observation of hydrides formation in cavity-grade niobium

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Niobium is an important technological superconductor used to make radio frequency cavities for particle accelerators. Using laser confocal microscopy we have directly investigated hydride precipitates formation in cavity-grade niobium at 77 and 140 K. We have found that large hydrides were usually formed after chemical or mechanical treatments, which are known to lead to a strong degradation of the quality factor known as Q disease. From our experiments we can conclude that hydrides causing Q disease are islands with a characteristic thickness of ≈ 100 nm and in-plane dimensions 1–10 μ m. Our results show that mechanical polishing uploads a lot of hydrogen into bulk niobium while electropolishing leads to a mild contamination. Vacuum treatments at 600–800°C are demonstrated to preclude large hydride formation in line with the absence of Q disease in similarly treated cavities.

I. INTRODUCTION

The phenomenon of drastic drop of niobium superconducting radio frequency (SRF) cavities quality factor Q_0 at low accelerating fields below 2 MV/m without field emission or multipacting was encountered more than 20 years ago $[1-4]$ $[1-4]$ $[1-4]$. It was found that a fast cooldown $(< 1$ hour from room temperature to 4.2 K) could in some cases solve the problem. A conclusion was made that the precipitation of the dissolved hydrogen as lossy hydrides at the rf surface at temperatures 100–150 K was the reason for the increase of the surface resistance. These hydrides are normal conducting at the typical cavity operating temperature of about 2 K. The effect was named ''hydrogen Q disease.'' It was supposed that the extent of Q degradation depends on the quantity of dissolved hydrogen, the cooldown rate, and the amount of atomic size defects and interstitial impurities serving as nucleation centers for forming hydrides [\[3](#page-5-2)–[6\]](#page-5-3). Degassing in a vacuum of better than 10^{-6} Torr at 800° C for 2–3 hours or at 600° C for 10 hours cures hydrogen contamination, restores low resistance state, and removes Q disease [[1](#page-5-0),[7\]](#page-5-4).

The typical niobium cavity fabrication procedure consists of multiple steps, which include half-cell stamping, weld preparation, light etch, and electron beam welding of the half-cell cups and beam tubes together. Manufactured cavities are then put through a sequence of different treatments developed primarily by empirical methods. Typically, bulk chemical removal by buffered chemical polishing (BCP), electropolishing (EP) or, more recently, centrifugal barrel polishing, are applied to remove \geq

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120 μ m from the surface. In the current processing sequence for International Linear Collider it is followed by 800°C vacuum bake for 2 hours, light EP (\sim 20 μ m), high pressure water rinsing, and 120° C in situ bake for 48 hours.

In general, whenever the naturally formed oxide layer is not present on the surface of the niobium, there is a possibility for hydrogen uptake into the material. Such opportunities for hydrogen to enter niobium exist along the processing steps. Chemical polishing (EP or BCP) is suggested to be one of the pathways whereby niobium becomes contaminated with hydrogen [\[8](#page-5-5)]. In the case of BCP it was found that keeping the acid temperature below 15°C helped minimizing hydrogen upload [[9\]](#page-5-6). Mechanical polishing using a solution, which is not hydrogen free, generally results in a strong hydrogen contamination [[10](#page-5-7)]. Finally, reabsorption from the furnace upon cooldown is another possible hydrogen pickup route [\[7,](#page-5-4)[11](#page-5-8)[,12\]](#page-5-9).

Nb-H systems were studied extensively in the 1970s [\[13–](#page-5-10)[15\]](#page-5-11). A complete equilibrium phase diagram is presented in Fig. [1](#page-1-0). Niobium can dissolve large amounts of hydrogen in interstitial states. The α phase is the solid solution in which hydrogen atoms are distributed on random tetrahedral sites. At room temperature the solubility limit of this phase extends up to 4 atomic percent of hydrogen. But this limit is drastically reduced with temperature decrease and is about 5×10^{-4} atomic percent at 100 K. When hydrogen concentration is above the solubility limit, different stoichiometric hydrides can be formed.

Despite the fact that the formation of hydrides is considered to be a well-established reason for the Q disease, they were never observed in cavities or cavity-grade material so far with only hypotheses put forward regarding their size, shape, and distribution [\[6](#page-5-3)]. This motivated us to carry out the presented metallographic studies of SRF

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FIG. 1. Phase diagram of Nb-H system (from Ref. [[15\]](#page-5-11)).

cavity-grade niobium samples. We have tried to investigate by direct observations how different treatments used for cavity processing affect the formation of hydrides.

