

Femtosecond Spectroscopy of Molecular Autoionization and Fragmentation

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Femtosecond laser pulses are applied to the study of the dynamics and the pathways of multiphoton-induced ionization, autoionization, and fragmentation of Na_2 in molecular-beam experiments. In particular, we report on first results obtained studying electronic autoionization (leading to $\text{Na}_2^+ + e^-$) and autoionization-induced fragmentation (leading to $\text{Na}^+ + \text{Na} + e^-$) of a bound doubly excited molecular state. The final continuum states are analyzed by photoelectron spectroscopy and by measuring the mass and the released kinetic energy of the corresponding ionic fragments with a time-of-flight arrangement.

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In contrast to the well studied two-electron excited states of various atomic systems, the electronic autoionization of doubly excited molecular states is still a field where very little is known. Doubly excited states are crucial to the understanding of the dynamics of diatomic molecules at energies where electronic autoionization, ionic fragmentation, and neutral fragmentation compete. The study of the coupling of the different continua of fragmentation and ionization is also of great theoretical interest, because of the need for an extension of multichannel quantum-defect theory (MQDT) to treat doubly excited electronic states, dissociating states, and autoionizing Rydberg states together.

Interesting studies of the dissociation and autoionization of unbound doubly excited states in H_2 have recently been reported by several groups.¹ The first spectroscopic study of a bound doubly excited state, the $^1\Sigma_g^-(3d+4p)$ state of K_2 , has also only recently been reported.² However, there has yet been no report on the interaction of a particular bound doubly excited state with different continua and the competition between the various decay channels of such a state. Electronic autoionization and in particular autoionization-induced molecular fragmentation have been neither experimentally nor theoretically explored as much as the vibrational and rotational autoionization of Rydberg states. Beside experimental difficulties, the main reason for this lack of information and understanding is that often neither the final continuum states are analyzed nor the dissociative ionization clearly distinguished from neutral fragmentation with subsequent excited-fragment photoionization when using nanosecond or even picosecond laser pulses. Of course, both processes may lead to formation of the same ionic fragment. Femtosecond laser pulses offer an experimental tool fast enough to distinguish between these two processes. If the laser pulse durations applied are considerably shorter ($\tau \approx 100$ fs) than the fragment separation time, no further laser-induced excitation or ionization of fragments can occur. For Na_2 , the time needed for the fragments to separate to 10 Å varies from 0.5 to 1 ps depending on the recoil

energy W . All observations therefore have to be related to excitation processes taking place at small internuclear distances.

Pioneering work in the field of femtosecond photochemistry and transition-state spectroscopy has been performed by Zewail and co-workers.³ In a series of beautiful experiments in the gas phase and in molecular beams, they have demonstrated the enormous advantage of applying femtosecond lasers to the study of molecular dynamics.³

The multiphoton ionization and fragmentation of alkali-metal dimers and, in particular, of Na_2 has been the subject of considerable current interest. In many experiments it has been found that in conjunction with the formation of Na_2^+ ions, the atomic ions Na^+ are also produced. It is well established⁴ that resonance-enhanced multiphoton-ionization (REMP) processes via the $A^1\Sigma_u^+$ or the $B^1\Pi_u$ states are responsible for this observation, whereas nonresonant multiphoton processes only play a minor role. Ionization and fragmentation of Na_2 Rydberg states involving vibrational-rotational autoionization have also been studied in recent years.⁵ However, very few experiments have reported a clear identification of the fragmentation pathways. Two examples of such detailed studies are (i) the fragmentation of the neutral molecule Na_2^* into $\text{Na}^*(3p) + \text{Na}^*(3p)$, i.e., the reverse process to laser-induced associative ionization, applying Doppler spectroscopy to measure the angular and energy distributions of the neutral fragments,⁶ and (ii) the two-photon ionization and dissociation of Na_2 resonantly enhanced via the $(2)^1\Sigma_u^+$ double-minimum state with kinetic-energy analysis of the ionic fragments.⁷

To study the dynamics of multiphoton-induced excitation, autoionization, and fragmentation of Na_2 , we have used a combination of different experimental techniques. Femtosecond laser pulses are used to induce the transitions and time-of-flight (TOF) spectroscopy to determine both the mass and released kinetic energies of the fragments and the energies and angular distributions of ejected electrons.

