ERRATA

Electroweak String Configurations with Baryon Number [Phys. Rev. Lett. 73, 373 (1994)]

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Equation (6) should have a factor of 2 on the right hand side [since $\epsilon_{ijk}Z^{jk} = 2(\vec{B}_Z)_i$]. This changes Eq. (11) by a factor of 2 to $N_{CS}(\mathrm{in}) = 2N_F\cos(2\theta_W)$ and similarly Eq. (16) to $Q_B = 2N_F\cos(2\theta_W)$. However, the helicity of the twisted loop in Fig. 1(b) is still $N_F\cos(2\theta_W)$—half that of the loops in Fig. 1(a)—as shown by R. I. Ricca and H. K. Moffatt [in *Topological Aspects of the Dynamics of Fluids and Plasma,* edited by H. K. Moffatt *et al.,* NATO ASI Ser. E, Vol. 218 (Kluwer Academic Publishers, Dordrecht, 1992)], and hence the rest of the arguments in the paper remain unchanged.

We are grateful to Jaume Garriga for spotting the factor of 2 in Eq. (6) and to Hoi-Kwong Lo for the factor of $\frac{1}{2}$ in the helicity of the twisted loop.

Suppression of Weak Localization due to Magnetic Flux in Few-Channel Ballistic Microstructures [Phys. Rev. Lett. 73, 2115 (1994)]

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Because of a missing factor of π in the original calculation of the factor k, the low-energy value of this constant (corresponding to $n \approx 400$ electrons), which is discussed in the first column of p. 2116, is in fact around $k \approx 3.9$. This changes the number for the theoretical FWHM to be compared with the experimental value. In this comparison, which appears in the second column of p. 2118, we should now say the following: Estimating the number n of electrons in the stadium to be about two thousand, we find $k \approx 5.8$. For a stadium of area 0.5 $(\mu\mathrm{m})^2$, this yields 0.43 for the height of the Lorentzian and 2.7 mT for the FWHM. These values differ somewhat from what we read off the Lorentzian shown as a dotted line in Fig. 3(b) of the second Ref. [1], namely 1/3 and 5 mT, respectively. The discrepancy is consistent with the presence in the experimental data of effects, such as finite temperature and inelastic scattering, that will tend to wash out coherent phenomena and manifest themselves as a flattening of the peak and broadening of the width.

The expression for x_c given in the second column of p. 2116 should read $x_c = \pi N v_c^2/\lambda$. In the line above Eq. (7), one should note that k is no longer close to unity. In the first column of p. 2118, the expression for τ_{mix} should read: $\hbar/\tau_{\mathrm{mix}} = (2\pi)\,\overline{(H_{\mathrm{mag}})^2}/d$, where H_{mag} is the second term on the right hand side of Eq. (1).

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