II. EXPERIMENTAL

Our samples were cut by wire electrical discharge machining from fine grain and large grain residual resistivity ratio (RRR) \geq 300 niobium sheets from the same material batch as was used to manufacture 9-cell and 1-cell cavities with state-of-the-art performance. Samples were embedded in either Bakelite or epoxy pucks and mechanically polished first by the increasing grit size (120–600) sandpaper. This was followed by automated polishing on Texmet polishing cloth with $6 \mu m$ size diamond solution for several hours. Final polishing was done on a vibrational polishing

machine with a colloidal water-based SiC solution. Electropolishing (EP) was performed using the in-house sample electropolishing setup with the standard acid solution used for cavities. 800° baking was accomplished in the same vacuum furnace that is used for cavities. 120° baking was done using a continuously pumped vacuum tube with resistive heating on the outside walls.

To study hydride formation, we either immersed the samples in liquid nitrogen or, when a temperature above 77 K was needed, we kept the samples in a nitrogen vapor above liquid level. The farther away the sample was from the liquid, the higher the equilibrium sample temperature. By adjusting the distance between a sample and liquid nitrogen level it was possible to stabilize temperature within 10 K. After warming the samples up to the room temperature we used a laser confocal optical microscope with a lateral resolution of order 1 μ m to observe the obtained surface structures.

III. RESULTS

First we studied mechanically polished samples without any subsequent chemical treatment. Presented in Fig. [2](#page-1-1) is a typical pattern obtained on a fine grain sample after placing it into liquid nitrogen for 1 h. Lenslike structures of a characteristic size of about $1-3 \mu m$ are clearly seen. No such structures were observed before placing the sample into liquid nitrogen. There is no difference between the patterns obtained after 1 h and after keeping samples in liquid nitrogen for 24 and 48 hours.

FIG. 2. (a) Lenslike structures after 77 K hold on a fine grain mechanically polished sample; (b) zoom-in from the selected area, arrows indicate surface relief due to hydrides; (c) smaller structures observed at another location.

Then we looked at the samples after mechanical polishing, 800°C baking for 2 hours, and keeping in liquid nitrogen for 24 hours. No apparent structures were observed. It is in good agreement with a commonly accepted concept that 800° C baking reduces bulk hydrogen concentration significantly which prevents the formation of large hydrides upon cooldown.

To investigate the possible 120° C baking effect, we polished one of the samples mechanically and baked it at 120° C for 48 hours before putting the sample into liquid nitrogen. The surface structures obtained are similar to those without 120° C treatment.

To check if the surface structures we observed are purely morphological or represent room-temperature hydrides, we baked one of the samples at 800° C for 2 hours after structures had been already formed. Structures did not disappear, and our interpretation of this fact consists in the following. The lattice constant in a hydride phase is about 10% larger than in pure niobium. The formation of

FIG. 3. Lenslike structures after 77 K hold on mechanically polished (a) fine grain sample; (b) single grain sample.

sufficiently large hydrides leads inevitably to mechanical deformation of the crystal around the precipitate. And it is the residual surface relief (not hydrides themselves) that we observe under a microscope. This explains why we can observe structures at room temperature where hydrides must have already disappeared according to the phase diagram, and also why these structures remain after 800°C bake.

It is interesting to note that hydride relief structures in adjacent grains are oriented in different ways, suggesting that their alignment is affected by crystallography. In line with this, a pattern obtained on the single grain sample (see Fig. [3\)](#page-2-0) shows lenslike structures parallel to each other.

Another apparent feature is that hydrides decorate surface irregularities such as scratches. This may mean that hydrogen is segregated at the irregularities before the cooldown or that such features provide nucleation centers for the hydride to form, or both.

Figure [4](#page-2-1) shows a typical surface profile across a groove attributed to large hydride formation. Characteristic height difference is of the order of 100 nm.

We produced analogous experiments at a higher temperature of 140 K by keeping samples in nitrogen vapor as described above. Figure [5](#page-2-2) shows much larger structures of a characteristic size of 10 μ m formed under these conditions. Some of them even have different ''starlike'' shape.

