

Universal High Energy Anomaly in the Angle-Resolved Photoemission Spectra of High Temperature Superconductors: Possible Evidence of Spinon and Holon Branches

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A universal high energy anomaly in the single particle spectral function is reported in three different families of high temperature superconductors by using angle-resolved photoemission spectroscopy. As we follow the dispersing peak of the spectral function from the Fermi energy to the valence band complex, we find dispersion anomalies marked by two distinctive high energy scales, $E_1 \approx 0.38$ eV and $E_2 \approx 0.8$ eV. E_1 marks the energy above which the dispersion splits into two branches. One is a continuation of the near parabolic dispersion, albeit with reduced spectral weight, and reaches the bottom of the band at the Γ point at ≈ 0.5 eV. The other is given by a peak in the momentum space, nearly independent of energy between E_1 and E_2 . Above E_2 , a bandlike dispersion reemerges. We conjecture that these two energies mark the disintegration of the low-energy quasiparticles into a spinon and holon branch in the high T_c cuprates.

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Understanding how doped oxygen holes are transported in the environment of antiferromagnetically coupled copper spin is one of the most fundamental problems in the field of high temperature superconductivity. It was proposed early on that the doped holes in the oxygen $2p\sigma$ orbitals form singlets with the spins of the neighboring coppers [1]. The resulting charge- e and spin-0 object is called the Zhang-Rice singlet (ZRS). As the ZRS moves through the CuO_2 plane, the copper spins get rearranged. As a result, the ZRS couples very strongly to the antiferromagnetic environment. Remarkably, as a consequence of such strong coupling, quasiparticles emerge at low energies. This is evidenced by the sharp nodal quasiparticle peaks seen in angle-resolved photoemission (ARPES) of almost all cuprate compounds [2,3]. In simple physical terms a quasiparticle is a composite object made of a ZRS and a $S = \frac{1}{2}$ copper spins. It is widely believed that, at sufficiently low temperatures, superconducting pairing occurs between these quasiparticles giving rise to the high temperature superconducting state. Thus a microscopic understanding of the pairing mechanism of high T_c superconductors requires an in-depth understanding of how a ZRS is dressed into a quasiparticle.

Here we present the first systematic study of the evolution of the ARPES spectral function from the Fermi level ($E_F \equiv 0$) to the valence band complex (at energy ≈ 1 eV [4]) for three families of high temperature superconductors. Our results provide a surprising new experimental

understanding on the quasiparticle formation process discussed above. Specifically, by covering a much broader energy range than typically studied [2,3], we have identified anomalies in the ARPES spectra occurring at two universal high energy scales, $E_1 \approx 0.38$ eV and $E_2 \approx 0.8$ eV from E_F . We conjecture that these two energies mark the threshold for the disintegration of the low-energy quasiparticles at two different binding levels.

ARPES data have been collected at the Advanced Light Source, beam lines 7.0.1, 10.0.1, and 12.0.1 for three different families of hole-doped cuprates: single layer $\text{Bi}_2\text{Sr}_{1.6}\text{La}_{0.4}\text{Cu}_2\text{O}_{6+\delta}$ (Bi2201), double layer $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ (Bi2212), and Pb-doped Bi2212 (Pb2212) and for several doping values. The data presented here were measured at least in both the first and the second Brillouin zones (BZs) and along the two polarization p_a and p_b as shown in Fig. 1. The double layer Bi2212 data were collected at 52 eV for the underdoped (UD) ($T_c = 64$ K) and optimally doped (OPT) ($T_c = 91$ K) sample, and 33 and 65 eV for the overdoped (OD) ($T_c = 80$ K) sample. The OD Pb2212 ($T_c = 65$ K) sample was measured at 55, 60, and 75 eV and the OPT Bi2201 ($T_c = 32$ K) sample at 33 eV. Unless specified otherwise, all the data reported were measured at 25 K.

