

Compact 810 kA linear transformer driver cavity

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We are performing experiments with a 92-kV, 810-kA, 74.6-GW linear transformer driver (LTD) cavity. This cavity generates a ~ 100 ns power pulse from DC-charged capacitors in a single step. Our experiments start with an existing 100-kV, 490-kA LTD cavity and are making a number of improvements to it that are aimed at increasing the cavity's peak output power and better understanding its operation. We are making improvements to the gas switches, the capacitors, and the magnetic toroids as well as heavily instrumenting the cavity. These experiments have increased the cavity's output current into a matched load by 65% without increasing its volume.

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I. INTRODUCTION

Linear transformer drivers (LTDs) are a new pulsed-power architecture that may dramatically reduce the cost and size of future pulsed-power accelerators. In LTDs, DC-charged capacitors are discharged in parallel to provide a ~ 100 -ns wide power pulse to a load in one step, without any of the pulse compression hardware used in traditional pulsed-power accelerators [1–9]. We are testing a number of changes to the existing LTD-II [7] cavity to increase its peak power. We are also adding diagnostics to the cavity to allow us to better understand its operation. We are calling the new cavity LTD-III.

In our first set of experiments, we replaced the baseline Russian gas switches with improved low-inductance switches and have also used an improved capacitor design developed specifically for LTD's. Both of these changes have decreased the system inductance and increased the peak output power. In the second set of experiments we replaced the silicon-steel magnetic toroids in the cavity with lower loss Metglas® toroids, which lowered the eddy-current losses in the cores and also increased the cavity's

output power. Since the new switches and capacitors were significantly smaller than the ones used in the original cavity, we performed a third experiment with 45% larger capacitors in the cavity, which again increased the output power and slightly lengthened the pulse.

Figure 1 shows a schematic of a basic linear transformer driver “brick” that consists of two capacitors charged to opposite polarity, a triggered switch and a set of magnetic toroids. When the switch is triggered, current attempts to flow around the outside of the LTD cavity, but is blocked by the magnetic toroids until the toroids saturate. This causes a fast-rising high-voltage pulse in the center of the

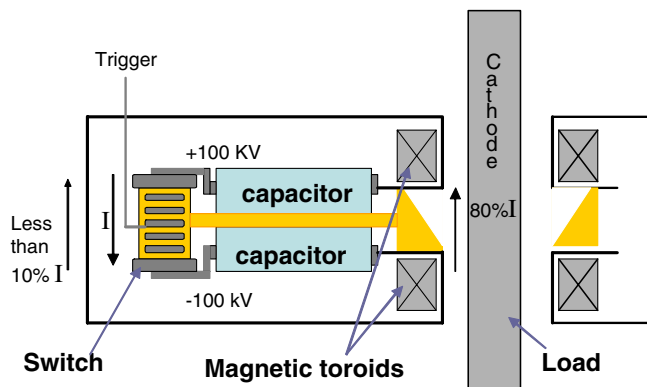


FIG. 1. Schematic of an LTD brick showing DC-charged capacitors, a triggered switch, and magnetic toroids.

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LTD cavity that may be coupled to a variety of loads, including a coaxial transmission line with a cathode stalk on the center line as shown in Fig. 1.

II. APPARATUS

A. Capacitors and switches

Figure 2 compares the baseline gas switch designed by the High Current Electronics Institute (HCEI) at Tomsk, Russia [10] with the switch jointly designed by L3 pulse sciences and Sandia Laboratories [11,12] and used in LTD-III. The L3-Sandia switch is significantly smaller than the HCEI switch and, hence, adds less inductance to the brick circuit. The L3-Sandia switch is a two-stage switch with midplane triggering and a UV preionization gap in the midplane to lower the switch's command jitter. The L3-Sandia switch has two each 6.4 mm electrode gaps and is insulated with dry air at a pressure of 158 psi (absolute). Figure 3 compares the single-ended general

