

Electronic Instability at High Flux-Flow Velocities in High- T_c Superconducting Films

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At high flux-flow velocities in type-II superconductors the nonequilibrium distribution of the quasiparticles leads to an electronic instability and an abrupt switching into a state with higher electric resistivity, as predicted by Larkin and Ovchinnikov (LO). We report the first observation of this effect in a high-temperature superconductor, namely in epitaxial c -axis oriented films of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. Using the LO theory, we have extracted from our results the inelastic quasiparticle scattering rate τ_{in}^{-1} , which strongly decreases with decreasing temperature below T_c .

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Vortex motion and energy dissipation in the mixed state of high- T_c superconducting films has been intensively studied during recent years. These studies focused predominantly on the mechanisms controlling vortex motion in the limit of low values of the driving Lorentz force or thermal force. Hence, in the past the onset of dissipation and the vortex dynamics in the limit of low vortex velocities occupied the center of interest. In the present paper, we report for the first time on experiments dealing with the opposite limit, namely the maximum possible vortex velocity established at high values of the driving force. In this regime an interesting electronic instability has been predicted about 20 years ago by Larkin and Ovchinnikov (LO) [1], resulting in a discontinuity in the current-voltage characteristic (IV) and an abrupt switching of the sample into another state of higher electric resistivity. The experimental observation of this effect for low- T_c superconductors has been reported some years ago [2,3], and reasonable agreement with the LO theory has been found. In this phenomenon, the relaxation rate of the quasiparticle energy plays a central role, and it is this feature which makes an experimental study of this effect in the high-temperature superconductors particularly interesting. Recently, inelastic scattering between quasiparticles has played an increasing role in the theoretical discussions of the electronic structure and the mechanism underlying high-temperature superconductivity [4,5]. Furthermore, inelastic quasiparticle scattering and its temperature dependence has become an important issue in the interpretation of recent experiments on the high-frequency conductivity [6,7], infrared optics [8], and thermal conductivity [9] of high-temperature superconductors.

We report on the first experimental observation in a high-temperature superconductor of the electronic instability due to quasiparticle excitation in the normal vortex core at high flux-flow velocities, as predicted by LO [1]. The effect manifests itself in an abrupt increase of the electric resistance of the sample. Our measure-

ments were performed with epitaxial c -axis oriented $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ films and clearly indicate that the inelastic electron-electron scattering rate strongly decreases with decreasing temperature below T_c . This temperature dependent quasiparticle scattering plays an increasing role in recent discussions of the mechanism for high-temperature superconductivity and of the underlying electronic structure.

The main point of the LO theory is the following. In the presence of the electric field resulting from current-induced vortex motion the nonequilibrium distribution of the quasiparticles in the superconductor is shifted upwards in energy compared to the equilibrium distribution. As a consequence, quasiparticles leave the potential well of the vortex core, the vortex core shrinks, viscous damping of the vortex motion is reduced, and the vortex velocity normalized to the driving force becomes larger. According to LO, the viscous damping coefficient η at vortex velocity v_φ is

$$\eta(v_\varphi) = \frac{\eta(0)}{1 + (v_\varphi/v_\varphi^*)^2}, \quad (1)$$

where $\eta(0)$ is the damping coefficient at $v_\varphi = 0$ and v_φ^* a critical vortex velocity at which the damping force ηv_φ attains a maximum. Because of this effect the IV characteristic is curved upwards and at the velocity v_φ^* an instability point is reached where the sample switches into another state. This instability point manifests itself as a distinct discontinuity in the IV characteristic. The nonequilibrium quasiparticle distribution and, hence, the value of v_φ^* depends sensitively on the inelastic electron-electron scattering rate τ_{in}^{-1} which tends to move the quasiparticle system back to its equilibrium distribution. A qualitative discussion of this point based on simple physical concepts has been given elsewhere [3]. For the critical velocity v_φ^* LO obtain the expression

$$v_\varphi^* = \frac{D[14\zeta(3)]^{1/2}(1 - T/T_c)^{1/2}}{\pi \tau_{\text{in}}}. \quad (2)$$

