

### Comment on "Transverse Electromagnetic Waves with $\mathbf{E}\parallel\mathbf{B}$ "

Chu and Ohkawa<sup>1</sup> (CO) have proposed that a class of transverse electromagnetic (TEM) waves with  $\mathbf{E}\parallel\mathbf{B}$  exists. This paper has provoked critical reaction<sup>2-6</sup> and a rebuttal.<sup>7</sup> Most of this discussion appears to be due to the failure of CO<sup>1</sup> to define their terminology carefully and explain their assumptions. Consequently, the respondents assumed instinctively that all TEM waves propagate and attempted to prove that  $\mathbf{E}\parallel\mathbf{B}$  waves could not propagate and hence, could not exist. The example of CO<sup>1</sup> implied that their proposed class consisted of only standing waves and Chu<sup>7</sup> was explicit in his rebuttal. This Comment derives the general conditions under which TEM standing waves with  $\mathbf{E}\parallel\mathbf{B}$  exist and remedies these deficiencies.

It is useful to classify TEM wave solutions of Maxwell's equations according to whether their Poynting vector  $\mathbf{S}(\mathbf{r},t)=0$  or  $\neq 0$ . The former are  $\mathbf{E}\parallel\mathbf{B}$  TEM standing waves and the latter  $\mathbf{E}\perp\mathbf{B}$  TEM traveling or standing waves.<sup>8</sup> CO<sup>1</sup> made implicit use of the Coulomb gauge in their derivation. The Coulomb gauge introduces the constraints  $\nabla\cdot\mathbf{A}=0$  and  $\Phi(\mathbf{r},t)=C$ . This gauge, which is consistent with the Lorentz gauge required to obtain independent wave equations for  $\mathbf{A}$  and  $\Phi$ , is used in the following analysis. However, the fields calculated are independent of the gauge.

The most general TEM solution of the vector wave equation for  $\mathbf{A}$  obtained from Maxwell's equations is<sup>9</sup>  $\mathbf{A}(\mathbf{r},t)=\mathbf{A}_+(\eta)+\mathbf{A}_-(\zeta)$ , where  $\eta\equiv\mathbf{k}\cdot\mathbf{r}-\omega t$  and  $\zeta\equiv\mathbf{k}\cdot\mathbf{r}+\omega t$ . Then  $\mathbf{B}(\mathbf{r},t)=\nabla\times\mathbf{A}=\mathbf{k}\times(\mathbf{A}'_++\mathbf{A}'_-)$ , where  $\mathbf{A}'_+\equiv[d\mathbf{A}_+(\eta)/d\eta]$  and  $\mathbf{A}'_-\equiv[d\mathbf{A}_-(\zeta)/d\zeta]$ , and so  $\mathbf{B}$  is transverse since  $\mathbf{k}\perp\mathbf{B}$ .  $\nabla\cdot\mathbf{A}=\mathbf{k}\cdot(\mathbf{A}'_++\mathbf{A}'_-)=0$ , so that  $\mathbf{B}\neq 0$  if  $\mathbf{A}'_+\neq-\mathbf{A}'_-$ . If  $\nabla\Phi=0$ , then  $\mathbf{E}(\mathbf{r},t)=-\partial\mathbf{A}/\partial t=\omega(\mathbf{A}'_+-\mathbf{A}'_-)$ , and so  $\mathbf{E}\parallel(\mathbf{A}'_+-\mathbf{A}'_-)$ .  $\mathbf{E}$  is transverse if  $\mathbf{k}\cdot(\mathbf{A}'_+-\mathbf{A}'_-)=0$ . It can be shown that the Poynting vector  $\mathbf{S}(\mathbf{r},t)\propto\mathbf{E}\times\mathbf{B}=\mathbf{k}[(\mathbf{A}'_++\mathbf{A}'_-)\cdot(\mathbf{A}'_+-\mathbf{A}'_-)]$  for TEM waves.  $\mathbf{S}(\mathbf{r},t)=0$  if  $(\mathbf{A}'_++\mathbf{A}'_-)\cdot(\mathbf{A}'_+-\mathbf{A}'_-)=0$ , which means that  $(\mathbf{A}'_++\mathbf{A}'_-)\perp(\mathbf{A}'_+-\mathbf{A}'_-)$ . It follows that  $|\mathbf{A}'_+|=|\mathbf{A}'_-|$ . If  $\mathbf{A}'_+\neq\mathbf{A}'_-$ , then  $\mathbf{E}\neq 0$ ,  $\mathbf{B}\neq 0$ , and  $\mathbf{E}\cdot\mathbf{B}\neq 0$ . Moreover,  $\mathbf{k}\cdot(\mathbf{A}'_++\mathbf{A}'_-)=0$  and  $\mathbf{k}\cdot(\mathbf{A}'_+-\mathbf{A}'_-)=0$  for these TEM waves. Consequently, TEM standing waves exist with  $\mathbf{E}\parallel\mathbf{B}$  and  $\mathbf{S}(\mathbf{r},t)=0$ . The example of CO<sup>1</sup> satisfies these conditions.

