

## Smectic-*A* Phase with Two Collinear Incommensurate Density Modulations

B. R. Ratna, R. Shashidhar, and V. N. Raja  
*Raman Research Institute, Bangalore 560080, India*  
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The first observation of a smectic-*A* phase with two collinear incommensurate density modulations is reported. X-ray diffraction studies of this incommensurate phase reveal reflections corresponding to both partially bilayer ( $A_d$ ) and bilayer ( $A_2$ ) periodicities, the relative intensities of the two reflections being strongly temperature dependent.

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A smectic-*A* liquid crystal may be described as an orientationally ordered fluid with a one-dimensional mass-density wave along the optic axis.<sup>1</sup> When the constituent molecules possess a strongly polar end group several types of smectic-*A* phases exist, which can be characterized unambiguously by x-ray diffraction.<sup>2</sup> The monolayer  $A_1$  phase exhibits a peak at a wave vector  $2q_0 = 2\pi/l$ , where  $l$  is approximately equal to the molecular length; in addition there may be diffuse scattering centered at a wave vector intermediate between  $q_0$  and  $2q_0$ . The bilayer phase ( $A_2$ ) is characterized by two reflections, the fundamental at  $q_0$  and its second harmonic at  $2q_0$ . In the case of the partially bilayer  $A_d$  phase, there is a reflection at  $q'_0 = 2\pi/l'$ , where  $l < l' < 2l$ , and, generally, a diffuse maximum centered at  $2q_0$ . There is also the fluid antiphase  $\bar{A}$  whose characteristic x-ray pattern consists of a spot at  $2q_0$  and two spots displaced from the  $Z$  axis (optic axis) in a perpendicular direction situated symmetrically about the  $q_0$  position.<sup>3</sup> Recent high-resolution x-ray studies<sup>4</sup> have shown the existence of two types of  $\bar{A}$  phases with rectangular lattices of different symmetries.

Although theoretically predicted,<sup>5,6</sup> a smectic-*A* phase with two incommensurate collinear periodicities has never been observed so far. The only case reported to date of a smectic with two coexistent incommensurate density modulations is the three dimensionally ordered smectic-*E* phase of 4-octyl-4'-cyanoterphenyl.<sup>7</sup> We present here the results of our x-ray studies on binary mixtures of 4-octyloxy-4'-cyanobiphenyl (8OCB) and 4-*n*-heptyloxyphenyl-4'-cyanobenzoyloxybenzoate (DB7OCN). The studies have revealed a new type of *A* phase (referred to here as  $A_{ic}$ ) which intervenes between the  $A_d$  and  $A_2$  phases and which has *two collinear incommensurate density modulations*, one of wavelength  $2\pi/q'_0$  and the other of wavelength  $2\pi/q_0$ . The amplitude of the former modulation decreases with decreasing temperature while that of the latter increases, leading finally to a lockin transition to the  $A_2$  phase.

The phase diagram of the binary system (obtained by a combination of optical and x-ray diffraction techniques) in the region of existence of the  $A_{ic}$  phase is shown in Fig. 1. For 8OCB molar concentrations

( $X < 24\%$ ), there is only the  $A_d$ - $A_2$  transition which is clearly seen as an abrupt change in the slope of the curve of layer spacing ( $d$ ) versus temperature (Fig. 2). We have taken x-ray diffraction photographs<sup>8</sup> for several concentrations as a function of temperature in the  $A_d$ ,  $A_{ic}$ , and  $A_2$  phases. The sample had to be cooled extremely slowly (less than 1°C per hour) in order to obtain a monodomain of the  $A_{ic}$  phase. Microdensitometer traces of a series of representative photographs scanned along the  $Z$  axis (parallel to the director) for the  $X = 34.8\%$  mixture are given in Figs. 3(a)–3(f). Starting from the  $A_d$  phase at 119°C [Fig. 3(a)], we see a sharp peak at  $q'_0$ . On cooling, a second sharp peak is seen at  $q_0$ , in addition to that  $q'_0$  [Fig. 3(b)]. This signifies the onset of the incommensurate phase. On further decrease of temperature the intensity of the reflection at  $q'_0$  decreases while that at  $q_0$  increases with an accompanying increase in the intensity of the second harmonic at  $2q_0$ . The switchover of the relative strengths of the  $q'_0$  and  $q_0$  reflections is clearly seen in Figs. 3(c) and 3(d). Finally at 108°C the peak at  $q'_0$  has disappeared altogether [Fig. 3(f)] leaving a

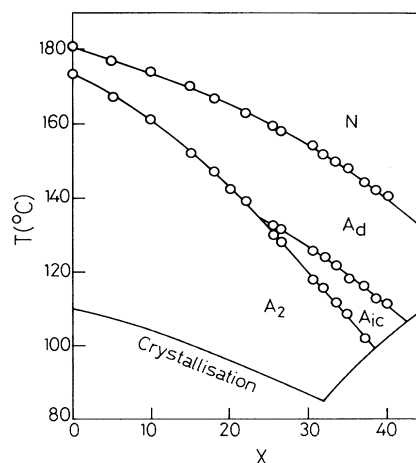


FIG. 1. Partial temperature-concentration ( $T$ - $X$ ) diagram for mixtures of 8OCB and DB7OCN.  $X$  is the mole percent of 8OCB in the mixture. The incommensurate  $A_{ic}$  phase intervenes between the partially bilayer ( $A_d$ ) and bilayer ( $A_2$ ) phases.