Figure 1 shows the ARPES intensity map as a function of energy and momentum in the $(0, 0)$ - (π, π) direction for an (a) UD, (b) OPT, and (c) OD Bi2212. In all panels two main features are apparent: a high intensity feature at low

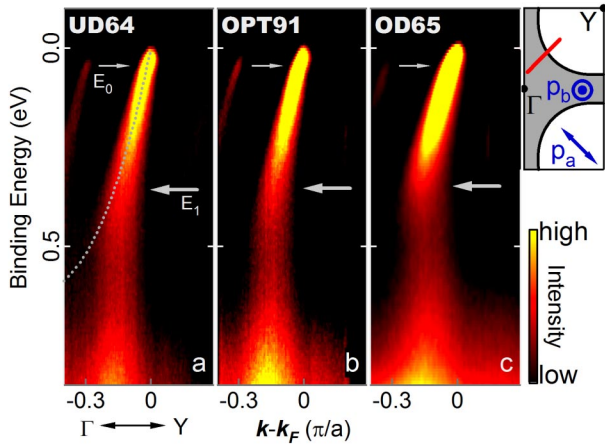


FIG. 1 (color online). ARPES intensity maps of Bi2212 samples for three different doping values. Data for the overdoped sample are in the normal state, 100 K. The location of the cut is shown in the BZ diagram on the right side, along with the different ARPES geometries used here. The gray dotted line is the dispersion obtained from a tight binding fit [5] up to energy 0.3 eV.

energy (widely studied in the literature [2,3]), and a weaker intensity feature at high energy. The high energy feature, “waterfall”-like feature, is the main focus of this Letter. Given the large energy span of Fig. 1, the “kink” at $E_0 \approx 0.06$ eV (gray arrows pointing to the right) is a very subtle feature. Aside from the kink, the low-energy dispersion can be well fitted by a single tight binding band [dotted gray line in panel (a)] [2,5]. Surprisingly, as the low-energy dispersion reaches $E_1 = 0.38 \pm 0.07$ eV, at momentum around $(\pi/4, \pi/4) \frac{1}{a}$ (gray arrows pointing to the left), the dispersion suddenly undergoes a steep downturn accompanied by a substantial drop of the ARPES intensity. The overall feature of this anomaly is nearly independent of doping.

In Fig. 2 we present selected raw EDCs (energy distribution curves, energy cuts at constant momentum; panel (a)) and MDCs [momentum distribution curves, momentum cuts at constant energy; panel (b)] for the OD Bi2212. We show results for the OD sample, making a strong case that pseudogap [6,7], disorder and inhomogeneity [8–10] can be ruled out as possible origin. However, similar behaviors are observed for all doping values we studied.

Panels (a) and (b) show that the behaviors of MDC and EDC peak become completely different as E_1 is reached, exposing the full view of the anomaly. The EDC peaks shown in panel (a) disperse in a simple manner and, as the momentum moves away from the Fermi momentum, the peak gets broader and weaker losing rapidly its strength as it approaches E_1 . This is consistent with the sudden decrease of the ARPES intensity observed at E_1 in Fig. 1. In contrast, the raw MDCs in panel (b) show a well-defined peak over the full energy range. For energy $\lesssim E_1$, the MDC peak disperses in a consistent fashion with the EDC dis-

