

Enhanced doping efficiency of ultrawide band gap semiconductors by metal-semiconductor junction assisted epitaxy

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Ultrawide bandgap semiconductors are important for a broad range of electronic and photonic devices, but their practical application has been limited by poor current conduction. Here, we demonstrate that with controlled tuning of the Fermi level by an *in situ* metal-semiconductor junction during epitaxy, efficient *p*-type conduction, which was otherwise impossible, can be achieved for large bandgap aluminum gallium nitride (AlGa_N). During epitaxy, the Fermi level is pinned away from the valence band edge, which fundamentally reduces the formation energy for substitutional Mg-dopant incorporation while simultaneously increasing the formation energy for compensating defects. We have demonstrated that Mg-doped Al_{0.9}Ga_{0.1}N can exhibit free hole concentration $\sim 4.5 \times 10^{17} \text{ cm}^{-3}$, with a resistivity $< 5 \Omega \cdot \text{cm}$, which is nearly three orders of magnitude lower compared to previous reports. Ultraviolet 280-nm light emitting diodes grown using this method exhibited nearly an order of magnitude improvement in external quantum efficiency compared to those grown using conventional molecular beam epitaxy. Such a unique technique can be extended for the epitaxy/synthesis of a broad range of wide bandgap semiconductors to achieve efficient current conduction that was not previously possible.

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I. INTRODUCTION

Wide bandgap semiconductors such as GaN, AlN, and their alloys have emerged as the materials of choice for high power and high frequency electronic devices [1,2], as well as a broad range of photonic devices including ultraviolet (UV) light emitting diodes (LEDs) [3], lasers [4], and solar blind photodetectors. Critical for the operation and performance of these devices is a precise control of the doping level in the different layers of the structure. To date, however, it has remained extremely challenging to achieve efficient *p*-type conduction of AlN and AlGa_N with relatively high Al content, which has been identified as one of the major obstacles to realizing high performance optoelectronic devices operating in the mid and deep UV spectra. Magnesium (Mg) has been established as the only viable *p*-type dopant of group III-nitride semiconductors [5]. However, it exhibits very large activation energy (up to 600 meV) in Al-rich AlGa_N [6–8], severely limiting the doping efficiency and the realization of large hole concentration at room temperature. Extensive studies have been performed to realize low-resistivity *p*-type AlGa_N, including the use of a high V/III ratio to suppress formation of compensating nitrogen vacancies, superlattices consisting of alternating AlGa_N layers, metal modulation epitaxy, Mg δ doping, indium as a surfactant, and polarization-induced doping, but with very limited success [9–31]. For example, the lowest resistivity reported for *p*-type Al_{0.85}Ga_{0.15}N epilayers

[21] to date is well over $10^3 \Omega \cdot \text{cm}$, which is more than three orders of magnitude larger than that of Mg-doped GaN.

To achieve *p*-type AlGa_N with large hole concentration and low resistivity, it is essential to incorporate large densities of Mg-dopant atoms. At very large concentrations ($\sim 10^{19} - 10^{20} \text{ cm}^{-3}$), an Mg impurity band is expected to form, which enables hole hopping conduction [32,33]. Moreover, the significantly broadened acceptor energy levels at large Mg-doping concentrations, together with the band tailing effect, also reduces the ionization energy for a fraction of Mg dopants. In practice, however, it becomes more difficult to incorporate Mg into AlGa_N with increasing Al concentration, due to the larger formation enthalpy (lower solubility) for Mg substituting Al in the AlGa_N lattice sites, compared to Ga [6,34]. Theoretical calculations in the Supplemental Material [35] confirm the difficulty in substitutional Mg incorporation, especially in Al sites. The formation energy for various compensating point defects also depends critically on the position of the Fermi level [5,6]. During conventional epitaxy, with Mg-dopant incorporation, the Fermi level shifts towards the valence band edge, illustrated in Figs. 1(a) and 1(b), which significantly reduces the formation energy for carbon, oxygen, and nitrogen vacancies [6,9,36,37]. These defects have a strong compensating effect and further degrade the structural, electronic, and optical properties of Mg-doped AlGa_N. In addition, the formation energy of N-substitutional and interstitial Mg incorporation decreases drastically with Mg incorporation, and becomes comparable to that of Al(Ga)-substitutional Mg incorporation when the Fermi level is positioned close to the valence band edge, further limiting the doping efficiency and the achievement of large hole concentrations. Evidently,

