



Competition between fractional quantum Hall liquid and Wigner solid at small fillings: Role of layer thickness and Landau level mixing

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What is the fate of the ground state of a two-dimensional electron system at very low Landau level filling factors (ν) where interaction reigns supreme? An ordered array of electrons, the so-called Wigner crystal, has long been believed to be the answer. It was in fact the search for the elusive Wigner crystal that led to the discovery of an unexpected, incompressible liquid state, namely the fractional quantum Hall state at $\nu = 1/3$. Understanding the competition between the liquid and solid ground states has since remained an active field of fundamental research. Here we report experimental data for a new two-dimensional system where the electrons are confined to an AIAs quantum well. The exceptionally high quality of the samples and the large electron effective mass allow us to determine the liquid-solid phase diagram for the two-dimensional electrons in a large range of filling factors near $\simeq 1/3$ and $\simeq 1/5$. The data and their comparison with an available theoretical phase diagram reveal the crucial role of Landau level mixing and finite electron layer thickness in determining the prevailing ground states.

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The ground state of an interacting, low-disorder two-dimensional electron system (2DES) in a large perpendicular magnetic field (B) has been of continued interest for decades. This is a regime where the kinetic (Fermi) energy of the electrons is quenched as they all occupy the highly degenerate, lowest Landau level (LL), and the interaction (Coulomb) energy dominates. As the 2DES enters deep into this extreme quantum limit, the LL filling factor ($\nu = nh/eB$) becomes very small; n is the 2DES density. At sufficiently small fillings $\nu \ll 1$, the ground state should be a magnetic-field-induced Wigner solid (WS), where the electrons arrange themselves in an ordered (triangular) array to minimize their mutual Coulomb potential energy [1,2]. Experiments on a relatively low disorder GaAs 2DES, however, revealed a new phase of matter at $\nu = 1/3$, namely the fractional quantum Hall state (FQHS), an incompressible liquid state with a vanishing longitudinal resistance and a quantized Hall resistance [3]. The search for the WS then moved to even lower fillings, and culminated in the observation of insulating phases flanking a well-developed FQHS at $\nu = 1/5$ in very high quality 2DESs [4,5]. The nonlinear I - V characteristics of these insulating phases [5], as well as resonances in their microwave spectrum [6], were interpreted as evidence for the formation of a WS, which is pinned by the ubiquitous disorder that is present in all real samples. This interpretation has been corroborated

by numerous theoretical calculations [7–10] and a variety of experiments [11–21].

Shortly after the observation of insulating phases near $\nu = 1/5$ in GaAs 2DESs [4,5], Santos *et al.* reported similar insulating phases in a dilute GaAs 2D *hole* system (2DHS) but there they flanked the $\nu = 1/3$ FQHS [22,23]. Another system, where insulating phases next to the $\nu = 1/3$ FQHS have been reported, is the ZnO 2DES [21]. These reports provided experimental evidence for the importance of LL mixing (LLM) in favoring the WS states over the FQH liquid states, and moving the stability of the WS to fillings as large as $\simeq 1/3$. Note that the 2D holes in GaAs and 2D electrons in ZnO have an effective mass (m^*) which is larger than m^* of GaAs 2D electrons by a factor of ~ 6 [24], leading to a larger LLM parameter κ , defined as the ratio of the Coulomb energy and the LL separation: $\kappa = (e^2/4\pi\epsilon_0\epsilon l_B)/(\hbar eB/m^*)$, where $l_B = \sqrt{\hbar/eB}$ is the magnetic length, and ϵ is the dielectric constant. The importance of LLM in the competition between the liquid and solid states at low ν has indeed been highlighted by a number of theoretical [25–32] and additional experimental studies [33–40].

