

## Random Signal Fluctuations Can Reduce Random Fluctuations in Regulated Components of Chemical Regulatory Networks

Johan Paulsson\* and Måns Ehrenberg†

*Department of Cell and Molecular Biology, Uppsala University, BMC, Box 596, Uppsala, 75124 Sweden*

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Many intracellular components are present in low copy numbers per cell and subject to feedback control. We use chemical master equations to analyze a negative feedback system where species  $X$  and  $S$  regulate each other's synthesis with standard intracellular kinetics. For a given number of  $X$ -molecules,  $S$ -variation can be significant. We show that this signal noise does not necessarily increase  $X$ -variation as previously thought but, surprisingly, can be necessary to reduce it below a Poissonian limit. The principle resembles Stochastic Resonance in that signal noise improves signal detection.

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Intracellular processes are regulated by signal molecules that often are present in a few to a few hundred copies and display significant internal noise [1]. Since a noisy signal only represents the underlying state of the cell in a probabilistic sense, this has generally been assumed to randomize control [2–4]. In this Letter we use chemical master equations [5,6] to demonstrate that signal noise instead can attenuate the concentration noise in a regulated component. Our minimal regulatory network consists of two components,  $X$  and  $S$ , that regulate each other's synthesis, and performance is defined by the capacity to reduce random  $X$ -variation. We chose this particular copy number control (CNC) system because its kinetic mechanisms constitute simple intracellular standards, but also because it is virtually identical to models of CNC of bacterial plasmids [7]. Plasmid CNC systems are comparatively lucid and have evolved primarily to attenuate copy number fluctuations [4].

The macroscopic equations describing the feedback system are

$$\begin{cases} \dot{[x]} = \frac{k[x]}{1+[s]/K} - [x], \\ \dot{[s]} = k_s[x] - k_d[s], \end{cases} \quad (1)$$

where  $[x]$  and  $[s]$  are continuous concentration variables. For plasmids, the autocatalysis of  $X$ -molecules comes from the constant frequency with which each plasmid molecule tries to replicate itself [8].  $S$ -molecules inhibit  $X$ -synthesis trials through so-called hyperbolic inhibition [Eq. (1)] so that the probability that a trial is successful depends on the concentration of  $S$ -molecules at the time of the trial [8].

Normalizing the variables by their nonzero steady states  $[x_r] = [x]/[\bar{x}]$  and  $[s_r] = [s]/[\bar{s}]$  gives  $[\dot{x}_r] = k[x_r][1 + [s_r](k - 1)]^{-1} - [x_r]$  and  $[\dot{s}_r] = k_d([x_r] - [s_r])$ . The parameters  $k_s$  and  $K$  in Eq. (1) thus determine two characteristic concentration scales, and  $k_d$  determines how rapidly  $[s_r]$  adjusts to changes in  $[x_r]$ . When  $k_d$  is small, relaxation to steady state is oscillatory, and when  $k_d \rightarrow \infty$ ,  $[s_r]$  is strictly proportional to  $[x_r]$  at

all times. The sensitivity with which the rate of  $[x_r]$  synthesis responds to changes in  $[s_r]$  increases with  $k$ , but for  $k \gg 1$ , sensitivity approaches an asymptote, where the rate of synthesis of  $[x_r]$  per unit of  $[x_r]$  is inversely proportional to  $[s_r]$ , i.e.,  $[\dot{x}_r] \approx [x_r][s_r]^{-1} - [x_r]$ . The basic control principle is that  $[s_r]$  follows changes in  $[x_r]$ , thereby boosting the relative rate of  $[x_r]$  synthesis at low concentrations and restraining it at high.

