

Evaporation of Neutral Strange Particles in \bar{p} -Ta at 4 GeV/c

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The reaction $\bar{p} + \text{Ta} \rightarrow V^0 + X$ ($V^0 = \Lambda, \bar{\Lambda}, \text{or } K_s^0$) at 4 GeV/c has been studied by comparing with the results in $\bar{p}p$ interaction. Strong enhancement of Λ production was observed, while $\bar{\Lambda}$ production is suppressed. The $\bar{\Lambda}/\Lambda$ production ratio is obtained as $(2.0 \pm 1.0) \times 10^{-2}$. Λ reveals evaporationlike behavior with a nuclear temperature of 97 MeV. Evaporationlike behavior with a temperature of 135 MeV, but via a short-range process, was observed also in K_s^0 production.

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Recently the effects of collective multinucleon processes¹ and phenomena of high nuclear temperature² have been deduced by the behavior of kaons produced in relativistic heavy-ion collisions. Λ production mechanisms are studied also in nucleus-nucleus collisions,^{3,4} but their results contradict each other in the dominant production channel $NN \rightarrow \Lambda KN$. It has been emphasized that strange particles could be a better probe for a study of high-energy nuclear reaction dynamics.⁵ Therefore, the strange-particle production mechanism has been theoretically studied from various points of view,⁶⁻¹¹ but details of the dynamics of the strange-particle production are still not clear.

In this report we present the result of an experiment in which we study the production mechanism of the neutral strange particles in \bar{p} -Ta collisions, and compare with the data of $\bar{p}p$ interaction.^{12,13} We chose \bar{p} as a projectile because of its large cross section for strange-particle productions in $\bar{p}p$ collisions. In addition, there has been only very scarce data of strange-particle production from \bar{p} -nucleus interactions.

The data come from an exposure of the KEK 1-m hydrogen bubble chamber to the beam of 4 GeV/c at the National Laboratory for High Energy Physics (KEK). Two 4.4-mm-thick Ta plates as a nuclear target are installed in the bubble chamber. The results presented here come from 107×10^3 pic-

tures. Details of the experimental procedure will be described elsewhere.

Λ , $\bar{\Lambda}$, and K_s^0 are observed through their charged decay modes and corrections for the unseen modes are carried out. The V^0 events are analyzed by use of the program TVGP and are weighted with a correction factor (~ 1.3 on the average) for the limited detection volume of the bubble chamber. Corrections are made for scanning losses (3%) and for unidentified V^0 events (10%). The cross sections are corrected for absorption of V^0 by Ta plates by use of its longitudinal momentum spectra (16%).

The results are summarized in Table I, together with those in $\bar{p}p$ at 4 GeV/c.^{12,13} The production cross section for Λ is much larger than that for $\bar{\Lambda}$ in \bar{p} -Ta, while those in $\bar{p}p$ are approximately equal to each other. If the strange particle is produced only at the primary collision of the \bar{p} , the cross section should be $32.0 (A^{2/3})$, due to A dependence) times larger than that in $\bar{p}p$. The reason is that the inelastic cross section is just the same size as the geometrical cross section of the Ta nucleus. However, the measured cross sections are 364 ± 41 (Λ), 8.9 ± 4.2 ($\bar{\Lambda}$), and 43.2 ± 3.5 (K_s^0) times as large as the ones in $\bar{p}p$. The Λ production is surprisingly enhanced, while the $\bar{\Lambda}$ production is suppressed by the $\Lambda N \rightarrow K\bar{N}N, \bar{N}N$, annihilation, . . . processes in the nucleus. The remarkable difference in these cross-section ratios is suggestive of different dynamics in

TABLE I. Summary of the cross sections.