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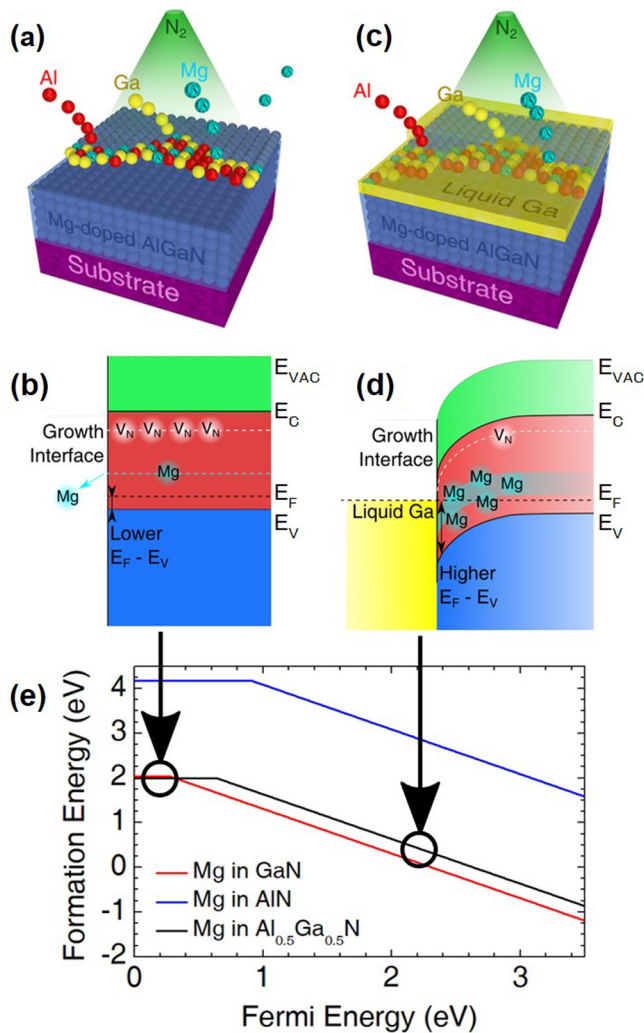


FIG. 1. (a) Schematic of conventional epitaxy. (b) Energy band diagram of the Mg-doped AlGaN layer during conventional epitaxy. (c) Schematic of metal-semiconductor junction assisted epitaxy, with the presence of a liquid Ga layer on the surface during epitaxy. (d) Energy band diagram at the growth front of Mg-doped AlGaN during metal-semiconductor junction assisted epitaxy, showing the pinning of the surface Fermi level away from the valence band edge. (e) Calculated formation energy for Mg substitution in GaN, AlN, and $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}$ as a function of the separation between the Fermi level and the valence band with substitutional Mg formation energies for the different growth processes indicated by their respective arrows.

some of the critical issues for achieving efficient *p*-type conduction of AlGaN can be well addressed, if the Fermi level at the growth front can be tuned away from the valence band during the epitaxy of *p*-type (Mg-doped) AlGaN.

To overcome these fundamental challenges, we investigate a unique epitaxial growth process—metal-semiconductor junction assisted epitaxy—of ultrawide bandgap AlGaN. Illustrated in Fig. 1(c), the epitaxy of AlGaN is performed in metal (Ga) rich conditions by using plasma-assisted molecular beam epitaxy (MBE). The excess Ga layer leads to the formation of a metal-semiconductor junction during the epitaxy of Mg-doped AlGaN, which pins the Fermi level away from