If  $X$ -synthesis were not inhibited by  $S$ -molecules, Eq. (1) would predict that any initial condition is perpetuated if  $k = 1$ . However, due to the random nature of chemical reactions, this would correspond to an accelerated unbiased random walk.  $X$ -variation would then be limited only by  $X$ -extinctions or by physical restrictions, such as a depletion of the cell's resources, when the number of  $X$ -molecules becomes too large. Both these effects are evolutionary unfavorable for the cell and negative feedback CNC has therefore evolved to attenuate the fluctuations [4]. Since the evolutionary rationale for CNC is to reduce fluctuations, it cannot be properly inspected using macroscopic equations. Here we instead use the chemical master equation [5,6].

Four types of events are included in our mesoscopic model of CNC. When there are  $m$   $X$ -molecules and  $n$   $S$ -molecules,  $S$ -molecules are synthesized with rate  $k_s m$  and degraded with rate  $k_d n$ .  $X$ -molecules are synthesized with rate  $g_{m,n} = km/(1 + n/K)$ , where  $K$  now includes the reaction volume, and degraded with rate  $m$ . For plasmids, this straightforward mesoscopic version of the inhibition mechanism is a good approximation since the duration of an  $X$ -synthesis event is only a few seconds while the half-life of  $S$ -molecules is about a minute [8]. If the number of  $S$ -molecules instead changed significantly over the duration of an  $X$ -synthesis trial, so that the trial cannot be considered instantaneous, the behavior is more complicated.

The rates are transition probabilities per time unit in the master equation for the probability  $p_{m,n}$  of being in a state with  $m$   $X$ -molecules and  $n$   $S$ -molecules. With the step operator [6]  $E_n^j f(n) = f(n + j)$ , where  $j = \pm 1$ , this can

be written as

$$\begin{aligned} \dot{p}_{m,n} = & (E_m^{-1} - 1)g_{m,n}p_{m,n} + (E_m - 1)mp_{m,n} \\ & + k_s m(E_n^{-1} - 1)p_{m,n} + k_d(E_n - 1)np_{m,n} \\ & + p_{m,n} \sum_{n=0}^{\infty} p_{1,n} \end{aligned} \quad (2)$$

for  $\{m > 0, n \geq 0\}$ . The last term comes from conditioning the distribution on  $m > 0$ . This is done because the state with zero  $X$ - and  $S$ -molecules is absorbing and all others are transient. When  $X$  is essential for the survival of the cell, the conditioning has its natural counterpart in the death of  $X$ -free cells.

When  $k_d$  is so high that the number of  $S$ -molecules rapidly adjusts to the current number of  $X$ -molecules, the number of  $S$ -molecules is a fast variable that can be removed from the equations. Equation (2) then simplifies to a master equation for the probability  $p_m$  of being in a state with  $m$   $X$ -molecules

$$\dot{p}_m = (E_m^{-1} - 1)g_m p_m + (E_m - 1)m p_m + p_1 p_m. \quad (3)$$

We will analyze Eq. (3) both for noise-free and noisy signals, i.e., when conditional  $S$ -variation is insignificant and significant, respectively.

When the conditional  $S$ -variation for a given value of  $X$  is negligible, then

$$g_m = km/[1 + (mk_s/k_d)/K]. \quad (4)$$

The most efficient control is then obtained when  $k$  is so high that  $g_m$  is approximately constant. The number of  $X$ -molecules then only deviates from the Poisson distribution at very low averages,  $\langle m \rangle$ . For a fixed average, lower values of  $k$  inevitably broaden the distribution. In fact, when conditional  $S$ -variation is negligible, the  $X$ -distribution can never be narrower than (approximately) Poisson even if all parameters in Eq. (2) can be chosen freely. Lower  $X$ -variation requires more efficient control kinetics, for instance,  $g_m \propto m^{1-i}$ . High  $i$  means high sensitivity amplification, defined as the percentage change in the response,  $g_m$ , over the percentage change in the signal,  $m$  [9]. In this case it means that the  $X$ -synthesis rate is high below the average and negligible above so that fluctuations are insignificant. Throughout this Letter we use  $\langle m \rangle = 10$ . For  $i \geq 1$ , the  $X$ -variances are then  $\sigma_m^2 \approx \langle m \rangle/i$ .

