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<sup>4</sup>R. K. Nesbet, Quarterly Progress Report, Solid State and Molecular Theory Group, Massachusetts Institute of Technology, October 15, 1955 (unpublished), pp. 4-8.

<sup>5</sup>R. K. Nesbet, Revs. Modern Phys. **32**, 272 (1960).

<sup>6</sup>See, for example, S. B. Welles, Phys. Rev. **62**, 197 (1942).

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## INELASTIC SCATTERING OF 18.9-Mev NUCLEONS FROM THE 9.6-Mev STATE OF C<sup>12</sup>

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Recently it has been suggested<sup>1</sup> that the 9.6-Mev state of C<sup>12</sup> has spin and parity  $J^\pi=3^-$  and not  $1^-$  as previously supposed.<sup>2</sup> This Letter reports a new analysis of Peelle's<sup>3</sup> 18.9-Mev inelastic proton scattering data, which definitely favors  $3^-$  for the 9.6-Mev state. The calculation uses a direct volume interaction with spin-dependent distorted waves.<sup>4</sup>

Peelle's<sup>3</sup> analysis of his data is based upon the direct surface interaction theory, which predicts a  $j_K^2(kR)$  angular distribution. Using an interaction radius  $R=3.3$  fermis, he found a best fit for  $K=1$  giving  $J^\pi=0^-, 1^-,$  or  $2^-$ . The same theory applied to inelastic scattering from the 4.4-Mev  $2^+$  state of C<sup>12</sup> gave very poor agreement with experiment.

Levinson and Banerjee<sup>5</sup> showed that the direct-surface-interaction theory is probably inadequate for a nucleus as small as C, and Robson and Robson<sup>4</sup> found spin-orbit effects to be important for 12-Mev nucleons inelastically scattered from the 4.4-Mev level of C<sup>12</sup>. Thus it was considered essential to take both these effects into account before using Peelle's data to determine the spin and parity of the 9.6-Mev state. In the following, only  $J^\pi=1^-$  or  $3^-$  are considered because there exists a reasonable amount of evidence<sup>1</sup> for eliminating other values.

The distorting potential is taken to be the usual Woods-Saxon optical potential plus a Thomas-Fermi spin-orbit potential,<sup>6</sup>

$$V_D(r) = -(V + iW)f(r) + (V_S + iW_S) \left( \frac{\hbar}{m c} \right)^2 \frac{1}{r} \frac{df}{dr} \vec{\sigma} \cdot \vec{L},$$

$$f(r) = \{1 + \exp[(r - R)/a]\}^{-1}. \quad (1)$$

The values of  $V$ ,  $W$ ,  $V_S$ , and  $W_S$  used are 45, 12,

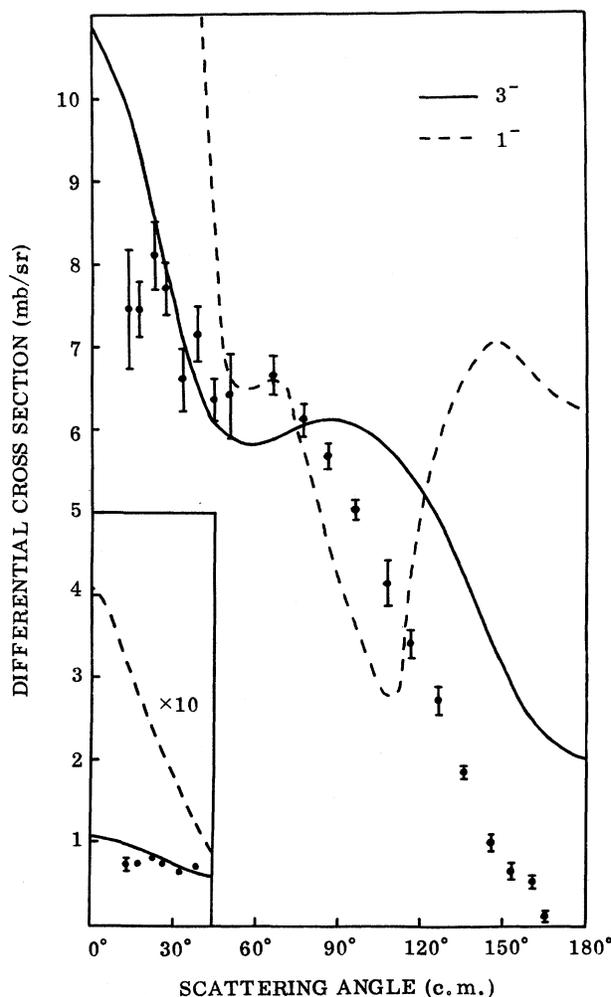


FIG. 1. Inelastic nucleon scattering from the 9.6-Mev state of C<sup>12</sup>. The points are the experimental results of Peelle for 18.9-Mev protons. The theoretical curves are for 18.9-Mev neutrons. The normalization is arbitrary and both curves are fitted at 75°.

