
ERRATA

Microscopic Theory of Periodic Conductance Resonances in Narrow Channels
[Phys. Rev. Lett. 70, 1513 (1993)]

N. F. Johnson and M. C. Payne

On p. 1515, the following sentence “However, we emphasize that the N -electron energies calculated in this paper do not include *all* exchange-correlation effects arising from our model electron-electron interaction” should have read “However, we emphasize that the N -electron energies calculated in this paper do include *all* exchange-correlation effects arising from our model electron-electron interaction.” A paraphrased version of this sentence appeared correctly earlier in the text on p. 1513.

The correct page number in Ref. [16] by F. Calogero is p. 2197.

Rigorous Results on the Ferromagnetism and Metal-Insulator Transition of the Anderson Lattice at Quarter Filling
[Phys. Rev. Lett. 70, 2024 (1993)]

Takashi Yanagisawa

Our propositions (Theorems 1 and 2) are valid for the one-dimensional Anderson lattice with open boundary conditions. Two- or three-dimensional models are not resolved. The theorem should read as follows.

Theorem 1.—Let us assume that $N_e = N + 1$ and U is infinitely large. If the level of f electrons are sufficiently deep, then the ground state of the 1D Anderson lattice Hamiltonian with open boundary conditions has the total spin $S = (N - 1)/2$ and is unique apart from S^Z degeneracy.

A proof for $N_e < N$ is not correct (lines 4–8 in the right column on page 2025), thus it should be omitted.

The author expresses his sincere thanks to Professor K. Ueda for pointing out vague places in the paper.

From Sideward Flow to Nuclear Compressibility
[Phys. Rev. Lett. 70, 2062 (1993)]

Qiubao Pan and Paweł Danielewicz

After this paper was published, we discovered discrepancies between the constraints imposed on the momentum-dependent potentials, the given parameter values, and the values used in the calculations. The proper values for the SM potential are $\mathcal{A} = 362$ MeV, $\mathcal{B} = -305$ MeV, $\mathcal{C} = -71$ MeV, $\gamma = 0.9$, and $\Lambda = 400$ MeV, and for the HM potential are $\mathcal{A} = 76.4$ MeV, $\mathcal{B} = 51.7$ MeV, $\mathcal{C} = -104$ MeV, $\gamma = 1.9$, and $\Lambda = 550$ MeV. The potentials yield compressibilities $K = 176$ and 323 MeV, respectively. At $\rho = \rho_0$ and high p the two potentials approach values $U \simeq 0$ and $U \simeq 35$ MeV, respectively, and they further give rise to an effective mass $m^* = 0.70m$ at $p \simeq p_F$.

If the two potentials were required to reach the same value at high p , the differences between the predictions for the Nb+Nb system in the paper would be somewhat reduced. The SM potential is closer, at high momenta, to the AV14+UVII potential of Ref. [13] (or to GBD potential), than to UV14+TNI (Fig. 1). Conclusions are not changed.