Here we report the observation of FQH and WS phases in 2DESs confined to AIAs quantum wells (QWs). These samples, which have exceptionally high quality, provide an ideal platform to explore the liquid-solid transitions for a large range of κ and ν . Electrons in AIAs QWs have larger m^* ($\simeq 0.46$ in units of free electron mass) and smaller ϵ ($\simeq 10$) compared to GaAs 2DESs ($m^* \simeq 0.067$ and $\epsilon \simeq 13$). As a result, the κ parameter for AIAs 2DESs, for a given density and filling factor, is $\simeq 9$ times larger than in GaAs 2DESs. Furthermore, the LLs are simple and their energies vary linearly with B , making the interpretation of κ straightforward. This is in contrast to the much more complex LLs in the 2DHS

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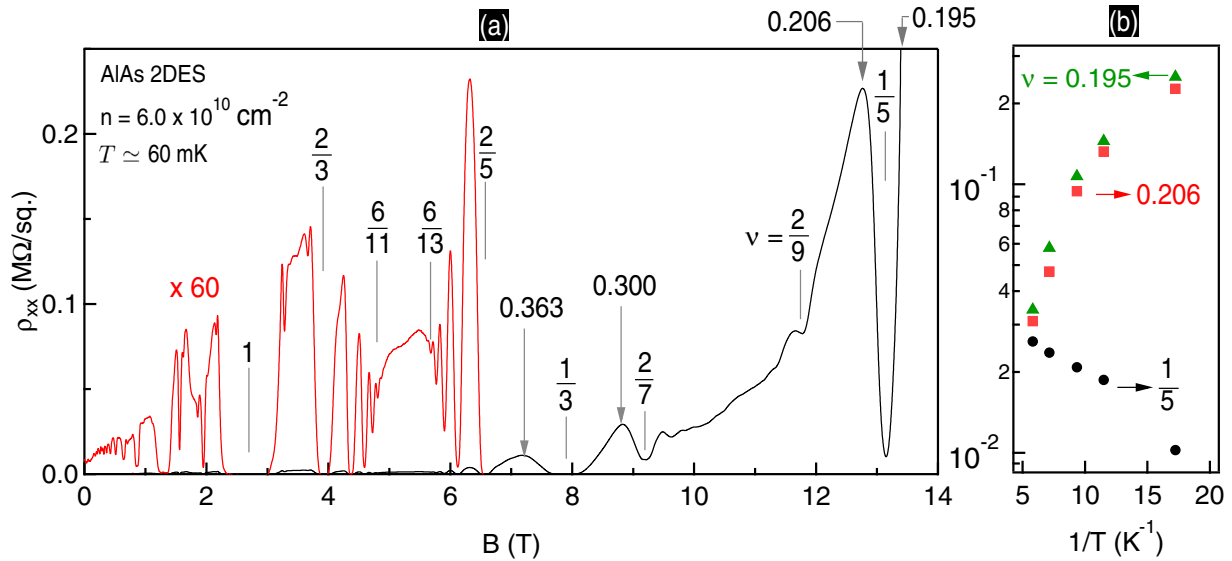


FIG. 1. (a) Longitudinal resistivity ρ_{xx} vs perpendicular magnetic field B for an AlAs 2DES with $n = 6.0 \times 10^{10} \text{ cm}^{-2}$ at $T \simeq 60 \text{ mK}$. The red trace is expanded by a factor of 60. Numerous FQHSs, marked by vertical lines, attest to the extremely high quality of the sample. (b) Arrhenius plots of ρ_{xx} at $\nu = 0.206$, $1/5$, and 0.195 . The insulating phases on the flanks of $\nu = 1/5$ signal a disorder-pinned Wigner solid.

which are nonlinear and have crossings as a function of B [39,41]. Our data provide a κ vs ν phase diagram for the liquid-solid states in a large range of fillings ($1/5 \lesssim \nu \lesssim 1/3$) and LLM ($4 \lesssim \kappa \lesssim 12$); note that the typical value for κ in GaAs 2DESs is $\simeq 1$. We compare this diagram quantitatively to the theoretically available diagrams near $\nu = 1/3$ and $1/5$ [31]. The overall agreement is quite good and demonstrates the role of LLM, *as well as* finite-layer thickness; however, some subtle discrepancies remain. We also report a thermal melting phase diagram, near $\nu = 1/3$, for the pinned WS.