FIG. 4. Surface profile across the AB line from Fig. [3.](#page-2-0)

FIG. 5. Patterns obtained after keeping a mechanically polished sample at 140 K for 1 h.

One of the important questions in cavities manufacturing is if and when chemical treatments lead to additional hydrogen upload. This is especially important since light electropolishing to remove 20–40 μ m of the material is customary at many labs after 600–800°C hydrogen degassing. We took 800 $^{\circ}$ C baked sample and removed a 40 μ m layer by electropolishing. After keeping the sample in liquid nitrogen we did not observe any structures. But after keeping the sample at 140 K some microscale structures appeared. Figure [6](#page-3-0) shows the evolution of hydrides formation with time. There was no apparent growth of the hydrides over three different times of observation. This may mean that precipitation reached the saturation determined by hydrogen availability and one can make a conclusion that 40 μ m electropolishing does upload some hydrogen but not as much as the mechanical polishing.

FIG. 6. Patterns obtained after keeping electropolished sample at 140 K for (a) 1; (b) 2; and (c) 4 hours.

IV. DISCUSSION

Appearance of surface damage due to hydrides was recently extensively investigated in steel and iron [\[16,](#page-5-12)[17\]](#page-5-13) and seems to be a general phenomenon for metal-hydrogen systems. In our studies we observe surface damage in niobium left behind by hydrides, which allows us to make conclusions about the original hydride size, morphology, and density on the niobium surface and correlate it with the performance of superconducting cavities.

Hydrogen is extremely mobile in niobium and thus interaction with defects is very important for hydride precipitation.

Interstitial hydrogen in metals interacts with microstructural defects including point defects (vacancies and interstitials), linear defects (e.g. dislocations), grain boundaries, surfaces, and interfaces (see Ref. [\[18\]](#page-5-14) for the latest review).

Typically, surface and interface binding are the strongest, which leads to hydrogen segregation near such inhomogeneities. In particular, for cavity-grade niobium such hydrogen enrichment at the niobium/oxide interface was reported [\[4,](#page-5-1)[19,](#page-5-15)[20\]](#page-5-16). In most of the studies the near-surface hydrogen concentration was found to be only weakly affected by vacuum heat treatments at 600–1000 $^{\circ}$ C typically applied on cavities to mitigate Q disease. In our experiments hydride-induced dents are absent after a cryohold of the mechanically polished sample annealed at 800° C for 3 hours in vacuum, which confirms the effectiveness of such a treatment for suppressing hydride formation. Even though the near-surface concentration may still be high, the amount of hydrogen in the bulk is very low inhibiting the growth of large hydrides in the first \sim 100 nm where rf currents flow in cavities.

Strong hydrogen interaction with dislocations, especially of edge character, was found to lead to the formation of hydrogen-rich Cottrell clouds around dislocation cores. Such objects were specifically identified in niobium [\[21](#page-5-17)[,22\]](#page-5-18). Mechanical deformation steps in cavity processing may therefore have a strong effect on hydrogen distribution and on the formation of hydrides upon cooldown. Our studies did not specifically address the effect of dislocations but we observed similarity in the density of surface dents between fine and large (single) grain samples, which have different bulk dislocation densities due to difference in mechanical deformation. Such similarity hints that predominant hydride nucleation centers in our samples may either form independently of mechanical deformation such as, for example, misfit dislocations at the niobium/oxide interface, or that nucleation centers are defects other than dislocations. Such a speculation though should be addressed by future investigations.

Other interstitial impurities in cavity niobium—oxygen, nitrogen, and carbon—are also good trapping centers for hydrogen. Strong correlation between the concentrations of oxygen and hydrogen near the surface was reported for several metals [[23](#page-5-19)[–26\]](#page-5-20) and, in particular, for niobium

Treatment	$77-140$ K hold outcome	Typical cavity performance
Mechanical polish (MP)	Large hydrides	Q disease
$+800C$ 3 hrs	No hydrides seen	No Q disease
$+BCP_5$ min	No hydrides seen	No Q disease
$+120C$ 48 hrs	No hydrides seen	No Q disease
$MP + 800C3$ hrs + EP40 min	Small hydrides after 140 K	No Q disease

TABLE I. Effect of different treatments.