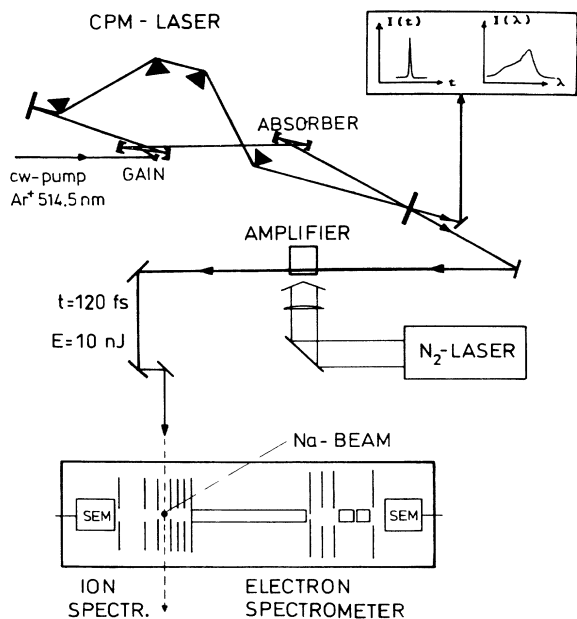


FIG. 1. Femtosecond-laser-molecular-beam arrangement. The molecular beam, the CPM laser, and the TOF axis are all perpendicular to each other.

Femtosecond pulses are generated in a home-built colliding-pulse mode-locked (CPM) ring dye laser.⁸ The emission peak of the CPM dye laser has been shifted to 616 nm by adjusting the diethyloxadicarbocyanine iodide (DODCI) absorber concentration. The spectral width (FWHM) was about 55 Å. The output of the CPM dye laser was amplified at a rate of 100 Hz in a N₂-laser-pumped dye amplifier to produce pulses of 10 nJ energy and $t = 120$ fs time duration measured by auto-correlation. The laser-pulse energy was intentionally kept at this low level. The laser system and the schematic experimental arrangement of the ion and electron time-of-flight spectrometers are shown in Fig. 1. The laser-molecular-beam interaction region is placed between parallel plates and since the ions are extracted perpendicular to the beam with a low electric field (10 V/cm), parent ions can be distinguished from fragment ions having initial kinetic energy. This experimental technique is widely known.⁹ The ion TOF spectrum displayed in Fig. 2(a) is obtained when our femtosecond laser ($I \sim 10^8$ W/cm²) interacts with the sodium molecular beam. The spectrum clearly shows the observation of Na₃⁺, Na₂⁺, and "slow" Na⁺ as well as "fast" Na⁺ ionic fragments resulting from the femtosecond laser excitation at $\lambda = 616$ nm.

The small signal of the metal cluster Na₃⁺ is due to an absorption of two photons of 616 nm, which are sufficient to ionize the Na₃ cluster.¹⁰ The intense Na₂⁺ signal is due to a REMPI process as is discussed below in detail. Predissociation of Na₂^{*} and photoionization of Na^{*} as the origin of the observed Na⁺ ions can be ruled out considering the time duration of our fs laser pulse.