15, and -4 Mev for the incident energy of 18.9 Mev and 45, 8, 17, and -4 Mev for the emergent energy of 9.3 Mev, with  $a = 0.4$  fermi and  $R = 2.75$  fermis in both cases.

The direct-interaction two-body potential is assumed to be zero ranged and spin independent, and all forms of exchange are neglected.

Following Barker *et al.*,<sup>1</sup> the  $C^{12}$  ground state is taken as

$$\psi_0 = \psi(1s^4 1p^8 [4, 4] 000, 0), \quad (2)$$

where the numbers following the configuration are values of  $[\lambda] TSL, J$ . The 9.6-Mev state is taken as

$$\psi_3 = \psi((1s^4 1p^7 [4, 3] \frac{1}{2} \frac{1}{2} 1, 1d) 003, 3) \quad (3)$$

for  $J^\pi = 3^-$ , and

$$\psi_1 = \sum_{i=1}^3 \alpha_i \psi_{1i} \quad (4)$$

for  $J^\pi = 1^-$ , where

$$\psi_{11} = \psi((1s^4 1p^7 [4, 3] \frac{1}{2} \frac{1}{2} 1, 2s) 001, 1), \quad (4a)$$

$$\psi_{12} = \psi((1s^4 1p^7 [4, 3] \frac{1}{2} \frac{1}{2} 1, 1d) 001, 1), \quad (4b)$$

$$\psi_{13} = \psi((1s^3 [3] \frac{1}{2} \frac{1}{2} 0, 1p^9 [441] \frac{1}{2} \frac{1}{2} 1) 001, 1). \quad (4c)$$

For the Elliott and Flowers interaction, Barker *et al.*,<sup>1</sup> found  $\alpha_1 = 0.916$ ,  $\alpha_2 = -0.399$ , and  $\alpha_3 = -0.028$ . For simplicity, the values used in the present calculation are  $\alpha_1 = 0.869$ ,  $\alpha_2 = -0.494$ , and  $\alpha_3 = 0$ . The  $1p$ ,  $1d$ , and  $2s$  wave functions are taken for an ideal harmonic oscillator with length parameter  $b = 2.0$  fermis, obtained by fitting the  $1p$

harmonic-oscillator wave function to the  $1p$  square-well solution of the Schrödinger equation for a radius  $R = 3.66$  fermis, which gives a reasonable fit to the 4.4-Mev data.<sup>4</sup>

Figure 1 shows the result of the calculation. It is seen that the theoretical curve for the  $3^-$  assumption is a much better fit than for the  $1^-$  assumption (both curves are normalized at  $75^\circ$ ). Greater emphasis should perhaps be placed on fitting the scattering for angles  $\theta < 90^\circ$ , since other processes like heavy-particle stripping<sup>7</sup> may become important for  $\theta > 90^\circ$ . On account of the very lengthy calculations involved, no attempt was made to vary the parameters, which were selected by a consideration of the elastic scattering and inelastic scattering from the 4.4-Mev state data.<sup>4,8</sup>

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<sup>2</sup>F. Ajzenberg-Selove and T. Lauritsen, *Nuclear Phys.* **11**, 1 (1959).

<sup>3</sup>R. W. Peelle, *Phys. Rev.* **105**, 1311 (1957).

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<sup>5</sup>C. A. Levinson and M. K. Banerjee, *Ann. Phys.* **3**, 67 (1958).

<sup>6</sup>The Riesenfeld-Watson notation is used here. W. B. Riesenfeld and K. M. Watson, *Phys. Rev.* **102**, 1157 (1956).

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<sup>8</sup>J. S. Nodvik and D. S. Saxon, *Phys. Rev.* **117**, 1539 (1960).