We studied 2DESs confined to modulation-doped AlAs QWs grown on GaAs (001) substrates. In the absence of symmetry-breaking, in-plane strain, our samples, which have QW widths (w) ranging from 20 to 50 nm, host electrons occupying two anisotropic in-plane valleys. The valleys have longitudinal and transverse effective masses $m_l = 1.1$ and $m_t = 0.20$, leading to an effective cyclotron mass $m^* = (m_l \times m_t)^{1/2} = 0.46$ [42]. The samples have 2DES densities (n) ranging from 3.3 to 33, in units of 10^{10} cm^{-2} , which we will use throughout the paper. They have a van der Pauw geometry, with alloyed InSn contacts at the corners and edge midpoints. We made measurements primarily in a dilution refrigerator, and used low-frequency ($\leq 7 \text{ Hz}$) lock-in techniques: constant-current mode for magnetoresistance and constant-voltage mode for differential resistance. We provide further details of sample parameters and device fabrication in the Supplemental Material (SM) [43].

In our samples, in the absence of uniaxial in-plane strain and electron-electron interaction, the two in-plane valleys are degenerate and should be equally occupied at zero magnetic field. The application of uniaxial strain, either intentional or unintentional, breaks this degeneracy and results in unequal occupation of the valleys [42,44]. If the uniaxial strain is sufficiently large and the 2DES density is small, the strain can lead to total transfer of electrons into one valley. This is the case in our experiments. Two of our samples (with $n = 33$ and 10) were glued to piezoelectric actuators which were biased

to ensure that all the electrons occupy only in-plane valley [42]. The other samples were attached to linear-dual-pin (DIP) headers using vacuum grease. Consistent with our past experience, these samples also experience some unintentional strain (e.g., see [44]). Thanks to this strain, and the very small 2DES density in our experiments, in these samples too the electrons occupied only one valley; our evidence for this conjecture is that we did not see any signs of valley degeneracy or partial valley occupation in the magnetoresistance data. We would like to add that our study focuses on the liquid to solid transitions in the lowest LLs and at small filling factors. Therefore, even for a minute amount of uniaxial strain all electrons occupy the lowest LL of the lowest energy valley.

Figure 1(a) shows the longitudinal resistivity (ρ_{xx}) vs B at $T \simeq 60 \text{ mK}$ for a sample with $n = 6.0$. The expanded red trace shows several FQHSs, manifested by deep minima in the ρ_{xx} trace. At high magnetic fields, the $\nu = 1/5$ FQHS exhibits a strong minimum. The numerous FQHSs and especially the $\nu = 1/5$ state, which has only been observed in the highest quality GaAs 2DESs [4,5,15,45] and more recently in graphene samples [46,47], attest to the exceptionally high quality of the sample. On the flanks of $\nu = 1/5$ (e.g., at $\nu = 0.206$ and 0.195), ρ_{xx} has very high values. Figure 1(b) shows Arrhenius plots of ρ_{xx} at $\nu = 1/5$ and its flanks. At $\nu = 1/5$, ρ_{xx} decreases as we lower the temperature, consistent with the formation of a FQHS. In contrast, on its flanks, ρ_{xx} increases by about one order of magnitude when we lower the temperature from 200 to 60 mK. Such insulating phases are generally believed to signal a disorder-pinned WS state [4,5,21–23,33,48–51]. Our data and analysis, as we describe below, are consistent with this interpretation. Moreover, they allow us to determine a quantitative quantum melting phase diagram for the magnetic-field-induced WS at low fillings, and compare this diagram to a recently calculated diagram, as reproduced in Fig. 2.