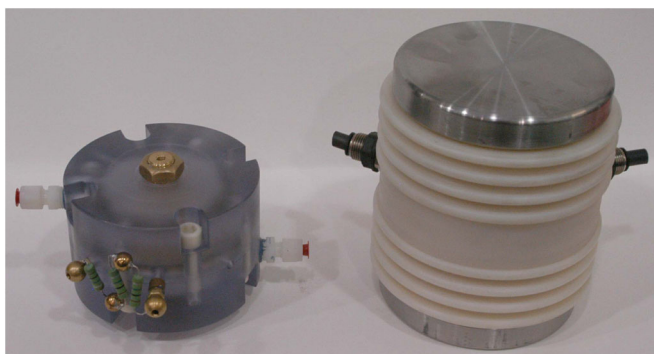


FIG. 2. Comparison of the 7.6-cm-tall L3-Sandia switch (left) used in the LTD-III experiments with the 15-cm-tall baseline switch (right) developed by the High Current Electronics Institute in Tomsk, Russia.



FIG. 3. Comparison of the double-ended 40-nF capacitors used in LTD-III with the single-ended 40-nF capacitors used in previous experiments.

atomics 40-nF capacitors used in LTD-II with the 40-nF double-ended capacitors used in LTD-III. The double-ended capacitors allow the brick circuitry to be packed more tightly, also resulting in a lower system inductance.

B. Magnetic toroids

Two sets of ferromagnetic toroidal cores were used in our experiments. Both types of cores were fabricated of a continuous winding of iron alloy with Mylar or polyethylene between adjacent windings to prevent shorting from turn to turn. The cores had inner diameters of 86.4 cm and were approximately 5.1-cm thick [13].

In the first experiments, we used the silicon-steel cores supplied with LTD-II. These cores were produced by researchers at the High-Current Electronics Institute in Tomsk, Russia, and were fabricated of ET3425 transformer iron [6,14]. The metal tapes used for the windings in the ET3425 cores were 80-microns thick and 18 mm tall with 12-micron thick layers of Mylar between metal turns. ET3425 has a saturation flux intensity of ~ 2.0 Tesla. The cavity contained six of the ET3425 cores, which had for a total Volt-second product from $-B$ saturation to $+B$ saturation of approximately 22 mV-seconds.

For the second two experiments, we used toroidal cores of 2605CO metallic glass that were supplied by Metglas®, Inc. [15]. The metal tapes used for the windings of the Metglas® cores were 23-microns thick and 3.56-cm tall. Four-micron-thick layers of polyethylene were wound between adjacent layers of Metglas®. The Metglas® in each of the four cores had a cross-sectional area of 15.5 cm^2 [6,8]. 2605CO has a saturation intensity of 1.7 Tesla. The cavity contained four of the Metglas® cores, which had a total volt-second product of approximately 18 mV-seconds. We expected the 23- μm -thick Metglas® cores to have much lower losses than the 80 μm -thick ET3425 cores since the losses in the cores are inversely proportional to the square of the winding thickness.

The magnetic domains in the cores were “set” just before each LTD-III shot by discharging a unipolar current pulse around the metal exterior casing of the LTD-III cavity in the opposite polarity to the main high-voltage pulse generated by the LTD. This “premagnetization pulse” had a peak current of 2 kA and a full width at half maximum of 140 microseconds.

C. LTD-III cavity

Figure 4 shows a picture of the LTD-III cavity. Twenty bricks consisting of 40-nF or 58-nF capacitors and L3-Sandia switches [11] surround magnetic toroids and an annular liquid resistive load cavity. During normal operation, the top capacitors are charged to positive voltage and the bottom capacitors are charged to negative voltage. The overall cavity diameter is 199 cm. The outer diameter of the annular resistive load is 77 cm.

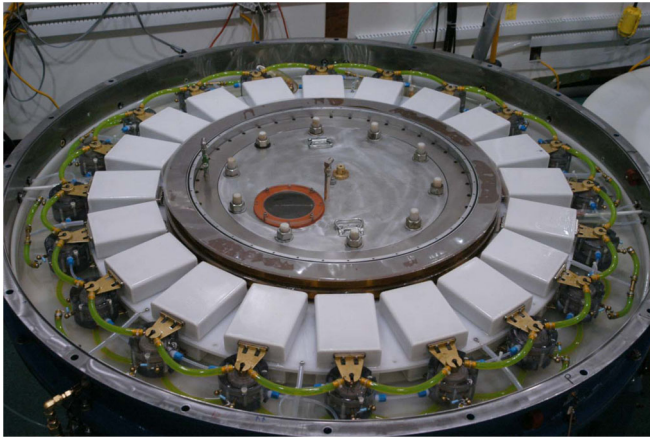


FIG. 4. Overview of the LTD-III cavity with 20 pairs of 40-nF capacitors, 20 switches and liquid charge and trigger resistors.

