## Measurement of the $\Lambda_c^+$ Lifetime

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A precise measurement of the  $\Lambda_c^+$  lifetime using approximately 1340 fully reconstructed  $\Lambda_c^+ \rightarrow$  $pK^{-}\pi^{+}$  and charge conjugate decays is presented. The data were accumulated by the Fermi-lab high energy photoproduction experiment E687. The lifetime of the  $\Lambda_{c}^{+}$  is measured to be  $0.215{\pm}0.016{\pm}0.008$  ps.

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This paper reports a new measurement of the  $\Lambda_c^+$  lifetime using approximately 1340 fully reconstructed  $\Lambda_c^+ \rightarrow$  $pK^{-}\pi^{+}$  decays. Throughout this paper, the charge conjugate state is implied when a decay mode of a specific charge is stated. Previous measurements have been limited to samples of  $\approx$  100 or fewer  $\Lambda_c^+$  decays. Because of the relatively large sample of  $\Lambda_c^+$  decays, extensive consistency checks of the results and a detailed systematic study can be made. Comparisons of accurate measurements of the  $\Lambda_c^+$  lifetime with those of other charm baryons and mesons provide information on the relative sizes of the different decay contributions: spectator decay, W exchange, and the interference of identical light quarks in the final state [1].

The data for this analysis were collected in 1990 and 1991 in the Fermilab wideband photoproduction experiment E687. The E687 detector is described in detail elsewhere [2].

The  $\Lambda_c^+$  decays were reconstructed by a *candidate* driven method [2]. The efficiency of this algorithm in finding vertices is essentially independent of the primary and secondary vertex separation, and so this method should not create a bias in the lifetime measurement.

Information from the Čerenkov counters is used to select protons, kaons, and pions. The confidence level at which the three microstrip tracks of the  $pK^{-}\pi^{+}$  combination form a vertex will be labeled CLD. The confidence level at which any of the three  $pK^-\pi^+$  tracks extrapolate back to the primary vertex is labeled CL1. The confidence level at which other microstrip tracks not already assigned to either the primary or secondary vertices point back to the secondary vertex is labeled CL2. The number of background  $pK^{-}\pi^{+}$  combinations can be greatly reduced by cuts on CL1 and CL2. The drawback is that these cuts can induce a small proper time dependence on the efficiency of vertex reconstruction. As a measure of the significance of detachment of the primary and secondary vertices we use the variable  $\ell/\sigma_{\ell}$ .  $\ell$  is the signed three-dimensional separation between the primary and secondary vertices, and  $\sigma_{\ell}$  is the error on  $\ell$  computed on an event-by-event basis.

Figure 1 shows  $pK^-\pi^+$  invariant mass plots for various cuts on CLD, CL1, CL2, and  $\ell/\sigma_\ell$ . It can be seen that the signal-to-noise ratio can be improved by cuts on CLD, CL1, and CL2. The  $\Lambda_c^+$  yields presented with the figures are determined from fits to the mass distributions with a Gaussian peak over a linear background. Note that this fit is not used in the extraction of the lifetime.

The method used to measure the  $\Lambda_c^+$  lifetime is very similar to the one we used to measure the  $D^0$  and  $D^+$ lifetimes [4]. We fit with the *reduced proper time*. The reduced proper time is given by  $t' = (\ell - N\sigma_\ell)/\beta\gamma_c$ , where N represents the significance of detachment cut  $(\ell/\sigma_\ell > N)$ , and  $\beta\gamma$  is the laboratory frame Lorentz boost of the  $\Lambda_c^+$ . To the extent that  $\sigma_\ell$  is independent of  $\ell$  (as data and Monte Carlo studies show), the t' distribution for  $\Lambda_c^+$  decays will be of the form  $\exp(-t'/\tau)$ , where  $\tau$  is the lifetime of the  $\Lambda_c^+$ .

A fit is made to the t' distribution for events within  $\pm 2\sigma$  of the  $\Lambda_c^+$  mass (approximately  $\pm 20 \text{ MeV}/c^2$ ), using a binned maximum likelihood method. Two reduced proper time histograms are made: one for events within  $\pm 2\sigma$  of the  $\Lambda_c^+$  mass (the signal histogram), and one for events within two sidebands, each  $4\sigma$  wide, above and below the  $\Lambda_c^+$  mass (the sideband histogram). The observed number of events within a reduced proper time bin (centered at  $t'_i$ ) for the signal histogram is labeled  $s_i$ , and for the sideband histogram is labeled  $s_i$ . The expected number of events  $(n_i)$  in the signal histogram for reduced proper time bin i is given by

$$n_{i} = S \frac{f(t_{i}')e^{-t_{i}'/\tau}}{\sum f(t_{i}')e^{-t_{i}'/\tau}} + B \frac{b_{i}}{\sum b_{i}},$$
(1)



FIG. 1.  $\Lambda_c^+$  signals for various cuts on CLD, CL1, CL2, and  $\ell/\sigma_\ell$  used in the lifetime analysis. The fit masses are within 2.5 MeV/ $c^2$  of the current world averages [3] and the widths of the signals are consistent with our experimental resolution.

where  $S = \sum s_i - B$  is the total number of signal events in the signal histogram, B is the total number of background events in the signal histogram, and  $f(t'_i)$  is a correction function. Fifty time bins are used to span five nominal lifetimes. The fit parameters are  $\tau$  and B.