Here D is the quasiparticle diffusion coefficient ($D = v_F l/3$, with v_F the Fermi velocity and l the electron mean free path); $\zeta(x)$ is the Riemann zeta function, and τ_{in} is the inelastic quasiparticle scattering time. The flux-flow voltage corresponding to v_ϕ^* is

$$\underline{V}^* = -(\underline{v}_\phi^* \times \underline{B})L, \quad (3)$$

where \underline{B} is the magnetic field and L the sample length between the voltage contacts. As discussed by LO, at voltages $V \leq V^*$ the curved IV characteristic is expected to follow the equation

$$I - I_c = \left[\frac{V}{1 + (V/V^*)^2} + V \left(1 - \frac{T}{T_c}\right)^{1/2} \right] \frac{1}{R_f}, \quad (4)$$

where I_c and R_f are the critical current and the flux-flow resistance, respectively. In an experimental study of this electronic nonequilibrium phenomenon great care must be exercised for eliminating heating effects.

Our sample films had a thickness of 60–100 nm and were deposited on single-crystalline MgO substrates by laser ablation. Microfabrication of the four-point sample geometry was performed by standard photolithography. For attaching the current and voltage leads, silver contact pads as large as $2 \times 2 \text{ mm}^2$ area were placed on top of the superconducting film in this way minimizing sample heating due to the contact resistance. The length (between voltage leads) and width of the sample film was typically 100 and 20 μm , respectively. The MgO substrate was in good thermal contact with a temperature-regulated copper block of the sample holder, allowing the variation of the sample temperature between 10 and 300 K. The samples studied in this paper had a sharp resistive transition within typically 2 K and zero resistance critical temperature T_{c0} of 89 K. At 100 K the resistivity was 70–80 $\mu\Omega \text{ cm}$. A magnetic field could be applied parallel to the c axis.

For substrates we have selected MgO, since the thermal conductivity in this material is about 10 times higher compared to single-crystalline SrTiO₃ [10]. The heat transfer coefficient between the YBa₂Cu₃O_{7- δ} film and MgO is about 2000 W/cm²K [11], thereby limiting the temperature rise due to the power dissipation in the sample films measured by us to values not larger than about 0.5 K. In addition, in order to reduce effects due to substrate heating to an acceptable low level, we found that it was necessary to measure the IV characteristic beyond the instability point using a rapid single-pulse method. The current was swept up and down in a single triangular pulse, and the voltage was recorded simultaneously. The time scale finally adopted for an up-sweep was typically 20 ms. Our instrumentation consisted of a HP 8165 Programmable Signal Source and a Tektronix model 11201A Digitizing Oscilloscope combined with a PAR model 113 preamplifier. We have studied three different epitaxial YBa₂Cu₃O_{7- δ} films, all showing similar results.

A typical set of experimental results is shown in Fig. 1. Here the applied current I and the voltage V measured during a single pulse is plotted versus time. The maximum current was selected such that the discontinuity in the voltage could clearly be observed. At $V = 225 \text{ mV}$ the temporal dependence of the voltage abruptly changes its slope turning into a nearly vertical branch. Note that the triangular current pulse slightly extends beyond this point, before the down-sweep of the current begins. The voltage $V = 225 \text{ mV}$ is interpreted as the flux-flow voltage V^* , corresponding to the critical flux-flow velocity v_ϕ^* according to Eq. (3). We have also confirmed that the functional form of Eq. (4) is well satisfied by the experimentally observed IV characteristics for $V \leq V^*$. The velocity v_ϕ^* is plotted as a function of magnetic field B for three temperatures in Fig. 2. As expected from Eq. (2), we see that v_ϕ^* is nearly independent of B . This clearly allows an accurate determination of v_ϕ^* . In Fig. 3 we plot the velocity v_ϕ^* versus temperature.