Considering the impact of a noisy signal, i.e., where conditional  $S$ -variation cannot be ignored, an adiabatic elimination of the fast variable  $S$  instead gives

$$g_m = km \sum_{n=0}^{\infty} \bar{p}_{n|m} / (1 + n/K). \quad (5)$$

The quasistationary conditional probabilities of  $n$   $S$ -molecules given  $m$   $X$ -molecules,  $\bar{p}_{n|m}$ , are Poissonian with conditional average  $\langle n \rangle_m = mk_s/k_d$ , since all synthesis and degradation events are independent. The number of  $S$ -molecules at any given time thus only represents the number of  $X$ -molecules in a probabilistic sense. For instance, when  $\langle n \rangle_m = m$ , a sample drawn

from  $\bar{p}_{n|10}$  is lower than a sample drawn from  $\bar{p}_{n|9}$  with probability 0.365. Since regulation is nonlinear, this will affect  $X$ -variation. In the notation of Haken [10] and Gardiner [5], the fast variable  $S$  can be said to be a noisy slave of the slow variable  $X$ .

To efficiently eliminate  $X$ -fluctuations, a higher  $m$  should correspond to a higher  $n$  so that the  $X$ -synthesis rate decreases efficiently. However, increasing the conditional  $S$ -variance increases the probability that a high  $m$  instead corresponds to a low  $n$ . It is tempting to assume that conditional  $S$ -fluctuations reduce performance by randomizing the critical step in regulation. However, having overlapping conditional  $S$ -distributions means that the probabilities of avoiding inhibition may become *more* separated than was possible without fluctuations. The reason is that the nonlinear regulatory mechanism receives a disproportional contribution from the tail of the distribution, and the probability mass in the tail in turn responds sensitively to changes in the average. This is a general principle for sensitivity amplification (Paulsson *et al.* [11]) for which the name Stochastic Focusing (SF) has been suggested.

Calculating the stationary distribution of Eq. (3) using Eq. (5) shows (Fig. 1) that  $X$ -fluctuations can be efficiently eliminated without conventional sensitivity amplification. In fact,  $X$ -variation can be reduced *indefinitely* if the rate constants can be chosen freely. Efficiency requires a high  $k$ , but there is an upper limit in  $k$  for given  $\langle m \rangle$  and  $\bar{p}_{n|m}$  since  $g_m \geq km\bar{p}_{0|m}$  [Eq. (5)].

We also studied the impact of  $k_d$  using the two-dimensional master equation (2) when  $k = 100$ , both for  $\langle n \rangle_{10} = 8$  and for insignificant conditional  $S$ -variation [12]. Decreasing  $k_d$  impairs CNC much more when conditional  $S$ -variation is significant (Fig. 2). This is not an effect of the higher sensitivity amplification. Using conventional sensitivity amplification in Eq. (2), such as  $g_{m,n} = km/(1 + n^2/K)$ , CNC works better at all values of  $k_d$  when conditional  $S$ -variation is negligible (Fig. 2). The difference between conventional sensitivity

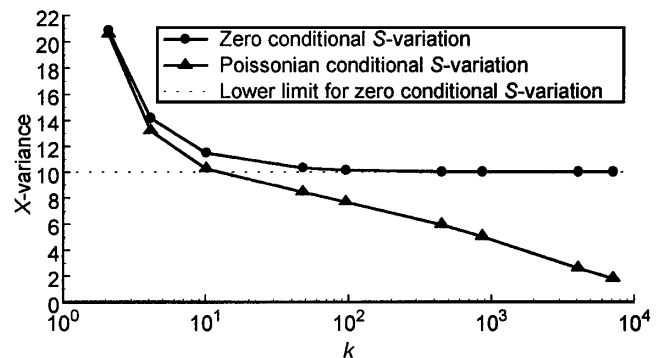


FIG. 1.  $X$ -variance as a function of  $k$  for hyperbolic inhibition when Eq. (5) is used in Eq. (3). The averages were  $\langle m \rangle = 10$  and  $\langle n \rangle_m = m$ . The inhibition constant  $K$  changes along the curve to keep  $\langle m \rangle = 10$ .

