## Comment on "Amplification under the Standard Quantum Limit"

Quantum limits on optical amplification have been widely discussed theoretically in the last decade [1], and several experiments [2–5] have crossed the standard quantum limit (SQL). In a recent Letter [6], D'Ariano, Machiavello, and Paris reported on Monte Carlo simulation of a saturated laser amplifier, and claimed to beat the SQL in this system.

This claim, however, is based on a binary "on"-"off" *large-signal* noise-figure calculation in a saturated, *nonlinear* device. Here we point out that the 3 dB SQL for a phase-insensitive amplifier only has a well defined meaning in a *linear* regime [1], i.e., either for a small signal linearized regime or for a truly linear device. For a nonlinear device which modifies the statistics between input and output, the noise figure based on a large signal, as in [6], must be used with caution since the noise figure is no longer easily related to the information degradation as expressed in symbol error probability [or bit error rate (BER)]. For digital coding, BER, not the signal-to-noise ratio (SNR), is the quantity of relevance to evaluate the information transfer efficiency [7].

In Ref. [6], the gain G and the noise figure  $\mathcal{R}$  are, respectively, defined by

$$G = \frac{S_{\text{out}}}{S_{\text{in}}}, \quad \mathcal{R} = \frac{(S^2/\mathcal{N})_{\text{in}}}{(S^2/\mathcal{N})_{\text{out}}},$$
 (1)

where S and  $\mathcal{N}$  denote signal and noise at the input (in) and at the output (out) of the amplifier. In [6], the quantities S and  $\mathcal{N}$  are defined for on-off modulation as

$$S = \langle \hat{O} \rangle^{\text{on}} - \langle \hat{O} \rangle^{\text{off}}, \quad \mathcal{N} = \frac{1}{2} (\langle \Delta \hat{O}^2 \rangle^{\text{on}} + \langle \Delta \hat{O}^2 \rangle^{\text{off}}),$$
(2)

where  $\hat{O}$  is the detected observable, and the brackets  $\langle \cdots \rangle$  denote an ensemble average.

Let us show by an example that, for a nonlinear device, the noise figure  $\mathcal{R}$  defined as above [6] does not account for the efficiency of the information transfer of the system, and therefore is not relevant to the SQL. Let us consider an optical repeater based on a detector plus electronic decision circuitry triggering a powerful laser. The transfer function of this device is described as  $\hat{O}_{\text{out}} = \{P_o, \text{ if } \hat{O}_{\text{in}} \geq P_{\text{th}}, \text{ and 0, if } \hat{O}_{\text{in}} < P_{\text{th}} \}$ , where  $P_{\text{th}}$  is the decision threshold. Assume now equal probability (equal average numbers of zeros and ones) on-off keying with zero photons in the off state and a coherent state, with  $\langle n \rangle$  photons on the average in the on state. The

input SNR is  $(S^2/\mathcal{N})_{in} = 2\langle n \rangle$ . Setting, for example, the decision threshold  $P_{\rm th}$  to one input photon, the output SNR is given by  $(S^2/\mathcal{N})_{\text{out}} = 2(e^{\langle n \rangle} - 1)$ , leading to a noise figure  $\mathcal{R} = \langle n \rangle / (e^{\langle n \rangle} - 1)$ . Here,  $P_o$  was assumed large enough that the relative output laser fluctuations could be neglected. As shown by this example,  $\mathcal{R}$  (as it is defined in [6]) not only can be less than 3 dB, but also less than 0 dB (unity), corresponding to an improvement of SNR. Yet, there is of course no improvement of the BER since symbol errors are introduced when a digital "one" is interpreted as a "zero," which is not reduced by a threshold function. This illustrates that for a nonlinear device the noise figure  $\mathcal{R}$  of [6] cannot be used directly to assess the information degradation, and hence  $\mathcal{R}$  cannot be compared to the 3 dB SQL for a linear phase-insensitive amplifier. Thus, the conclusion in Ref. [6] on "amplification below the standard quantum limit" is unwarranted.

If the noise figure instead is defined for a small signal input, which should be a small linearizable modulation  $S = \delta \langle \hat{O} \rangle$  around a given working point, the model used by the authors of [6] would likely show that a saturated laser amplifier does not go below the SQL.

We acknowledge useful discussions with Philippe Grangier and Jean-François Roch.

Olle Nilsson, Anders Karlsson, Jean-Philippe Poizat, and Eilert Berglind

<sup>1</sup>Ellemtel Utveckling AB, 125 25 Älvsjö, Sweden <sup>2</sup>Laboratory of Photonics and Microwave Engineering Department of Electronics, Royal Institute of Technology S-164 40 Kista, Sweden

<sup>3</sup>Institut d'Optique Théorique et Appliquée B.P. 147, 91403 Orsay, France

Received 16 March 1995 PACS numbers: 03.65.Bz, 42.50.Dv, 42.65.Ky

- [1] C. M. Caves, Phys. Rev. D 26, 1817 (1982).
- [2] Z. Y. Ou, S. F. Pereira, and H. J. Kimble, Phys. Rev. Lett. 70, 3239 (1993).
- [3] E. Goobar, A. Karlsson, and G. Björk, Phys. Rev. Lett. **71**, 2002 (1993).
- [4] J.-F. Roch, J.-Ph. Poizat, and P. Grangier, Phys. Rev. Lett. 71, 2006 (1993).
- [5] J. A. Levenson, I. Abram, Th. Rivera, and Ph. Grangier, J. Opt. Soc. Am. B 10, 2233 (1993).
- [6] G. M. D'Ariano, C. Machiavello, and M. G. A. Paris, Phys. Rev. Lett. 73, 3187 (1994).
- [7] This point arose already in the comments by J. H. Shapiro, IEEE Phot. Tech. Lett. 4, 647 (1992) and P. A. Humblet, *ibid.* 4, 650 (1992), on a paper by K. Kikuchi, *ibid.* 4, 195 (1992).