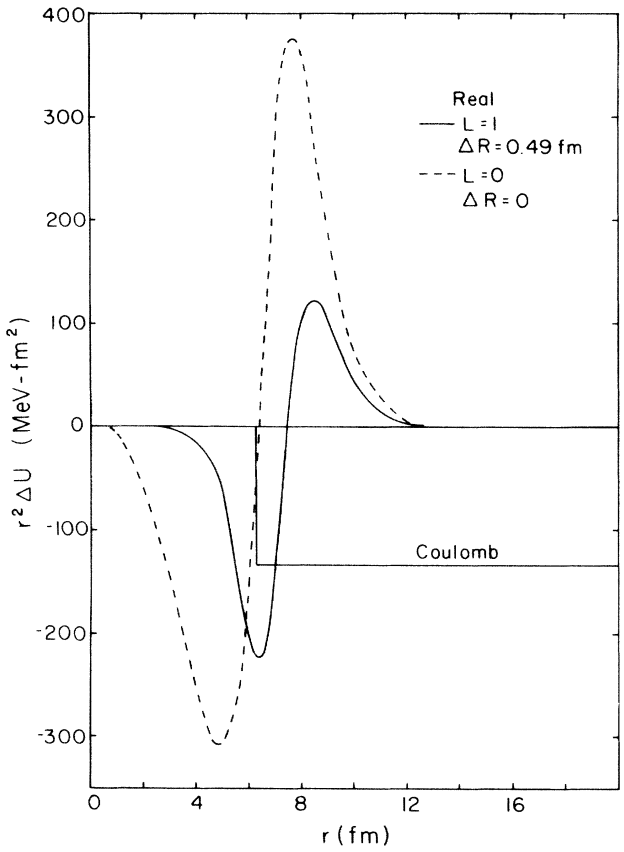

FIG. 2. Comparison of the nuclear real transition potentials, weighted by r^2, for isoscalar monopole and isovector dipole scattering of 129-MeV α particles by ^{116}Sn. Cross sections as shown in Fig. 1(a) are predicted with these. A neutron excess radius is used for the dipole mode, $r_{0n}=r_{0p}+0.1$ fm. The $L=1$ Coulomb potential as used in the DWBA calculations is also shown, weighted also by r^2.

Complete Integrability in a Quantum Description of Chaotic Systems. K. NAKAMURA and M. LAKSHMANAN [Phys. Rev. Lett. **57**, 1661 (1986)].

The equality

$$\langle n|(d|n\rangle/dt)=(d\langle n|/dt)|n\rangle=0,$$

which is used to prove Eqs. (1c) and (1d) on page 1662, is not correct in general. Instead, $\langle n|(d|n\rangle/dt)=i\phi_n(t)$ and its complex conjugate should be employed, where $\phi_n(t)$ is a t-dependent real parameter. The correction seems to yield additional terms on the right-hand side of Eqs. (1c) and (1d). But this problem can be resolved immediately as follows: Let us define $\Phi_n(t)$ which satisfy $d\Phi_n/dt=\phi_n$; if $\{|n\rangle e^{-i\Phi_n}\}$ and their Hermitian conjugates are regarded as $\{|n\rangle\}$ and $\{\langle n|\}$, respectively, together with the corresponding redefinition of $\{V_{nm}\}$, then the original forms of Eqs. (1a)–(1d) and other succeeding statements in the text prove not to be altered at all.

Nucleation and Growth of Colloidal Crystals. D. J. W. AASTUEN, N. A. CLARK, L. K. COTTER, and BRUCE J. ACKERSON [Phys. Rev. Lett. **57**, 1733 (1986)].

All values of the colloidal particle density, n, should be rescaled by a factor of 5.5. That means that the melting density, n_m, which was reported as $n_m\cong 1\ \mu m^{-3}$, should be $n_m\sim 5.5\ \mu m^{-3}$ and the values of n along the abscissa in Fig. 4 should be 5.5 times greater. Because of this rescaling the fitting parameter A in Eq. (2) and reported in the caption of Fig. 4 is actually 0.77.