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Mean-Free-Path Estimators for Finite-Range Measurements and Low Statistics. S. GARPMAN, I. LUND, I. OTTERLUND, and K. SÖDERSTRÖM [Phys. Rev. Lett. **53**, 2195 (1984)].

On pg. 2197, Eq. (15) reads

$$\frac{(e^{D/\lambda} - e^{-ND/\lambda})}{(e^{D/\lambda} - 1)(N + 1)}.$$

It should read

$$\frac{(e^{D/\lambda} - e^{-ND/\lambda})}{(e^{D/\lambda} - 1)(N + 1)}.$$

Measurements of Gamow-Teller Strength Distributions in Masses 13 and 15. C. D. GOODMAN, R. C. BYRD, I. J. VAN HEERDEN, T. A. CAREY, D. J. HOREN, J. S. LARSEN, C. GAARDE, J. RAPAPORT, T. P. WELCH, E. SUGARBAKER, and T. N. TADDEUCCI [Phys. Rev. Lett. **54**, 877 (1985)].

Table I, line 1, column 3 should read 0.46 ± 0.02 .

An incorrect version of Fig. 1 was published. The correct version is given here.

Vector-Meson Mass Generation by Chiral Anomalies. R. JACKIW and R. RAJARAMAN [Phys. Rev. Lett. **54**, 1219 (1985)].

It has been brought to our attention, after our Letter was published, that a two-dimensional model of a vector meson carrying an explicit bare mass term and coupled to chiral fermions was solved by C. R. Hagen [Ann. Phys. (N.Y.) **81**, 67 (1973)]. We thank Dr. A. Guha for pointing this out to us.

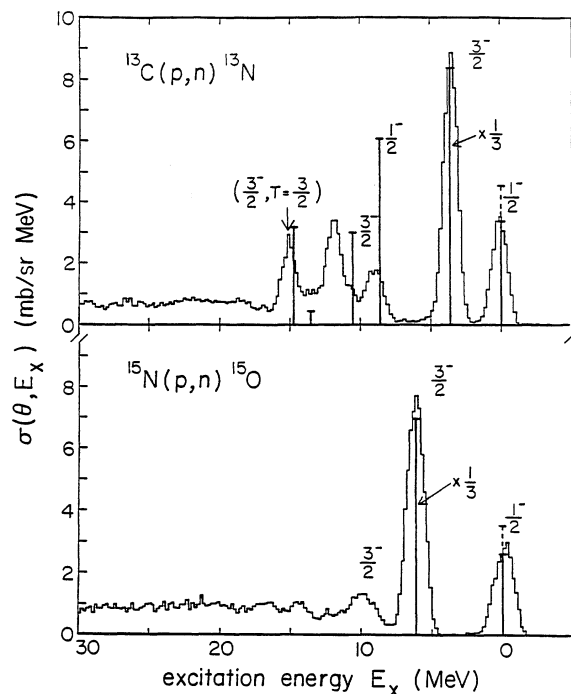


FIG. 1. Differential cross section vs excitation energy for (top) $^{13}\text{C}(p,n)^{13}\text{N}$ and (bottom) $^{15}\text{N}(p,n)^{15}\text{O}$ at $\theta = 0^\circ$ and $E_p = 160$ MeV. The $^{15}\text{N}(p,n)$ spectrum was obtained from measurements with a Melamine ($\text{C}_3\text{H}_6^{15}\text{N}_6$) target and has had the carbon contribution subtracted. The $^{12}\text{C}(p,n)$ contaminant transitions would appear in the ^{15}N spectrum at excitation energies greater than 14.6 MeV. The vertical bars illustrate the calculated GT strength distributions of Ref. 5. The dashed portions of the ground-state bars represent the Fermi contributions to the cross sections. The scale is chosen so that the ground-state bars would match the data peaks in height if the beta-decay $B(\text{GT})$ values were used instead of the calculated values. The bars for the major $\frac{3}{2}^-$ peaks are reduced by a factor of 3 to keep them on scale. The spin assignments of the unlabeled experimental peaks cannot be made on the basis of this experiment alone.