

## Observation of the Acceleration of a Single Bunch by Using the Induction Device in the KEK Proton Synchrotron

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A single rf bunch in the KEK proton synchrotron was accelerated with an induction acceleration method from the injection energy of 500 MeV to 5 GeV.

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Four years ago, the concept of an induction synchrotron employing induction accelerating devices was proposed by Takayama and Kishiro [1] for the purpose of overcoming the shortcomings, such as a limitation of the longitudinal phase space available for the acceleration of charged particles in a rf synchrotron, which has been one of the indispensable instruments for nuclear physics and high-energy physics since the invention by McMillan [2] and Veksler [3]. Accelerating devices in a conventional synchrotron, such as a rf cavity, are replaced with induction devices in the induction synchrotron. A gradient focusing force, seen in the rf waves, is not indispensable for the longitudinal confinement of particles. Pulse voltages, which are generated at both edges of some time period with opposite sign, as shown in Fig. 1, are capable of providing longitudinal focusing forces. A pair of barrier-voltage pulses work in a similar way to the rf barrier, which has been demonstrated at Fermilab and BNL [4]. The acceleration and longitudinal confinement of charged particles are independently achieved with induction step voltages in the induction synchrotron. This notable property of the separated function in the longitudinal direction brings about a significant freedom of beam handling never seen in a conventional rf synchrotron, in which radio-frequency waves in a resonant cavity simultaneously take both roles of acceleration and longitudinal confinement. Associated with this separated function, various figure of merits are expected. The formation of a superbunch, which is an extremely long bunch with a uniform line density, and the use of which is considered in a proton driver for the second-generation neutrino physics and a future hadron collider [5], is most attractive. In addition, transition crossing without any longitudinal focusing seems to be feasible

[1], which could substantially mitigate undesired phenomena, such as bunch shortening due to nonadiabatic motion and microwave instabilities [6].

The time duration between barrier-voltage pulses determines the size of a superbunch. For accelerating a superbunch, a long accelerating step voltage with a pulse length on the order of  $\mu$  sec is required. The voltage has to be generated at the revolution frequency of the beam in a ring. In the case of a ring-circumference of 300 m, a repetition rate of 1 MHz is required. So far, there has been no induction accelerating system capable of meeting these parameters. Recently, prototype devices, which can generate a 250 nsec flat-top voltage at a repetition rate of cw 1 MHz, have been assembled at KEK after the 3-year research and development stage, and combined with the existing rf accelerating system. A single proton bunch trapped in a rf bucket was accelerated with the induction accelerating system from 500 MeV to 5 GeV within 1.3 sec. In this Letter, we report on the first experimental result of the induction acceleration of a single rf bunch in

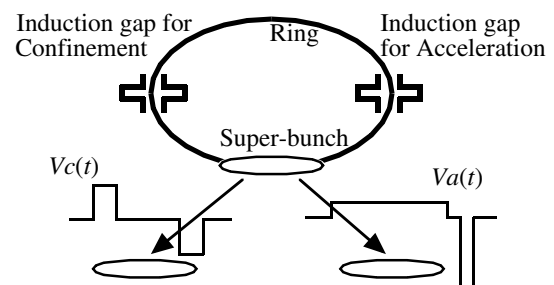


FIG. 1. Schematic view of the induction synchrotron and step-voltage profiles for confining and acceleration.

the KEK proton synchrotron (KEK-PS) as well as a brief description of the induction accelerating system.

The key devices required to realize an induction synchrotron are an induction accelerating cavity [7] and a pulse modulator driving the cavity [8]. These devices are notably different from similar devices employed in modern linear induction accelerators [9]. Remarkable characteristics are its switching characteristics, repetition ratio, and duty factor. Another different feature is that the pulse modulator has to be kept far from the induction cavity placed in the accelerator tunnel, because the solid-state power switching elements obtainable at present cannot survive an extremely high radiation dose. Thus, the pulse modulator is connected with the accelerating cavity through a long transmission cable. In order to reduce reflection from the load, a matching resistance has been installed at the end of the transmission cable. The induction accelerating system consists of an induction cavity with a matching load, a transmission cable, a pulse modulator, and a dc power supply. A typical system capable of generating a step-pulse of 2 kV output voltage and 18 A output peak current at 1 MHz with 50% duty has been demonstrated at KEK. An equivalent circuit for the induction accelerating system is shown in Fig. 2.