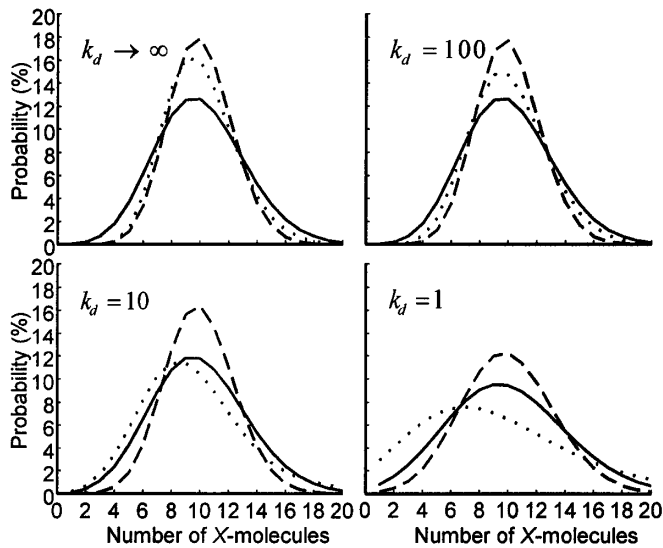


FIG. 2. Stationary  $X$ -distributions calculated by numerical integration of Eq. (2). The solid curve is hyperbolic inhibition without conditional  $S$ -variation [12], and the dotted curve is the same with Poissonian conditional  $S$ -fluctuations. The broken curve was obtained without conditional  $S$ -variation when  $g_{m,n} = km/(1 + n^2/K)$ .  $K$  changes to keep  $\langle m \rangle = 10$  and  $\langle n \rangle_{10} = 8$  in all curves.

amplification and fluctuation enhanced sensitivity (SF) is that low  $k_d$  and significant conditional  $S$ -fluctuations result in correlations between subsequent events. If  $n$  by chance is low at  $t = 0$ , it is likely to remain low for some time. Figure 2 shows that as time correlations become more and more significant, the beneficial effects of noise become overshadowed by stronger correlations between outcomes of subsequent  $X$ -synthesis trials.

The results presented in Figs. 1 and 2 are not restricted to Poisson fluctuations or exceptionally low conditional averages  $\langle n \rangle_m$ . We performed the same analysis when  $S$ -molecules were produced in instantaneous, geometrically distributed bursts. Geometric bursts are common in intracellular processes since they arise for Poisson processes in exponentially distributed time windows [3]. With  $k'_s$  as the rate constant for initiation of  $S$ -synthesis and  $G_j = q^j(1 - q)$  as the probability for a burst of  $j$  molecules, adiabatic elimination of the fast variable  $S$  results in Eqs. (3) and (5) with  $\bar{p}_{n|m}$  as the stationary distribution of

$$\dot{p}_{n|m} = k'_s m \left( \sum_{j=1}^n G_j p_{n-j|m} - q p_{n|m} \right) + k_d (E_n - 1) n p_{n|m}. \quad (6)$$

This gives the same type of macroscopic equation as before [Eq. (1)], and  $\bar{p}_{n|m}$  is the negative binomial (NB)

$$\bar{p}_{n|m} = \frac{\beta^n}{(1 + \beta)^{\lambda m + n}} \frac{\Gamma(\lambda m + n)}{\Gamma(\lambda m) n!}, \quad (7)$$

where  $\beta = q/(1 - q)$  is the average burst size and  $\lambda = k'_s/k_d$ . The corresponding conditional  $S$ -average and vari-

ance are  $\langle n \rangle_m = m\lambda\beta$  and  $\sigma_{n|m}^2 = m\lambda\beta(\beta + 1)$ , respectively. Fluctuations can thus be significant also at high averages. For plasmids, NB distributions can arise when  $S$ -synthesis depends on a transcription activator or repressor. The NB also arises as the stationary distribution for a large number of other simple birth and death processes of biochemical relevance, for instance autocatalysis [4].

