

Comment on "Moriya's Anisotropic Superexchange Interaction, Frustration, and Dzyaloshinsky's Weak Ferromagnetism"

In a recent Letter [1], Shekhtman, Entin-Wohlman, and Aharony have shown that the single-bond anisotropic superexchange interaction [2] does not provide anisotropy, and that frustration is a necessary condition for yielding weak ferromagnetism in La_2CuO_4 . Their arguments have been based entirely on the single-band Hubbard model. In this Comment, we would like to point out that, in a realistic model with multiorbitals on the ligand oxygen ions, the interaction between the two spins on the neighboring Cu ions is already anisotropic, and that this single-bond anisotropy, rather than frustration among the bonds, is in fact the origin of the weak ferromagnetism in La_2CuO_4 .

A relevant model, which one starts with, must contain in-plane oxygen orbitals coming into play due to lattice distortions. Let us write the one-hole Hamiltonian

$$\begin{aligned}
 H = & \sum_{\mathbf{R}'} \sum_n \sum_{\sigma} \varepsilon_n(\mathbf{R}') p_{n\sigma}^{\dagger}(\mathbf{R}') p_{n\sigma}(\mathbf{R}') \\
 & + \sum_{\mathbf{R} \neq \mathbf{R}'} \sum_n \sum_{\sigma} [b_n(\mathbf{R} - \mathbf{R}') d_{\sigma}^{\dagger}(\mathbf{R}) p_{n\sigma}(\mathbf{R}') + \text{H.c.}] \\
 & + \sum_{\mathbf{R} \neq \mathbf{R}'} \sum_n \sum_{\sigma\sigma'} \{d_{\sigma}^{\dagger}(\mathbf{R}) [\mathbf{C}_n(\mathbf{R} - \mathbf{R}') \cdot \boldsymbol{\sigma}]_{\sigma\sigma'} p_{n\sigma'}(\mathbf{R}') + \text{H.c.}\}
 \end{aligned} \quad (1)$$

in an extended notation of Ref. [1]. \mathbf{R} runs over the Cu sites and \mathbf{R}' runs over the O sites. $d_{\sigma}^{\dagger}(\mathbf{R})$ creates a hole with spin σ in the ground-state $3d$ orbital at site \mathbf{R} , and $p_{n\sigma}^{\dagger}(\mathbf{R}')$ creates a hole with spin σ in the in-plane oxygen $2p_n$ ($n=x,y,z$) orbital at site \mathbf{R}' , where the z axis is taken perpendicular to the CuO_2 plane. It is then proved, by following a recipe of Ref. [1], that the one-bond part of the Hamiltonian (1) is isotropic *only if* $\mathbf{C}_n(\mathbf{R} - \mathbf{R}')/b_n(\mathbf{R} - \mathbf{R}')$ is independent of n , i.e., in this case the unitary transformation of operators

$$\begin{aligned}
 d_{\sigma}(\mathbf{R}_1) & \rightarrow \hat{\mathbf{d}}_{\sigma}(\mathbf{R}_1) = \sum_{\sigma'} [e^{-i(\theta/2)\hat{\mathbf{e}} \cdot \boldsymbol{\sigma}}]_{\sigma\sigma'} d_{\sigma'}(\mathbf{R}_1), \\
 d_{\sigma}(\mathbf{R}_2) & \rightarrow \hat{\mathbf{d}}_{\sigma}(\mathbf{R}_2) = \sum_{\sigma'} [e^{i(\theta/2)\hat{\mathbf{e}} \cdot \boldsymbol{\sigma}}]_{\sigma\sigma'} d_{\sigma'}(\mathbf{R}_2),
 \end{aligned} \quad (2)$$

transforms the one-bond Hamiltonian into the isotropic one, where $\hat{\mathbf{e}}$ denotes a unit vector in the direction of $\mathbf{C}_n(\mathbf{R}_1 - \mathbf{R}')/b_n(\mathbf{R}_1 - \mathbf{R}')$, and \mathbf{R}_1 and \mathbf{R}_2 are the neighboring Cu sites between which the O site \mathbf{R}' locates. In the real situations, however, $\mathbf{C}_n(\mathbf{R}_1 - \mathbf{R}')/b_n(\mathbf{R}_1 - \mathbf{R}')$ cannot be independent of n . We thus have an *anisotropic single bond*. In Ref. [3], we have derived the full spin Hamiltonian including the oxygen $2p_z$ orbitals and found that, in the low-temperature orthorhombic phase of La_2CuO_4 , the single-bond anisotropy is characteristic of

the weak-ferromagnetic spin canting.

Shekhtman, Entin-Wohlman, and Aharony [1] argued that their Hamiltonian can be mapped onto the Dzyaloshinsky thermodynamic potential via Eqs. (12a) and (12b) of Ref. [1], and thus the weak-ferromagnetic ordered state is realized in the presence of the frustration. We have reexamined the ground-state spin structure of the system by a finite-size exact diagonalization technique [3]. It is found that for their Hamiltonian (i.e., in the system without the effect of the $2p_z$ orbitals), the frustration works to suppress the correlation for the uniform magnetization rather than to raise it. Thus, the ground state of this two-dimensional quantum system seems not to follow the behavior suggested in Ref. [1]. However, for the Hamiltonian including the $2p_z$ -orbital contribution, a strong enhancement of the correlation for the uniform magnetization is observed, which straightforwardly leads to the emergence of the weak ferromagnetism. The single-bond anisotropy due to the oxygen multiorbitals is thus essential. The frustration among different bonds seems insufficient. The weak ferromagnetism recently observed in the low-temperature tetragonal phase of $\text{La}_{2-x}\text{Nd}_x\text{CuO}_4$ [4] can also be explained in terms of the single-bond anisotropy due to the $2p_z$ -orbital contribution [3]; it should be noted that the frustration among different bonds has no possibility for yielding the weak ferromagnetism in this phase [5].

Finally, let us note that our arguments are not only relevant for the La_2CuO_4 -type crystals, but are to imply generally that when one studies the magnetism which arises from spin-orbit coupling under distorted lattices, it is of essential importance to examine contributions from multiorbitals of the ligand ions.

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