The cavity was filled with Luminol [16] transformer oil during operation. The oil was filled and drained through filters that removed particles and moisture. Most LTD cavities are used in a vertical orientation and a small (~ 2 psi) negative pressure is applied to the top of the cavity to draw off any air bubbles in the cavity. Since our cavity was used in a horizontal position, debubbling the cavity was more challenging. We took several steps to remove air bubbles from the cavity once it was filled and the cover was in place. A ~ 10 liter transparent tank that was half-filled with oil was connected to the top of our cavity and we applied a ~ 2 psi negative pressure to this tank to draw bubbles out of the cavity. We also tilted the cavity 2° from horizontal and recirculated oil from the bottom of the transparent tank back through the cavity to further flush out any remaining bubbles.

Figure 5 shows a cross section of the LTD-III cavity. Locations of the capacitors, switches, and resistive load cavity are indicated. The locations of ten load current B-dot detectors are indicated in a small cavity immediately

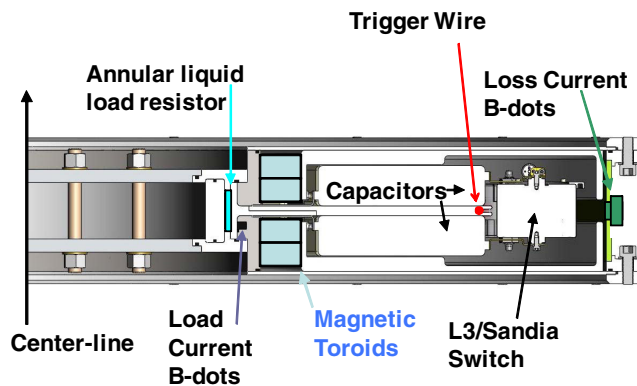


FIG. 5. Cross section of LTD-III cavity with 58-nF capacitors, showing location of capacitors, switches, magnetic toroids, and annular KBr resistive load along with load and loss current B-dot monitors.

outside the resistive load. Loss current B-dots are located on the outer diameter of the cavity. The load current B-dots were calibrated *while mounted in the LTD cavity hardware* in the calibration laboratory for the Z facility. The loss current B-dots were calibrated in separate fixtures in the same calibration laboratory. Figure 6 shows details of connections between the switches and capacitors in the initial experiments with 40 nF capacitors and solid-state resistors. Figure 7 shows similar details of the connections in the final experiments using 58-nF capacitors and all-liquid resistors. The charging resistors are 1-k Ω each and the trigger resistors are 500- Ω each.

Three types of resistors were used in these experiments: Small 0.63-cm diameter 500- Ω liquid resistors initially used as trigger resistors and 1-k Ω solid-state resistors initially used as charge resistors. (Fig. 6) 1.9-cm diameter style of liquid resistors were eventually used for both 500- Ω charge resistors and 1-k Ω charge resistors (Fig. 7). The small diameter liquid resistors were replaced because they developed air bubbles too frequently. The solid-state resistors were replaced because they occasionally shattered. All of the liquid resistors were copper sulfate solutions in yellow, oil-resistant plastic tubing.

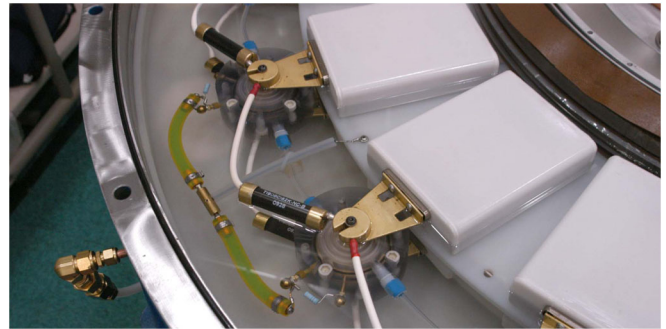


FIG. 6. Details of LTD-III connections with 40-nF capacitors, solid-state charge resistors, and liquid trigger resistors.