The function f(t'), derived from Monte Carlo simulation, corrects the signal lifetime distribution for effects of acceptance, analysis cut efficiencies, hadronic absorption, and decay of the charm secondaries. Figure 2 shows plots of f(t') for the two  $\Lambda_c^+$  samples shown in Fig. 1. Comparison of the measured lifetimes for these samples will show the reliability of the f(t') correction function used. As a consistency check of the background lifetime evolution  $(b_i)$  used in the fit, different sidebands are used, at  $4\sigma$ ,  $6\sigma$ , and  $8\sigma$  from the  $\Lambda_c^+$  mass.

In order to tie the value of B to the background level expected from the mass sidebands, we include an additional factor in the likelihood function. The final likelihood function is given by

$$L = \prod \frac{n_i^{s_i} e^{-n_i}}{s_i!} \frac{(2B)^{N_b} e^{-2B}}{N_b!},$$
(2)

where  $N_b = \sum b_i$ .

Simulations where care was taken in modeling the background lifetime evolution revealed both the presence of a small bias in the lifetime fitting procedure and a slight underestimation of the true statistical error due to the neglected fluctuations in  $b_i$ . The fitting bias is found to be  $\approx 0.005$  ps due to the presence of a non-negligible long-lived background component, and the statistical error is found to be underestimated by  $\approx 15\%$ . The fitting bias is smaller for samples with better signal-to-noise ratio, but is always small compared to the statistical error. The lifetime results are corrected for this fitting bias and the statistical error bars from the fits are corrected to include contributions from fluctuations in  $b_i$ .

For a consistency check of the measured lifetime, two variations of this method are also used for comparison. In the first, we use an *event-by-event* method where the likelihood function is calculated from the product of the probabilities for each event [5]. In the second variation of the maximum likelihood fit method, the *absolute proper time* is used instead of the reduced proper time. The



FIG. 2. f(t') correction function as described in the text. The fits shown are to a linear form.



FIG. 3. (a),(b) Results of lifetime fits as a function of the  $\ell/\sigma_{\ell}$  cut for different sets of cuts on CLD, CL1, and CL2. (c) Results of lifetime fits using the absolute proper time. (d) Results of lifetime fits using the event likelihood fit method.

absolute proper time is given by  $t = \ell/\beta \gamma c$ .

Figures 3(a) and 3(b) show the measured  $\Lambda_c^+$  lifetimes plotted versus the  $\ell/\sigma_\ell$  cut for the two different sets of cuts on CLD, CL1, and CL2 shown in Fig. 1. The measured lifetimes using the *absolute* proper time, and using the *event likelihood* method, are shown in Figs. 3(c) and 3(d), respectively. It can be seen that all the results are very consistent with each other and no significant variation with respect to the  $\ell/\sigma_\ell$  cut is seen.

Figures 4(a) and 4(b) show the background subtracted and f(t') corrected reduced proper time distributions for the two  $\Lambda_c^+$  samples shown in Fig. 1. The superimposed curve is a pure exponential function using the  $\Lambda_c^+$  lifetime found by the fit. Figures 4(c) and 4(d) show, for the two data samples shown in Fig. 1, the raw lifetime distribution for events under the  $\Lambda_c^+$  peak (points) and the raw lifetime distribution for events in the sidebands (line). Note that events under the  $\Lambda_c^+$  peak include both signal and background, and the background (sideband) lifetime distribution has both a short-lived component ( $\tau \approx$ 0.06 ps) as well as a long-lived component ( $\tau \approx 0.6$  ps). In choosing the  $\ell/\sigma_{\ell}$ , CLD, CL1, and CL2 cuts at which to quote the lifetime, we are guided by the principle of maintaining a reasonable signal to noise while keeping the statistical error to a minimum. We will quote the lifetime at the cuts used for Fig. 1(a). The lifetime is 0.215±0.016 ps.

Further consistency checks are made on the measured lifetime. The fitted lifetime could show a possible dependence on variations of the background proper time distribution. Different sidebands are used in the fit to check this. Background from the decays  $D^+$  and  $D_s^+ \to K^+ K^- \pi^+$ , where the  $K^+$  is misidentified as a proton, could cause problems in the lifetime measurement. As a check on this, the lifetime was measured for two  $\Lambda_c^+$  samples where this background contamination is significantly reduced. In one sample, the proton had to be identified by the Čerenkov counters as definitely a proton (i.e., not allowing protons identified as



FIG. 4. (a),(b) The background subtracted and Monte Carlo corrected lifetime evolutions measured for the two  $\Lambda_c^+$  samples shown in Fig. 1. (c),(d) The raw lifetime distributions for events under the  $\Lambda_c^+$  peak (points) and for events in the sidebands (line).

