

## An Uncontrolled Walking Toy That Cannot Stand Still

Michael J. Coleman and Andy Ruina

*Department of Theoretical and Applied Mechanics, Cornell University, Ithaca, New York 14853-7501*

(Received 16 June 1997)

We built a simple two-leg toy that can walk stably with no control system. It walks downhill powered only by gravity. It seems to be the first McGeer-like passive-dynamic walker that is statically unstable in all standing positions, yet is stable in motion. It is one of a few known mechanical devices that are stable near a statically unstable configuration but do not depend on spinning parts. Its design is loosely based on simulations which do not predict its observed stability. Its motion highlights the possible role of uncontrolled nonholonomic mechanics in balance. [S0031-9007(98)05783-4]

PACS numbers: 87.45.Dr, 46.10.+z

Human walking on level ground involves dynamic balance which, if viewed in a coarse-grained way, is presumably asymptotically stable. This observed stability of walking must depend on some combination of neurological control and mechanical features. The common view is that neuromuscular control is responsible for this balance. To what extent is neuromuscular coordination of animal locomotion, say human walking, really necessary? The bold proposal of McGeer [1–3] is that much of the stabilization of walking might be understood without control.

That asymptotically stable balance might be achieved without control is somewhat nonintuitive since top-heavy upright things tend to fall down when standing still or, more generally, since dynamical systems often run away from potential energy maxima. Two mechanics issues that bear on such stability considerations are that (i) Hamiltonian (conservative and holonomic) dynamical systems cannot have asymptotic stability, and (ii) conservative nonholonomic systems can have asymptotically (exponentially) stable steady motions in some variables, as recalled in Zenkov *et al.* [4].

Since before the clever patent of Fallis in 1888 [5], there have been two- and four-leg passive-dynamic walking toys that can walk downhill. All such toys that we know about are statically stable when they are not walking. While their motion is engaging to watch, their dynamic stability is perhaps not so great a surprise.

*McGeer's passive-dynamic walkers.*—Inspired by a double pendulum simulation of swinging legs [6] and by simple walking toys, McGeer found two-dimensional, straight-legged and kneed walking models that displayed graceful, stable, humanlike walking on a range of shallow slopes with no actuation (besides gravity) and no control. McGeer termed these motions *passive-dynamic* walking. All of McGeer's successful designs, as well as those of his imitators thus far [7–9], have been more or less constrained against falling over sideways so that their dynamic balance is fore-aft only. These machines cannot stand stably upright except when their legs are spread fore and aft. The dynamic stability of these devices could be dependent on the static stability of this spread-leg configuration which is visited momentarily.

While human walking motion is mostly in the sagittal (fore-aft and vertical) plane, the stability of out-of-plane (sideways) motions is also important. McGeer's [2] numerical 3D studies only led to unstable periodic motions. Fowble and Kuo [10] numerically simulated a passive-dynamic 3D model of walking but also did not find stable passive motions.

Our recent investigations of walking balance have been based on attempts to design mechanisms that vaguely mimic human geometry and walk without control. This paper describes one such primitive design (first reported in [11]) which extends to three dimensions, at least experimentally, McGeer's remarkable two-dimensional walking mechanisms.

*Spinning parts and nonholonomic constraints.*—Humans are notably lacking in gyros, flywheels, or other spinning parts. Things with spinning parts, like tops and gyros, are well known to be capable of balancing near a potential energy maximum. The common model of an energy conserving point-contact gyro, however, does not have asymptotic stability since it is Hamiltonian. Adding a rounded tip to the top, with nonholonomic rolling contact, is not stabilizing. A spinning top with dissipation, however, can be asymptotically stable in a transient sense in that, over a limited time until the spinning rate has slowed too much, vertical motion is approached exponentially. The observed asymptotic stability of rolling coins and the like also depends on dissipation.

We know of only a few uncontrolled three-dimensional devices that can have asymptotically stable steady motions at or near a potential energy maximum, without depending on fast spinning parts. These devices are all nonholonomically constrained and conservative: (i) a “no-hands” bicycle with massless wheels (say skates) and a special mass distribution [12,13], (ii) a no-hands tricycle (where gyroscopic terms do not affect the dynamics) with a mildly soft decentering (negative spring constant) spring on the steering [14,15], (iii) a rigid rider attached appropriately to a moving skateboard [16], and (iv) a statically unstable boat with an ideal keel that is steered by the boat lean. Certain gliding aircraft might also be considered as an example, but defining a potential energy maximum is

less clear for planes since no well defined reference for measuring potential energy exists.

