

## Oscillatory Structures in GaAs/(AlGa)As Tunnel Junctions

Recently, Hickmott *et al.*<sup>1</sup> reported a remarkable oscillatory structure with period  $e\Delta V = \hbar\omega_{LO}$ , the longitudinal-optic phonon energy, in the low-temperature reverse-bias  $J(V)$  characteristics of  $n^+$ GaAs/(AlGa)As/ $n^-$ GaAs/ $n^+$ GaAs devices. The structure was observed only in magnetic fields ( $B$ ) above 4 T which, they claimed, reduced ionized-impurity scattering by inducing freezeout onto shallow donors. This permitted sequential ballistic acceleration of electrons up to the LO phonon emission threshold.

We have studied similar molecular-beam-epitaxy-grown, 200- $\mu$ m-diam mesas with layer thicknesses and doping as follows: substrate and buffer, 200  $\mu$ m,  $2 \times 10^{18}$   $\text{cm}^{-3}$ ;  $n^-$  layer, 1  $\mu$ m,  $1 \times 10^{15}$   $\text{cm}^{-3}$ ; Al<sub>0.35</sub>Ga<sub>0.65</sub>As barrier,  $168 \pm 10$  Å, undoped; top layer, 1  $\mu$ m,  $2 \times 10^{18}$   $\text{cm}^{-3}$ . Figure 1 shows typical  $J(V)$  plots. Mesas of type A have  $J(V)$  of similar overall form to those in Ref. 1, but of much larger amplitude ( $\sim 10^3$  larger at  $V = -0.1$  V). Mesas of type B have an additional low-bias ( $V > -0.2$  V) contribution to  $J$  which may arise from imperfections (e.g., microchannels) in the (AlGa)As barrier. Even at  $B=0$ , derivative plots of both mesa types show oscillatory structure (Fig. 2) of comparable relative amplitude (e.g.,  $\sim 2\%$  at  $-0.2$  V) to that in Ref. 1. The periodicity of the main peaks is  $\hbar\omega_{LO}/e$ . Structure is clearly visible (Fig. 2, curves *c* and *d*) even at temperatures (30 K) at which the majority of donors are ionized. The weaker intermediate extrema, seen only for  $T < 4$  K, probably arise from higher harmonic structure in  $J(V)$ ; at 4 K some peaks in  $d^2J/dV^2$  are only 2 mV wide. Note that our barriers are  $\sim 20\%$  narrower than those in Ref. 1 and that  $J$  is exponentially sensitive<sup>2</sup> to barrier height and width. Calculation shows that  $J(V)$  of Fig. 1 (type A) is fully consistent with the barrier-tunneling characteristics. The

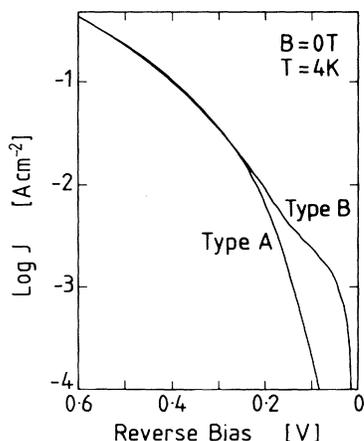


FIG. 1.  $\text{Log } J$  vs  $-V$  for two of our devices.

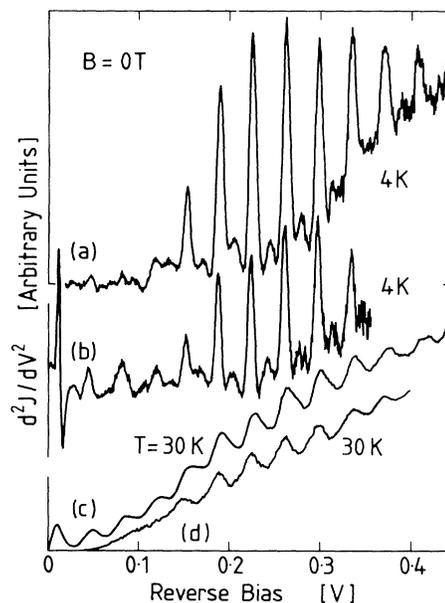


FIG. 2.  $d^2J/dV^2$  vs  $-V$ : Curves *a* and *d*, type-A mesas; curves *b* and *c*, type-B mesas.

larger  $J(V)$  passed by our devices may be crucial for observation of oscillatory structure even at  $B=0$  and up to 30 K and for understanding of its origin.<sup>3,4</sup> The low reverse-bias shoulder on  $J(V)$  in type-B mesas is removed by modest ( $\approx 5$  T) transverse magnetic fields ( $J \perp B$ ), suggesting that it may be due to microchannels in the barrier.<sup>5</sup> However, oscillations for  $V > 0$  are not observed.<sup>5</sup> Oscillatory structure in both mesa types for  $V < -0.2$  V is essentially unaffected by magnetic fields ( $\parallel$  and  $\perp J$ ) up to 12 T.

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<sup>1</sup>T. W. Hickmott *et al.*, Phys. Rev. Lett. **52**, 2053 (1984).

<sup>2</sup>The exponential argument is  $\sim -30$ .

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<sup>4</sup>C. Hanna and R. B. Laughlin, Bull. Am. Phys. Soc. **30**, 631 (1985).

<sup>5</sup>P. F. Lu, D. C. Tsui, and H. M. Cox, Phys. Rev. Lett. **54**, 1563 (1985).