K/p ambiguous), and in the other sample  $D^+$  and  $D_s^+$  decays are eliminated by eliminating events where the  $K^+K^-\pi^+$  mass is within  $\pm 24$  MeV of the  $D^+$  mass or within  $\pm 30$  MeV of the  $D_s^+$  mass. The contamination due to  $D^+ \rightarrow K^-\pi^+\pi^+$  decays, where one of the pions is misidentified as a proton, is negligible. The fitted lifetime may also be sensitive to the assumed momentum distribution of the  $\Lambda_c^+$  particles. To check this the  $\Lambda_c^+$  sample is split into a low (p < 100 GeV/c) and a high ( $p \geq 100 \text{ GeV}/c$ ) momentum sample, and the lifetime is measured separately for these two samples. Additionally, the total  $\Lambda_c^+$  sample is split into particle and antiparticle



FIG. 5. A comparison of measured lifetimes for different sidebands used in the binned maximum likelihood fit, and for different  $\Lambda_c^+$  samples described in the text.

samples, and into data taken in 1990 and 1991 as further consistency checks of the measured lifetime. Figure 5 shows a comparison of the measured lifetimes for all these different samples for a single  $\ell/\sigma_{\ell}$  cut of  $\ell/\sigma_{\ell} > 4$ . Again, it is evident that the results are all very consistent with each other and that any systematic errors in the measured lifetime are small compared to the statistical error.

Although the systematic error on the measured lifetime is seen to be small compared to the statistical error, an upper limit on the systematic error can be estimated.

Because of uncertainties in the target absorption corrections, we include a systematic error of 0.002 ps for the  $\Lambda_c^+$  lifetime. Two effects are present: hadronic absorption of the decay secondaries, if not corrected for, will tend to give a larger fitted  $\Lambda_c^+$  lifetime, and absorption of the  $\Lambda_c^+$  by the target material prior to decay will tend to give a lower  $\Lambda_c^+$  lifetime. Uncertainties in the secondary absorption correction arise because we are uncertain of the extent to which elastic scattering of the secondaries causes severe mismeasurement of the parent  $\Lambda_c^+$ . The  $\Lambda_c^+$  absorption cross section is unknown. By varying the  $\Lambda_c^+$  absorption cross section between 0 and 2 times the proton-nuclear cross section, the variation of the fitted  $\Lambda_c^+$  lifetime is only 0.003 ps. We take the  $\Lambda_c^+$  nuclear cross section to be the proton nuclear cross section for the absorption correction.

We ascribe a systematic error of 0.005 ps for the uncertainty in the f(t') correction function used. This is obtained by looking at variations of the fitted lifetime when the f(t') correction function is varied within its statistical error as given by the Monte Carlo simulation, and by looking at the variations when different cuts on CLD, CL1, CL2, and  $\ell/\sigma_{\ell}$  are used.

A systematic error of 0.005 ps is included for uncertainties in the background lifetime distribution under the signal peak. This is determined by looking at variations in the fitted lifetimes using different background sidebands, and by looking at differences in the fitted lifetimes for the two  $\Lambda_c^+$  samples, mentioned previously, where backgrounds from  $D^+$  and  $D_s^+$  decays have been significantly reduced.

The acceptance of the charm secondaries depends partly on the  $\Lambda_c^+$  momentum. Higher momentum charm particles will, on the average, decay nearer the microstrip system than low momentum ones, and thus have a larger acceptance. Also, charm particles produced near the edge of the experimental target have a slightly smaller acceptance than those produced at the center of the target because their decay secondaries have a larger probability of laying outside the transverse fiducial volume of the first microstrip station. These two effects depend on the assumed momentum spectrum and the beam profile, and due to uncertainties in these we add an additional systematic error of 0.002 ps.

For uncertainties in the correction for fit biases in the fitting method used, we include a systematic error of 0.003 ps. This includes uncertainties due to the choice of the number of time bins used and the maximum lifetime cut used. This is obtained by looking at variations in the lifetime using the other two fitting methods, variations in the fitted lifetime when different numbers of time bins are used, and variations in the fitted lifetime when different maximum reduced proper time cuts are used.

Combining all sources of systematic errors incoherently, we obtain a final  $\Lambda_c^+$  lifetime measurement of  $0.215\pm0.016(\text{stat})\pm0.008(\text{syst})$  ps.

In summary, we report on a new measurement of the  $\Lambda_c^+$  lifetime based on a sample of  $\approx 1340$  fully reconstructed  $\Lambda_c^+ \to pK^-\pi^+$  decays. The measured lifetime is  $0.215\pm0.016(\text{stat})\pm0.008(\text{syst})$  ps. The data satisfy many consistency checks, and extensive systematic studies were made.

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