Comment on "Electronic Structure of Ideal Metal/GaAs Contacts"

In an interesting paper [1], van Schilfgaarde and Newman have analyzed theoretically the electronic structure of ideal metal/GaAs contacts by means of a selfconsistent local-density approximation (LDA), and have concluded that "models which invoke intrinsic interface states to explain Fermi-level pinning are not consistent with experimental observations. This strongly suggests the nonideal nature of the experimentally observed interfaces."

Contrary to these remarks, I intend to show that ideal interfaces and the induced-density-of-interface-states (IDIS) model (or the intrinsic-interface-states model in Ref. [1]) can explain the observed behavior of metalsemiconductor interfaces.

The conclusions of Ref. [1] are based on the theoretical results obtained for the barrier heights of different metal/GaAs interfaces that show a lack of correlation with the experimental evidence and, apparently, with the charge-neutrality level of the IDIS model [2]. Leaving apart the problems associated with the error introduced [3] in the LDA, which makes the conclusions of Ref. [1] at least questionable, I would like to concentrate on the actual characteristics of the IDIS model. First of all, I should mention that it has already been known for some years that the IDIS model would yield different interface change-neutrality levels depending on the interface conditions [4,5]. Accordingly, intrinsic and extrinsic chargeneutrality levels have been introduced [6]. The intrinsic level (or the midgap energy [2]) only appears in very idealized conditions: (i) where the metal has a broad and featureless electronic density of states [7]; and (ii) for (110) III-V semiconductor surfaces, when the adsorbed metal atom is very large, and its valence electrons interact similarly with the cationlike and the anionlike surface states [7,8].

Even Al on $Si(111)$ can give $[4]$ different chargeneutrality levels depending on the position of the last Al layer on the Si surface. Theoretical results yield values of ϕ_{Bp} ranging between 0.4 and 0.9 eV, depending on the surface geometry. These values fluctuate around 0.67 eV, the level that one can define as the Si intrinsic chargeneutrality level [6].

K on $GaAs(110)$ is a very ideal case [8] due to the K large size. The charge-neutrality level and the Fermi energy are almost independent of the metal adsorption site, yielding a Schottky barrier ϕ_{Bp} of \sim 0.7 eV.

For a thick Ag-GaAs(110) interface it has been found [5] that the Schottky barrier ϕ_{Bp} varies between 0.54 and 0.93 eV, depending on the metal adsorption sites on the semiconductor. The intrinsic charge-neutrality level was also found [5] to be around 0.7 eV.

The important point that one can deduce from these results $[2-4]$ is that the extrinsic charge-neutrality level fluctuates around the intrinsic one, depending on the metal-atom properties (as shown very elegantly in Ref. [I]) and its adsorption site at the semiconductor interface.

The results of Ref. [I] seem to confirm this conclusion, since the calculated values of ϕ_{Bp} fluctuate around 0.7 eV, the GaAs intrinsic charge-neutrality level [6]. On the other hand, those results have been calculated for a fixed interface geometry: One can expect, however, for the metal adatoms different adsorption sites from the ones assumed in Ref. [I], changing accordingly the interface Fermi level.

In conclusion, the actual theoretical evidence [4-6] does not support the claim of Ref. [1] mentioned at the beginning of this Comment and suggests that more theoretical calculations minimizing the total interface energy are necessary before a definitive conclusion about the properties of ideal metal/GaAs contacts can be reached. Let us remark, however, that LCAO calculations along these lines have been performed for the alkali-metal-GaAs(110) interfaces [8,9]; the results of these calculations support unambiguously the IDIS model. Moreover, LDA calculations $[3]$ for the NiSi₂-Si junction for which the interface geometry is well established also support the same model. All these results suggest the ideal nature of at least some of the experimentally observed interfaces.

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- [I] M. van Schilfgaarde and N. Newman, Phys. Rev. Lett. 65, 2728 (1990).
- [2] C. Tejedor, F. Flores, and E. Louis, J. Phys. C 10, 2163 (1977); J. Tersoff, Phys. Rev. Lett. 52, 465 (1984).
- [3] G. P. Das, P. Blöck, O. K. Andersen, N. E. Christensen, and O. Gunnarsson, Phys. Rev. Lett. 63, 1168 (1989).
- [4] H. I. Zhang and M. Schluter, Phys. Rev. B 18, 1923 (1978); G. Platero, J. A. Vergés, and F. Flores, Surf. Sci. 168, 100 (1986).
- [5] J. Sánchez-Dehesa, J. Ortega, F. Flores, and J. Dow, Surf. Sci. 200, 424 (1988); J. Ortega, J. Sánchez-Dehesa, and F. Flores, in Proceedings of the Nineteenth International Conference on the Physics of Semiconductors, Warsaw, l988, edited by W. Zawadski (Polish Academy of Sciences, Warsaw, 1989), p. 641.
- [6] F. Flores et al., Appl. Surf. Sci. 41/42, 144 (1989).
- 7] M. Lannoo et al., in Metallization and Metal Semiconductor Interfaces, edited by I. P. Batra, NATO Advanced Study Institutes, Ser. B, Vol. 195 (Plenum, New York, 1989), p. 259; F. Flores and C. Tejedor, J. Phys. C 20, 145 (1987).
- [8]J. Ortega and F. Flores, Phys. Rev. Lett. 63, 2500 (1989).

[9] J. Ortega et al., Surf. Sci. (to be published).

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