## Total-Scattering Measurements and Comparisons for Collisions of Electrons and Positrons with  $N_2O$

Ch. K. Kwan, Y.-F. Hsieh, W. E. Kauppila, Steven J. Smith, T. S. Stein, and M. N. Uddin Department of Physics and Astronomy, Wayne State University, Detroit, Michigan 48202

and

## M. S. Dababneh

Department of Physics, Yarmouk University, Irbid, Jordan (Received 12 December 1983)

Total cross sections ( $O_T$ ) for 1–500-eV positron and electron scattering by N<sub>2</sub>O are measured by a beam-transmission method. Comparisons of these results with prior electron and positron measurements of  $Q_T$  for  $CO_2$ , N<sub>2</sub>, and CO, made by our group, reveal several remarkable similarities in the shapes and magnitudes of the  $Q<sub>T</sub>$  curves for the triatomic and also for the diatomic molecules. Small differences in the absolute  $Q_T$  values for N<sub>2</sub>O and  $CO<sub>2</sub>$ , and for N<sub>2</sub> and CO, may be due to the weak permanent dipole moments of N<sub>2</sub>O and CO.

PACS numbers: 34.80.-i

It has been known for some time that certain molecules, because of their similar structure and geometry, scatter electrons in intriguingly similar ways. An example is the isoelectronic pair  $N_2$  and CO. A number of authors have commented on a variety of processes in which  $N_2$  and CO give remarkably similar results for electron scattering: Hake and Phelps' have observed that the transport properties of electrons in CO resemble those in  $N_2$ ; Schulz<sup>2</sup> has commented on their similarities for resonance formation and vibrational excitation; Onda and Truhlar<sup>3</sup> have discussed their similar behavior in rotational excitation; and Bromberg<sup>4</sup> and Dubois and Rudd<sup>5</sup> have recognized similarities in their elastic differential cross sections at higher energies. Recent theoretical calculations for electron scattering from these molecules by Jain<sup>6</sup> and Tayal et al.<sup>7</sup> have found interesting similarities in the elastic differential, integral, and momentum transfer cross sections in the 40-800 eV energy range. Measurements of  $Q_T$  for these gases performed in our laboratory<sup>8,9</sup> have not only confirmed the findings of Jain regarding electron  $Q_T$ 's but also have shown that many of the conclusions concerning the  $Q_T$ 's for electrons can be extended to those of positron  $Q_T$ 's on these gases as well (see Ref. 9 for a detailed discussion). In order to investigate further the possibility that other molecules may have similar characteristics, we choose for comparison the following pair of molecules in the N-0-C family:  $N_2O$  and  $CO_2$ . They are isoelectronic, linear in their ground states, and formed by adding an oxygen atom to  $N_2$  and CO, respectively.

In this paper we present measurements of the to-

tal cross sections for electron and positron scattering by  $N_2O$  in the energy range of 1-500 eV. This is the first time that total cross sections have been reported for positron impact on nitrous oxide. In addition, no  $Q_T$  data have been available for the  $e^-$ -N<sub>2</sub>O system at energies beyond 50 eV.