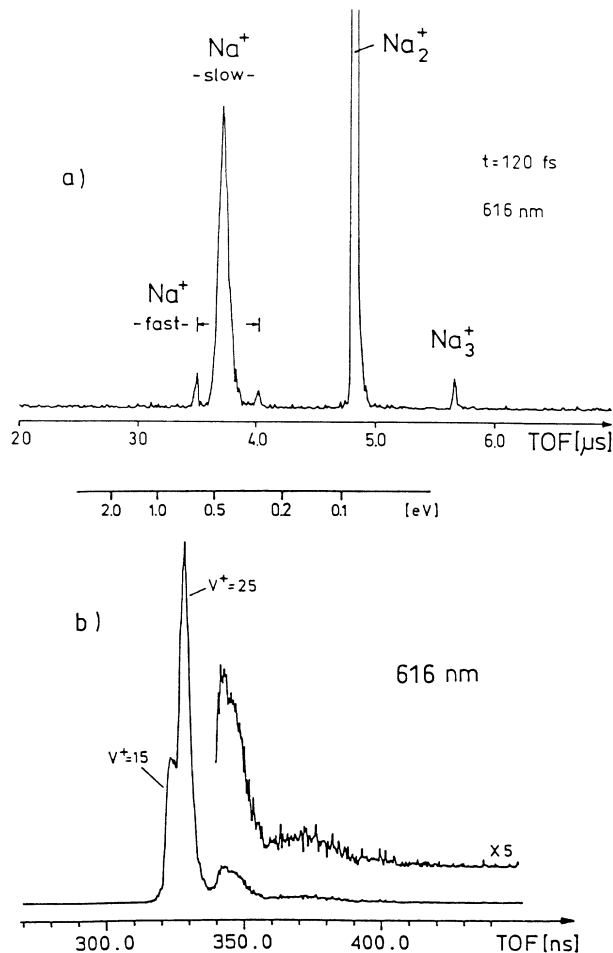


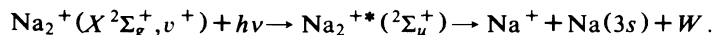
FIG. 2. (a) Time-of-flight spectrum of ions formed by the interaction of femtosecond laser pulses with a sodium molecular beam and (b) time-of-flight spectrum of ejected electrons.

Therefore the fast and slow Na⁺ ions originate from fragmentation processes occurring at small internuclear distances of Na₂. Based on this result, which is rather difficult to obtain from other experiments, and with the known molecular-potential curves, we completely determined for this model case the multiphoton excitation and fragmentation pathways.

The one-color REMPI spectrum between 610 and 620 nm obtained with a 2.5-ns pulsed tunable dye laser shows an isolated and strong peak at 616.08 nm in both the Na₂⁺ and the Na⁺ channels. Based on the known spectroscopy of molecular Rydberg states of Na₂,¹¹ it is evident that the peak in the Na₂⁺ spectrum is due to resonance-enhanced three-photon ionization of Na₂ (Fig. 3). The ionization-enhancing intermediate electronic state ¹Π_g(3s+3d) is populated by the two-photon process Na₂(X¹Σ_g⁺, v''=0) + 2hν → Na₂(Ryd¹Π_g, v* = 20). The Rydberg molecule is then photoionized by absorption of a third photon leading to the observed strong Na₂⁺ signal. This direct ionization leads to observed electrons having kinetic energies of $E = 810 \pm 10$ meV

and $E = 940 \pm 10$ meV. From this we conclude that the dimer ions are preferentially formed in the $v^+ = 24, 25$ and $v^+ = 14, 15$ vibrational states of the electronic ground state $X^2\Sigma_g^+$, which is in accord with the calculated Franck-Condon factors. Exactly the same excitation and ionization process happens when we use our femtosecond laser with its central wavelength of 616 nm. Considering the predominant production of the Na_2 mol-

ecules in the $v'' = 0$ level in a supersonic molecular-beam expansion, the two-photon transition probability, and the upper $^1\Pi_g$ -Rydberg-state vibrational spacing $\Delta G \approx 88$ cm^{-1} , we estimate that 90% of the excited Na_2 molecules are in the $v^* = 20$ level, despite the broad spectral distribution (≈ 140 cm^{-1} FWHM) of the fs laser pulse. The created dimer ions may now undergo a bound-free transition by absorption of one more photon from still the same fs laser pulse:



Taking into account the populated v^+ levels and the known potential curves for the ionic ground and first-excited states,¹² the corresponding bound-free transitions lead to recoil energies W between 9910 and 11037 cm^{-1} . The energy obtained from the analysis of the ion TOF spectrum, namely, $W = 10500 \pm 500$ cm^{-1} , surely agrees with this prediction. We therefore conclude on the basis of the observed ion and electron kinetic energies (Fig. 2) and the restriction imposed by using 120-fs laser pulses that fast Na^+ ions are produced by the ionization and fragmentation process shown in Fig. 3.