[\[27–](#page-5-21)[30\]](#page-5-22). A direct experimental evidence for their role in cavities comes from the absence of Q disease in the cavities made of low RRR (≤ 50) material. Our samples are of high RRR (\sim 300) with the O, N, and C content on the level of ≤ 20 ppm as compared to bulk hydrogen concentrations of more than 100 ppm routinely obtained in mechanically polished samples. Thus, other interstitials are probably of secondary importance for hydrogen trapping in our studies.

Among point defects, interaction of hydrogen with vacancies in metals is an important issue, which gained significant attention lately due to the discovery of so-called "superabundant vacancies" [\[31\]](#page-5-23). It was found out that niobium is among the metals where the presence of hydrogen may lead to the enormous enhancement of vacancy concentration [\[32](#page-5-24)[,33\]](#page-5-25) due to the formation of vacancyhydrogen complexes. There are experimental hints [\[34](#page-5-26)[,35\]](#page-5-27) that these objects may be involved in a 120° C baking effect on cavities, which remains poorly understood for more than a decade. We speculate that in our experiments such complexes may be abundant in the near-surface region and serve as primary trapping and nucleation centers for hydride growth upon cooldown.

It is useful to perform simple diffusion-controlled precipitation estimates for any kind of above-mentioned defects serving as a nucleation center. Final hydride size at $T = 2$ K after a particular cooldown process depends on the amount of trapped hydrogen at the nucleation center, which, in turn, is proportional to hydrogen concentration and to a in turn, is proportional to hydrogen concentration and to a volume $(\sqrt{Dt})^3$ from which more hydrogen can diffuse towards a nucleation center. Here, $D = D(T)$ is the diffusion constant of hydrogen in niobium. There are two different possibilities for hydrides to stop growing. The first one is ent possibilities for hydrides to stop growing. The first one is
when \sqrt{Dt} becomes comparable with the distance between adjacent nucleation centers, all hydrogen becomes trapped, and hydrides stop growing due to the lack of hydrogen supply. Another one can be realized due to the lack of hydrogen mobility when the temperature becomes low. In this case hydrogen distribution remains ''frozen'' with the size of hydrides determined by the time spent at higher temperatures. We believe in our experiments we have encountered both of these growth stopping scenarios.

In particular, in cases of 77 and 140 K hold of similarly treated samples there is no difference among pictures obtained after keeping samples for different durations at these temperatures. At the same time there is a significant difference in sizes of hydrides obtained at 140 (Fig. [5\)](#page-2-2) and 77 K (Fig. [2](#page-1-1)). Therefore there is more hydrogen available for segregation at 77 K but hydrides are not growing. A logical interpretation of this fact could be lack of hydrogen mobility. Similarly, the 140 K result appears to be determined by hydrogen supply. Hence, one can make a conclusion that only when we are passing the $100-150$ K interval in a process of cooling down are hydrides really growing.

Based on the above, the surface relief picture obtained after 77 K holds can mimic the distribution of large hydrides that is realized in niobium cavities upon ''fast'' cooldown typical of cavity operations. Similarly, 140 K hold can be used as being representative for a ''slow'' cooldown. A summary of our findings for different treatments is shown in Table [I.](#page-4-0) We observe a strong correlation between our findings and the presence or absence of Q disease in cavity performance after the same treatments.

Finally, it should be noted that since our experimental technique relies on the formation of surface relief, we cannot exclude the presence of hydrides at cryogenic temperatures in the samples where we found no surface relief after cryohold. Lateral resolution of our microscope is limited to about 1 μ m and hence smaller hydrides could have been present during the hold leaving surface relief which is beyond the resolution available to us. There may also in principle exist hydride phases which do not leave a signature surface relief at all, for example, due to the smaller mismatch between the lattice parameters with host niobium. Therefore our observations can only provide a ''lower bound'' on the occurrence of hydride precipitation after a particular surface treatment.

V. CONCLUSIONS

We have developed a procedure and directly observed surface relief caused by the formation of niobium hydrides on the surface of the cavity-grade material. We correlated the formation of hydrides with the Q disease observed in SRF niobium cavities and explicitly described the morphology and characteristic sizes of culprit hydrides. Possible hydrogen trapping mechanisms and their effect on the formation of surface hydrides were discussed.

