

Comment on "Does Antimatter Fall with the Same Acceleration as Ordinary Matter?"

In a series of Letters [1,2] the Eöt-Wash Collaboration has reported on their equivalence-principle experiment. In a summary article [3] they stated that their experiment limited an anomalous acceleration of antimatter to $10^{-5}g$ from quantum-gravity models. Reference [2] expanded on Ref. [3], now claiming a limit of $2 \times 10^{-6}g$.

However, what they have actually shown is that, since their [2] vector coupling is proportional to a conserved charge (e.g., baryon number), a precise cancellation by a scalar requires its coupling to be effectively proportional to this charge to a high accuracy; cancellations more precise than Ref. [2] envisions can occur naturally in many theories (see below and Ref. [4]). Further, they give no detailed long-range topographic or geologic error analysis [1-3]. Independently, it is important to measure gravity on antimatter [4,5] just as it is important to test the principle of equivalence [1].

In Ref. [2] there is much discussion about the experimental conclusion that Witteborn and Fairbank (WF) observed the Schiff-Barnhill "electron-sag" effect in their electron gravity experiment. However, for a quarter century this result has remained an object of experimental and theoretical controversy, involving the Dessler-Michel-Rorschach-Trammell-Herring "ion-sag" and patch effects [4,6]. (Eventually, Schiff came to accept ion-sag dominance, concluding that WF "cannot be understood" given the "usual picture of the metallic surface" [7]). In any event, surfaces have been developed which should allow gravity to be measured on ions [8]. This would culminate in the original experimental proposal [4] to measure $g(\bar{p})/g(H^-)$.

The analysis of Ref. [2] considers a possible anomalous acceleration of antimatter due to "quantum gravity." However, of the three cases discussed, only the first two are based upon a Lagrangian scalar interaction density. The other case is a phenomenological parametrization. Further, the analysis, especially in case (1), depends heavily on details of only one example in the quantum-gravity framework [4]. One is simply *not* restricted to the forms in Ref. [2]. For example, all three cases in Ref. [2] assume $\mathcal{L}_V = (a)^{1/2} V_\mu J^\mu$, with J^μ a dimension-three operator. Such a dimension constraint is found in renormalizable gauge-field theories, but does not generally apply to gravitational theories [4].

The scalar coupling of case (1) of Ref. [2] also assumes $\mathcal{L}_S = (b)^{1/2} \phi T^\mu_\mu$. However, there exist other couplings of the graviscalar field to either the gauge fields and/or other sources that reduce the differences between the vector and scalar couplings above. For example, $\mathcal{L}_{S2} = (b)^{1/2} \phi U$, where $U = \sum_j (b_j^1)/4 [F_j^{\mu\nu} F_{\mu\nu}^j]$, has a $j=0$ QED component which is closely related (by a virial) to a constant fraction of the binding energy, whatever its Z dependence. When added to \mathcal{L}_S , the resulting effective coupling of the scalar is more closely proportional to the

vector coupling. (A similar argument can be made for the $j=1, \dots, 8$ QCD couplings [4].)

Case (2) is superfluous as it is equivalent to changing the vector coupling in case (1) from B to $(1-\epsilon)B$ (for [2] $\lambda_V = \lambda_S$).

Case (3) generates an *ad hoc* phenomenological dependence of the scalar coupling on $|B|$ and B^2 (and perhaps μ —see Fig. 3 of Ref. [2]). However, again the vector coupling is assumed proportional to B/μ , so cancellation simply requires the scalar to vary similarly. A more revealing assumption would have been a case "orthogonal" to case (1), that is, to assume $q_S = \epsilon\mu + |B|$. However, then the Eöt-Wash constraint is satisfied for reasonable values of ϵ .

Thus, while we agree that the precise cancellation described in Ref. [2] is ruled out for the very specific cases discussed there, it is neither excluded nor unnatural in the more general case. A precise cancellation is "natural" only if it is due to some (approximate) symmetry; however, that symmetry may not be immediately apparent. Matter experiments can only lead us to conclude that *either* there are no new effects *or* they cancel as if enforced by a symmetry. An antimatter experiment can decide between the two. Therefore, although strong constraints can be set on some theories by present experiments [1-3], they alone cannot rule out all anomalous gravitational forces.

Finally, and most importantly, independent of any particular theoretical motivation [4], be it quantum gravity, *CPT* violation, or something else, the measurement of gravity on the antiproton would be unique among all fundamental gravity and particle-physics experiments. Fortunately, a consensus now exists to this effect [5]. Even if one obtains the Newtonian value one expects, one will have learned something new and important.

T. Goldman, Michael Martin Nieto, M. H. Holzscheiter,
T. M. Darling, Martin Schauer, and Jay Schecker
Los Alamos National Laboratory
Los Alamos, New Mexico 87545

Received 19 March 1991

PACS numbers: 04.80.+z, 04.60.+n, 04.90.+e

- [1] C. W. Stubbs *et al.*, Phys. Rev. Lett. **58**, 1070 (1987); **61**, 2409 (1988); **62**, 609 (1989); E. G. Adelberger *et al.*, *ibid.* **59**, 849 (1987); **59**, 1790(E) (1987); B. R. Heckel *et al.*, *ibid.* **63**, 2705 (1989).
- [2] E. G. Adelberger *et al.*, Phys. Rev. Lett. **66**, 850 (1991).
- [3] E. G. Adelberger *et al.*, Phys. Rev. D **42**, 3267 (1990).
- [4] M. M. Nieto and T. Goldman, Phys. Rep. **205**, 221 (1991), and references therein.
- [5] K. Pretzel, in CERN Report No. PSCC/90-28, p. 125 (unpublished).
- [6] L. I. Schiff, Phys. Rev. B **1**, 4649 (1970).
- [7] *Near Zero: Now Frontiers of Physics*, edited by J. D. Fairbank *et al.* (Freeman, New York, 1988).
- [8] J. B. Camp, T. W. Darling, and R. E. Brown, J. Appl. Phys. **69**, 7126 (1991).