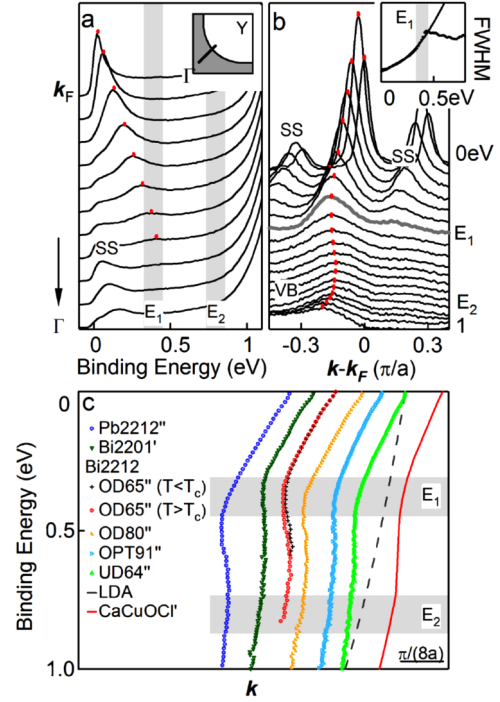


FIG. 2 (color online). (a),(b) EDCs from k_F (top curve) toward Γ and MDCs for the OD Bi2212. The spectra are vertically shifted for an easy view. The small peaks on the right and left of the main MDC peak are due to the superstructure (SS). At ≈ 1 eV, the valence band can also be distinguished. The inset shows the FWHM of the MDC peak as a function of energy. (c) MDC dispersion, horizontally shifted for an easy view, for several compounds and doping. The $\text{Ca}_2\text{CuO}_2\text{Cl}_2$ dispersion [19] is shifted in energy by 0.45 eV to account for the energy gap. The OD-Bi2212 data are shown both above, 100 K, and below T_c , 25 K. The prime and double prime stand for data taken in the first and second BZ, respectively. The dashed line is the LDA band dispersion [20].

persion, and as it reaches E_1 it suddenly becomes almost energy independent all the way up to E_2 . As the energy increases beyond E_2 , the MDC peak starts dispersing again. It is surprising that a well-defined MDC peak can still be identified within this large energy range. A further tracking of the MDC dispersion, well above E_2 , is made difficult due to the strong valence band complex dominating the ARPES signal, as seen by the strong rise of the intensity on the right end side of panel (a). Interestingly, the MDC peak width stops increasing at E_1 and shows a small decrease [inset of panel (b)].

From now on, we will refer collectively to the anomalies at E_1 and E_2 as “high energy anomaly.” In Fig. 2(c) we report MDC dispersions for different materials, various doping values, and different temperatures. It is clear that the overall features of the high energy anomaly and their energy and momentum locations are a universal feature in all the materials studied, from heavily overdoped to undoped compound, and do not show any substantial change going from the superconducting to the normal state.

Interestingly, beyond E_2 the MDC dispersion is in a reasonable agreement with the LDA prediction [panel (c)]. This gives an estimate of the energy scale of the measured band to be $\approx 1.3\text{--}1.4$ eV.

In Fig. 3 the full momentum space information about the high energy anomaly is summarized. Panels (a)–(h) show the momentum space distribution of the ARPES intensity at different energies from E_F to 1.3 eV. Representative data from OD Pb-Bi2212 are shown. From E_F to $\approx E_1$ [panels (a)–(c)] the ARPES contour and the tight binding fit (white solid line) are in very good agreement. This region corresponds to the low energy region, below E_1 , of Figs. 1 and 2. As the energy increases, the E_1 anomaly is marked unambiguously by the strong departure of the data from the tight binding fit and, over a wide energy region from E_1 to E_2 [panels (d)–(f)], the main ARPES intensity is “pinned” at the boundary of a diamond (dashed line) whose four corners are located near $(\pm\pi/2a, 0)$ and $(0, \pm\pi/2a)$. We caution readers that despite the color scale the spectral weight within this diamond is not strictly zero (Fig. 2). When the energy increases beyond E_2 [panels (g),(h)], the ARPES intensity starts moving again toward Γ , consistently with the resumed dispersion of the MDC peak discussed in Fig. 2. This is a common feature of all the materials reported here and of optimally doped $\text{La}_{1.64}\text{Eu}_{0.2}\text{Sr}_{0.16}\text{CuO}_4$ (Eu-LSCO) [11].

We have searched for, but did not find, a similar anomaly in more conventional materials such as GaAs (a simple band insulator), $\text{K}_{0.9}\text{Mo}_6\text{O}_{17}$, SmTe_3 [12] and CeTe_2 [13] (quasi-two-dimensional metals with conventional charge density wave orders), and graphite (a quasi-two-dimensional semimetal) [14]. However, a similar high

energy anomaly is also seen in several ruthenate compounds [15] suggesting that it might be an intrinsic feature of Mott physics.

