

**Vager and Naaman Reply:** In their Comment [1], Hernando and García correctly compared the treatment propose in Ref. [2] to the well-known system of “non-interacting electrons on a circular ring.” The Hamiltonian in this case is given by [3]

$$H = \frac{1}{2m} \left( p - \frac{e}{c} A \right)^2 \quad (1)$$

where  $p$  is the momentum operator along the path and  $A$  is the vector potential along the same path.  $A$  is commonly chosen as

$$\vec{A} = \frac{1}{2} \vec{H} \times \vec{r} \quad (2)$$

where  $H$  is the magnetic field and  $r$  originate from the center of the ring of a radius  $R$ . Therefore  $A = \frac{RH}{2}$  and the angular momentum component along the magnetic field is  $\hbar L = pR$ .

The Hamiltonian can be written now as

$$H = \frac{\hbar^2 L^2}{2mR^2} + \mu LH + \frac{e^2}{8mc^2} R^2 H^2 \quad (3)$$

where  $\mu$  is the Bohr magneton.

The matrix elements resulting from the last term in (3) are diagonal and constant. Therefore the form of the Hamiltonian used in Ref. [2], which includes only the first two terms of the Hamiltonian shown in (3), has the same spacing between energy levels and the same eigenstates. Therefore, it is justified for a circular ring. [4]

The Hamiltonian shown in (3) can be expressed as

$$H = \frac{\hbar^2}{2mR^2} (L + q)^2 \quad (4)$$

where  $q = \frac{\pi R^2 H}{\Phi_0}$  is a measure of the external flux in units of  $\Phi_0 = hc/e$ .

The eigenstates associated with this Hamiltonian are

$$\psi = \exp[i(M - q)\phi] \equiv \exp[i l \phi] \quad (5)$$

where  $\phi$  is the angular variable and  $M$  is an eigenvalue of  $L$ .

Uniqueness of  $\psi$  requires that  $l = M - q = \dots, -2, -1, 0, 1, 2, \dots$

The ground state angular momentum obeys  $|M| = \min|l + q|$  and carries with it a magnetic moment  $\mu M$  which is periodic in  $q$ .

Near zero field,  $l = 0$  is the ground state and is paramagnetic, namely  $\mu M = \mu q$ . This happens for fields obeying  $-1/2 < q < 1/2$ .

In the system described in Ref. [2], due to its small area, paramagnetism can persist till external magnetic fields approaching 1 T. Hence the statement regarding the paramagnetism of the system made in Ref. [2] is justified.

Noninteracting electrons theory as described above predicts magnetization which does not exceed half of a Bohr magneton per participating electron. This prediction is about 2 orders of magnitude too low as compared to the experimental observations [5]. The same discrepancy exists between the theory and experiments performed on the magnetic response of quasi-one-dimensional rings [6]. Therefore, the comment by Hernando and García is correct *vis a vis* the size of the magnetism. Clearly a noninteracting boson model, such as in Ref. [2], cannot explain the size of the giant magnetization observed. Recently, modified theories [7] that assume interacting electrons and disorder partially succeeded in reducing the discrepancy between theory and experiments, but the mystery is not completely solved yet.

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