The core material of the induction cavity employed for the first acceleration experiment is a nanocrystalline alloy, called Finemet (Hitachi Metal). Heat generated due to core loss is cooled down by insulation oil. The electrical parameters of each unit cell are given by the capacitance of 260 pF, the inductance of 110  $\mu$ H, and the resistance of the 330  $\Omega$ , which determine the properties of pulse rising and falling. Three unit cells with a 2 kV output voltage per unit are mechanically combined into a single module for the convenience of installation. Since the inner conductor with three ceramic gaps is common to three unit cells, but both sides of each gap are electrically connected to the outer edges of each cell, a particle is accelerated with the same voltage passing these gaps.

A full-bridge switching circuit in the pulse modulator, which is depicted in Fig. 2, was employed because of its simplicity. The pulse modulator is capable of generating bipolar rectangular shaped voltage pulses. The full-bridge

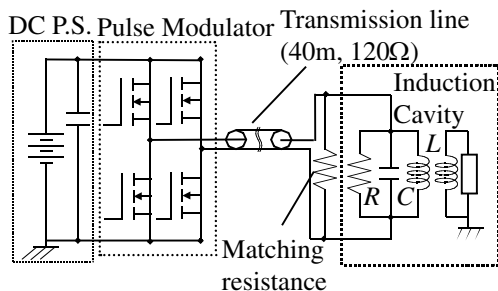


FIG. 2. Equivalent circuit for the induction accelerating system.

type pulse modulator consists of four identical switching arms. Each switching arm is composed of 7 metal-oxide-semiconductor field-effect transistors, which are arranged in series. Their gates are driven by their own gate-driving circuits. The gate signals are generated by converting light signals provided from the pulse controller, which is a part of the accelerator control system, to electronic signals. Its details will be presented elsewhere [8].

The entire system employed for the current experiment is schematically shown in Fig. 3. The generation of a 2 kV voltage pulse is directly controlled by trigger pulses for the switching elements of the pulse modulator, the master signal of which is created in the digital signal processing (DSP) synchronizing with ramping of bending magnets, and gate-driving signal patterns initiated by this master trigger signal are generated by the following signal-pattern generator to be sent to the gate controller of the pulse modulators through a long coaxial cable. The DSP counts the  $B$ -clock signal, and achieves the desired revolution frequency. Any delay between the accelerating pulse and a bunch monitor signal is always corrected by the DSP. The system is connected to the existing rf system through the rf signal, which shares the  $B$ -clock signal. As a result, synchronized induction acceleration is guaranteed. Here, the rf does not contribute to acceleration of the beam bunch, but gives the focusing force in the longitudinal direction; the beam bunch is in principle trapped around the phase of zero. The machine parameters of the KEK-PS employed for the experiment are as follows; transition energy  $\gamma_t$  is 6.63, injection and extraction energy are 500 MeV/8 GeV, revolution frequency  $f_0$  is 668–877 kHz, rf voltage  $V_{rf}$  is 40 kV and harmonic number  $h$  is 9.

In the experiment, the signals of the bunch monitor and three current transformers (CTs), which always observed the current flow through the matching resistances, were monitored on a digital oscilloscope located at the central

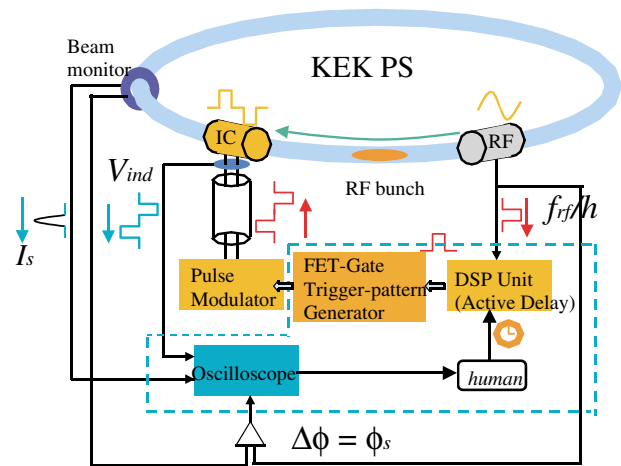


FIG. 3 (color). Schematic view of the hybrid accelerating system.

accelerator control room. Before the experiment, an actual induced voltage at the ceramic gap, an output voltage of the transmission cable, and the CT signal were simultaneously measured and the correspondence between each was well calibrated. In addition, a delayed timing of the master gate signal triggering the pulse modulator is adjusted by the DSP so that the bunch signal would stay around the center of the induction voltage pulse through the entire accelerating period. Typical waveforms of the CT signals are shown in Fig. 4 together with the bunch signal.