FIG. 7. Details of LTD-III connections with 58-nF capacitors, and all liquid charge and trigger resistors.

In a simple resistor-inductor-capacitor (RLC) circuit, the maximum energy is transferred to the resistive load under matched-load conditions; i.e., when

$$R = (L/C)^{0.5}. \quad (1)$$

Under matched-load conditions in a simple RLC circuit, there is a 16% reversal of the current. In our experiments, we defined matched load as a 16% current reversal even though our circuit, which had magnetic toroidal cores in it, was not strictly a simple RLC circuit. Thus, our results are “near matched load.”

Because LTD-III is a sealed cavity with high pressure gas switches inside it, mechanical rupture of a switch or a gas line could subject the cavity to the full switch gas pressure, which is more than the cavity is designed to withstand. We took two precautions to prevent the cavity from being overpressured in this manner. First, we put four pressure-relief valves [17] set for 1.5 psi (gauge) on the outer diameter of the cavity. Hoses from these valves ran to a common dump tank. Second, since failures of the gas system are usually caused by an electrical arc inside the cavity that occurs during the charging of the cavity, we programmed the control system to abort the charge and lower the switch pressure to 2 psi in the event of a charge imbalance in the \pm charging system.

D. Trigger system

The trigger pulse for LTD-III was generated by discharging two 20-nF capacitors in parallel through a modified version of the HCEI switches described previously [8,10]. The trigger generator was outside the LTD-III cavity and the trigger pulse was conveyed to the cavity via two high-voltage 67- Ω cables. The two cables entered the cavity 180° apart from each other and attached to a wire located in a groove in the insulator plate separating the capacitors (see Fig. 5). This trigger pulse was then

connected through twenty 500- Ω resistors to each of the 20 switch midplanes.

III. RESULTS

Table I summarizes the results of our three experiments. Figure 8 shows currents through matched resistive loads during three different experiments. The peak current in the first experiment of 615 kA is \sim 25% higher than the 490 kA peak currents achieved in LTD-II using HCEI switches and single-ended capacitors. This increase in peak current is due solely to the decrease in system inductance due to the improved capacitors and switches. In the second experiment, the Metglas® cores replaced the silicon-steel cores, reducing the losses and yielding a 650 kA peak current for a 32% increase over LTD-II results. In the third experiment, 58-nF capacitors replaced the 40 nF capacitors yielding a 810 kA peak current with a 60 ns 10% to 90% rise time. While the use of 58 nF capacitors increases the rise time of the pulse, the 100 ns zero-to-peak rise time we obtain with the 58 nF capacitors is still a good match to 100-ns wire-array implosion times on the Z accelerator [18,19].

Figure 9 shows the voltage and power waveforms for the tests with the 58 nF capacitors and Metglas® toroids when the capacitors were charged to \pm 100 kV and discharged into a matched load. The peak load power was 74.6 GW.

Figure 10 shows the loss currents that leaked past the toroidal cores and traveled around the outside of the cavity. In the first experiment, the loss currents peaked at 59 kA or about 9.6% of the load current. In the second experiment with the Metglas® cores, the loss currents peaked at about 15 kA or 2.3% of the load current. In the experiment with the Metglas® cores and the 58 nF capacitors, the loss current peaked at 51 kA or 6.3% of the load current. In the experiments with the silicon-steel cores, the loss currents peaked about 20 ns later than the load current,

TABLE I. Summary of results for LTD-III shots into matched loads with capacitors charged to \pm 100 kV.

Experiment	1	2	3
Capacitors	40 nF	40 nF	58 nF
Toroidal cores	Silicon steel	Metglas®	Metglas®
Matched-load peak current (kA)	615	650	810
Peak power (GW)	52	57	74.6
Peak voltage (kV)	84.4	87	92
10%–90% rise time (ns)	48	50	60
0-peak rise time (ns)	80	84	100
Energy to load by peak of pulse (J)	2947	2951	4147
Energy to load by 100 ns (J)	3769	3781	4147
Energy to load by end of pulse (J)	5471	5742	9217
Percent of capacitor bank energy to load by end of pulse	68.4	71.8	79.5
Volt-second product across load by end of pulse (mV-sec)	11.5	11.5	14.25
Load resistance from V_{load}/I_{load} (Ω)	0.144	0.136	0.113
Peak value of loss currents (kA)	59	15	51