We note that this peculiar high energy behavior described in Figs. 1–3 cannot be explained by the well-known ARPES matrix element effect [16]. The matrix element is strongly sensitive to the experimental settings in contrast to the high energy anomaly, e.g., predicting [17] that the intensity near the Γ point is enhanced in the second BZ, while it does predict an intensity depression in the first BZ. Instead, we find a similar high energy behavior independently of the BZ, the photon energy (25, 33, 43, 52, 55, 59, 75, 90, 100, 130, and 150 eV), or the light polarization settings (along the Cu-O bonds, along the Cu-Cu bonds, and normal to the CuO_2 planes).

One possible explanation of the observed high energy anomaly is in terms of coupling to a bosonic mode along a similar line as the low-energy kink. However, we believe this picture is unlikely since so far there is no known mode occurring at these energies and robust over the whole doping range, from zero doping to high overdoping.

A far more general possibility is that the data presented here show the disintegration of the low-energy quasiparticle. In this view, we propose that the dispersive band in the energy range from E_0 to E_1 is the signature of a composite object, as schematically represented in Fig. 3(j), made of a ZRS bound to a Cu spin $\frac{1}{2}$. This composite particle has quantum number $S = \frac{1}{2}$ and charge e and is consistent with that of a photohole. At lower energies this composite object is further dressed by phonons and low-energy collective spin excitations to become the ultimate quasiparticle in the energy range between E_F and E_0 . The

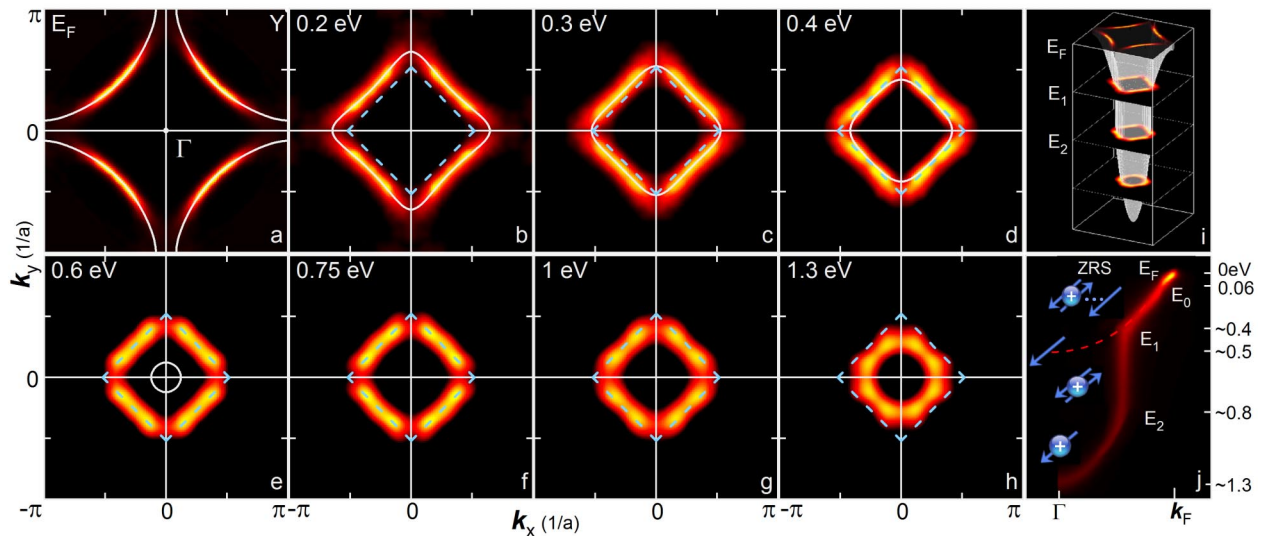


FIG. 3 (color online). (a)–(h) Maps of the ARPES intensity in the momentum space at increasing energies for the Pb-Bi2212 sample. Data were taken in the second BZ and symmetrized according to the tetragonal symmetry. The color scale is normalized independently for each cut. The white solid lines correspond to the tight binding fit and the dashed line diamond indicates the characteristic geometry of the high energy anomaly. (i) Three dimensional plot of the ARPES intensity as a function of energy and in-plane momentum. (j) Our proposed scenario for the high energy anomaly.