All of these devices differ from walking mechanisms in that they are constrained against fore-aft tipping (the walking devices have fore-aft dynamics), they conserve energy (the walkers lose energy at joint and foot impacts and use up gravitational potential energy), and they are nonholonomically constrained (most of the walkers are well modeled as piecewise holonomic).

*Intermittent contact and nonholonomicity.*—Asymptotically stable mechanical systems must be non-Hamiltonian. Two mechanisms for losing the Hamiltonian structure of governing equations are dissipation and nonholonomic constraints. The primary examples of nonholonomic constraint are rolling contact and skatelike sliding contact. For these two smooth constraints, and other less physical nonholonomic constraints, the set of allowed differential motions is not integrable. That is, the constraints are not equivalent to a restriction of the space of admissible configurations. For smooth nonholonomic systems, the dimension of the configuration space accessible to the system is greater than the dimension of the velocity space allowed by the constraints.

An intermittent nonslipping contact constraint can also cause the dimension of the accessible configuration space to be greater than the dimension of the accessible velocity space. As suggested by one simple example [18], this discrete nonholonomicity may account for exponential stability of some systems. The walking models we study are all nonholonomic in this intermittent sense (and also in the conventional sense if they have rounded feet). They can, for example, translate forwards by walking although the contact constraint does not allow forward sliding.

*Dynamical modeling.*—Figure 1 shows a 3D model which probably captures the essential geometric and mass-distribution features of the physical model presented here. The device, at least at the level of approximation which we believe is appropriate, is a pair of symmetric rigid bodies (leg 1: stance leg; leg 2: swing leg) that have mass  $m$ , symmetrically located (in the rest state) centers of mass  $G_{1,2}$ , and mirror-symmetry-related moment of inertia matrices with respect to the center of mass  $\mathbf{I}_{1,2}$ . The

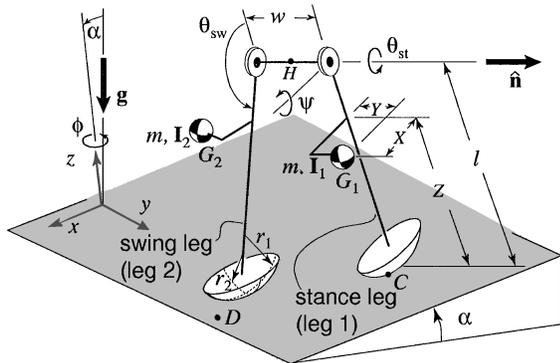


FIG. 1. A rigid body model of the simple walker. Parameters and state variables are described in the text.

legs are connected by a frictionless hinge at the hip with center point  $H$  and orientation  $\hat{\mathbf{n}}$  normal to the symmetry plane of the legs. Each of the two legs can make rolling and collisional contact with the ground (slope =  $\alpha$ ). The gravitational acceleration is  $g$ .

The (reduced) dynamical state of the model is determined by the orientations and angular velocities of the legs. The stance leg orientation is determined by standard Euler angles  $\psi$ ,  $\theta_{st}$ ,  $\phi$  for lean, pitch, and steer, respectively. The configuration of the swing leg is described by the angle  $\theta_{sw}$ . The absolute position of the walker on the plane does not enter into the governing equations. The instantaneous point of contact of the stance leg with the ground is  $C$  and the point of the impending contact is  $D$ . We assume ground collisions are without bounce or slip.

The unreduced accessible configuration space is six dimensional (the above angles plus position on the slope) whereas at any instant in time the accessible velocity space is four dimensional (the four dynamical state variables), hence the overall nonholonomicity ( $6 > 4$ ) of this system. The model is also dissipative due to kinetic energy loss at the collisions.

The model is well posed since the governing equations for rigid bodies in hinged, rolling, and plastic-collisional contact are well established. The equations which govern the evolution of the state of the system  $\mathbf{q} = \{\phi, \dot{\phi}, \psi, \dot{\psi}, \theta_{st}, \dot{\theta}_{st}, \theta_{sw}, \dot{\theta}_{sw}\}$  follow from angular momentum balance (or other equivalent principles). Between collisions, we have angular momentum balance for the whole system about the contact point  $C$

$$\sum_{i=1,2} \mathbf{r}_{G_i/C} \times m\mathbf{g} = \sum_{i=1,2} [\mathbf{r}_{G_i/C} \times m\mathbf{a}_i + \boldsymbol{\omega}_i \times (\mathbf{I}_i \boldsymbol{\omega}_i) + \mathbf{I}_i \dot{\boldsymbol{\omega}}_i], \quad (1)$$

