

# Comment on “Yield of Singlet Excitons in Organic Light-Emitting Devices: A Double Modulation Photoluminescence-Detected Magnetic Resonance Study”

In a recent Letter [1] Lee *et al.* studied the spin 1/2 photoluminescence-detected magnetic resonance (PLDMR) of a soluble derivative of poly-phenylene vinylene [MEH-PPV] using a double modulation scheme, where both the laser excitation intensity and the microwave power were modulated at  $\omega_{\text{ex}}$  and  $\omega_{\mu}$  ( $\omega_{\mu} \ll \omega_{\text{ex}}$ ), respectively. The obtained results were used to attempt differentiating between the quenching and the spin dependent recombination (SDR) models for PLDMR [2]. Also, an alternative model to the SDR was introduced [1] in which triplet excitons interact with polarons, (*T-P* model) leading to a reduction of both species. It was concluded [1] from the  $\omega_{\mu}$  dependence of PLDMR that the polaron recombination time is  $\sim 30 \mu\text{s}$ . This conclusion is unjustified, however, since the authors overlooked [already in Eq. (1)] the role of the polaron lifetime distribution, and spin subsystem dynamics on the PLDMR frequency response. In addition, the *T-P* model [1] is not viable in  $\pi$ -conjugated polymers in general, and in MEH-PPV, in particular.

In textbook introduction to PLDMR [3] the spin-lattice relaxation and microwave ( $\mu w$ ) power are shown to play an important role in determining the PLDMR dynamics and magnitude, especially for SDR [4]. Since SDR is inherently dispersive with a very broad distribution of recombination times  $\tau$  [3], then the spin-lattice relaxation time,  $T_{\text{SL}}$ , separates pairs with  $\tau \ll T_{\text{SL}}$  that contribute strongly to PLDMR [5] from those pairs with  $\tau \gg T_{\text{SL}}$  that contribute very weakly to PLDMR [3–5]. Thus the  $\omega_{\mu}$ -dependence PLDMR is sensitive to the *spin dynamics* as well as recombination kinetics of the spin-carrying excitations.

We have compared the  $\omega_{\text{ex}}$  dependence of polarons in MEH-PPV at 10 K using the photoinduced absorption (PA) technique at 0.4 eV with the spin 1/2 PLDMR  $\omega_{\mu}$  dependence at different  $\mu w$  powers (Fig. 1). These modulation frequency dependencies are not similar to each other for both in-phase and quadrature components, thus revoking the claim in [1] that the spin 1/2 PLDMR  $\omega_{\mu}$  dependence “yields the lifetimes of the spin-carrying species responsible for the resonance.” In particular, it is clear from Fig. 1 that the recombination mechanism of the polaron PA, unlike the PLDMR, is both *slower* and *dispersive* due to a very broad distribution of lifetimes. In addition, the PLDMR response depends on the  $\mu w$  power, whereas the PA response does not. We thus conclude that the apparent time constant in the PLDMR  $\omega_{\mu}$  dependence is *not* the spin-carrying recombination lifetime as claimed in [1]; but instead is determined by  $T_{\text{SL}}$ ,  $\mu w$  power, and  $\tau$  distribution. Consequently, the analysis of the flat double modulation  $\omega_{\text{ex}}$  dependence in [1] that is based on Eq. (1) is questionable.

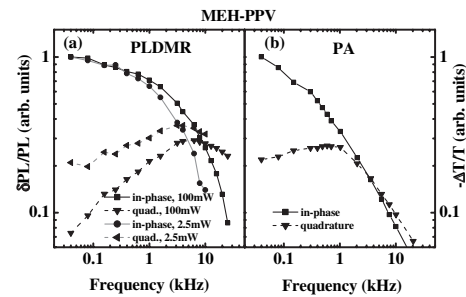


FIG. 1. (a) The  $\mu w$  modulation frequency dependence of spin 1/2 PLDMR in MEH-PPV film measured at 10 K at two  $\mu w$  powers. (b) The laser excitation frequency dependence of the polaron PA at 0.4 eV. The PA response is dispersive and does not depend on the  $\mu w$  power.

Second, the *T-P* model advanced in [1] cannot explain the existence of spin 1/2 PLDMR when the triplet exciton density is negligibly small. This occurs in MEH-PPV at elevated temperatures ( $\sim 230 \text{ K}$ ) where the triplet PA is negligible whereas the spin 1/2 PLDMR dynamics does not change much, as well as in many  $\pi$ -conjugated polymers where the triplets are unstable. Two examples of such polymers include: (i)  $\text{C}_{60}$ -doped MEH-PPV, in which the triplet excitons dissociate at the  $\text{C}_{60}$  sites [6], yet we measured  $\delta L/L \sim 4 \times 10^{-4}$  at  $g \cong 2$ . (ii) The degenerate ground state polymer poly(di-phenyl-acetylene), in which the triplet excitons are unstable against dissociation into solitons pairs [7], yet we measured  $dL/L \sim 2 \times 10^{-4}$ . The underlying mechanism of the spin 1/2 PLDMR in these examples *cannot be associated with triplet excitons* and thus the SDR model, which does not explicitly involve triplets [2–5], can uniquely explain the results. Consequently, if the SDR model is viable when the triplet density is small, then this mechanism cannot be turned off in other situations and polymers.

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