To study the impact of increased variation using Eq. (7) in Eqs. (3) and (5), we numerically compared different combinations of  $\beta$  and  $\lambda$  for fixed averages  $\langle m \rangle = 10$  and  $\langle n \rangle_{10} = 50$ . The rate constant  $k$  was fixed to 100, and  $K$  was changed to keep  $\langle m \rangle = 10$ . The results are shown in Fig. 3.

When  $\beta \rightarrow 0$ , the NB distribution converges to a Poissonian, and the conditional  $S$ -fluctuations have no significance because of the high average, resulting in an approximately Poisson distributed number of  $X$ -molecules (Fig. 3). When  $\beta$  increases, so does the conditional  $S$ -variation, but the  $X$ -variation instead decreases. In fact, the uncertainty in the number of  $X$ -molecules can only be decreased below the Poissonian limit by simultaneously increasing the conditional uncertainty in the number of  $S$ -molecules. This “kinetic uncertainty principle” is obviously not universal but rather depends on the kinetic mechanisms of the individual systems.

In conclusion, intrinsic noise in one component of a regulatory chemical network can be exploited to reduce intrinsic noise in another component, in direct contradiction to what has been previously assumed [2]. Signal noise arising from simplistic birth and death processes combined with simple and realistic biochemical mechanisms can, in fact, work as efficiently as a threshold mechanism combined with insignificant signal noise. This principle

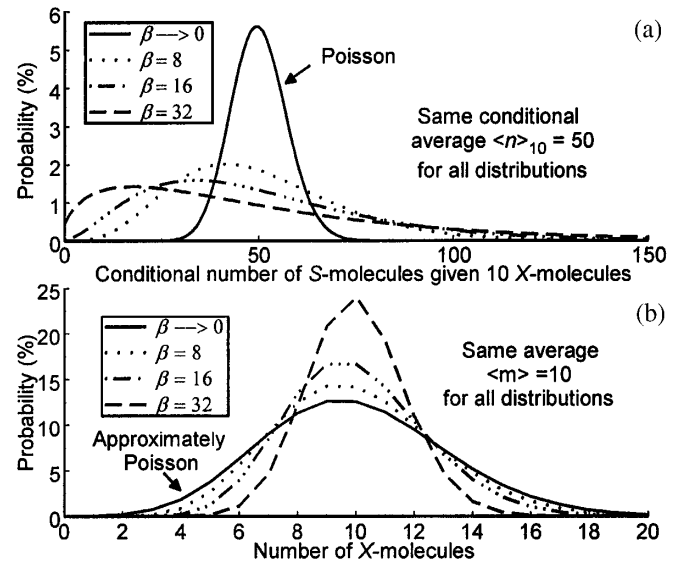


FIG. 3. Copy number distributions. (a) Conditional NB  $S$ -distribution [Eq. (7)] corresponding to ten  $X$ -molecules, i.e.,  $\bar{p}_{n|10}$ . (b)  $X$ -distributions using Eqs. (5) and (7) in Eq. (3). See main text for parameters.

resembles Stochastic Resonance (SR) [13] in that signal noise can improve performance of a nonlinear system. However, signal, noise and performance are all different from SR standards. Instead, the results are more closely related to Renyi's classic paper [14] showing that average rates of bimolecular reactions are affected by copy number variances.

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\*Electronic address: Johan.Paulsson@icm.uu.se

†Electronic address: Mans.Ehrenberg@icm.uu.se

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