where  $\mathbf{r}_{G_i/C} \equiv \mathbf{r}_{G_i} - \mathbf{r}_C$ , the center of mass velocities and accelerations are  $\mathbf{v}_{1,2}$  and  $\mathbf{a}_{1,2}$ , and the angular velocities are  $\boldsymbol{\omega}_{1,2}$ . Angular momentum balance for the swing leg about the hip axis  $\hat{\mathbf{n}}$  is

$$\hat{\mathbf{n}} \cdot \{\mathbf{r}_{G_2/H} \times m\mathbf{g} = \mathbf{r}_{G_2/H} \times m\mathbf{a}_2 + \boldsymbol{\omega}_2 \times (\mathbf{I}_2 \boldsymbol{\omega}_2) + \mathbf{I}_2 \dot{\boldsymbol{\omega}}_2\}. \quad (2)$$

The eight collisional jump conditions come from continuity of configuration through the collision, conservation of angular momentum of the system about the new contact point  $D$ ,

$$\sum_{i=1,2} \mathbf{r}_{G_i/D} \times m\mathbf{v}_i + \mathbf{I}_i \boldsymbol{\omega}_i \Big|_- = \sum_{i=1,2} \mathbf{r}_{G_i/D} \times m\mathbf{v}_i + \mathbf{I}_i \boldsymbol{\omega}_i \Big|_+ \quad (3)$$

and conservation of angular momentum for the swing leg about the swing hinge axis

$$\hat{\mathbf{n}} \cdot \{\mathbf{r}_{G_1/H} \times m\mathbf{v}_1 + \mathbf{I}_1 \boldsymbol{\omega}_1 \Big|_- = \mathbf{r}_{G_2/H} \times m\mathbf{v}_2 + \mathbf{I}_2 \boldsymbol{\omega}_2 \Big|_+\}, \quad (4)$$

where the respective sides are to be evaluated just before (−) and after (+) foot collision with the ground. The second jump condition Eq. (4) is applied to the same leg as it switches from stance (subscript 1) to swing (subscript 2). Equations (3) and (4) also assume no collisional impulse from the ground to the leg which is just leaving the ground.

The governing equations and jump conditions above are expressed in terms of positions, velocities, and accelerations, which are functions of the state variables. The governing equations are massive expressions (pages long). We assembled the kinematic expressions and governing differential equations using symbolic algebra software (Maple®).

The no-slip rolling condition is that the velocity of the material point in contact at  $C$  is zero. So far, we have studied only a simplification with point-contact feet ( $r_1 = r_2 = 0$ ) and no hip spacing ( $w = 0$ ). In this case, when a foot is on the ground, the contact acts like a ball-and-socket joint and the only nonholonomy is that of intermittent contact. At all times in between collisions, this point-foot system is smooth and holonomic.

In order to study the stability of such systems, following McGeer, we represent an entire gait cycle by a Poincaré map

$$\mathbf{f}(\mathbf{q}_k) = \mathbf{q}_{k+1} \quad (5)$$

from the state of the system  $\mathbf{q}_k$  just after a foot collision to the state  $\mathbf{q}_{k+1}$  just after the next collision of the same foot (two leg swings and two foot collisions per map iteration). We evaluate  $\mathbf{f}$  using numerical integration of Eqs. (1) and (2) between collisions and applying the jump conditions Eqs. (3) and (4), at each foot collision. For this model, the map is seven dimensional ( $8 - 1$ ), but we treat it as eight dimensional for numerical convenience.

Fixed points of the return map  $\mathbf{f}$  [ $\mathbf{q}$  with  $\mathbf{f}(\mathbf{q}) = \mathbf{q}$ ] correspond to periodic gait cycles (not necessarily stable). We find fixed points by numerical root finding on the function  $\mathbf{f} - \mathbf{q}$ , sometimes using fixed points from models with nearby parameter values to initialize searches.

We determine the stability of periodic motions by numerically calculating the eigenvalues of the linearization of the return map at the fixed points. If the magnitudes of some of the eigenvalues are less than one (with all others equal to one), then the fixed point is asymptotically stable in those variables. Because there is a family of limit cycles at different headings one eigenvalue is always one. Because we use eight instead of seven dimensions in our map, one eigenvalue is always zero.