the valence band at the growth front, illustrated in Fig. 1(d). In this unique epitaxy process, the Fermi level position is decoupled from Mg-dopant incorporation, i.e., the surface band bending allows the formation of a nearly *n*-type growth front despite *p*-type dopant incorporation, which is in direct contrast to the fixed Fermi level position near the valence band edge during the conventional epitaxy of Mg-doped AlGaN. As such, the formation energy for substitutional Mg is dramatically reduced, even when very large densities of Mg-dopant atoms are incorporated, which is accompanied by a significant suppression of the formation of compensating defects. Epitaxy of AlGaN under metal-rich conditions was reported previously. However, the role of Fermi-level tuning at the growth front and the resulting effect on the enhanced substitutional incorporation of Mg and suppressed compensating defect formation was not identified [38–41]. Figure 1(e) illustrates the variation of the theoretically calculated substitutional formation energy, for Mg in GaN, AlN, and $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}$, with the energy separation between the Fermi level and valence band maximum. These calculations are explained in the next section, and details of the band structure calculations for AlGaN during growth using the two methods is discussed in the Supplemental Material [35]. By increasing the separation between the Fermi level and valence band, the formation energy is reduced to only ~ 0.43 eV for a metal-semiconductor junction assisted epitaxy of AlGaN, which is ~ 1.6 eV lower than that during growth using conventional epitaxy. This powerful method simply relies on the spontaneous formation of a metal-semiconductor junction at the growth interface by excess Ga and does not involve any modification of the system under use. In this growth regime, the Al composition of AlGaN can be controllably varied by tuning the Al material flux while keeping the nitrogen flow rate constant. This is because the bond strength of Al-N is much stronger compared to that of Ga-N [42,43] and, as a consequence, Al-N will preferably form, whereas any excess Ga will accumulate on the surface to form the metal-semiconductor junction during epitaxy. Utilizing this technique, we have demonstrated that Mg incorporation in AlGaN can be enhanced by nearly one order of magnitude compared to the conventional growth process: Mg concentration $\sim 2 \times 10^{20} \text{ cm}^{-3}$ was measured in $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ for a moderate Mg flux of $\sim 7 \times 10^{-8}$ Torr. A significant reduction of carbon impurity incorporation was also confirmed through detailed secondary ion mass spectrometry (SIMS) measurements. Significantly, a free hole concentration of $\sim 4.5 \times 10^{17} \text{ cm}^{-3}$ was measured for $\text{Al}_{0.9}\text{Ga}_{0.1}\text{N}$, with resistivity values $< 5 \Omega \cdot \text{cm}$, which is nearly three orders of magnitude lower compared to previous reports [21]. Ultraviolet 280-nm light-emitting diodes fabricated using the two methods demonstrated significant improvements in the device characteristics and efficiency for the samples grown by metal-semiconductor junction assisted epitaxy, as compared to the sample grown using the conventional growth mode.

II. METHODOLOGY

A. Theoretical calculations

Density functional theory calculations were performed with the HSE06 hybrid functional [44] and the projector

augmented wave method [45,46], as implemented in the Vienna *Ab initio* Simulation Package (VASP) [47]. GW-compatible pseudopotentials were used, where Mg 3*s*, Al 3*s*3*p*, Ga 4*s*4*p*, and N 2*s*2*p* states were treated as valence electrons. The fraction of nonlocal Hartree-Fock exchange in the hybrid functional is 0.30 for GaN and Al_{0.5}Ga_{0.5}N, and 0.33 for AlN, giving their band gap values of 3.51, 4.65, and 6.20 eV, respectively, in close agreement with the experiments [48–50]. Defect calculations were performed for GaN, Al_{0.5}Ga_{0.5}N and AlN using 96-atom orthorhombic supercells [51] with $2 \times 2 \times 2$ Γ -centered Brillouin-zone sampling and a plane wave energy cutoff of 500 eV. The special quasi-random supercell of Al_{0.5}Ga_{0.5}N was generated by the Alloy Theoretic Automated Toolkit [52]. All structures were relaxed until the force on ions is less than $0.02 \text{ eV } \text{\AA}^{-1}$, with spin-polarization included for unpaired electrons. Corrections to the charged defects were employed, based on the Freysoldt scheme [53], which eliminates the electrostatic interactions between image charges in the supercell approach.

B. Molecular beam epitaxial growth

The samples used in this work were grown on $\sim 1\text{-}\mu\text{m}$ AlN-on-sapphire substrates from DOWA Holdings Co. Ltd. using a plasma-assisted Veeco Gen 930 MBE system. The Mg-doped AlGa_n epilayers were grown at a temperature $\sim 700^\circ\text{C}$ (thermocouple reading), and the Al flux was varied for samples with different Al compositions. The nitrogen flow was 0.4 standard cubic centimeter per min, with a forward plasma power of 350 W. The reflection high energy electron diffraction (RHEED) pattern was closely monitored during epitaxy to determine the growth condition, which was achieved by varying the Ga flux. A relatively dim and streaky RHEED pattern indicated that the growth was under metal-rich conditions, whereas a brighter pattern with segmented streaks suggests the growth is in slightly N-rich conditions.

The growth was initiated with an undoped AlN buffer layer, followed by $\sim 30\text{-nm}$ undoped AlGa_n, before the growth of the Mg-doped AlGa_n layer. The thickness of the AlGa_n layers for Hall measurement was approximately 430 nm. For the samples grown for Hall measurements, the growth is terminated with a $\sim 2\text{-nm}$ *p*-Ga_n capping layer, which was subsequently etched in the fabrication process for Hall measurements.

