

### Comment on "Superconducting PrBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>"

Recently, Zou *et al.* [1] reported the observation of bulk superconductivity (SC) for a PrBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> (Pr123) single crystal grown by the traveling-solvent floating-zone (TSFZ) method. The SC of Pr123 itself and also the increase of the  $T_c$  from 85 to  $\approx 105$  K under pressure are of general interest. These unexpected results (see also [2]) are in sharp contrast with the generally accepted view that Pr123 is the only nonsuperconducting compound among the orthorhombic  $RBa_2Cu_3O_{7-y}$  ( $R = Y$ , rare earth) cuprates. More detailed knowledge of TSFZ crystal properties is required to resolve this discrepancy.

Zou *et al.* reported only slight differences in the crystal structure between TSFZ Pr123 and crystals grown by the flux method: An elongation of the  $c$ -axis parameter connected with the expansion of the distance between two CuO<sub>2</sub> planes was found. We agree with Zou *et al.* that it is hard to attribute the hole delocalization and the occurrence of SC in the TSFZ crystal to the small elongation of the Pr-O(2) distance. Because of the "strong sample inhomogeneity" the substitution of Ba for Pr might create mobile superconducting holes [1]. However, the effective magnetic moment of Pr  $\mu_{\text{eff}}$  was found to be  $2.92\mu_B$  [1], i.e., close to that of their flux crystal. Thus, given this value of  $\mu_{\text{eff}}$  it is difficult to imagine a substantial substitution of nonmagnetic Ba for Pr.

The aim of this Comment is to show the *inconsistency* of the value of  $\mu_{\text{eff}}$  reported by Zou *et al.* with their magnetic susceptibility  $\chi(T)$  data (Fig. 3 of Ref. [1]). For clarity the data of Ref. [1] are shown here together with recent results of our group for a high quality Al-free Pr123 single crystal grown in a Pt crucible [3]. For this Pr123 crystal the  $\chi^{-1}(T)$  curves are shown in Fig. 1 for the field parallel ( $H \parallel c$ ) and perpendicular ( $H \parallel ab$ ) to the  $c$ -axis. (The field direction for the TSFZ crystal was not mentioned in [1].) The values of  $\mu_{\text{eff}} = 2.9\mu_B$  and  $3.1\mu_B$  were obtained for our crystal for  $H \parallel ab$ -plane and  $H \parallel c$ -axis, respectively, from the best fits of points at  $50 \leq T \leq 300$  K to the modified Curie-Weiss law including a temperature independent  $\chi_0$  (shown in Fig. 1 by solid lines). These values are in good agreement with previously reported ones; see, e.g., [4]. Since it is impossible to have very closely similar values of  $\mu_{\text{eff}}$  from quite different "flux" and "TSFZ" curves, we reestimated the value of  $\mu_{\text{eff}}$  from the TSFZ data shown in Fig. 1.

The first estimate from the *linear* approximation by Zou *et al.* to the  $\chi^{-1}(T)$  data (dotted line in Fig. 1 here) gives  $\mu_{\text{eff}} = 2.32\mu_B$ . According to [1] their fit was obtained with  $\chi_0 = 4.5 \times 10^{-4}$  emu/mol. But the *straight* line in Fig. 3 of [1] representing the fit *cannot* be reproduced with that  $\chi_0$  value. The second estimate made directly from their data points fitted to the modified Curie-Weiss law including  $\chi_0$  gives an even smaller value of  $\mu_{\text{eff}} =$

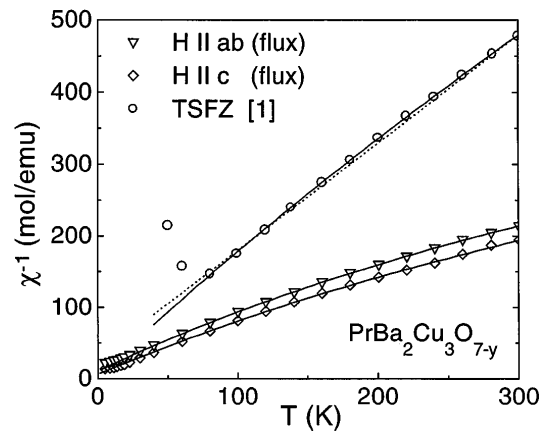


FIG. 1.  $\chi^{-1}$  vs  $T$  for flux grown [3] and TSFZ [1] Pr123 single crystals. Solid lines: fits to the Curie-Weiss law. Only some representative points are shown. For details, see text.

$2.09\mu_B$  and a Curie constant  $C = 0.546$  emu K/mol. The  $C$  value for the TSFZ crystal is about one-half of that for our flux crystal (1.04 and 1.19 emu K/mol for  $H \parallel ab$  and  $H \parallel c$ , respectively). This suggests that Pr occupies only about a half of the  $R$  sites (assuming for the TSFZ crystal nearly the same Pr local moment as for the flux grown one). The other half of the  $R$  sites is occupied most probably by the nonmagnetic Ba. Noteworthy, SC with  $T_c \approx 43$  K was observed for Pr<sub>0.5</sub>Ca<sub>0.5</sub>Ba<sub>2</sub>Cu<sub>3</sub>O<sub>7-y</sub> thin films [5]. Ba<sup>2+</sup> as well as Ca<sup>2+</sup> on  $R$  site dopes additional mobile holes and compensates for the localization of holes by the Pr-O(2,3) hybridization. Ba<sup>2+</sup> has a larger ionic radius than Pr<sup>3+</sup> and so the substitution of Ba for Pr could give a natural explanation not only for the SC in TSFZ Pr123 but also for the elongation of the distance between the CuO<sub>2</sub> planes observed in [1].

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