

Comment on “Wave Field Determination Using Three-Dimensional Intensity Information”

Recently, Nugent has shown [1] that, within the framework of the paraxial approximation, the cross-spectral density function (CSDF) of the light field may be recovered via a direct transformation of the three-dimensional intensity distribution. In this Comment I will show that, actually, the transformation suggested in Ref. [1] does not yield the detailed CSDF but rather an integral of the CSDF along a certain line.

The transformation is given in Eq. (14) of Ref. [1], which reads

$$W(\mathbf{r}, \mathbf{x}, 0) = \frac{x_1 x_2}{\lambda_0^2} \int \int I(\mathbf{r}', a) e^{-2\pi i \mathbf{x} \cdot (\mathbf{r} - \mathbf{r}') / \lambda_0} d\mathbf{r}' da. \quad (1)$$

In this equation λ_0 is the wavelength, \mathbf{r} , \mathbf{r}' , and \mathbf{x} are vectors in the plane $z = \text{const}$, x_1, x_2 are the two components

of \mathbf{x} , and $a = 1/z$.

$$W(\mathbf{r}, \mathbf{x}, z) = \langle E(\mathbf{r} + \mathbf{x}/2, z, \lambda_0) E^*(\mathbf{r} - \mathbf{x}/2, z, \lambda_0) \rangle$$

is the cross-spectral density function, namely, the correlation between the component of the field with wavelength λ_0 at the points $(\mathbf{r} + \mathbf{x}/2, z)$ and $(\mathbf{r} - \mathbf{x}/2, z)$. $I(\mathbf{r}, z) \equiv W(\mathbf{r}, 0, z)$ is the intensity.

Note that according to Eq. (1), by changing the origin of the z coordinate, the CSDF at every z may be recovered. In other words Eq. (1) implies that the CSDF $\langle E(\mathbf{r} + \mathbf{x}/2, z, \lambda_0) E^*(\mathbf{r} - \mathbf{x}/2, z, \lambda_0) \rangle$ at every point (\mathbf{r}, z) in the volume and for every difference vector \mathbf{x} may be recovered from the intensity $I(\mathbf{r}, z) = \langle E(\mathbf{r}, z, \lambda) E^*(\mathbf{r}, z, \lambda) \rangle$. This is, of course, a very strong statement. In the rest of this Comment I will show that actually only part of the information about the CSDF may be recovered.

My derivation goes as follows: The 3D distribution of the CSDF is related to the CSDF at the $z = 0$ plane through the formula [2]

$$W(\mathbf{r}', \mathbf{x}, z) = \int K(\mathbf{r}' + \mathbf{x}/2, \boldsymbol{\rho} + \boldsymbol{\xi}/2, z) K^*(\mathbf{r}' - \mathbf{x}/2, \boldsymbol{\rho} - \boldsymbol{\xi}/2, z) W(\boldsymbol{\rho}, \boldsymbol{\xi}, 0) d\boldsymbol{\rho} d\boldsymbol{\xi}, \quad (2)$$

where K is the free-space propagator of the paraxial wave equation:

$$K(\mathbf{u}, \mathbf{v}, z) = (1/\lambda_0 z) e^{2\pi i (\mathbf{u} - \mathbf{v})^2 / 2\lambda_0 z}. \quad (3)$$

In the case $\mathbf{x} = 0$ Eq. (2) is a formula for the 3D distribution of the intensity in terms of the CSDF at the $z = 0$ plane. In order to get the CSDF at $z = 0$ in terms of the 3D intensity distribution one has to invert this formula. A partial inversion is obtained by multiplying this formula by $(1/z^2) \exp[-2\pi i \mathbf{x} \cdot (\mathbf{r} - \mathbf{r}') / \lambda_0 z]$, and then integrating over \mathbf{r}' and z . The result is

$$\frac{x_1 x_2}{\lambda_0^2} \int W(\mathbf{r}', \mathbf{x}, 0) \delta\left(\frac{\mathbf{x} \cdot (\mathbf{r}' - \mathbf{r})}{\lambda_0}\right) d\mathbf{r}' = \frac{x_1 x_2}{\lambda_0^2} \int \int I(\mathbf{r}', a) e^{-2\pi i \mathbf{x} \cdot (\mathbf{r} - \mathbf{r}') / \lambda_0} d\mathbf{r}' da. \quad (4)$$

The right-hand side of this equation is the same as in the transformation derived in Ref. [1] [Eq. (1) above]. The left-hand side says that this transformation when applied to the intensity does not yield the detailed CSDF but rather an integral of this function along a line (parametrized by l) $\mathbf{r}'(l)$ on which $\mathbf{r} - \mathbf{r}'$ is perpendicular to \mathbf{x} .

The difference between the present result [Eq. (4) above] and the result of Ref. [1] [Eq. (1) above] is traced to the improper treatment of the delta function in the paragraph between Eqs. (13) and (14) in Ref. [1], which leads to the incorrect conclusion that $W_m = W$. A careful treatment shows that actually Eq. (13) in Ref. [1] leads to the conclusion that W_m equals a line integral of W .

In summary, the central statement of Ref. [1] should be replaced by a weaker statement saying that *some*

characteristics of the cross-spectral density function of the light field may be recovered via a direct transformation of the three-dimensional intensity distribution.

G. Hazak
 Physics Department
 Nuclear Research Center
 P.O. Box 9001, Beer Sheva, Israel

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- [2] See, for example, Eq. (4) in A. Gamliel and G. P. Agrawal, J. Opt. Soc. Am. A **7**, 2184 (1990).