For the samples on which SIMS was performed, the grown structure consisted of several AlGa_n layers grown with different Mg fluxes, with undoped layers in between. The Mg flux used in the different layers of the samples on which SIMS was performed was measured at different temperatures before every growth, using a beam flux monitor equipped in the MBE system.

The LED heterostructure consists of a 250-nm thick Si-doped Al_{0.7}Ga_{0.3}N layer, multiple Al_{0.45}Ga_{0.55}N/Al_{0.7}Ga_{0.3}N quantum wells, 60-nm-thick Mg-doped AlGa_n layer, and 3-nm *p*-Ga_n contact layer. The first 30-nm Mg-doped AlGa_n layer was graded from an Al composition $\sim 70\%$ to 50% to make use of polarization induced doping to maximize hole injection into the active region [54], followed by a 30-nm-thick Mg-doped Al_{0.5}Ga_{0.5}N. An AlGa_n electron blocking

layer was also incorporated to reduce electron overflow. The device active region was calibrated for emission at $\sim 280 \text{ nm}$.

C. Characterization

The composition of AlGa_n epilayers was derived from Vegard's law based on x-ray diffraction measurements, which were performed using a Rigaku SmartLab x-ray diffractometer. The sample surface was examined by scanning electron microscopy (SEM) using a Hitachi SU8000 cold field emission SEM. Atomic force microscopy (AFM) scans of the sample surface was taken by a Bruker Dimension Icon AFM, operated under tapping mode in air, to determine the surface roughness. SIMS was performed by EAG Laboratories to derive Mg-dopant atom concentration and C-impurity atom concentration. Photoluminescence spectroscopy measurements were performed at room temperature using a 193-nm ArF excimer laser as the excitation source.

D. Device fabrication

To fabricate the samples for Hall measurements, metal contacts consisting of Ni (20 nm)/Al (100 nm)/Au (20 nm) were defined using photolithography and deposited by e-beam evaporation, followed by annealing at 500°C for 5 min in the presence of air. The thin *p*-Ga_n contact layer was subsequently isolated and removed. The UV LED fabrication process involves the use of standard photolithography, dry etching, and contact metallization techniques. The device mesa has an area $50 \mu\text{m} \times 50 \mu\text{m}$. A Ti (40 nm)/Al (120 nm)/Ni (40 nm)/Au (50 nm) metal stack was deposited on *n*-AlGa_n and annealed at 750°C for 30 sec in nitrogen ambient to form the *n*-metal contact. A Ni (20 nm)/Al (100 nm)/Au (20 nm) metal stack was deposited on *p*-AlGa_n and annealed at 500°C for 5 min in air to form the *p*-metal contact.

III. RESULTS AND DISCUSSION

A. Effect on Mg incorporation

We have studied the extraordinary effect of a metal-semiconductor junction on Mg-dopant incorporation during the epitaxy of AlGa_n with Al composition $\sim 75\%$. Sample A was grown under Ga-rich conditions, using metal-semiconductor junction assisted epitaxy, to ensure coverage of the substrate surface with metallic Ga during growth, while Sample B was grown under nearly stoichiometric conditions using the conventional growth mode but with otherwise identical conditions, e.g., the same growth rate and same Mg fluxes, for corresponding layers. Subsequently, Mg concentrations were obtained using SIMS measurements. Details about the Ga coverage during epitaxy are presented in the Supplemental Material [35]. Figure 2(a) shows the Mg atom concentration profile for the two samples, wherein the different Mg-doped AlGa_n layers are separated by undoped layers. It is seen that, with the use of metal-semiconductor junction assisted epitaxy, Mg concentrations are significantly higher than those grown using the conventional growth method. Variations of the Mg incorporation vs Mg flux, measured as beam equivalent pressure (BEP) for the two samples are further plotted in Fig. 2(b). For both samples, Mg concentration increases with Mg BEP.