Our interpretation of the values of V^* and v_ϕ^* in terms of an electronic instability of the type discussed by LO (in contrast to an instability dominated by heating) is supported, among others, by the following experimental facts. (1) v_ϕ^* was found to be independent from B as expected from Eq. (2). On the other hand, the voltage V^* increased proportional to B , whereas the current I^* , at which V^* is reached, remained approximately constant. Although in some cases by changing B the power dissipation I^*V^* varied by a factor as large as 10, the independence of v_ϕ^* from B remained unaffected. As shown theoretically [12] and demonstrated experimentally [13] sample heating leads to distinct deviations from the proportionality between V^* and B and from the B independence of I^* . Both kinds of deviations we did not observe. However, using SrTiO₃ instead of MgO as substrate material, both kinds of deviations clearly appeared indicating effects due to substrate heating. (2) Variation of the time scale of the up-sweep of the current from 20 ms upwards by a factor of 3 or downwards by

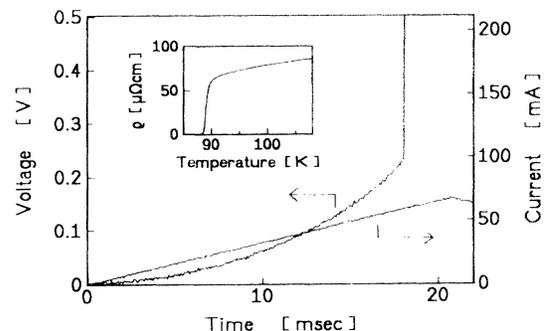


FIG. 1. Voltage (left vertical axis) and current (right vertical axis) versus time. $T = 72 \text{ K}$, $B = 1.8 \text{ T}$. The inset shows the resistivity versus temperature near the superconducting transition.

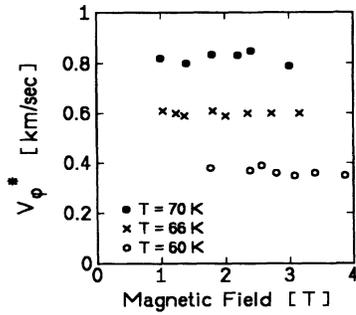


FIG. 2. Critical flux-flow velocity v_{ϕ}^* versus magnetic field for three temperatures. Same sample as in Fig. 1.

a factor of 50, left the values of V^* and I^* unchanged within 1%. However, by increasing this time scale beyond about 100 ms, we noted the first tendency for an increase of V^* and a decrease of I^* , as expected when sample heating starts to play a role. (3) By limiting the maximum current of the triangular current pulse to values $I \leq I^*$, at which the voltage instability is reached, the voltage could be reversibly traced backwards during the downturn of the current. However, if this maximum current exceeded I^* , the large increase in voltage and power dissipation caused thermal hysteresis and a deviation of this reversible behavior.

Because of the arguments given above, we feel confident that our interpretation of the observed break of the IV characteristic at V^* or v_{ϕ}^* in terms of an electronic instability and not a thermal one is justified. Using Eq. (2) derived by LO for calculating the inelastic scattering rate τ_{in}^{-1} , we obtain the plot of τ_{in}^{-1} versus temperature shown in Fig. 4. In this calculation we have used the values $v_F = 1 \times 10^7$ cm/s [14] and $l = 5$ nm [15].

As seen in Fig. 4, again we clearly find the rapid decrease of τ_{in}^{-1} with decreasing temperature below T_c , as observed already in Refs. [6–9]. Our determination of τ_{in}^{-1} could not be extended appreciably above 82 K, since at these temperatures the scattering rate τ_{in}^{-1} quickly approaches the value for electron-phonon scattering of typically $10^{12} - 10^{13}$ s $^{-1}$ [6,7,15], thereby leaving the validity range of the LO theory. Experimentally, above

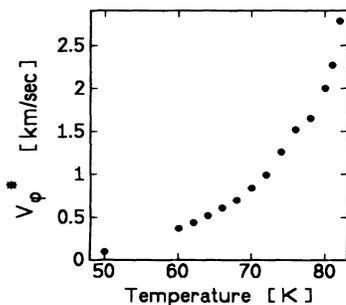


FIG. 3. Critical flux-flow velocity v_{ϕ}^* versus temperature. Same sample as in Fig. 1.