Figure 2 summarizes our experimental results (symbols) in κ vs ν phase diagrams near $\nu = 1/3$ and $1/5$. The gray lines

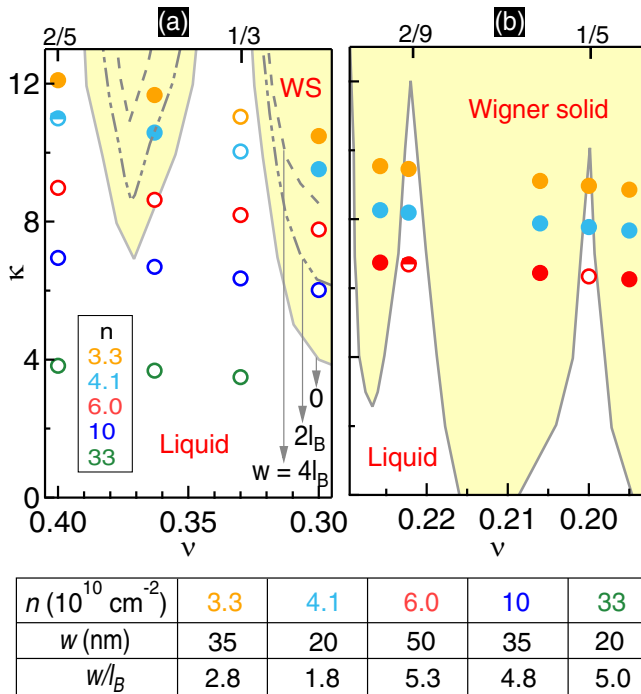


FIG. 2. Liquid-Wigner solid *quantum melting* phase diagrams around (a) $\nu = 1/3$ and (b) $\nu = 1/5$; κ is the Landau level mixing parameter. The gray solid lines represent zero-layer-thickness calculations by Zhao *et al.* [31] for the boundary between the WS and the liquid phases. The closed and open circles indicate the experimentally deduced WS and liquid phases, respectively. The half-filled circles are used to imply a close competition between WS and liquid phases. The filling factors at which the experimental data are presented are $2/5$, 0.363 , $1/3$, 0.300 , $2/9$, 0.206 , $1/5$, and 0.196 . In (a), the dotted-dashed and dashed lines represent the theoretical boundaries between the liquid and WS states for $w = 2l_B$ and $w = 4l_B$, respectively [31]. The table shows the density, QW width, and w/l_B (at $\nu = 1/3$) for the different samples that we studied.

and yellow and white regions are from theoretical calculations [31]. The experimental data are shown by symbols. The open, closed, and semiopen symbols represent a liquid, a disorder-pinned WS, and a close competition between liquid-solid, respectively. Data for $n = 6.0$ are shown in Fig. 2 by red symbols. Around $\nu = 1/5$, in Fig. 2(b), there is overall agreement with the theoretical calculations except for $\nu = 2/9$. At $\nu = 2/9$, as the temperature decreases, ρ_{xx} increases, yet it remains a local minimum [see Fig. 1(a)]. This behavior suggests a competition between the FQH liquid and a disorder-pinned WS. It is likely that disorder is partly responsible for the competition. It is worth remembering that early GaAs 2DESs, which had lower quality, showed a competition between a FQHS and an insulating behavior at $\nu = 1/5$ [52]. When samples of much better quality became available, ρ_{xx} at $\nu = 1/5$ displayed a clear FQHS with a vanishing resistance at the lowest temperatures [4]. Also, previous experimental studies in samples with a controlled (minute) amount of alloy disorder report an insulating behavior at and around $\nu = 1/5$ [53,54]. Local compressibility measurements in a GaAs 2DES showed that disorder creates puddles of slightly different densities, and as a result of different filling factors with a typical varia-