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- [1] S. Isagawa, [J. Appl. Phys.](http://dx.doi.org/10.1063/1.327523) **51**, 6010 (1980).
- [2] B. Aune, B. Bonin, J. M. Cavedon, M. Juillard, A. Godin, C. Henriot, P. Leconte, H. Safa, A. Veyssiere, and C. Zylberajch, in Proceedings of LINAC'90 (National Lab, Los Alamos, NM, 1990), pp. 253–255 [[http://www.jacow](http://www.jacow.org) [.org\]](http://www.jacow.org).
- [3] B. Bonin and R.W. Roth, in Proceedings of the 5th Workshop on RF Superconductivity, edited by D. Proch (DESY, Hamburg, Germany, 1992), pp. 210–244 [\[http://](http://www.jacow.org) [www.jacow.org\]](http://www.jacow.org).
- [4] C. Z. Antoine, B. Aune, B. Bonin, J. Cavedon, M. Juillard, A. Godin, C. Henriot, P. Leconte, H. Safa, A. Veyssiere, A. Chevarier, and B. Roux, in Proceedings of the Fifth Workshop on RF Superconductivity (DESY, Hamburg, Germany, 1991), pp. 616–634.
- [5] J. Halbritter, P. Kneisel, and K. Saito, in Proceedings of the 6th Workshop on RF Superconductivity (Gordon and Breach, New York, 1994), p. 196 [[http://www.jacow.org\]](http://www.jacow.org).
- [6] R. W. Roth, in Proceedings of the 6th Workshop on RF Superconductivity (Ref. [[5](#page-5-28)]) [[http://www.jacow.org\]](http://www.jacow.org).
- [7] S. Isagawa, [J. Appl. Phys.](http://dx.doi.org/10.1063/1.328267) 51, 4460 (1980).
- [8] T. Higuchi, K. Saito, Y. Yamazaki, T. Ikeda, and S. Ohgushi, in Proceedings of the 10th Workshop on RF Superconductivity, Tsukuba, Japan (2001), pp. 427–430 [\[http://www.jacow.org\]](http://www.jacow.org).
- [9] H. Padamsee, J. Knobloch, and T. Hays, RF Superconductivity for Accelerators (John Wiley and Sons, New York, 1998).
- [10] T. Higuchi and K. Saito, in Proceedings of the 11th Workshop on RF Superconductivity (Travemunde, Lubeck, Germany, 2003), pp. 572–578 [[http://www.jacow](http://www.jacow.org) [.org\]](http://www.jacow.org).
- [11] K. Faber and H. Schultze, Scr. Metall. **6**[, 1065 \(1972\).](http://dx.doi.org/10.1016/0036-9748(72)90188-3)
- [12] A. Grassellino, A. Romanenko, A. Rowe, L. D. Cooley, C. Ginsburg, and V. P. Yakovlev, in Proceedings of IPAC'12, New Orleans (2012), WEPPC115 [[http://www.jacow.org\]](http://www.jacow.org).
- [13] T. Schober, Scr. Metall. 7[, 1119 \(1973\).](http://dx.doi.org/10.1016/0036-9748(73)90025-2)
- [14] H.K. Birnbaum, M.L. Grossbeck, and M. Amano, [J. Less-Common Met.](http://dx.doi.org/10.1016/0022-5088(76)90048-5) 49, 357 (1976).
- [15] J.-M. Welter and F.J. Johnen, Z. Phys. B 27[, 227 \(1977\).](http://dx.doi.org/10.1007/BF01325532)
- [16] D. P. Escobar, C. Minambres, L. Duprez, K. Verbeken, and M. Verhaege, Corros. Sci. 53[, 3166 \(2011\)](http://dx.doi.org/10.1016/j.corsci.2011.05.060).
- [17] I.M. Robertson, [Eng. Fract. Mech.](http://dx.doi.org/10.1016/S0013-7944(01)00011-X) **68**, 671 (2001).
- [18] A. Pundt and R. Kirchheim, [Annu. Rev. Mater. Res.](http://dx.doi.org/10.1146/annurev.matsci.36.090804.094451) 36, [555 \(2006\)](http://dx.doi.org/10.1146/annurev.matsci.36.090804.094451).