This work constitutes an extension of the experimental studies being pursued in this laboratory on electron and positron total-scattering cross sections by gases using a beam-transmission method; detailed descriptions of the apparatus, experimental procedure, and error analysis have previously been procedure, and error analysis have previously been<br>reported.<sup>10,11</sup> The primary difference in the electron and positron measurements of  $Q_T$  is that the positron beam originates from an  $\overline{1}C$  positron source (produced by bombardment of a polycrystalline  $^{11}$ B target with a 4.75-MeV proton beam from a Van de Graaff accelerator), while the electron beam is produced by a thermionic electron source. The energy widths for both the  $e^+$  and the  $e^-$  beams are about 0.<sup>1</sup> eV. Using the method of Kauppila are about  $0.1$  eV. Using the method of Kauppila *et al.*, <sup>11</sup> we estimate that the experimental errors for our present absolute electron  $Q_T$  results range from 5% at 1.2 eV to 3% at 500 eV and for the  $e^+$ -N<sub>2</sub>O  $Q_T$  results, range from 6% to 4%, respectively. The errors for comparisons of our electron and positron data are smaller than the absolute  $Q_T$  errors because of the fact that several of the error components affect the electron and positron measurements in the same way. For the same reason, this is also true when comparing our results for different gases. We estimate that the total errors in the comparisons of our electron and positron measurements (and in our comparisons for different gases) is  $3\%$  for electrons and 4% for positrons. However, the above estimates do not include the potential source of error associated with incomplete discrimination against projectiles scattered at small angles in the forward direction. Depending on the extent that forward-angle scattering occurs, our measured results may be lower than the actual total cross sections. Following the procedure elaborated by Kauptions. Following the procedure elaborated by Kaup<br>pila *et al.*, <sup>11</sup> the discrimination angle ranges from  $20^{\circ}$  at 2 eV to 8° at 500 eV for the present  $e^{+}$  measurements, and ranges from  $11^{\circ}$  at 2 eV to 5 $^{\circ}$  at 500 eV for the  $e^-$  work. In order to estimate the extent that small-angle scattering could affect our  $Q_T$ results, detailed information on the differential scattering cross section would be needed. We note also that with the use of a retarding element after the scattering region, there is complete discrimination against projectiles scattered with an energy loss of more than a few tenths of an electronvolt.

In Fig. 1 we compare our present low-energy  $e^-$ -N<sub>2</sub>O  $Q_T$  measurements with earlier measurements by Bruche,<sup>12</sup> Ramsauer and Kollath,<sup>13</sup> and Zecca et al.  $14$  Our measurements agree quite well with the measurements of Bruche, except at the peak (2.3 eV) of the shape resonance where our  $O<sub>T</sub>$ values are about 25% higher. The results of Zecca er al. (which were normalized to those of Bruche at 4 eV) show a much less pronounced shape resonance than the present measurements and those of Bruche.

The complete set of present  $Q_T$  measurements for  $e^+$  and  $e^-$  scattering by N<sub>2</sub>O are shown in Fig.



FIG. 1. Low energy  $e^-$ -N<sub>2</sub>O total–cross-section measurements. The present results are shown with the measurements of Bruche (Ref. 12), Ramsauer and Kollath (Ref. 13), and Zecca et al. (Ref. 14). The "n" refers to normalized measurements.

2 along with prior  $(e^+, e^-)$ -CO<sub>2</sub>  $Q_T$  measurements. <sup>8,9</sup> It is seen that above 5 eV the shape and magnitude of the  $Q_T$  results for both N<sub>2</sub>O and CO<sub>2</sub> are very similar for each respective projectile. In fact, the electron  $Q_T$  curves are particularly close to each other (within a few percent) above 5 eV. Meanwhile, the  $e^+$ -N<sub>2</sub>O  $Q_T$  results tend to be slightly, but noticeably, higher than those for the  $e^+$ -CO<sub>2</sub> system with the discrepancy ranging from 8% in the vicinity of 30 eV to 1% at 500 eV. At the lowest energies  $(< 5 eV$ ) the internal structure of the target molecule plays a more prominent role in the collision process, and it would be expected that for different gases the differences in the  $Q_T$ 's will be more pronounced there. Nevertheless, the comparison in Fig. 2 of the present results with those of Hoffmann et al.<sup>8</sup> for  $(e^+,e^-)$ -CO<sub>2</sub> scattering at these low energies show that all the features discussed there are also present for  $N_2O$  scattering. For electron scattering there are prominent shape



FIG. 2. Comparison of  $(e^+, e^-)$ -N<sub>2</sub>O and  $(e^+, e^-)$ - $CO<sub>2</sub>$  total cross sections up to intermediate energies. The present measurements of  $(e^+, e^-)$ -N<sub>2</sub>O total cross sections are compared with the  $CO<sub>2</sub>$  results of Hoffman er al. (Ref. 8) and Kwan er al. (Ref. 9). The threshold energies for formation of positronium in the ground and first excited states are indicated by the arrows labeled Ps and Ps", respectively (inset).