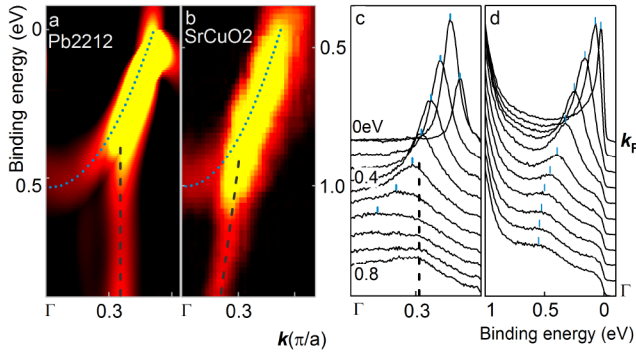


FIG. 4 (color online). Second derivatives of ARPES intensity maps along the nodal direction of Pb2212 (a) and SrCuO₂ (b) [18]. Only negative values of second derivatives are shown, to trace peaks but not dips. Sum of the second momentum-derivative and the second energy-derivative is shown at each point. (c),(d) MDCs and EDCs from the data shown in panel (a). The MDCs and EDCs are vertically shifted for an easy view. The dotted line and the dashed line highlight the proposed spinon and holon dispersions, respectively.

characteristics of the fermionic composite particle existing between E_0 and E_1 is the broad EDC and MDC peaks and their mutually consistent dispersions. At E_1 this composite object breaks down into a ZRS and a copper spin. The former is referred as a “holon” and the latter a “spinon” in the theory literature of high T_c . Experimentally this is seen as the sudden loss of spectral intensity of the broad dispersive ARPES peaks at E_1 . In the energy range between E_1 and E_2 the photoemission spectrum is the convolution of those of a spinon and a holon. In principle, the spectral features associated with both excitations can be observed by ARPES [18,21].

By tuning the energy (60 eV) and the polarization (out-of-plane) of the photon, we have observed for the first time a new dispersive feature [dotted line in Fig. 4(a)] whose bottom occurs at ≈ 0.5 eV at the Γ point. The corresponding peak of this new branch can be observed both in the MDCs and also in the EDCs spectra [panels (c),(d)]. We interpret this new feature as the spinon branch while the waterfall feature as the holon branch. The two peaks in the MDCs [panel (c)] in the energy range above 0.4 eV and below ≈ 0.5 eV represent the two branches. This interpretation is supported by the striking similarity between our data [panel (a)] and the recent photoemission result of the spinon and holon branches of the one dimensional cuprate ([panel (b)] [18]. Note, however, that, while the comparison between panels (a),(b) is quite appealing, not only qualitatively but also quantitatively, in our case two peaks can be observed only in MDCs but not in EDCs, due to the puzzling near-vertical nature of the holon dominated branch (waterfall). Finally, E_2 is the energy where the ZRS disintegrates into a bare oxygen hole and a copper spin. This reappearance of the fermionic oxygen hole explains the reemergence of a bandlike dispersion. Obviously, more

investigations would be necessary to test our conjecture further.

In conclusion, we have reported for the first time a universal high energy anomaly in the ARPES spectra of different families of high temperature superconductors, identified by a sudden change in the dispersion of the main spectral peak. This phenomenon is robust under the change of doping, as well as chemical composition. We conjecture that the high energy anomaly provides the long-sought-after evidence of spin charge separation in the high T_c compounds.

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