To date, like McGeer [2] and Fowble and Kuo [10] who studied similar simulations, we have found only unstable periodic motions, though less unstable than theirs. A nearly stable case from our numerical studies has maximum eigenvalue modulus of about 1.15, one of exactly one, and the other six less than one. Fore-aft balance has already been achieved with two-dimensional walking models whose stable fixed points we use as starting points for the 3D analysis. Thus the eigenvector associated

with the maximum eigenvalue corresponds to falling over sideways (i.e., is dominated by  $\psi, \dot{\psi}$  components) as expected. The most stable mass distributions we have found do not have very humanlike parameters; each leg has a center of mass closer to the foot than the hip, and laterally displaced at about 90% of the leg length.

In this almost-stable case, the walker’s legs have a mass distribution corresponding roughly to laterally extended balance bars, like what might be used for tightrope walking. In the limit, as the lateral offset of the center of mass gets very large, the device approaches, for sideways balance, an inverted pendulum with large rotational inertia. The step periods remain bounded. Negligible falling acceleration can thus occur in one step and the modulus of the maximum eigenvalue of the linearized step-to-step map asymptotically approaches one, or apparent neutral stability, from above. Thus, the closeness of the largest map eigenvalue modulus to one is not a complete measure of closeness to stability. However, when averaged over a step cycle, this model falls more slowly than a corresponding inverted pendulum and the low eigenvalue is not just a result of slowed falling due to large rotary inertia.

*The toy.*—As a nonworking demonstration of the kinematics and mass distributions in our simulations, and not for walking experiments, we assembled a device similar to the one shown in Fig. 2. It has two straight legs, separated by simple hinges at the hips, laterally extending balance mass rods, and rounded feet. Playing, with no hopes of success, we placed the toy on a ramp. Surprisingly, it took a few serendipitous, if not very steady or stable, steps. After some nonquantifiable tinkering, we arrived at the functioning device shown.

Our physical model is constructed from a popular American child’s construction toy, brass strips to round the feet bottoms, and various steel nuts for balance masses. The walking ramp has about a 4.5° slope and is narrow enough to avoid making contact with the balance

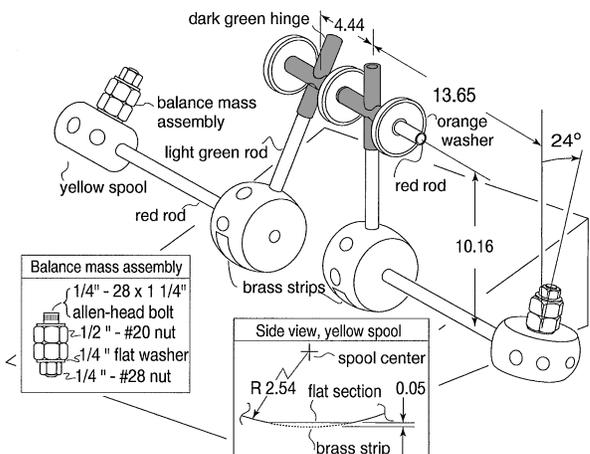


FIG. 2. The 3D Tinkertoy® walking model with hardware description and dimensions (in cm, not drawn to scale). The balance masses and the brass strips are fastened with tape (not shown).

masses as the walker rocks side to side. Another more complex assembly of similar toy parts (not described here) walks on a wide ramp.

*Construction details.*—The device is built using the Playskool® Tinkertoy® Construction System: Colossal Constructions™, 1991 set. One leg is made from a yellow spool, a light green rod, and a dark green hinge (plus “+” shaped) glued together. Then, we slid the legs onto a red rod (loose fit) which acts as an axle. The green hinges are separated and kept from sliding apart by three orange washers, friction fit to the red axle. The legs and axle can rotate independently.

To support the side weights, we glued a yellow spool rigidly to the end of a red rod and inserted the other end into the side of a yellow foot with a friction fit to allow for rotational adjustment.

We assembled each balance mass from two stacked steel nuts held together between two washers by a nut and bolt. Each nut assembly has a mass of about 50 grams. Then, each balance mass assembly was located on the yellow spools at the end of the balance rods and held in place with vinyl electrical tape. The balance mass assembly is tilted behind the leg. As a result, the legs have low mass centers located laterally at a distance comparable to the leg length, above the center of curvature of the feet, and just behind the leg axes. The mass of the fully assembled walking device is about 120 grams, only 20 grams more than the two balance masses. When the toy is in its unstable-equilibrium standing position the nominally vertical legs are approximately orthogonal to the ramp.