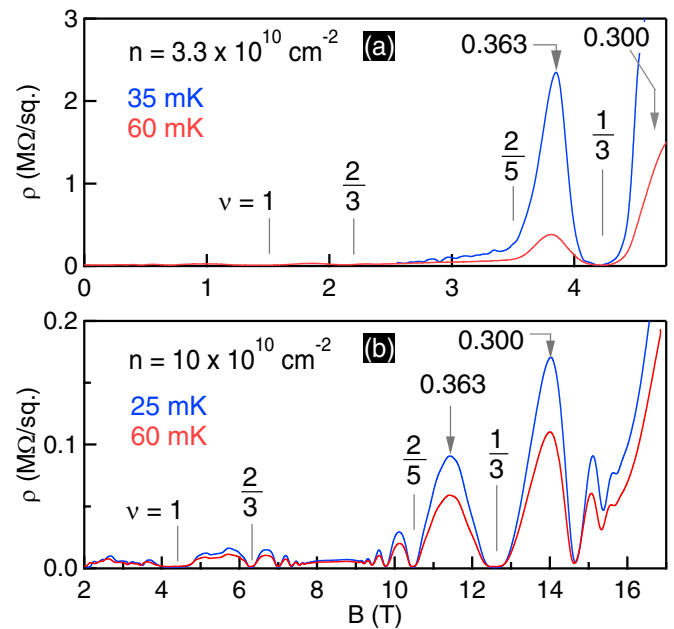


FIG. 3. ρ_{xx} vs B for AlAs 2DESs with (a) $n = 3.3 \times 10^{10} \text{ cm}^{-2}$ and (b) $n = 10 \times 10^{10} \text{ cm}^{-2}$.

tion $\Delta\nu$ [55,56]. In the theoretical phase diagram reproduced in Fig. 2, at large κ , variations in ν would make the WS (yellow region) overcome the $2/9$ FQH liquid (narrow white region), preventing its percolation. For a detailed discussion of the role of disorder in the competition between FQHSs and solids, see Ref. [32].

The red symbols in Fig. 2(a) show our experimental results for $n = 6.0$ near $\nu \simeq 1/3$. The ρ_{xx} trace [see Fig. 1(a)] exhibits fully developed $\nu = 1/3$ and $2/5$ FQHSs, which agree with the prediction of liquid phases [open red symbols in Fig. 2(a)]. As seen in Fig. 1(a), ρ_{xx} maxima on the flanks of $\nu = 1/3$ (at $\nu = 0.363$ and 0.300) are several times larger than the low-field resistance, but do not show strong insulating behavior. This behavior appears to be at odds with the calculations. However, the calculations that delimit the yellow and white regions in Fig. 2 are for a 2DES with zero thickness, but our 2DES is confined to a finite-width QW. Zhao *et al.* [31] reported that for a 2DES with finite layer thickness, the WS-liquid boundary moves to larger values of κ (see SM of Ref. [31]). In Fig. 2(a), we also include the results of finite-thickness calculations from Ref. [31], for $w = 2l_B$ and $4l_B$. As seen in the table shown in the lower part of Fig. 2, our AlAs 2DES with $n = 6.0$ has $w/l_B \simeq 5.3$ at $\nu = 1/3$. Comparing the experimental data with the theoretical $w = 4l_B$ liquid-solid boundary, we find agreement with the liquid states predicted in the calculations.

