

Carrier-Carrier Scattering in High- T_c Superconductors

Lee and Read¹ first proposed that the linear-in- T resistivity ρ of the high- T_c superconductors is due to carrier-carrier scattering. This scattering should also lead to a frequency dependence of the conductivity,² and the recent observations of Thomas *et al.*³ are very suggestive in this respect.

The frequency dependence of the electron-electron scattering rate Γ is closely related to its temperature dependence. The carriers are spread over a range $k_B T$ about the Fermi energy E_F , leading to $\Gamma \sim T^2$, and likewise $\Gamma \sim \omega^2$. Gurzhi and Kaganov² find $\Gamma \sim \omega_n^2$, where $\omega_n^2 = \omega^2 + (\pi k_B T/h)^2$, and that $\Gamma \rightarrow 0$ if $\omega \gg E_F B$ (with B the bandwidth). This cannot, however, be the whole story. Since the system is causal, a frequency-dependent Γ must lead to a frequency-dependent mass renormalization m^* , which is given by a Kramers-Kronig relation. This m^* cannot simply be a function of ω_n .

Figure 1 shows that Γ found by Thomas *et al.*³ scales as ω_1 (not as ω_2), except near the high-temperature saturation, for T between 70–270 K. The frequency-dependent m^* is also observed.³ Away from saturation, Γ appears to be quadratic in ω_1 . However, the data could also be fitted by a linear dependence, with a small quadratic correction. This latter form is consistent with the dc resistivity measured on these samples.³

This scaling has important consequences. First, this is very different from the frequency-dependent conductivity observed in heavy-fermion superconductors.⁴ Moreover, in extracting the gap function from the low- T optical properties, it is essential to compare the superconducting reflectivity to the normal-state reflectivity at the same temperature.

The frequency dependence of m^* is readily understood in a Fermi-liquid picture. The strong linear-in- T electron-electron scattering is evidence that the hole Fermi surface is near the Van Hove singularity (VHS) in these materials.¹ To explain the low carrier density observed in Hall measurements, these carriers must be separated into two groups: a high density of heavy holes near the Van Hove singularity, with a few light holes associated with the rest of the Fermi surface.⁵ The strong scattering of the heavy holes ensures that they become localized at low temperatures, contributing negligibly to the Hall effect. The T and ω dependence of m^* can be directly related to this localization since the effective carrier density is proportional to $\omega_p^2 m/m^*$. The role of ad-

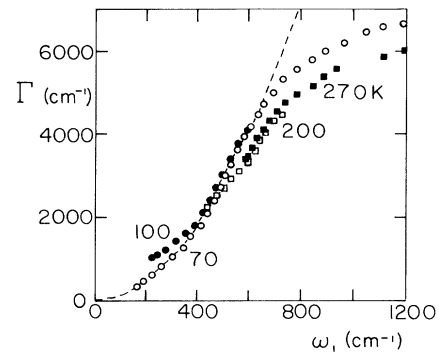


FIG. 1. Replot of data of Ref. 3, upper part of Fig. 4: Γ vs ω_1 . Dashed line: $\Gamma \propto \omega_1^2$.

ditional free carriers associated with intact Cu-O chains⁶ will probably complicate interpretation of $\sigma(\omega, T)$ in stoichiometric $\text{Ba}_2\text{YCu}_3\text{O}_7$.

It should be stressed that the VHS can coincide with the Fermi level only at a fixed carrier density. Hence, if the VHS is important for high- T_c superconductivity, there must be an electronic or structural instability which pins the Fermi level at the Brillouin zone boundary, similar to the Hume-Rothery phases in alloys. This appears to be the case.⁵

R. S. Markiewicz
Physics Department
Northeastern University
Boston, Massachusetts 02115

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