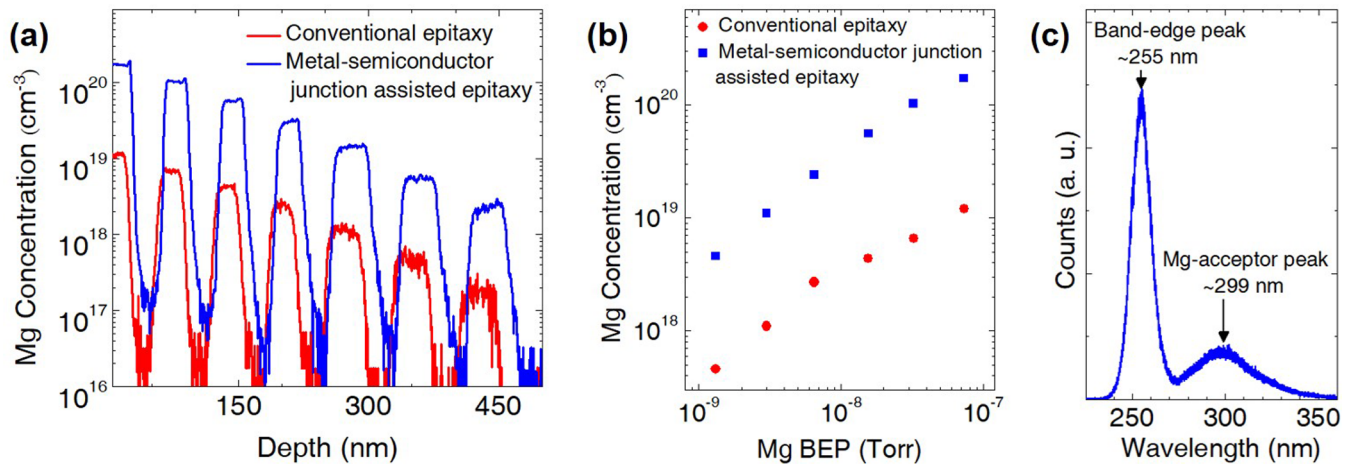


FIG. 2. (a) Mg atom concentration versus depth obtained from SIMS measurements on Mg-doped $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ grown using metal-semiconductor junction assisted epitaxy and conventional epitaxy. (b) Mg concentration versus the Mg flux for Mg-doped $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ grown using metal-semiconductor junction assisted epitaxy and conventional epitaxy. (c) Photoluminescence spectrum of a Mg-doped $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ sample grown using metal-semiconductor junction assisted epitaxy with band-edge peak and Mg-acceptor peak shown with arrows.

However, the Mg atom density is over an order of magnitude higher for Sample A grown with metal-semiconductor junction assisted epitaxy, as compared to Sample B using the conventional process. The maximum Mg incorporation achieved for Sample A was $\sim 2 \times 10^{20}$ Mg atoms cm^{-3} , without showing any sign of saturation. The significantly enhanced Mg incorporation is attributed to the reduced formation energy for Al(Ga)-substitutional Mg incorporation [5,6,37] when the Fermi level is pinned away from the valence band edge by utilizing the metal-semiconductor junction at the growth front. We would also like to point out that the C-impurity concentration was significantly reduced compared to metal-organic chemical vapor deposition (MOCVD). For the sample grown using metal-semiconductor junction assisted epitaxy, the carbon concentration is limited by the measurement background of SIMS ($\sim 1 \times 10^{16}$ cm^{-3}). For comparison, carbon concentrations $\sim 5 \times 10^{16}$ cm^{-3} to 2×10^{18} cm^{-3} have been commonly measured in Al-rich AlGaIn grown by MOCVD [55]. The pinning of the Fermi level at the growth front through metal-semiconductor junction assisted epitaxy also leads to a significant decrease in the formation of point defects, which explains the observation that undoped (Al)GaIn layers grown under Ga-rich conditions showed approximately three orders of magnitude higher resistivity compared to films grown under conventional conditions [6,38].

B. Photoluminescence characteristics

We further studied the photoluminescence properties of the sample grown using metal-semiconductor junction assisted epitaxy. A typical photoluminescence spectrum for a Mg-doped $\text{Al}_{0.75}\text{Ga}_{0.25}\text{N}$ sample is shown in Fig. 2(c). A strong peak near the band edge at ~ 255 nm (4.86 eV) was measured, accompanied with the Mg-acceptor related transition at ~ 298 nm (4.16 eV). It has been previously reported that this emission originates from a donor-acceptor pair transition in Mg-doped AlGaIn epilayers [39]. The broad linewidth of the Mg-acceptor related transition and its partial overlap with the band-edge luminescence emission indicates a very large

spread of the Mg-acceptor energy levels, extending nearly to the valence band edge of AlGaIn. Evidently, the significantly enhanced Mg dopant incorporation can not only lead to the formation of an impurity band for hole hopping conduction [31–33], but more importantly, results in substantially reduced activation energies for a portion of Mg-dopants, thereby enabling the presence of large hole carrier concentrations at room temperature in AlGaIn that were not possible otherwise [56].

