

Transverse Electromagnetic Waves with Nonzero $\vec{E} \cdot \vec{B}$

In a recent Letter,¹ Khare and Pradhan demonstrate electromagnetic wave equations with nonzero $\vec{E} \cdot \vec{B}$ values and obtain nonzero values of $q = -\int \vec{E} \cdot \vec{B} d^4x$ and $s = \int (E^2 - B^2) d^4x$. The first result is not actually that remarkable. The latter two seem to stem partly from an error in analysis and partly from introducing static uniform components of \vec{E} and \vec{B} in the long-wavelength limit that fill the universe and give the nonzero q and s . The transversality and mutual orthogonality of \vec{E} and \vec{B} for ordinary light waves are usually demonstrated for unidirectional, propagating waves. However, the quantities q and s need not be zero for static fields, and such nonzero values can be introduced artificially as the zero-frequency limit of a spectral representation. There are nevertheless finite-frequency waves for which $\vec{E} \cdot \vec{B}$ is nonzero.² Consider a single wave, here labeled (1), to have $E_y^{(1)} = B_z^{(1)}$ and propagation vector $k_x^{(1)}$. Superimposing a second wave labeled (2) propagating in the opposite direction, $k_x^{(2)} = -k_x^{(1)}$, but with $E_z^{(2)} = B_y^{(2)} = E_y^{(1)}$, then gives a standing wave with $\vec{E} \cdot \vec{B} = 2(E_y^{(1)})^2 [\cos(2kx) + \cos(2\omega t)]$ which has nonzero time-averaged $\vec{E} \cdot \vec{B}$.

To construct solutions of the above type with nonzero q , consider (instead of sine waves) isolated square-wave pulses having the polarization properties above and encountering one another. During the brief moment that the waves overlap, one has finite nonzero $\vec{E} \cdot \vec{B}$, and zero all other times. If one now considers an infinite train of such pulses going in each direction, then $\vec{E} \cdot \vec{B}$ will apparently have a finite time and spatial average. But such wave trains have dc components of E and B . What we are really doing is filling the universe with uniform \vec{E} and \vec{B} fields, in which case (at least for an empty Euclidean universe) there is no prohibition against having nonzero q and s . The condition for nonzero $\vec{E} \cdot \vec{B}$, in the above example, is that the wave forms have nonzero path integrals: $\int_{-\infty}^{+\infty} A(\eta) d\eta \neq 0$, where $A(kx \pm kt)$ is the wave form and A can be E ,

B , or the vector potential. Such waves simply have a static component, as does a single square wave rising from zero initial amplitude. In a spectral representation, the issue is then how the amplitude varies as $k \rightarrow 0$.

Khare and Pradhan also take a faulty limit. One sees in their Eqs. (15b) and (15c) that their q and s vary as $\sin 2\alpha$ and $\cos 2\alpha$, where α is the phase of the above standing wave, written as $\sin(Kt + \alpha)$, etc. [their Eqs. (5) and (8)]. If so, these two "gauge and Lorentz-invariant properties" depend on the choice of the zero of the arbitrary time coordinate, and are not even Galilean invariants. In the case of q , for example, this result seems to follow in Eq. (12b) from rewriting $2\sin(Kt + \alpha)\cos(Kt + \alpha)$ as $\sin(2\alpha)\cos(2Kt)$, apparently on the assumption that the time integral over the $\sin(2Kt)$ piece from $-\infty$ to $+\infty$ is symmetric and therefore identically zero. Such a "symmetry" is gauge noninvariant. The average of these two limits is indeterminant, and cannot be constrained to fall at $t=0$ regardless of the (arbitrary) choice of $t=0$! In any case, it seems obvious by inspection that the time average of $\sin(kt + \alpha)\cos(kt + \alpha)$ is zero if k is nonzero. In a separate erratum,³ the same authors still claim to obtain their original results ["... q , s , and E are all finite and identical to Eqs. 15(a) to 15(c)"] under even more general conditions.

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¹A. Khare and T. Pradhan, Phys. Rev. Lett. **49**, 1227, 1594(E) (1982).

²C. Chiu and T. Ohkawa, Phys. Rev. Lett. **48**, 837 (1982). The waves in this reference have nonzero $\vec{E} \cdot \vec{B}$ which vanishes on time averaging. They superimpose circularly polarized waves rather than linearly polarized waves as done here.

³A. Khare and T. Pradhan, Phys. Rev. Lett. **51**, 1108(E) (1983).