Under coexistence with the rf voltage and the induction voltage, a charged particle receives an energy gain per turn,

$$eV_{\text{acc}}(t) = e[V_{\text{rf}} \sin \phi(t) + V_{\text{ind}}], \quad (1)$$

where  $V_{\text{rf}}$  and  $V_{\text{ind}}$  are the rf voltage and the induction voltage, respectively, and  $\phi(t)$  is the position of the particle in the rf phase,  $\omega_{\text{rf}}t$ . The orbit and energy of a particle are dominated by the following equations: for the force balance in the radial direction

$$m\gamma \frac{(c\beta)^2}{\rho} = ec\beta B(t), \quad (2)$$

where  $B(t)$  is the bending field and  $\rho$  is the bending radius, and for the change in energy

$$mc^2 \frac{d\gamma}{dt} = \frac{ec\beta}{C_0} V_{\text{acc}}(t). \quad (3)$$

From Eq. (2) and (3), the accelerating voltage must satisfy the relationship  $V_{\text{acc}}(t) = \rho C_0 dB/dt$ , so that the particle is synchronously accelerated with ramping of the bending field. The bending field is linearly ramped over 1.7 sec, as shown in Fig. 5. In this linear ramping region, a  $V_{\text{acc}}$  of 4.7 kV is required. In the present experiment, the induction voltage was fixed to be 5.28 kV.

In order to confirm the induction acceleration, the phase signal, which shows the relative position of the bunch center to the rf, was measured through an entire accelerat-

ing region. Particularly, we focused on three cases: (1) with a rf voltage alone, (2) with a rf voltage and a positive induction voltage, and (3) with a rf voltage and a negative induction voltage. From Eq. (1), a theoretical prediction is  $\phi_s = \sin^{-1}(V_{\text{acc}}/V_{\text{rf}}) \sim V_{\text{acc}}/V_{\text{rf}} = 6.73$  degrees for case (1),  $\phi_s = -0.68$  degree for case (2),  $\phi_s = \sin^{-1}(2V_{\text{acc}}/V_{\text{rf}}) \sim 2V_{\text{acc}}/V_{\text{rf}} = 14.59$  degrees for case (3), where  $\phi_s$  represents the position of the bunch center in the rf phase. Since the induction voltage is devoted to the acceleration for case (2), the rf does not serve for acceleration, but takes a role of capturing alone; thus, the phase must be nearly zero. In case (3), the rf has to give a two-times larger energy to the bunch than that in case (1) from the energy-conservation law; the phase should increase by a factor of 2. Actually, the time evolution of the phases and a beam-current for each case has been observed. There were notable differences between three cases for the temporal evolution in phase. The beam itself behaved in the same manner for the three cases as shown in Fig. 5, where a temporal evolution of the beam intensity from 500 MeV to 5 GeV is shown.

Before we proceed to an analysis of the experimental results, we have to pay attention to the phase signal amplitude, because there were several ambiguities in our phase pick-up monitoring system, such as time-varying offset and gain. In order to avoid unknown factors, we have made a kind of calibration table, from which the relationship between the monitored phase pick-up voltage and the exact phase angle is readable, by measuring the phase pick-up voltages corresponding to the different rf voltages from 40 kV to 90 kV. In Fig. 6, the experimental results are shown along with corrections obtained from a calibration table. For case (1), the phase is very close to the theoretical prediction. For case (2), the phase is nearly 0 as a rough estimate. For case (3), the magnitude is higher by an about factor of 2 than that of case (1). At first glance, we can find a qualitative agreement with the theoretical pre-

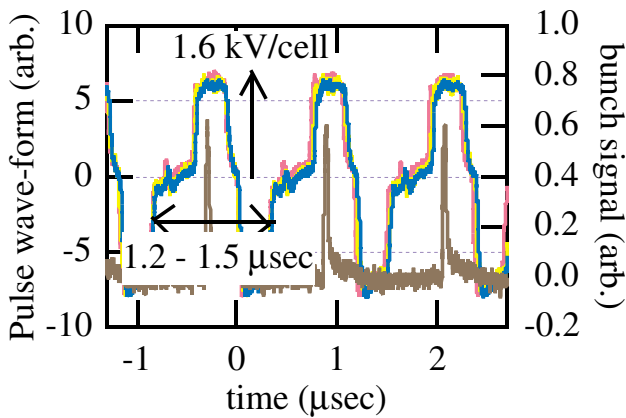


FIG. 4 (color). Induction pulse waveforms (pink, yellow, blue) and a bunch signal (brown). Revolution period at 500 MeV and 5 GeV is 1.5 and 1.2  $\mu$ sec, respectively.