resonances at 2.3 and 3.8 eV for N<sub>2</sub>O (of <sup>2</sup> $\Sigma$ <sup>+</sup> symmetry) and  $CO_2$  (of <sup>2</sup>II<sub>u</sub> symmetry), respectivelya feature which is absent for positron scattering from these gases. Beyond these shape resonances, the  $e^-$  total cross sections increase and reveal a broad maximum around 25 eV for  $N_2O$  and 30 eV for  $CO<sub>2</sub>$ . It is interesting that broad shape resonances at intermediate energies (10—40 eV) caused by higher excited states of the negative molecular ion are not uncommon for electron-molecule scattering, as has been discussed by Lynch et al.<sup>15</sup> for  $CO<sub>2</sub>$ , OCS, and CS<sub>2</sub>. They suggest a possible enhancement of vibrational excitation at these resonances. Total-cross-section measurements, however, may be too insensitive to detect these effects.

In the case of positrons the  $Q_T$  curves for N<sub>2</sub>O and  $CO<sub>2</sub>$  also reveal similar qualitative features. The  $Q_T$ 's increase rapidly below 2 eV. Both curves show an abrupt rise ("bump") at the positronium formation thresholds of 6.25 eV for  $N_2O$  and 7.0 eV for  $CO<sub>2</sub>$  suggesting a significant increase of inelastic scattering (due to positronium formation) at these energies, which is consistent with our earlier measurements for other gases. $8 - 10$  It is curious that the overall increases of the  $Q_T$ 's from the positronium formation thresholds for  $N_2O$  and  $CO_2$  are not as large as for most of the other gases that have been studied up to the present time. There appears to be a second statistically significant increase in the  $O<sub>T</sub>$ curves for  $N_2O$  and  $CO_2$  at an energy of about 5 eV above each positronium formation threshold (see inset of Fig. 2). Since the first excited state of positronium is 5.1 eV above the ground state, these latter increases in the  $Q_T$  curves may be due to the formation of positronium in the first excited state.

It is relevant to the present discussion that comparisons between prior electron and positron measurements of  $Q_T$  for N<sub>2</sub> and CO by our group,<sup>8,9</sup> which are shown together in Fig. 3, indicate a very similar situation as we have observed for  $N_2O$  and  $CO<sub>2</sub>$ , in that there are striking similarities and some small differences in the shapes and magnitudes of the respective  $Q_T$  curves. For electron scattering both the  $N_2$  and CO  $Q_T$  curves exhibit shape resonances at low energies and agree to within a few percent of each other for electron energies above 10 eV. For positron scattering both of the  $Q_T$  curves are increasing at the lowest energies and exhibit appreciable increases after the positronium formation threshold.

In comparing the electron and positron  $Q_T$  magnitudes for the diatomic and triatomic molecules in Figs. 2 and 3, it is intriguing that both of the polar molecules, CO and  $N<sub>2</sub>O$  (with dipole moments of



FIG. 3. Comparison of  $(e^+, e^-)$ -N<sub>2</sub> and  $(e^+, e^-)$ -CO total cross sections up to intermediate energies. The  $(e^+, e^-)$ -N<sub>2</sub> results are from Hoffman *et al.* (Ref. 8) and the  $(e^+, e^-)$ -CO results from Kwan *et al.* (Ref. 9). The threshold energies for formation of positronium in the ground state are indicated by the labeled arrows.

