

Universal Power Law for the Dimension of Strange Attractors near the Onset of Chaos

In a recent Letter¹ it was shown that the envelope of the Lyapunov exponent of one-dimensional, one-parameter families of dynamical systems that period double on their way to chaos rises like $(\mu - \mu_\infty)^\beta$ where μ is the stress parameter, μ_∞ is the point of accumulation of period doublings, and $\beta = \ln 2 / \ln \delta$, $\delta = 4.669$ Here we comment that (i) this result can be generalized to the (single) positive Lyapunov exponent in any dissipative multidimensional map $M_\mu: R^F \rightarrow R^F$ that undergoes period doubling, and (ii) that this implies a universal scaling law for the dimension of strange attractors which develop after the completion of a cascade of period doublings. By this we mean that the information dimension² D_1 depends on the stress parameter μ according to $(D_1 - 1) \sim (\mu - \mu_\infty)^\beta$ for values of μ that do not belong to a window of periodicity.

The vector of Lyapunov exponents, $\vec{\lambda}[M_\mu]$ obeys the identity $\vec{\lambda}[M_\mu] = 2^{-n} \vec{\lambda}[T^n M_\mu]$, where T is the doubling operator.³ For $\mu - \mu_\infty \ll 1$ one can linearize T and show that $\vec{\lambda}[M_\mu] = 2^{-n} \vec{\lambda}[G + \delta^n(\mu - \mu_\infty)aH] = (\mu - \mu_\infty)^\beta \vec{\lambda}[G + aH]$. Here G is the fixed-point map of T at μ_∞ , a is a number, and H is the eigenfunction in the (single) unstable direction. The map G has $F-1$ (negative) infinite Lyapunov exponents due to infinite contraction rates in $F-1$ directions.³ Accordingly, $F-1$ Lyapunov exponents of M_μ may reach a (negative) finite limit (a "regular part") when $\mu = \mu_\infty$ and only the one associated with the singled out direction in phase space³ vanishes like $(\mu - \mu_\infty)^\beta$ (i.e., is purely "singular").

To calculate the dimension, we invoke Kaplan-Yorke (KY) formula^{4,5} $D_1 = j + \sum_{i=1}^j \lambda_i / |\lambda_{j+1}|$, where j is the highest index for which $\sum_{i=1}^j \lambda_i \geq 0$. Since $|\lambda_{j+1}|$ would be dominated by its regular part, the considerations above lead directly to the universal scaling law of $(D_1 - 1)$.

We tested the prediction numerically in two examples, i.e., Helleman's standard map $(x', y') = (2\mu x + 2x^2 + by, x)$, $b = 0.3$, $\mu_\infty = -0.7639$, and the simple three-dimensional map $(x, y, z) = (1 - \mu x^2 + by + cz, x, y)$, $b = c = 0.15$, $\mu_\infty = 1.3892$. Excellent agreement with the scaling law of $(D_1 - 1)$ is found for all values of μ that do not belong to a window of periodicity; see Fig. 1.

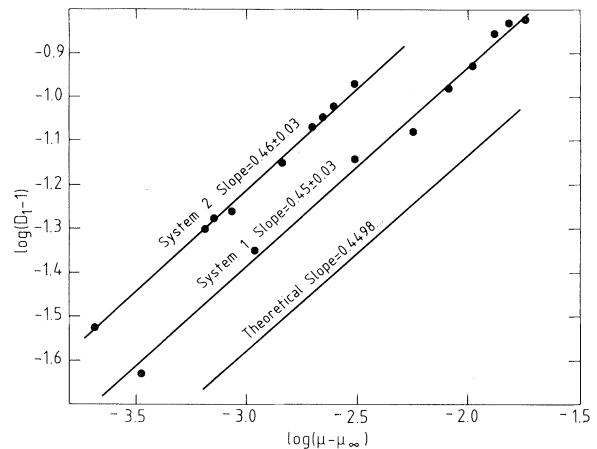


FIG. 1. $\ln(D_1 - 1)$ vs $\ln(\mu - \mu_\infty)$ for the examples discussed in the text.

An experimental verification of the scaling law of $(D_1 - 1)$ is particularly important since it would serve as a test of the KY formula in addition to providing support for the theory leading to this scaling law. Such a test is important since a rigorous proof of the KY formula as an equality for high-dimensional systems is still lacking, although insight on its validity was obtained in Ref. 5.

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