In order to explore larger values of κ , we studied samples with densities $n = 4.1$ and 3.3 (light-blue and yellow symbols in Fig. 2, respectively). Figure 3(a) shows ρ_{xx} vs B at $T \simeq 35$ and 60 mK for $n = 3.3$ (ρ_{xx} for $n = 4.1$ has similar characteristics; see SM). On the flanks of the $\nu = 1/3$ FQHS, ρ_{xx} has very large values and exhibits a strong insulating behavior, which we interpret as a signature of a disorder-pinned WS. The samples with $n = 3.3$ and 4.1 have $w/l_B \simeq 2.8$ and 1.8 , respectively. When compared to the solid-liquid calculated

boundary with $w = 2l_B$, we find agreement with the predicted WS on the flanks of $1/3$. The behavior at $\nu = 2/5$ is an exception to the agreement: in the experiments we find a WS for $n = 3.3$, and a competition between a FQHS and a solid for $n = 4.1$. For smaller ν , around $1/5$ in Fig. 2(b), there is a strong insulating phase without any hint of FQHSs. As discussed in the previous paragraph, we invoke the role of disorder to explain the above discrepancies with calculations.

We also studied samples with densities $n = 10$ and 33 , to explore smaller values of κ . The data are shown in Fig. 2 by dark-blue and green symbols, respectively. Figure 3(b) shows ρ_{xx} vs B at $T \simeq 25$ and 60 mK for the sample with $n = 10$. We find fully developed FQHSs at $\nu = 1/3$ and $2/5$. At $\nu = 0.363$ and 0.300 , ρ_{xx} does not show a strong insulating behavior, suggesting that there is no WS on the flanks of $\nu = 1/3$. Considering that the sample with $n = 10$ has $w/l_B \simeq 4.8$ at $\nu = 1/3$, the experimental data points agree with the theoretical boundary, once the finite thickness of the 2DES is taken into account. Note that the filling range near $\nu = 1/5$ was experimentally out of our reach for $n = 10$. At $n = 33$, the experimental data near $\nu = 1/3$ are in good overall agreement with calculations (the ρ_{xx} trace is shown in the SM); there is no sign of an insulating phase down to $\nu \simeq 0.32$, the smallest ν we could reach experimentally for this sample in a hybrid magnet.

Two remarks are in order here. First, the WS and FQH states near $\nu = 1/5$ are very close in energy, and the role of finite-layer-thickness could not be assessed reliably in calculations [31]. Second, as seen in Fig. 2(b), the range of ν near $\nu = 1/5$ (as well as near $\nu = 2/9$) where the FQHS is stable, at large κ , is extremely narrow according to the calculations. This may explain why FQHSs are not experimentally seen for the lower density samples: even a very small level of disorder can cause density inhomogeneity (of the order of 1%), which can lead to the disappearance of the $\nu = 1/5$ FQHS. In this context, our observation of a well-developed $\nu = 1/5$ FQHS, with a vanishing ρ_{xx} as the temperature approaches zero, in the $n = 6.0$ sample (Fig. 1) at κ exceeding 6 [Fig. 2(b)] attests to the exceptionally high quality and homogeneity of the sample.

Another property of interest for the WS is its *thermal melting* phase diagram [5,16,20,39,40,45,48–50]. To probe this thermal melting, we monitor the nonlinear behavior of the differential resistance (dV/dI) as a function of a DC current (I_{DC}) flowing through the 2DES with $n = 3.3$. This technique was used in the past to study the magnetic-field-induced WS in other 2D carrier systems [5,22,45,49,50]. Measurement details are outlined in the SM [43].

The Fig. 4 inset shows the temperature dependence of dV/dI vs I_{DC} at $\nu = 0.363$. At the lowest temperatures, the 2DES shows increased conduction for $I_{DC} \neq 0$, signaled by a decrease in dV/dI , and consistent with a disorder-pinned WS [5,22,45,48–50]. At the highest temperatures, where we expect the WS to have melted, dV/dI is independent of I_{DC} , consistent with ohmic conduction observed in liquid phases [5,21,22,38,45,48–50], and in our sample very near the $\nu = 1/3$ FQHS (data not shown). As we raise the temperature, there is a gradual change from the low-temperature, nonlinear behavior to the high-temperature, linear limit. This temperature dependence is qualitatively generic for all the filling factors where we observe a strong insulating behavior in this