Channel	Number of events	Cross section (mb)	Cross section in $\bar{p}p$ (mb)
Total inelastic	...	1628 ± 30	...
$K_s^0 + X$	445	82.0 ± 6.0	1.90 ± 0.07
$\Lambda + X$	929	193 ± 12	0.53 ± 0.05
$\bar{\Lambda} + X$	21	3.8 ± 2.0	0.44 ± 0.05

strange-particle productions. The $\bar{\Lambda}/\Lambda$ production ratio is obtained as $(2.0 \pm 1.0) \times 10^{-2}$. The value is larger than an upper limit of the $\bar{\Lambda}/\Lambda$ production ratio (4.5×10^{-3}) in nucleus-nucleus collision at $4.5A$ GeV/c.⁴ This is reasonable because the projectile is an antiparticle, \bar{p} , in the present experiment, while in Ref. 4, $\bar{\Lambda}$ cannot be produced by elementary process because of the lack of incident energy. Further studies are needed to determine the $\bar{\Lambda}$ production mechanism. In this paper we restrict ourselves to the analyses on Λ and K_s^0 because of insufficient statistics in the $\bar{\Lambda}$ data.

Figure 1 shows rapidity distributions of Λ and K_s^0 from \bar{p} -Ta in the laboratory system. The rapidity distribution of K_s^0 is symmetric. Its peak lies at $y = 0.6$ and the peak goes to $y^* = 0$ if one assumes the \bar{p} -3N c.m. frame, i.e., that the incident \bar{p} collides with an effective target of three-nucleon mass in the center-of-mass system. The frame is moving with $\beta = 0.54$ in the laboratory system. The rapidity distribution of Λ is roughly symmetric but has an

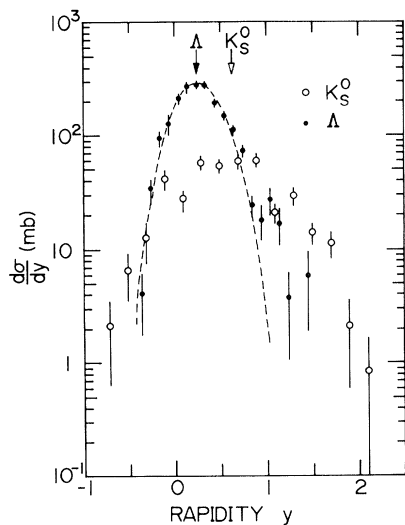


FIG. 1. Rapidity spectra of K_s^0 and Λ . The dashed curve shows a symmetric curve for the Λ distribution. The centers of the distributions of K_s^0 and Λ are indicated by arrows.

asymmetric tail on the higher-rapidity side. The tail lies just in the rapidity region of the Λ 's from the $\bar{p}p$ interaction.¹³ The peak of the rapidity distribution of Λ goes to $y^* = 0$ if one assumes the \bar{p} -13N c.m. frame, moving with $\beta = 0.24$.

In order to explore the source of the asymmetric distribution of Λ , a scatter plot of momentum p^* versus rapidity y^* of Λ 's from \bar{p} -Ta is shown in Fig. 2. Momentum p^* and rapidity y^* are defined in the \bar{p} -13N c.m. frame. In Fig. 2 the absence of data points around $y^* = -0.35$ is due to the Ta target plates. A dashed curve is drawn in order to show a shape of the distribution of Λ from $\bar{p}p$ in the moving system with $\beta = 0.24$. The data points concentrate almost symmetrically in the region of $p^* \leq 0.8$ GeV/c and $|y^*| \leq 0.8$ (assigned as a region S , shown by a solid curve). Data points outside of the region S overlap well with the distribution expected from the $\bar{p}p$ interactions, which is shown with a dashed curve. This means that these events come from elementary $\bar{p}N$ interactions and could be separated from the main symmetric region S . It is also worthwhile to notice that the number of these events and the angular distribution correspond to

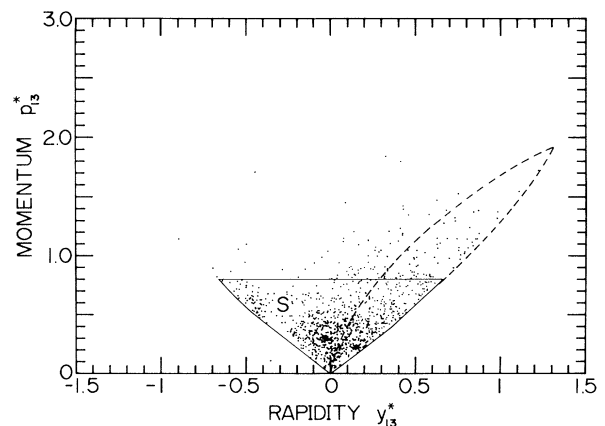


FIG. 2. Scatter plot of momentum p^* vs rapidity y^* which are defined in the \bar{p} -13N c.m. frame. A dashed curve shows the shape of the distribution of Λ 's in $\bar{p}p$ collisions.

those expected from the elementary $\bar{p}p$ cross section [see Fig. 3(b)].