The observation of slow Na^+ ions, however, cannot be explained within this framework since photoionization of the vibrational level $v^* = 20$ of the $^1\Pi_g$ Rydberg state directly into the $^2\Sigma_u^+$ continuum is energetically forbidden for the applied laser wavelength. The electron spectrum in Fig. 2(b) shows in addition to the two strong peaks less intense broad structures in the range between 300 and 500 meV and around 160 meV extending to

very low energies. Based on both the measured electron energy distribution and the distribution of recoil energies $W = 900 \pm 400$ cm^{-1} , obtained from the analysis of the slow Na^+ ionic fragments, a consistent interpretation is found by considering the excitation and the decay of doubly excited molecular states of Na_2 . Figure 4 shows theoretical¹² potential curves of the ground state $^2\Sigma_g^+$ and the repulsive state $^2\Sigma_u^+$ of the Na_2^+ ion and a schematic potential curve of the doubly excited $^1\Pi_u(4s+3p)$ state of Na_2^{**} (dashed curve). There is no other asymptotic limit of the bound doubly excited state in that energy range. The doubly excited states $\text{Na}_2^{**}(nl+3p)$ form a Rydberg series converging to the $^2\Pi_u$ state of Na_2^+ , whose potential curve is known

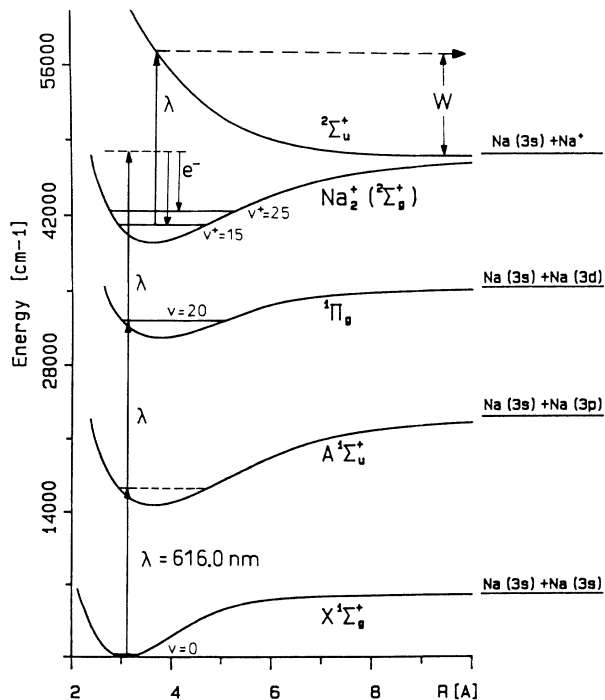


FIG. 3. Potential-energy diagram of Na_2 illustrating the process leading to fast Na^+ ionic fragments.

