

Hydrogen-Acceptor Pairs in Silicon

In a recent Letter, Assali and Leite (AL)¹ propose a model for the structure of a hydrogen-boron pair in silicon and the concomitant loss of shallow-acceptor activity. Using the *MS-X α* cluster method, they explore single-particle and total energies as functions of H displacement from the T_d interstitial position towards the substitutional B.

In this Comment, we would like to point out recent theoretical and experimental studies of the H-B pair which were not addressed in the Letter of AL. These results and further calculations that we report here strongly suggest that the AL model for the H-B pair (B_sH_i) is not the stable, minimum-energy configuration.

We have previously reported^{2,3} both *MS-X α* ⁴ and modified neglect of diatomic overlap (MNDO)⁵ computational results for the H-B pair. Using the MNDO-cluster method, we found that, with B on center, H did not remain along the line explored by AL.^{2,3} Furthermore, allowing both the H and B to relax, we found that the interstitial H moved into a position between the B and one of its bonding Si neighbors producing a Si—H bond, while the B moved off center in the opposite direction to attain threefold coordination with the remaining silicons.

This acceptor-deactivating configuration was originally suggested by Pankove and co-workers.⁶ Here, the B and H interact with each other only weakly; we have effectively a SiH- BSi_3 system. This can be contrasted with the AL model, which has an interstitial H interacting primarily with on-center substitutional B (a H- BSi_4 system).

Further support for the Pankove model comes from recent infrared studies^{7,8} and supporting theoretical work.³ In the experimental studies, an 1870- cm^{-1} absorption was correlated with the H-B pair defect. The Pankove (SiH- BSi_3) model was favored since the observed H vibrational frequency was consistent with a Si—H bond. This conclusion was subsequently supported by our MNDO computations, where we calculated a value of 1880 cm^{-1} .³ Any model which would involve H vibrations in a B—H bond was rejected by Pankove since it would be expected to produce an absorption at ~ 2560 cm^{-1} .

We have done follow-up MNDO calculations using a hydrogen-terminated version of the AL cluster ($BHSi_{10}H_{18}$ with frozen silicons). Exploring first the AL configuration, we find a minimum when H is 1.19 Å along the interstitial direction from an on-center B (compare with 1.8 Å from Fig. 2 of AL). For this configuration, we calculate a 2590- cm^{-1} H vibration,⁹ which is consistent with the argument by Pankove and co-workers.

We also calculate the total energies for the Pankove and AL orientations of the B-to-H pair axis $\{[111]$ vs

$[1\bar{1}\bar{1}]$, respectively, in a $(1\bar{1}0)$ plane}. For each orientation, the total energy is minimized with respect to independent H and B motion along the line. We find that the $[111]$ Pankove orientation is lower in energy by 1.3 eV. Therefore, the AL minimum is a saddle point according to our calculations.

As a final, related comment, we note that the acceptor-deactivating configuration suggested by AL has been proposed and well studied for the Li-B pair in silicon.¹⁰ In fact, we find in using the MNDO method that it is the $[11\bar{1}]$ AL configuration of the Li-B pair which is favored over the $[111]$ Pankove orientation by 1.9 eV. This affinity of Li to the interstitial site is consistent with our earlier studies.¹¹

In summary, we suggest that the potential-energy surface explored by AL was too limited. This we conclude on the basis of our MNDO calculations which indicate that (i) the AL configuration is a saddle point, with the Pankove configuration producing the true minimum, and (ii) the computed H vibrational frequency is in agreement with recent infrared studies for the Pankove model, but not for the AL model.

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¹L. V. C. Assali and J. R. Leite, Phys. Rev. Lett. **55**, 980 (1985).

²G. G. DeLeo and W. B. Fowler, J. Electr. Mater. **14a**, 745 (1985).

³G. G. DeLeo and W. B. Fowler, Phys. Rev. B **31**, 6861 (1985).

⁴K. H. Johnson and F. C. Smith, Jr., Phys. Rev. B **5**, 831 (1972).

⁵M. J. S. Dewar and W. Thiel, J. Am. Chem. Soc. **99**, 4899, 4907 (1977); W. Thiel, Quantum Chemistry Program Exchange **11**, 379 (1979).

⁶J. I. Pankove, D. E. Carlson, J. E. Berkeyheiser, and R. O. Wance, Phys. Rev. Lett. **51**, 2224 (1983); J. I. Pankove, R. O. Wance, and J. E. Berkeyheiser, Appl. Phys. Lett. **45**, 1100 (1984).

⁷J. I. Pankove, P. J. Zanzucchi, C. W. Magee, and G. Lucovsky, Appl. Phys. Lett. **46**, 421 (1985).

⁸We calculate a frequency of 2640 cm^{-1} when the motion of both H and B are considered about relaxed H and B positions for the AL $[111]$ orientation.

⁹M. Balkanski, R. J. Elliott, W. Nazarewicz, and P. Pfeuty, in *Lattice Defects in Semiconductors*, edited by R. R. Hasiguti (Univ. of Tokyo Press, Tokyo, 1968), p. 3.

¹¹G. G. DeLeo, W. B. Fowler, and G. D. Watkins, Phys. Rev. B **29**, 1819 (1984).