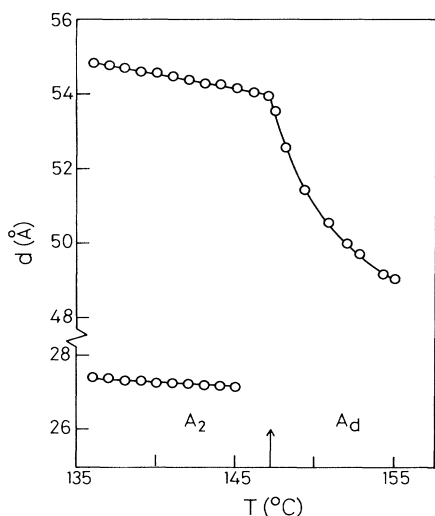


FIG. 2. Temperature variation of the layer spacing ( $d$ ) in the  $A_d$  and  $A_2$  phases of the  $X=18\%$  mixture. The arrow represents the  $A_d$ - $A_2$  transition temperature.

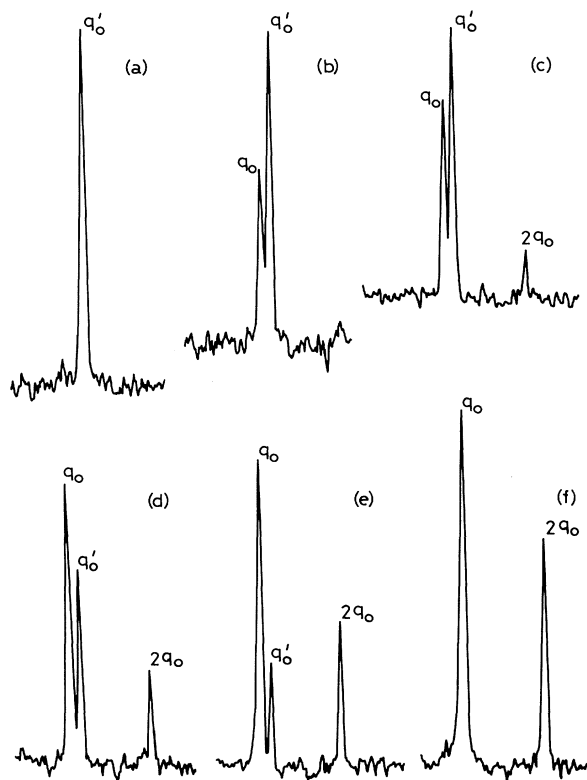


FIG. 3. Raw microdensitometer scans of the x-ray diffraction photographs taken for the  $X=34.8\%$  mixture at different temperatures: (a)  $A_d$  phase,  $119^\circ\text{C}$ ; (b)–(e)  $A_{ic}$  phase,  $117$ ,  $116$ ,  $114.5$ , and  $112^\circ\text{C}$ , respectively; (f)  $A_2$  phase,  $108^\circ\text{C}$ . The wave vector corresponding to each reflection is marked. The direction of the scan is along the  $Z$  axis (optic axis) for all the photographs.

clear signature of the  $A_2$  phase—strong reflections at  $q_0$  and  $2q_0$ . It must be emphasized that regardless of their amplitudes, the sharpness of these reflections remains the same throughout at all temperatures. No reflections corresponding to combinations of  $q_0$  and  $q'_0$  were recorded even with long exposures. [The setup did not allow very-low-angle reflections ( $\theta < 0.5^\circ$ ) to be recorded.] We have also verified from the high-angle diffraction maximum that the in-plane order is liquidlike in all the three phases.

Figure 4 gives the intensity contour diagram (obtained with an  $X$ - $Y$  microdensitometer—Joyce-Loebl Scandig 3—in conjunction with an on-line computer) of the photograph taken for the  $X=34.8\%$  mixture at  $115.5^\circ\text{C}$ . Typical widths of the spots are  $0.8 \times 10^{-2} \text{ \AA}^{-1}$  in the  $Z$  direction and  $1.7 \times 10^{-2} \text{ \AA}^{-1}$  in the  $X$  direction. The larger width in the  $X$  direction arises from the geometry of the x-ray monochromator setup. However, it is evident that any displacement of the reflections along the  $X$  axis arising from a lateral periodicity of several hundred angstroms would at once be revealed in the contours. We therefore conclude that the three wave vectors are collinear along the  $Z$  axis.