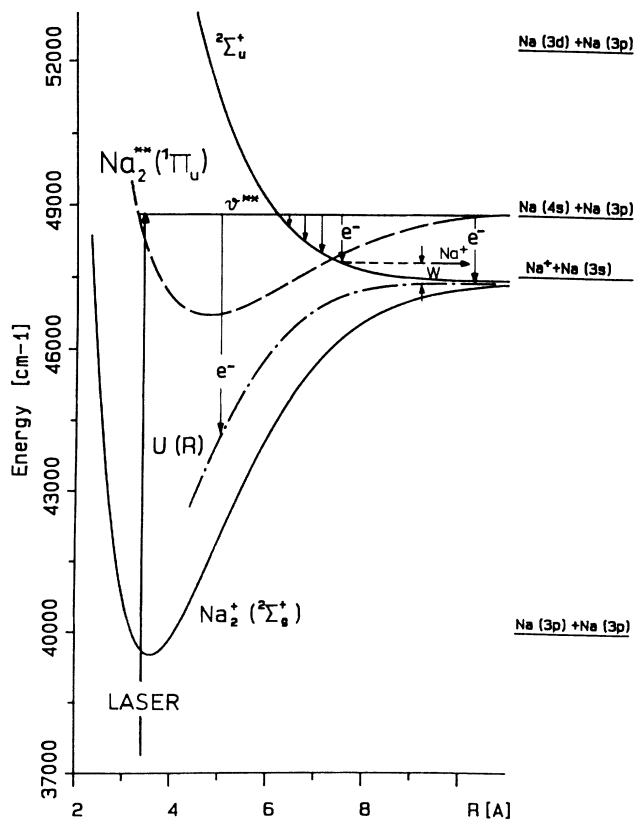
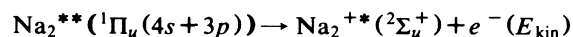


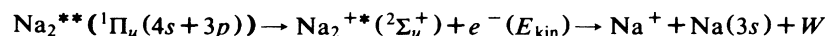
FIG. 4. Excitation and autoionization processes of the doubly excited $^1\Pi_u(4s+3p)$ state of Na_2 .

theoretically.¹² Assuming that the shape of a ${}^1\Pi_u$ -state potential curve, correlated to $\text{Na}(4s)+\text{Na}(3p)$, is similar to the ionic ${}^2\Pi_u$ curve and that it has a small potential barrier at large internuclear distances like the $B\ {}^1\Pi_u$ state built from $\text{Na}(3s)+\text{Na}(3p)$, we conclude that in the resonance-enhanced three-photon process vibronic levels close to the dissociation limit of the ${}^1\Pi_u(4s+3p)$ state are excited. The wave functions of these vibronic levels extend from 3 Å to ~ 10 Å. These doubly excited levels may now electronically autoionize into the $X\ {}^2\Sigma_g^+$ ground state of Na_2^+ . Taking into account Mulliken's difference potential $U(R)$,¹³ we expect for this particu-

lar process ejected electrons with energies between 260 and 500 meV. These electrons are definitely observed. For internuclear distances greater than 6 Å, the vibronic levels of this doubly excited state interact additionally with the continuum of the repulsive ${}^2\Sigma_u^+$ state of Na_2^+ . Therefore for $R \gtrsim 6$ Å a second channel is open for electronic autoionization. The autoionization process



is responsible for the observed electrons having energies in the range from 0 to 160 meV. The autoionization-induced fragmentation



leads to slow Na^+ ions whose kinetic energies depend on the internuclear distance R where the autoionization takes place. A calculation of the autoionization electron spectrum arising from bound-free transitions ${}^1\Pi_u(4s+3p) \rightarrow {}^2\Sigma_u^+$ shows essentially a broad distribution in agreement with the experimental electron energy spectrum. The observed range of released kinetic energies W of the slow Na^+ fragment ions is also in accord. This example of autoionization and autoionization-induced fragmentation shows that a unique assignment can only be given when the final continuum states are determined by measurements of the kinetic energy of the electrons and of the ionic fragments.

In conclusion, this is the first study where, in a molecular-beam experiment, a femtosecond laser is used in combination with ion and electron spectroscopy to clarify the ionization, autoionization, and fragmentation mechanisms of molecules excited by multiphoton processes into the energy region of doubly excited electronic states. The electronic autoionization of the ${}^1\Pi_u(4s+3p)$ bound state of Na_2 into the two ion continua (${}^2\Sigma_g^+$ and ${}^2\Sigma_u^+$ of Na_2^+) and the autoionization-induced fragmentation are processes which have not been observed before and which undoubtedly enhance our knowledge of doubly excited molecular states.

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