- [19] T. Tajima, R. L. Edwards, F. L. Krawczyk, J. Liu, D. L. Schrage, A. H. Shapiro, J. R. Tesmer, C. J. Wetteland, and R. L. Geng, in Proceedings of the 11th Workshop on RF Superconductivity (Ref. [\[10\]](#page-5-7)), THP19 [[http://www.jacow](http://www.jacow.org) [.org\]](http://www.jacow.org).
- [20] A. Romanenko and L.V. Goncharova, [Supercond. Sci.](http://dx.doi.org/10.1088/0953-2048/24/10/105017) Technol. 24[, 105017 \(2011\)](http://dx.doi.org/10.1088/0953-2048/24/10/105017).
- [21] J.A. Rodriguez and R. Kirchheim, [Scr. Metall.](http://dx.doi.org/10.1016/0036-9748(83)90091-1) 17, 159 [\(1983\)](http://dx.doi.org/10.1016/0036-9748(83)90091-1).
- [22] E. Krautz, [Int. J. Refract. Met. Hard Mater.](http://dx.doi.org/10.1016/0263-4368(93)90039-I) 12, 315 (1993).
- [23] G. Maire, L. Hilaire, P. Legare, F.G. Gault, and A. O'Cinneide, J. Catal. 44[, 293 \(1976\)](http://dx.doi.org/10.1016/0021-9517(76)90399-7).
- [24] M. H. Mintz, J. A. Schultz, and J. W. Rabalais, *[Phys. Rev.](http://dx.doi.org/10.1103/PhysRevLett.51.1676)* Lett. 51[, 1676 \(1983\).](http://dx.doi.org/10.1103/PhysRevLett.51.1676)
- [25] K. Ojima and K. Ueda, [Appl. Surf. Sci.](http://dx.doi.org/10.1016/S0169-4332(00)00371-8) 165, 149 (2000).
- [26] E. Nowicka and R. Dus, Vacuum 67[, 409 \(2002\).](http://dx.doi.org/10.1016/S0042-207X(02)00239-7)
- [27] P. Zapp and H. Birnbaum, Acta Metall. **28**[, 1275 \(1980\).](http://dx.doi.org/10.1016/0001-6160(80)90083-8)
- [28] P. Zapp and H. Birnbaum, Acta Metall. 28[, 1523 \(1980\).](http://dx.doi.org/10.1016/0001-6160(80)90053-X)
- [29] A. Shirley, C. Hall, and N. Prince, [Acta Metall.](http://dx.doi.org/10.1016/0001-6160(83)90193-1) 31, 985 [\(1983\)](http://dx.doi.org/10.1016/0001-6160(83)90193-1).
- [30] T. Pfiz, R. Messer, and A. Seeger, [Z. Phys. Chem.](http://dx.doi.org/10.1524/zpch.1989.164.Part_1.0969) 164, 969 [\(1989\)](http://dx.doi.org/10.1524/zpch.1989.164.Part_1.0969).
- [31] Y. Fukai and N. Okuma, *[Phys. Rev. Lett.](http://dx.doi.org/10.1103/PhysRevLett.73.1640)* **73**, 1640 (1994).
- [32] J. Cizek, I. Prochazka, R. Kuzel, F. Becvar, M. Cieslar, G. Brauer, W. Anwand, R. Kirchheim, and A. Pundt, [J.](http://dx.doi.org/10.1016/j.jallcom.2004.11.125) [Alloys Compd.](http://dx.doi.org/10.1016/j.jallcom.2004.11.125) 404–406, 580 (2005).
- [33] J. Cizek, I. Prochazka, S. Danis, G. Brauer, W. Anwand, R. Gemma, E. Nikitin, R. Kirchheim, A. Pundt, and R. Islamgaliev, Phys. Rev. B 79[, 054108 \(2009\)](http://dx.doi.org/10.1103/PhysRevB.79.054108).
- [34] A. Romanenko and H. Padamsee, [Supercond. Sci.](http://dx.doi.org/10.1088/0953-2048/23/4/045008) Technol. 23[, 045008 \(2010\)](http://dx.doi.org/10.1088/0953-2048/23/4/045008).
- [35] A. Romanenko, C. Edwardson, and P. Coleman, at TFSRF'2012 [\[www.jlab.org/indico/conferenceDisplay](www.jlab.org/indico/conferenceDisplay.py?confId=22) [.py?confId=22\]](www.jlab.org/indico/conferenceDisplay.py?confId=22).