

### Comment on "Observation of the Photorefractive Effect in a Polymer"

In a recent, very interesting Letter [1] Ducharme *et al.* make what they believe to be the first report of the photorefractive (PR) effect in a polymer. In this Comment we draw attention to electro-optic (EO) conjugated polymers, in which PR effects have also been observed, as reported for *trans*-polyacetylene [ $t$ -(CH) $_x$ ] in 1983 [2], although not described as such at that time.

Above-gap photoexcitation of  $t$ -(CH) $_x$  results, at the long times ( $\mu$ s to ms) probed in mechanically chopped cw laser experiments, in a photoinduced absorption (PA) spectrum characteristic of charged solitons [2,3]; peaks appear at "midgap" and in the infrared, together with an associated bleaching of the  $\pi$ - $\pi^*$  absorption band. In addition, an oscillatory structure (superposed on the bleaching), straddles the absorption edge. The latter reproduces the spectrum obtained in direct electroabsorption (EA) measurements, and is assigned [2] to EO modulation of the absorption due to local fields arising from soliton generation.

Similar photoinduced EA is found in other conjugated polymers, e.g., poly(2,5-thienylene vinylene) [4] (PTV), and is now appropriately recognized [5] as a "photorefractive effect" (N.B. while manifest in PA as a change in the imaginary part,  $\Delta k$ , of the refractive index, there is by necessity an accompanying change,  $\Delta n$ , in the real part of the refractive index). The relative ease with which this PR effect is detected relates to the large third-order EO response. The corresponding complex susceptibility,  $\chi^{(3)}(-\omega; 0, 0, \omega)$ , evaluated from the EA data for PTV has a peak  $\text{Re}\{\chi^{(3)}\}$  value of  $7.5 \times 10^{-9}$  esu and a peak  $\text{Im}\{\chi^{(3)}\}$  value of  $6.2 \times 10^{-9}$  esu. The close correspondence between photoinduced [4] and directly measured [5] EA spectra is demonstrated for PTV in Fig. 1. Comparison of the magnitudes of the two EA signals at their maxima at 1.78 eV [ $(-\Delta T/T)_{\text{max}} \approx 1 \times 10^{-4}$  for photoexcitation with  $200 \text{ mW/cm}^2$  of 488-nm (2.54-eV) light at 80 K and  $(-\Delta T/T)_{\text{max}} \approx 5.2 \times 10^{-4}$  for 100 kV/cm applied electric field at 130 K] allows an estimate of the average space-charge field that results from photoexcitation: We find a value of  $\approx 50 \text{ kV/cm}$ .

It should be emphasized that the results discussed here differ in two important ways from those reported by Ducharme *et al.* [1].

(i) The PR effect seen in  $t$ -(CH) $_x$  and PTV is found in a homopolymer with charge generation and trapping occurring without external bias fields and without the addition of carrier transport agents. Charge generation is assisted by photoexcitation of electron-hole pairs separated between chains, and intrinsic trapping may occur through self-localization arising from electron-phonon coupling [4]. Extrinsic factors including defects and disorder must, however, also play a role.

(ii) The EO effect in conjugated polymers is due to a

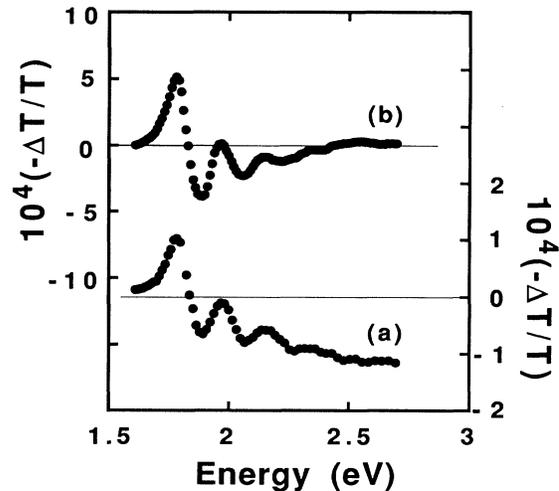


FIG. 1. (a) Photoinduced [4] (right scale, 80 K, 488-nm excitation) and (b) directly measured [5] (left scale, 130 K, 100 kV/cm) EA of PTV. The increasingly negative offset at higher energies for (a) is due to bleaching of the  $\pi$ - $\pi^*$  absorption by charge-carrier generation [4].

*third-order* nonlinearity, while that in the materials considered by Ducharme *et al.* [1] has a *second-order* origin. As a consequence, the PR response in conjugated polymers is potentially weaker, but this clearly depends on how large a space-charge field can be achieved. Note also that for second-order materials the polymer must at some stage be poled in order to achieve the EO response, while for conjugated polymers molecular alignment is not required.

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