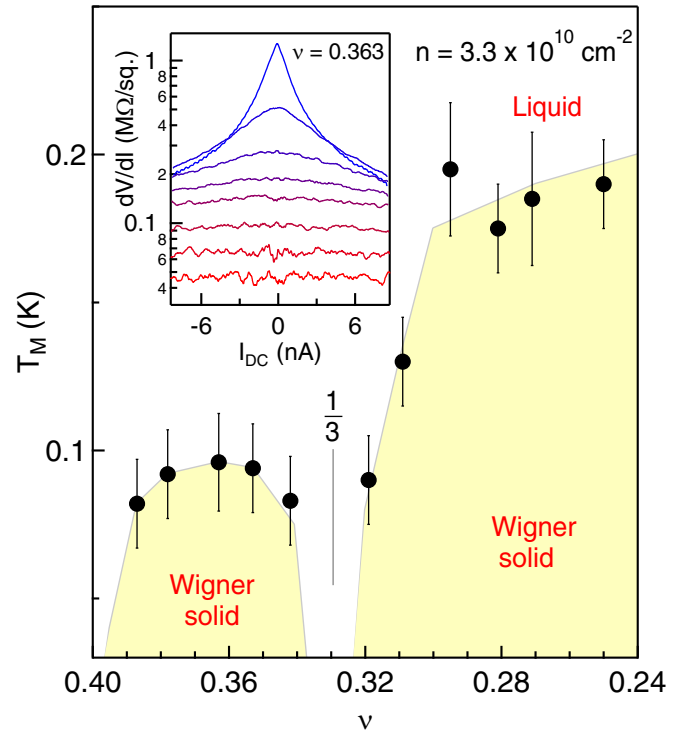


FIG. 4. Wigner solid *thermal melting* phase diagram. The light yellow and white regions indicate the solid and liquid phases, respectively. The gray line is a guide to the eye. The inset shows the temperature dependence of the differential resistance dV/dI vs I_{DC} traces at $\nu = 0.363$. The temperatures are (from top to bottom) 40, 53, 67, 80, 90, 113, 145, and 206 mK.

sample, namely near $\nu \simeq 0.36$ and for $\nu \lesssim 0.32$. We deduce the melting temperature (T_M) of the WS by estimating the temperature at which dV/dI changes from nonlinear to ohmic; see SM for details.

Figure 4, which contains T_M vs ν , displays our deduced WS thermal melting phase diagram of the AIAs 2DES with $n = 3.3$. As ν approaches 0.33, T_M decreases sharply and vanishes close to the $\nu = 1/3$ FQHS. When ν is larger than $1/3$, we observe a reentrant WS. The values of T_M in our thermal phase diagram are close to those recently reported for the melting temperatures of a WS reentrant around $\nu = 1/3$ in a GaAs 2D hole system of density 3.8 [39].

Our experimental data and their quantitative comparison with state-of-the-art calculations provide strong evidence for the validity of associating the insulating phases, observed in ultrahigh-quality 2DESs, at low fillings, on the flanks of FQHSs with disorder-pinned WS states. The data also highlight the importance of LLM, as well as the finite-layer-thickness, in the competition between the correlated solid-liquid phases. The role of finite-layer thickness was not assessed experimentally prior to our study. It is also important to note that the calculations [31] were done for isotropic 2DESs while our data are taken on samples which have an effective mass anisotropy of $\simeq 5$ [42]. The very good agreement between the calculations and the experimental data might thus sound surprising. However, as reported experimentally [57] and theoretically [58–63], it turns that FQHSs are quite

robust against anisotropy, and that (anisotropic) WS states only become prevalent in systems with much larger anisotropies than in our samples. Finally, we emphasize that our data described here were taken in AIAs 2DESs with only one conduction-band valley occupied. An exciting future study would be to experimentally and/or theoretically assess the stability of liquid and solid states in AIAs 2DESs where the electrons occupy two degenerate valleys.

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