C. Low-resistivity *p*-type AlGaIn with high Al content

By employing the metal-semiconductor junction assisted epitaxy, we have also grown a series of Mg-doped samples with Al compositions varying from $\sim 75\%$ up to $\sim 90\%$. Atomic force microscopy measurements indicated a smooth surface with roughness below 1 nm for all the samples. Hall measurements were performed on the samples, using the van der Pauw method, to determine the hole concentration, hole mobility, and resistivity of the AlGaIn layers for temperatures ranging from room temperature to 500 °C. The room temperature hole concentration is observed to monotonically decrease with increasing Al content, illustrated in Fig. 3(a). Figure 3(b) shows variations of hole mobility with the Al composition. The decrease in hole concentration and mobility with increasing Al composition can be explained by a reduction in Mg incorporation for alloys with a higher Al mole fraction due to the lower solubility of Mg and the increased formation enthalpy for Al-substitutional Mg [5,57]. Even for the incorporated Mg atoms, the increase in the activation energy of the Mg-acceptor with increasing Al content further shrinks the free hole concentration. This reduction in Mg concentration with increasing Al composition has been further verified through SIMS measurements of AlGaIn samples grown with different alloy compositions, but with the same Mg flux, shown in the Supplemental Material [35]. These factors result in an increase in the resistivity for Mg-doped AlGaIn layers with increased Al composition, shown in Fig. 3(c). However,

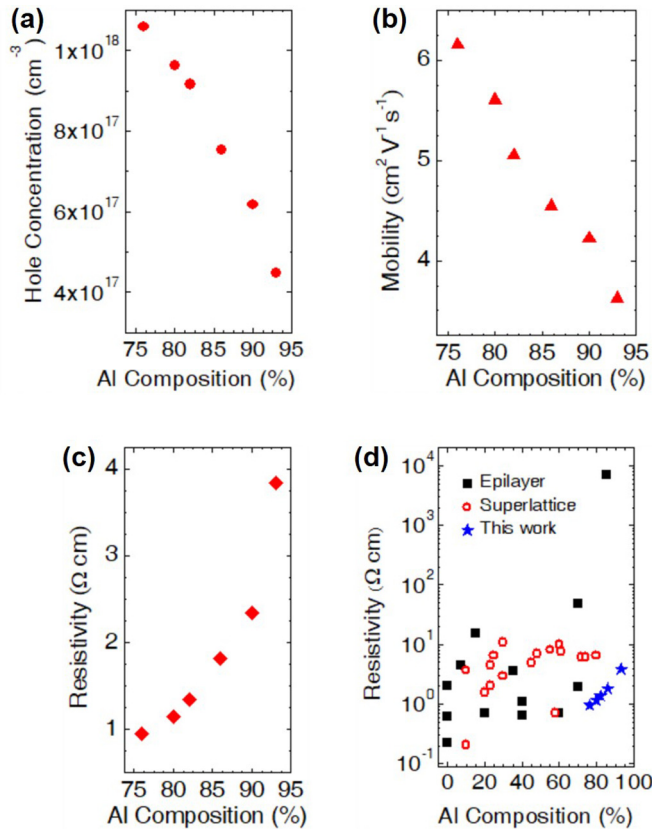


FIG. 3. Room-temperature Hall measurement data for Mg-doped AlGaN epilayers plotted against Al composition, showing (a) free hole concentration, (b) hole mobility, (c) resistivity, and (d) resistivities of Mg-doped AlGaN layers obtained from literature and this work.

even for Al_{0.9}Ga_{0.1}N, the measured resistivity remains below 5 Ω·cm. The resistivity values of some previously reported Mg-doped AlGaN layers [9–31] are plotted vs Al compositions in Fig. 3(d), along with the resistivity of the AlGaN layers obtained in this work. It is seen that the resistivities of Mg-doped AlGaN grown using metal-semiconductor junction assisted epitaxy is nearly one to three orders of magnitude lower compared to previously reported results.