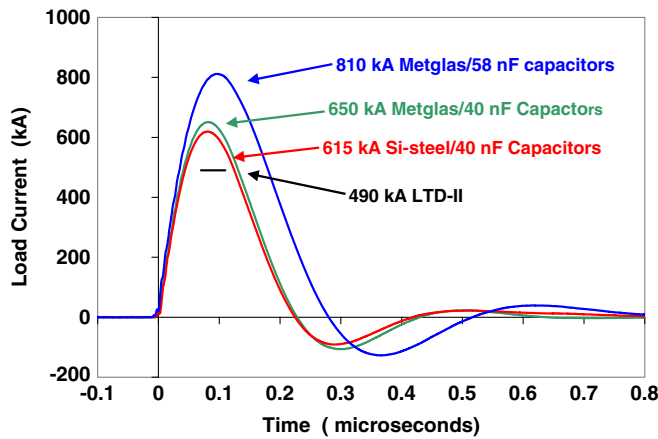


FIG. 8. LTD-III peak currents into matched loads with 40 nF capacitors and Si-steel toroids (red), 40 nF capacitors and Metglas® toroids (green), and 58 nF capacitors and Metglas® toroids (blue). Peak LTD-II current levels are shown for comparison.

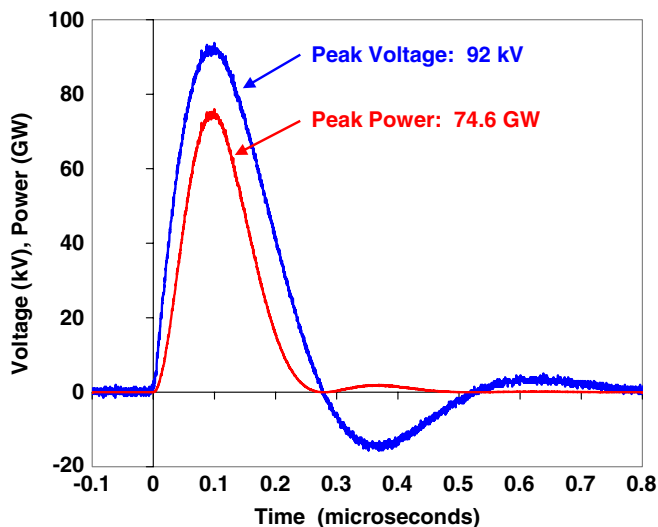


FIG. 9. Load voltage and power waveforms for LTD-III with 58 nF capacitors and Metglas® toroids.

indicating that the effective impedance of the silicon-steel cores decreases with time during the pulse, in agreement with predictions by Kim *et al.* [20]. In contrast, the loss currents with the Metglas® cores peaked before the peak of the power pulse.

It is sometimes useful to refer to an effective “resistance” of the magnetic cores [8] defined simply as the load voltage divided by the loss current. With the silicon-steel cores, the effective core resistance at peak power was 1.5Ω . With the Metglas® cores and 40 nF capacitors, the effective resistance at peak power was 7.2Ω . With the same Metglas® cores and 58 nF capacitors, the effective resistance was 1.9Ω . All of these values are much larger than the $0.11\text{--}0.15\text{-}\Omega$ resistance of the flowing resistive load.

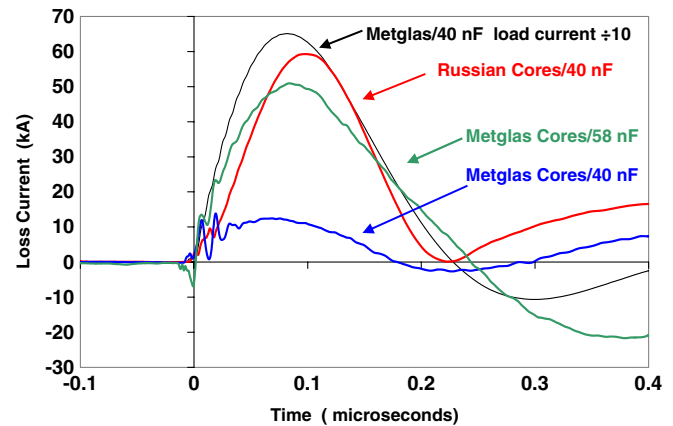


FIG. 10. Loss currents for the three experiments. A scaled current trace for one of the main current waveforms is shown for comparison.