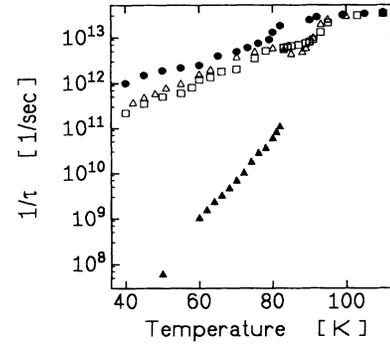


FIG. 4. Inelastic quasiparticle scattering rate τ_{in}^{-1} (\blacktriangle) obtained from Eq. (2) versus temperature. Same sample as in Fig. 1. For comparison we also show the scattering rate obtained from measurements of the microwave surface impedance and reported in Ref. [7] (\bullet) and Ref. [6] (sample A \square , sample B \triangle).

82 K the value of v_{ϕ}^* obtained from V^* shows a distinct B dependence, making an unambiguous determination of v_{ϕ}^* impossible. For comparison we also show in Fig. 4 the quasiparticle scattering rate obtained recently from measurements of the microwave surface impedance of YBaCuO films [7] and single crystals [6]. It is interesting that the decrease of the scattering rate τ_{in}^{-1} with decreasing temperature displayed by our data is much more rapid than that of the microwave data. This difference in behavior can possibly be explained by the fact that the microwave results are much more sensitive to extrinsic scattering mechanisms (defects, impurities, surface irregularities, etc.) than the nonlinear IV characteristic measured in our experiments. We emphasize that the IV characteristic is only determined by the flux-flow process. Furthermore, it is only the *inelastic* quasiparticle scattering affecting the critical velocity v_{ϕ}^* , whereas the microwave losses are also due to the elastic processes. Although a detailed theoretical treatment still needs to be done, it appears that the scattering rate τ_{in}^{-1} is dominated by inelastic electron-electron scattering and that the opening of a gap in the quasiparticle energy spectrum below T_c will be crucial for explaining the rapid suppression of τ_{in}^{-1} [16,17]. A crude estimate of the temperature dependence of the scattering rate τ_{in}^{-1} can be obtained as follows. Noting that the Cooper pair binding energy is twice the energy gap Δ , the thermally excited quasiparticle density n_{qp} is proportional to $\exp(-2\Delta/k_B T)$. Hence, we find the proportionality $\tau_{in}^{-1} \sim n_{qp}^2 \sim \exp(-4\Delta/k_B T)$. Using for approximation the zero-temperature limit $\Delta(0) \approx 2k_B T_c$, we obtain $\tau_{in}^{-1} \sim \exp(-8T_c/T)$ and, finally, in our case of $T_c = 89$ K $\tau_{in}^{-1} \sim \exp[(-710 \text{ K})/T]$. Fitting our data in Fig. 4 with the proportionality $\tau_{in}^{-1} \sim \exp(-T^*/T)$ we obtain $T^* = 950$ K, which is reasonably close to our estimated value of 710 K. On the other hand, electron-phonon scattering is clearly inconsistent with this rapid

suppression of τ_{in}^{-1} below T_c . Attempting to fit our data with a power law of the type $\tau_{in}^{-1} \sim T^n$, we obtained an unphysically high value $n \approx 14$.

The distance $v_\phi^* \tau_{in}$ calculated from our data presented in Figs. 3 and 4 increases with decreasing temperature, with a value of 30 nm at 80 K and 1.5 μm at 50 K. This range of values is clearly much larger than the diameter of the normal vortex core (given by twice the coherence length). Noting that at a magnetic field of 1 T the intervortex distance is about 20 nm, we see that, in the range of magnetic fields and temperatures we have studied, near the LO instability the nonequilibrium distribution of the quasiparticles extends over the whole volume of the superconductor.

To conclude, we have detected an interesting electronic instability in high- T_c superconducting films due to the nonequilibrium quasiparticle distribution at high flux-flow velocities. Nonequilibrium effects in low- T_c superconductors have been studied extensively some time ago, and summaries can be found in Refs. [18] and [19]. In high- T_c superconductors this subject is just beginning to be investigated. Here the possible influence of the doping level on the nonequilibrium effects appears particularly interesting.

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