D. Hole conduction mechanism

We have further measured variations of the hole concentration vs temperature for Mg-doped AlGaN samples with different Al compositions, illustrated in Fig. 4(a). A low activation energy (~10–20 meV) at temperatures less than 600 K is seen in these samples, which is characteristic of hole hopping conduction in the impurity band [13,14] and can also be partly explained that a portion of the Mg dopants have significantly reduced activation energy, evidenced by the photoluminescence spectrum shown in Fig. 2(c) [31,58]. At higher temperatures, more Mg-dopants get activated, contributing to holes in the valence band. This leads to a sharp increase in hole concentration at elevated temperatures (>650 K), which is characterized by a large activation energy (~300–400 meV). Such activation energy values, however, are somewhat lower than those theoretically expected [6] for AlGaN alloys with Al

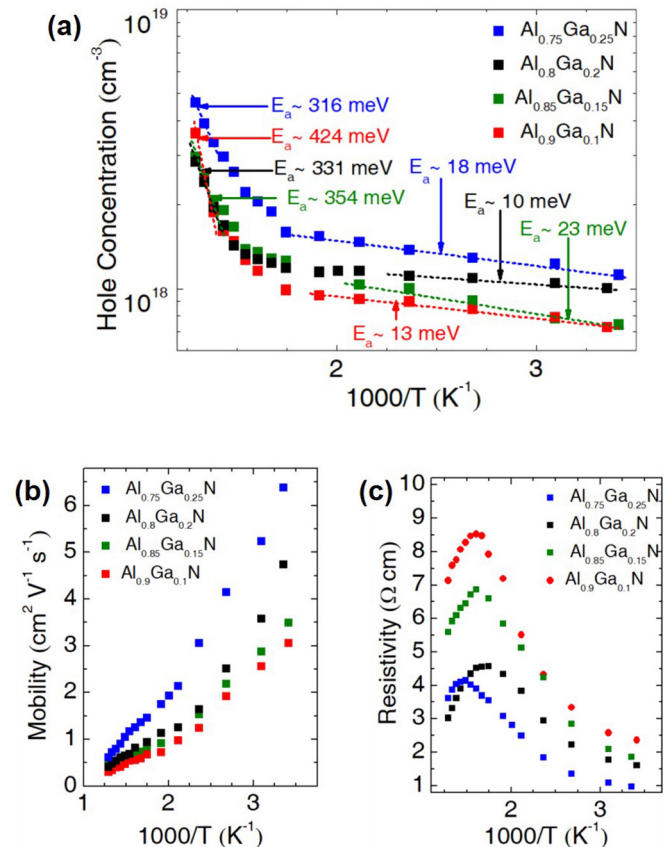


FIG. 4. Temperature dependent Hall measurement of Mg-doped AlGaN epilayers grown by metal-semiconductor junction assisted, with Al content between 75% and 90%, for (a) hole concentration, (b) hole mobility, and (c) resistivity plotted against the inverse of temperature.

compositions ~75–90%. This can be explained by the presence of band-tailing effects and the significantly broadened acceptor energy level distribution, which effectively reduces the activation energy for a portion of Mg acceptors. The measured hole mobility, shown in Fig. 4(b), has a monotonically decreasing trend with increasing temperature, as expected due to an increase in phonon scattering. The resistivity, shown in Fig. 4(c), is first observed to increase with temperature between 300 and ~650 K due to the decrease in hole mobility and relatively small change in hole concentration. At higher temperatures, when the Mg acceptors get thermally activated, the dramatic rise in hole concentration results in the observed decrease in resistivity.

E. UV LED characteristics

The significantly reduced resistivity of Mg-doped AlGaN, enabled by the unique metal-semiconductor junction assisted epitaxy, is crucial to improve the efficiency of optoelectronic devices operating in the mid- and deep-UV wavelengths. We have fabricated and characterized AlGaN UV LEDs, emitting at ~280 nm, grown using this unique approach, and compared the device characteristics with identical LEDs grown using conventional epitaxy. A typical electroluminescence spectrum is shown in Fig. 5(a), with a narrow linewidth ~11 nm. The