0.112 and 0.167  $D<sub>1</sub><sup>16</sup>$  respectively), have noticeably larger  $Q_T$ 's for positron scattering than the corresponding nonpolar molecules,  $N_2$  and  $CO_2$ . Meanwhile, the  $Q_T$ 's for electron scattering by these two pairs of molecules are much closer to each other than for positrons. These observations raise some interesting questions as to what role the permanent dipole moment plays in positron and electron scattering by these molecules and why there is a more pronounced effect for positrons. The  $Q_T$  differences for positrons may be related to the positronium channel (which is open only to positrons), although the  $Q_T$  differences also are seen to exist below the positronium formation thresholds. In electron differential elastic scattering cross section measurements for  $N_2$  and CO, Brom $berg<sup>4</sup>$  has found that in the energy range 300–500 eV there are no significant differences in the differential cross sections at a particular energy for scattering angles greater than 12°, while at smaller angles the CO cross sections rise at a slightly faster rate than for  $N_2$ . This observation led Bromberg to suggest that these small differences may be attributable to the permanent dipole moment of CO, which has been supported in the recent theoretical work has been supported in the recent theoretical work<br>by Tayal *et al.*<sup>11</sup> In view of this information, it is likely that the  $e^-$ -CO  $Q_T$  results may be slightly larger than the  $e^-$ -N<sub>2</sub> results because of the additional small-angle elastic scattering for CO. As a result of the lack of theoretical work and any more detailed experimental work (than  $Q_T$  measurements) on positron scattering by CO,  $N_2$  CO<sub>2</sub>, and  $N<sub>2</sub>O$ , it can only be speculated that the larger differences in the  $Q_T$  measurements for positron scattering than for electron scattering by  $N_2O$  and  $CO_2$ , and by  $N_2$  and CO, may result from the permanent dipole moments of  $N_2O$  and CO having a greater effect on positron scattering. It would seem that these latter comparisons and the possible evidence for formation of postronium in the first excited state for positron scattering by  $N_2O$  and  $CO_2$  should provide a stimulus for future experimental and theoretical work.

We are grateful to D. Jerius for his untiring help in various aspects of this project and to Dr. J. M. Wadehra for useful discussions concerning this work. This work was supported by the National Science Foundation under Grant No. PHY80-07984.

 $R$ . D. Hake, Jr., and A. V. Phelps, Phys. Rev. 158, 70

(1967).

 ${}^{2}G$ . J. Schulz, in *Principles of Laser Plasmas*, edited by George Bekefi (Wiley, New York, 1976), Chap. 2.

 $3K$ . Onda and D. G. Truhlar, J. Chem. Phys. 73, 2688 (1980).

4J. P. Bromberg, J. Chem. Phys. 52, 1243 (1970).

5R. D. Dubois and M. E. Rudd, J. Phys. B 9, 2657 (1976).

6Ashok Jain, J. Phys. B 15, 1533 (1982).

7S. S. Tayal, Ashok Jain, A. N. Tripathi, and M. K. Srivastava, J. Chem. Phys. 78, 3021 (1983).

8K. R. Hoffman, M. S. Dababneh, Y. F. Hsieh, W. E. Kauppila, V. Pol, J. H. Smart, and T. S. Stein, Phys. Rev. A 25, 1393 (1982).

9Ch. K. Kwan, Y. F. Hsieh, W. E. Kauppila, S. J. Smith, T. S. Stein, M. N. Uddin, and M. S. Dababneh, Phys. Rev. A 27, 1328 (1983).

10T. S. Stein, W. E. Kauppila, V. Pol, J. H. Smart, and G. Jesion, Phys. Rev. A 17, 1600 (1978).

<sup>11</sup>W. E. Kauppila, T. S. Stein, J. H. Smart, M. S. Dabaneh, Y. K. Ho, J. P. Downing, and V, Pol, Phys. Rev. A 24, 725 (1981).

 $12E$ . Bruche, Ann. Phys. (Leipzig) 83, 1065 (1927).

<sup>13</sup>C. Ramsauer and R. Kollath, Ann. Physik (Leipzig) 7, 176 (1930).

<sup>14</sup>A. Zecca, I. Lazzizzera, M. Krauss, and C. E. Kuyatt, J. Chem. Phys. 61, 4560 (1974).

<sup>15</sup>M. G. Lynch, D. Dill, J. Siegel, and J. L. Dehmer, J. Chem. Phys. 71, 4249 (1979).

<sup>16</sup>R. D. Nelson, Jr., D. R. Lide, Jr., and A. A. Maryott Selected Values of Electric Dipole Moments for Molecules in the Gas Phase, U. S. National Bureau of Standards, National Standard Reference Data Series—<sup>10</sup> (U. S. GPO, Washington, D. C., 1967).