To ensure that the walker is statically unstable (cannot stand on the flat sections or in any other way), a small (0.50 cm wide) strip of thin (0.013 cm) brass shim stock material was fastened over the flat section contacting the floor so as to ensure its curvature there.

*Observed motion.*—Because the center of mass is above the center of curvature of the round feet, we cannot stably stand this device with parallel or with splayed legs. When aimed downhill on a ramp, tipped to one side, and released, the device rocks side to side and, coupled with swinging of the legs, takes tiny steps. When a foot hits the ground, it sticks and then rolls, until the swinging foot next collides with the ground. Except at the moment of foot collision, only one foot is in contact with the ground at any time. When the swinging foot collides with the ground, the trailing leg leaves the ground. The gait is more or less steady; after small disturbances the toy either falls or stumbles a few steps while returning to near-periodic gait. At a slope of  $4.5^\circ$ , it takes a step about every 0.47 s and advances forward about 1.3 cm per step, where a step is measured from a foot collision to the next collision of that same foot. The side-to-side tilt is about  $4^\circ$ ; there is no visible variation in  $\phi$  during a step, but there is slight directional drift (one way or another) over many steps. The rounded metal strips at the feet bottom deform during foot collision in a way that may or may not be essential; we do not know.

In conclusion, we have constructed a device which can balance while walking but cannot stand in any configuration. Although our new machine does not have a very humanlike mass distribution, it does highlight the possibility that uncontrolled dynamics may not just contribute to fore-aft walking balance, as indicated by previous McGeer models, but also to side-to-side balance. The mechanism joins a small collection of statically unstable devices which dynamically balance without any rapidly spinning parts.

Our too-simple mathematical/computational model does not explain this behavior. We do not yet know what key modeling features need be included to predict the observed dynamic stability. An open and possibly unanswerable question is whether the stability of this intermittently dissipative system can be explained, in part, by the fact that its piecewise holonomic contact constraints act somewhat like nonholonomic constraints.

The authors thank Les Schaffer, Saskya van Nouhuys, and Mariano Garcia for editorial comments.

- 
- [1] T. McGeer, *Int. J. Robot. Res.* **9**, 62 (1990).
  - [2] T. McGeer, in *Proceedings of the Experimental Robotics II: The 2nd International Symposium*, edited by R. Chatila and G. Hirzinger (Springer-Verlag, Berlin, 1992), pp. 465–490.
  - [3] T. McGeer, *J. Theor. Biol.* **163**, 277 (1993).
  - [4] D. Zenkov, A. Bloch, and J. Marsden, University of Michigan report (unpublished).
  - [5] G. T. Fallis, U.S. Patent No. 376,588 (1888).
  - [6] S. Mochon and T. McMahon, *Math. Biosci.* **52**, 241 (1980).
  - [7] M. Garcia, A. Chatterjee, A. Ruina, and M. J. Coleman, *ASME J. Biomech. Eng.* (to be published).
  - [8] M. J. Coleman, Ph.D. thesis, Cornell University, Ithaca, NY, 1998.
  - [9] B. Thuillot, A. Goswami, and B. Espiau, in *Proceedings of the 1997 IEEE International Conference on Robotics and Automation* (IEEE, New York, 1997), pp. 792–798.
  - [10] J. V. Fowble and A. D. Kuo, in *Biomechanics and Neural Control of Movement, Proceedings of the Engineering Foundation Conferences* (EFC, Mount Sterling, OH, 1996), pp. 28, 29.
  - [11] M. J. Coleman *et al.*, in *Proceedings of the IUTAM Chaos '97 Symposium* (to be published).
  - [12] R. S. Hand, Master's thesis, Cornell University, Ithaca, NY, 1988.
  - [13] J. Papadopoulos (private communication) (unpublished).
  - [14] Y. Rocard, *General Dynamics of Vibrations* (Frederick Ungar Publishing Co., New York, 1960), 3rd ed.
  - [15] R. S. Sharp, in *Proceedings of the 8th IAVSD Symposium: The Dynamics of Vehicles on Roads and Tracks*, edited by J. K. Hedrick (Swets and Zeitlinger B. V., Lisse, 1983), pp. 564–577.
  - [16] M. Hubbard, *J. Appl. Mech.* **46**, 931 (1979).
  - [17] T. Cardanha and R. Bennet, Cornell University Engineering Undergraduate Research Project (unpublished).
  - [18] A. Ruina, *Rep. Math. Phys.* (to be published).