## Comment on "VO<sub>2</sub>: Peierls or Mott-Hubbard? A View from Band Theory"

In a recent Letter [1] Wentzcovitch, Schulz, and Allen reopened the discussion on the nature of the metalinsulator transition in vanadium dioxide (VO<sub>2</sub>). On the basis of LDA (local density approximation) calculations on the monoclinic  $M_1$  structure which gave a semimetal with a very small number of carriers, they concluded that this phase is an ordinary band (Peierls) insulator and not a Mott-Hubbard insulator in which onsite Coulomb correlations dominate. The authors, however, presented results only for one of the insulating phases of VO<sub>2</sub>: They and others have pointed out that all V<sup>4+</sup> ions are in singlet V-V pairs which makes a clear distinction impossible. In this Comment we wish to remark that this ambiguity does not apply to the other insulating phases of VO<sub>2</sub>, which are clearly of the Mott-Hubbard type.

There are two structural components to the lattice distortion from the high temperature rutile (R) phase, namely a pairing and a twisting of V-V pairs out of the rutile axis  $c_r$ . Around twenty years ago the nature of a second monoclinic insulating phase ( $M_2$ ) was established [2]. The metal-insulator transition is  $R \rightarrow M_2$ in  $V_{1-x}Cr_xO_2$  with x very small ( $x \ge 3 \times 10^{-3}$ ) [2] or in pure VO<sub>2</sub> [3] when a small uniaxial pressure is applied along (110)<sub>r</sub>. In  $M_2$  one-half of the V chains of the R phase pairs but does not twist and the other half twists but

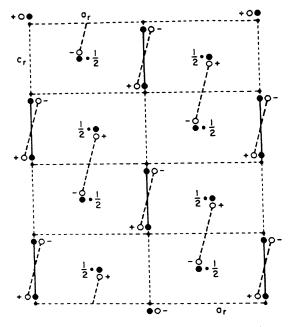


FIG. 1. Comparison of V-V pairing in the three phases  $(R, M_1, \text{ and } M_2)$ . In  $M_1$  (open circles) all the vanadium atoms both pair and twist from the rutile positions. In  $M_2$  (filled circles) one-half of the vanadium atoms pairs but does not twist and the other half forms unpaired zigzag chains. (The distortions are exaggerated by a factor of 2 for clarity.)

does not pair (see Fig. 1). There are simple electrostatic reasons that the pairing on one set of V chains induces a twist in the other V chains. Thus the  $M_1$  phase of VO<sub>2</sub> can be viewed as a simple superposition of two lattice distortions of the  $M_2$  type.

In  $M_2$  one-half of the V<sup>4+</sup> ions form equally spaced V chains, and NMR and EPR experiments show that they behave magnetically as  $s = \frac{1}{2}$  Heisenberg chains ( $J \approx$ 300 K). It is clear that these V chains in  $M_2$  are magnetic (or Mott-Hubbard) insulators. Since  $M_2$  can be stabilized by minimal perturbations, it is also clear that  $M_2$  is a local minimum for VO<sub>2</sub>, whose free energy is only very slightly higher than that of  $M_1$  at room temperature and pressure. Further cooling leads to a continuous  $M_2 \rightarrow M_1$ transition through an intermediate insulating triclinic (T) phase where pairing (or dimerization) on one set of V chains grows continuously [4].

The  $M_2$  phase is a Mott-Hubbard insulator and the  $M_1$  phase is a superposition of two  $M_2$ -type lattice distortions. Further a continuous  $M_2 \rightarrow M_1$  transition through the intermediate insulating T phase is observed with decreasing temperature. These experiments led us to conclude previously [2,3] that all the insulating phases of VO<sub>2</sub>,  $(M_1, M_2, \text{ and } T)$  were of the same type and should be classified as Mott-Hubbard and not band insulators.

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- [4] Note, however, the initial transition  $M_2 \rightarrow T$  is discontinuous.

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