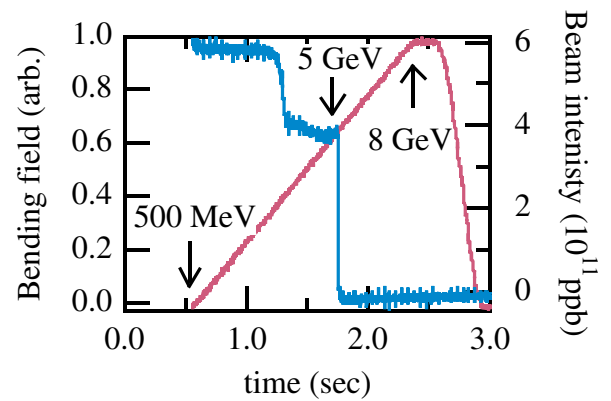


FIG. 5 (color). Ramping pattern of the bending field (purple) and beam intensity (sky blue). The beam is stopped just below a transition at 1.75 sec, and some beam loss due to transverse motion is observed around 0.5 sec earlier before the transition.

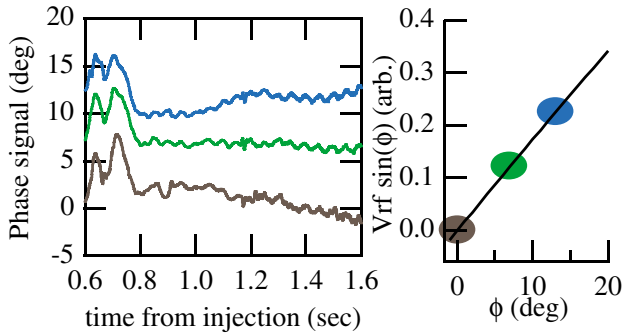


FIG. 6 (color). Phase signals vs time (sec) for case (1) (green), case (2) (brown), and case (3) (blue).

diction. Using the calibrated gain in the linear ramping region from 700 msec to 1.6 sec, we arrived at Table I, where the results are more quantitatively compared with the prediction.

The experimentally obtained phases are slightly different from the theoretically predicted values. A difference growing near the transition energy may be attributed to the fact that the rf bunch becomes shorter when approaching the transition energy; as a result, the effective acceleration voltage in this region sensitively depends on the relative position in the induction voltage with a droop. Meanwhile, we realize that a difference between the accelerating voltage pulse shape and the decelerating voltage pulse shape, which can be seen in Fig. 4, gives a slightly asymmetric evolution around Case (1). From the overall experimental results and their analysis, we conclude that a single rf bunch was accelerated from 500 MeV up to 5 GeV with the induction accelerating device, as expected.

The present induction voltage waveform has a droop of 10–15% over 250 nsec, although it has not been discussed. To achieve a uniform acceleration of the superbunch, flatness in the waveform is important. Several methods [10] to compensate the droop have been proposed. They will be tried at the KEK-PS. Following the original scenario of the POP experiment to realize an induction synchrotron [11], hereafter, we will proceed to its next step in the formation of a superbunch by the induction step-barriers.

It is emphasized that for the first time charged particles in a high-energy accelerator ring were accelerated with an induction accelerating system. The experimental fact of induction acceleration in a circular ring and the reality that the key devices developed for this purpose, such as a pulse modulator as a switching driver, are on our hands as big milestones for us to achieve an induction synchrotron in the near future, and a superbunch hadron collider in the no-so-far future. Last, it is emphasized that the induction

TABLE I. Results quantitatively compared with the prediction.

	experimental (deg)	theoretical (deg)
Case (1)	6.5 ~ 7	6.73
Case (2)	2.0 ~ -2.0	-0.68
Case (3)	10 ~ 13	14.6

system worked well through the entire operating period of 24 hours without any trouble.

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*Note added in proof.*—After the submission, a single rf proton bunch was accelerated beyond the transition energy to the flat energy of 8 GeV [12].

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