A similar analysis for TEM traveling or standing waves with  $\mathbf{E}\perp\mathbf{B}$  and  $\mathbf{S}(\mathbf{r},t)\neq 0$  yields the following results:  $\mathbf{E}\cdot\mathbf{B}=\omega\mathbf{k}\cdot[(\mathbf{A}'_++\mathbf{A}'_-)\times(\mathbf{A}'_+-\mathbf{A}'_-)]=0$  if either  $(\mathbf{A}'_++\mathbf{A}'_-)\parallel(\mathbf{A}'_+-\mathbf{A}'_-)$  or  $(\mathbf{A}'_++\mathbf{A}'_-)\times(\mathbf{A}'_+-\mathbf{A}'_-)\perp\mathbf{k}$ , and  $\mathbf{S}(\mathbf{r},t)\neq 0$  if  $(\mathbf{A}'_++\mathbf{A}'_-)\not\perp(\mathbf{A}'_+-\mathbf{A}'_-)$ .

The approach taken by CO<sup>1</sup> of defining another vector

potential  $\mathbf{F}_k(\mathbf{r})=\mathbf{A}_k(\mathbf{r})+k^{-1}\nabla\times\mathbf{A}_k(\mathbf{r})$ , which leads to  $\nabla\times\mathbf{F}_k(\mathbf{r})=k\mathbf{F}_k(\mathbf{r})$  if  $\nabla\cdot\mathbf{F}_k(\mathbf{r})=\nabla\cdot\mathbf{A}_k(\mathbf{r})=0$ , is insufficient to define those TEM standing waves with  $\mathbf{E}\parallel\mathbf{B}$  and  $\mathbf{S}(\mathbf{r},t)=0$ , unless  $\mathbf{A}_k(\mathbf{r})$  is constrained to be real since all standing waves must satisfy  $\mathbf{A}(\mathbf{r},t)=\text{Re}[\exp(-i\omega t)]\times\text{Re}[\mathbf{A}_k(\mathbf{r})]$ . Otherwise, it can be shown that the vector potential

$$\begin{aligned}\mathbf{A}(\mathbf{r},t) &= \mathbf{A}_0[a\cos(\mathbf{k}\cdot\mathbf{r}+\omega t)+b\sin(\mathbf{k}\cdot\mathbf{r}+\omega t)] \\ &= \text{Re}[\exp(-i\omega t)\mathbf{A}_k(\mathbf{r})],\end{aligned}$$

where  $\mathbf{A}_k(\mathbf{r})=\mathbf{A}_0\exp[-i(\mathbf{k}\cdot\mathbf{r}-\delta)]$  and  $\delta=\tan^{-1}(b/a)$ , can be used to obtain a derived vector potential  $\mathbf{F}_k(\mathbf{r})$  for which  $\mathbf{E}\perp\mathbf{B}$  and  $\mathbf{S}(\mathbf{r},t)\neq 0$ .

It has been shown that a class of TEM waves with  $\mathbf{E}\parallel\mathbf{B}$  exists that can be derived from a vector potential  $\mathbf{A}(\mathbf{r},t)=\mathbf{A}_+(\eta)+\mathbf{A}_-(\zeta)$ , satisfying  $\mathbf{k}\cdot[d\mathbf{A}_+(\eta)/d\eta]=0$  and  $\mathbf{k}\cdot[d\mathbf{A}_-(\zeta)/d\zeta]=0$ , and a scalar potential  $\Phi=C$ , if  $|d\mathbf{A}_+(\eta)/d\eta|=|d\mathbf{A}_-(\zeta)/d\zeta|$  and  $d\mathbf{A}_+(\eta)/d\eta\parallel d\mathbf{A}_-(\zeta)/d\zeta$ , where  $\eta\equiv\mathbf{k}\cdot\mathbf{r}-\omega t$  and  $\zeta\equiv\mathbf{k}\cdot\mathbf{r}+\omega t$ . These are the most general conditions for TEM waves with  $\mathbf{E}\parallel\mathbf{B}$  to exist. Those  $\mathbf{E}\parallel\mathbf{B}$  solutions obtained by the condition given by CO<sup>1</sup> can be obtained by use of the above formalism. These waves do not propagate since  $\mathbf{S}(\mathbf{r},t)=0$ , and should be described as TEM standing waves with  $\mathbf{E}\parallel\mathbf{B}$  to distinguish them from those classical TEM traveling and standing waves with  $\mathbf{E}\perp\mathbf{B}$  and  $\mathbf{S}(\mathbf{r},t)\neq 0$ .

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<sup>3</sup>K. K. Lee, Phys. Rev. Lett. **50**, 138 (1983).

<sup>4</sup>M. Salingaros, Am. J. Phys. **53**, 361 (1985), and J. Phys. A **19**, 101 (1986).

<sup>5</sup>F. C. Michel, Phys. Rev. Lett. **52**, 1351 (1984).

<sup>6</sup>K. R. Brownstein, J. Phys. A **19**, 159 (1986).

<sup>7</sup>C. Chu, Phys. Rev. Lett. **50**, 139 (1983).

<sup>8</sup> $\mathbf{S}(\mathbf{r},t)_{av}=0$  for all standing waves.  $\mathbf{S}(\mathbf{r},t)=0$  is a special case for  $\mathbf{E}\parallel\mathbf{B}$  standing waves.

<sup>9</sup>W. R. Smythe, *Static and Dynamic Electricity* (McGraw-Hill, New York, 1950), 2nd ed., p. 443.