In ± 100 kV shots, the total volt-second product impressed across the magnetic cores for both tests with 40-nF capacitors was 11.5 mV-seconds, well below the 22 mV-second product of the silicon-steel cores or the 18 mV-second product of the Metglas® cores. The volt-second product impressed across the Metglas® cores during the tests with 58 nF capacitors was 14.25 mV, which was still lower than the 18 mV-second theoretically available. The dramatic increase in loss current in changing from 40- to 58-nF capacitors, however, suggests that we are nearing saturation of the Metglas® cores in the 58-nF shots.

Since we were interested in understanding the inductance of the LTD-III cavity, we took several low-voltage shots of LTD-III into a short-circuit load. To form a short-circuit load, we drained the resistive load and then wrapped a thin brass foil around the entire circumference of the inner load insulator, providing a low-inductance short circuit between the top and bottom plates. Figure 11 shows

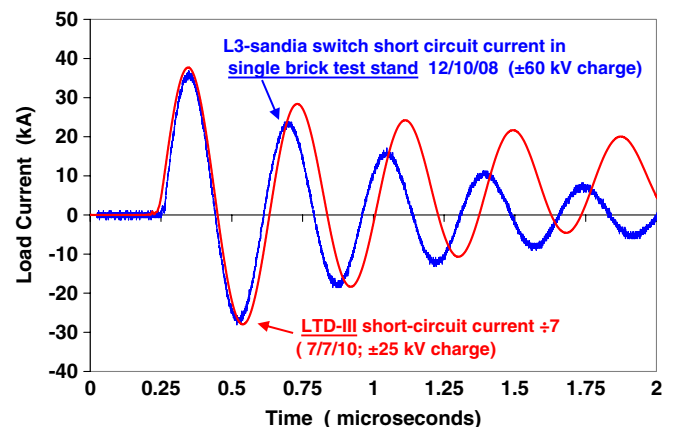


FIG. 11. Low-voltage, short-circuit currents in LTD-III compared with earlier data from our single brick test stand. Most of the increase in ringing period in LTD-III compared to the single brick test stand is caused by the thicker and longer spacer used to separate the capacitors on LTD-III.

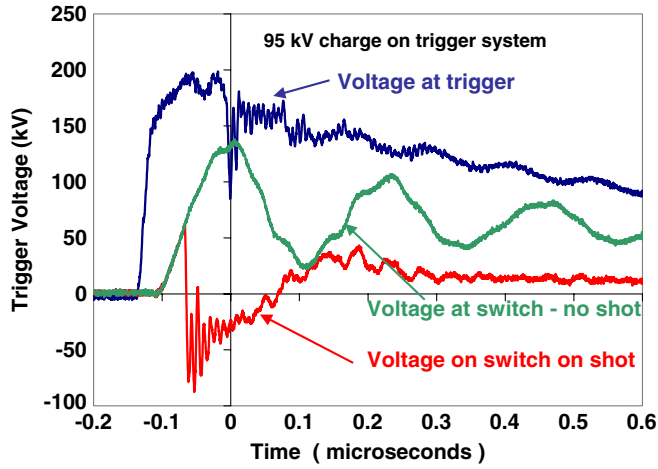


FIG. 12. Comparison of LTD-III trigger voltages measured at the trigger generator (blue), at a switch when no voltage was on the switch (green), and at a switch during an LTD-III shot (red).

the current into the short-circuit load with 40 nF capacitors in LTD-III charged to ± 25 kV. For comparison, Fig. 11 also shows a short-circuit trace from our earlier experiments with a single LTD brick [11]. Based on the ringing period of 0.38 microseconds, we estimate the total inductance of the shorted cavity to be 9.1 nH. In our single LTD brick experiments, we measured an inductance per brick to 155 nH. Dividing this inductance by 20 to simulate the effect of 20 single bricks in parallel gives a simulated inductance of 7.75 nH for the 20 bricks in parallel, which is significantly smaller than the 9.1 nH that we measure for LTD-III.