In order to study the dynamical characteristics of Λ 's inside the region S , we have to know how much contribution is made from the Λ 's produced in the elementary $\bar{p}N$ interactions. This contribution is estimated to be 8% from the $\bar{p}p$ data.¹³ Therefore, it is negligible for further analysis, such as angular distributions. Figure 3(a) shows the angular distribution of Λ 's inside the symmetric region S in the \bar{p} -13N c.m. frame. The fact that the distribution is clearly isotropic indicates an evaporation character of the process. The difference of the Λ production processes is clear when one compares the angular distribution with the one in Fig. 3(b), that is, for the events outside the region S . The angular distribution of K_s^0 is also almost flat in the \bar{p} -3N c.m. frame, as seen in Fig. 3(c). The fact suggests a similar evaporation character of the K_s^0 production process.

The inverse slopes of kinetic-energy spectra $[(1/p^*)d\sigma/dT^*]$ are analyzed for K_s^0 and Λ inside the region S and obtained as 135 MeV for K_s^0 in the \bar{p} -3N c.m. and 97 MeV for Λ in the \bar{p} -13N c.m. frame, respectively. The value of 135 MeV for K_s^0 is comparable with the experimental and theoretical ones for K^+ production in the nucleus-nucleus collisions.^{2,10} These temperatures may not be high enough to induce a phase transition to quark-gluon plasma, which requires a higher temperature, of the order of 200 MeV.¹⁴ K_s^0 is emitted through an evaporationlike process from the hot cluster consisting of three or four nucleons. This indicates a short-range correlation of K_s^0 . The main fraction of the Λ 's, which is responsible for the enhancement of Λ production, appears to be leaving from the hot \bar{p} -13N cluster through the evaporation mechanism. This cluster may break up during the traversal of the nucleus because of its high internal energy, $\frac{1}{3}$ of the kinetic energy of the cluster. An average Λ polarization was also measured to be consistent with zero ($P=0.02 \pm 0.09$), which is in agreement with the results in nucleus-nucleus collisions.^{3,4} The polarization of the Λ also indicates the evaporationlike behavior.

In summary, we have seen a very large Λ production cross section and have observed evaporationlike Λ 's from a hot cluster with $\beta=0.24$ and $T=97$ MeV. A similar observation is made also for the evaporation process of K_s^0 from a moving system with $\beta=0.54$ and $T=135$ MeV.

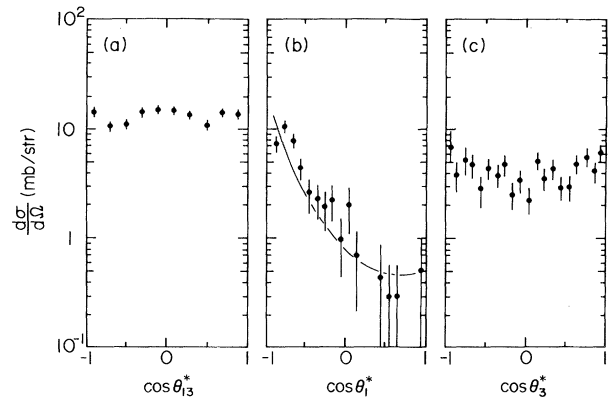


FIG. 3. Angular distributions of Λ and K_s^0 . (a) Λ 's inside the region S in the \bar{p} -13N c.m. frame. (b) Λ 's outside of the region S , which is the angular distribution in the $\bar{p}p$ c.m. system; the solid curve shows the angular distribution of Λ 's in the $\bar{p}p$ interaction, but multiplied by 32.0 ($= A^{2/3}$). (c) K_s^0 in the \bar{p} -3N c.m. frame.

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