The thermal evolution of the layer spacing corresponding to the different modulations in the  $A$  phases of the same mixture is shown in Fig. 5. The variations in the  $A_d$  and  $A_2$  phases are similar to those seen for the  $18\%$  mixture (Fig. 2). In the  $A_{ic}$  phase,  $2\pi/q'_0$  shows a marked decrease with decrease of temperature. Measurements of the layer spacing for a number of concentrations in the region of existence of the  $A_{ic}$

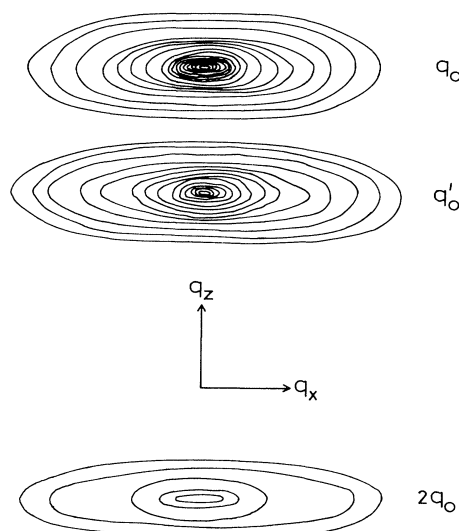


FIG. 4. Intensity contour diagram of a photograph taken for the  $X=34.8\%$  mixture at  $115.5^\circ\text{C}$ . Typical widths of the spots are  $q_z=0.8 \times 10^{-2} \text{ \AA}^{-1}$  and  $q_x=1.7 \times 10^{-2} \text{ \AA}^{-1}$ . The spot at  $2q_0$  has been displaced (along  $Z$ ) closer to the other spots for convenience of representation.

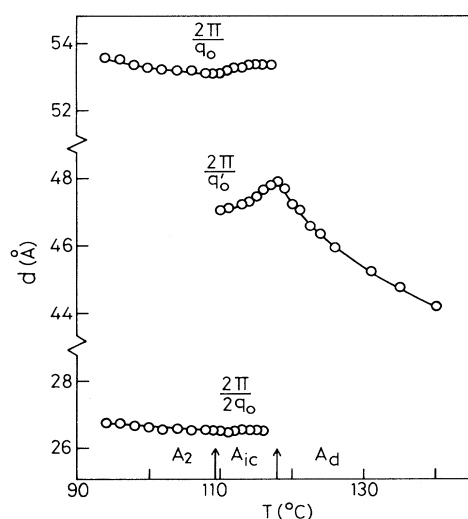


FIG. 5. Temperature variation of the layer spacing ( $d$ ) in the  $A_d$ ,  $A_{ic}$ , and  $A_2$  phases of the  $X=34.8\%$  mixture. The wave vectors corresponding to the different periodicities are discussed in the text. The arrows indicate the temperatures of transition between the  $A$  phases.

phase have confirmed this behavior. This decreasing trend is exactly opposite to what one gets from a simple calculation of the layer spacing variation in the  $A_{ic}$  phase assuming it to be a two-phase region. Also, differential scanning calorimetry runs (taken with a Perkin-Elmer DSC-4 in conjunction with the Thermal Analysis Data Station) of the  $A_d$ - $A_2$  transition show a rapid decrease in the strength of the signal with increasing 8OCB concentration, and no signal is observable for  $X \sim 20\%$  which is a few percent away from the concentration at which the  $A_{ic}$  phase makes its appearance. Even in the region where the  $A_{ic}$  phase exists, no signals were observed corresponding to the  $A_d$ - $A_{ic}$  and  $A_{ic}$ - $A_2$  transitions. Thus the possibility of the  $A_{ic}$  phase being a two-phase region is ruled out.

Cladis and Brand<sup>9</sup> have recently observed an inverted cholesteric phase which appears as an island surrounded by different types of smectic- $A$  phases in the temperature-concentration plane. These authors argued that their results imply that the coexistence of

two percolating collinear density waves with different periodicities in a smectic- $A$  phase is incompatible with the fluidity and order-parameter rigidity of the  $A$  phase. The present study shows that the  $A$  phase can support two incommensurate collinear periodicities over a range of temperature. Prost and Barois<sup>5</sup> have suggested two molecular models for the incommensurate  $A$  phase: Depending on the relative strengths of the elastic and lockin terms in the free-energy expansion, the incommensurate density modulations can coexist either by percolating through each other or as a multisoliton regime. Which of the two molecular models represents the  $A_{ic}$  phase remains to be settled. Clearly, further studies are needed for a complete understanding of this new phase.

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<sup>1</sup>See, e.g., J. D. Litster, in *Liquid Crystals of One- and Two-Dimensional Order*, edited by W. Helfrich and G. Heppke (Springer-Verlag, Berlin, 1980), p. 65.

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<sup>8</sup>For a description of the setup used in the x-ray diffraction experiments see B. R. Ratna, S. Krishna Prasad, R. Shashidhar, G. Heppke, and S. Pfeiffer, *Mol. Cryst. Liq. Cryst.* **124**, 21 (1985).

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