Figure 12 shows voltage waveforms of the trigger pulse measured at the trigger generator and *at a switch* inside LTD-III. The measurement at the switch shows two conditions: first, when no charge voltage was applied across the switches, allowing the trigger pulse to ring and, second, during a normal shot of the cavity. The trigger pulse measured at the generator has a 10%–90% rise time of ~ 20 ns, but this trigger rise time has increased to ~ 60 ns by the time the trigger actually reaches a switch. While the trigger pulse still has more than enough voltage to trigger the switches, this increase in rise time is likely to adversely

affect the simultaneity of the triggering of the twenty switches as well as the overall cavity jitter.

IV. DISCUSSION

In order to better understand the sources of inductance in the system, we assembled an “inductance budget” for the LTD-III cavity with 40 and 58 nF capacitors and compared it to estimates from our previous single brick experiments. In our inductance budget, we used known component inductances and zeroth-order estimates of the inductances caused by the spacing of power feeds in the cavity [21]. The inductances of the capacitors were obtained from the manufacturer. The inductances of the switches are from our previous work [11]. Table II summarizes our simulated inductances and then compares them to experimental results from short-circuit shots. Our calculated inductances are 5%–10% lower than the experimental values from the short-circuit shots. The increased thickness (1.4 cm vs 0.64 cm) and length (31.5 cm vs 19.7 cm) of the spacer between the capacitors accounts for most of the difference in inductance between our single brick experiments and the LTD-III cavity. Note that even with our improved switches and capacitors, over half of the circuit inductance is still in the switch. 45 nH of this inductance is caused by the switch arcs themselves and 40 nH is caused by the inductance of the current loop from the top capacitor, through the switch and back to the bottom capacitor.

The flowing resistive load was made up of a concentrated solution of potassium bromide in water. Typical resistivities ranged from 6 Ω -cm to 8 Ω -cm. These concentrated alkaline solutions rapidly degraded conductivity measuring probes [22] unless the measurement was taken in less than 10 seconds and then the probe was rinsed in deionized water then vinegar and then stored in deionized water. Load resistances determined from measuring the resistivity and then calculating resistance based on the load geometry typically were about 15% higher than load resistance values obtained using Ohm’s law with the load voltage and current traces.

Large pulsed power drivers with 100–1000 LTD cavities will need each cavity to operate for 1000 shots without

TABLE II. Inductance of components in LTD systems. All inductances are in nano-Henries.

Parameter for single brick	Single LTD brick experiments	LTD-III with 40 nF capacitors	LTD-III with 58 nF capacitors
Switch loop	40	40	40
Switch arcs	45	45	45
Two capacitors in series	30	30	30
Spacer between capacitors	9.6	30.4	36
Load resistor	22.6	20.6	20.6
Total inductance per brick	147	166	171.6
Inductance of 20 bricks in parallel	7.36	8.3	8.58
Experimental value	7.75	9.1	Not measured

maintenance [23]. In our proof-of-principle experiments with the LTD-III cavity, we have not demonstrated this level of reliability. There are several areas where more work needs to be done. First, replacing the liquid resistors used in LTD-III is a high priority since liquid resistors change their values over time whether or not they are used. Second, air bubbles in the oil were a constant problem in these experiments and led to our pumping on the oil and recirculating it. New cavity designs that carefully consider how all the oil bubbles can be removed from the system would be very helpful. Finally, it is not clear that we are using the best transformer oil for our experiments. While Luminol was chosen over Shell Diala-X because of its higher breakdown strengths in small areas under AC electric fields, more tests need to be done with larger areas and DC fields. More studies of the effects of dissolved air and water on the oil breakdown strengths would also be helpful.

V. SUMMARY

We have performed three experiments with an existing 490 kA LTD cavity to increase its peak current and better understand its operation. In the first experiment, we replaced the existing Russian switches and single-ended capacitors with improved L3-Sandia switches and double-end capacitors, significantly lowering the inductance of the cavity. This decrease in inductance raised the peak current to a matched load from 490 to 615 kA. In the second experiment, we replaced the existing silicon-steel magnetic cores with Metglas® cores, which lowered the core losses and increased the peak current to 650 kA. In the third experiment, we replaced the 40 nF capacitors with 58 nF capacitors, raising the peak current to a matched load to 811 kA. These experiments may lead to improved pulsed-power accelerators for flash radiography, dynamic materials experiments, and inertial-confinement fusion research.

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