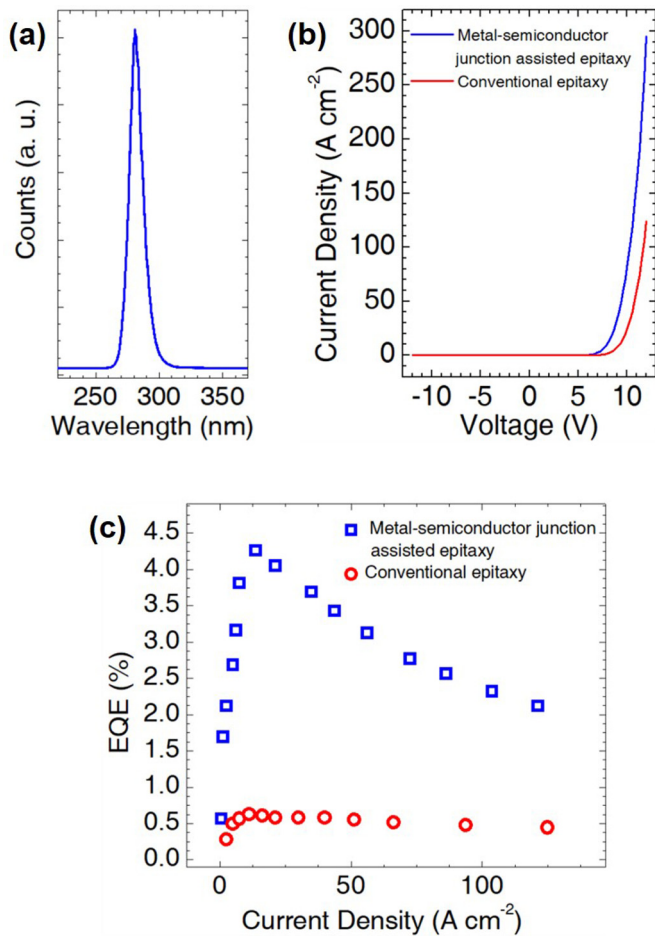


FIG. 5. (a) Electroluminescence spectrum for an UV LED grown using metal-semiconductor junction assisted epitaxy. (b) I - V characteristics for LEDs grown using metal-semiconductor junction assisted epitaxy and conventional epitaxy. (c) EQE versus current density for these devices.

current-voltage characteristic is shown in Fig. 5(b), which exhibits a turn-on voltage of ~ 7 V for the LED grown using metal-semiconductor junction assisted epitaxy, while we see a turn-on voltage ~ 9 V for the conventional epitaxy sample. The higher turn-on voltage is likely a result of the ineffective Mg-doping when using conventional epitaxy. The external quantum efficiency (EQE) was further measured directly on wafer, without any packaging, substrate removal, or cooling. Shown in Fig. 5(c), a maximum on-wafer EQE of $\sim 4.3\%$ was measured at room temperature, which is significantly better than the LED device grown using the conventional epitaxy process, which exhibits a maximum EQE of $\sim 0.6\%$. This device performance is also better than other previously reported AlGa_N UV LEDs grown by MBE around this wavelength [38,39,59,60]. The improved device characteristics seen in

the samples grown using metal-semiconductor junction assisted epitaxy highlights the significance of efficient p -type conduction on device performance. It is also worth noting that the measured EQE can be significantly increased through proper packaging of devices to increase the light extraction efficiency, and by employing more comprehensive techniques to capture all the emitted light, such as the use of an integrating sphere.

IV. CONCLUSION

In summary, we have demonstrated that, by tuning the surface Fermi level using metal-semiconductor junction assisted epitaxy, efficient p -type conduction can be achieved for Al-rich AlGa_N that was not previously possible. The presence of a metal-semiconductor interface at the growth front pins the Fermi level away from the valence band edge, which can significantly enhance Al(Ga)-substitutional Mg-dopant incorporation and further reduces the formation of compensating point defects, as demonstrated in this work, both theoretically and experimentally. It should also be noted that the presence of surface states, which are strongly affected by growth conditions, as has been previously described for both polar [61] and nonpolar [62] surfaces, may further play a role in pinning the Fermi level away from the valence band, although the surface state density structure at the elevated temperatures required for crystal growth has remained unknown. As such, large concentrations of Mg acceptors can be incorporated in Al-rich AlGa_N, which enables the formation of a Mg impurity band. Al-rich AlGa_N epilayers, with resistivity values below $1 \Omega \cdot \text{cm}$ for Al_{0.75}Ga_{0.25}N and $\sim 4 \Omega \cdot \text{cm}$ for Al_{0.9}Ga_{0.1}N have been measured, which are essentially required for achieving high efficiency mid and deep UV optoelectronic devices. Deep UV LEDs grown using this method showed a great improvement in external quantum efficiency, and lower turn-on voltage, as compared to devices grown using conventional epitaxy. Such a unique technique can be further extended for the epitaxy/synthesis of a broad range of semiconductor nanostructures and heterostructures to achieve controlled dopant incorporation and to fundamentally improve their structural, electronic, and optical properties.

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