is given by the point at which the relative slope  $dI/dV = 1$ . On this basis the gap width for lead is  $(4.2 \pm 0.1) kT_c$ .

The experiment has been repeated with tin and indium giving entirely similar results; the gap in each case is approximately  $4kT_c$ . These results are of a preliminary nature, and experiments at lower temperatures will make them more precise.

I wish to thank C. P. Bean and J. C. Fisher

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 $<sup>1</sup>J.$  C. Fisher and I. Giaever (to be published).</sup>

W. A. Harrison (private communication) has pointed out that the tunnel current is not proportional to the density of states except in the limiting case of a low density of states.

 $3J.$  Bardeen, L. N. Cooper, and J.  $\tilde{h}$ . Schrieffer, Phys. Rev. 108, 1175 (1957).

## CRITICAL FIELD FOR SUPERCONDUCTIVITY IN NIOBIUM- TIN

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It is well known<sup>1</sup> that  $Nb<sub>3</sub>Sn$  is a superconductor with a high critical temperature, 18'K. The measurements here reported show that it has also an exceptionally high critical field, about 70 000 oersteds at 4.2'K, necessary for the suppression of all superconductivity.

The material was prepared by melting together niobium and tin in the argon arc, and the button so obtained was formed by grinding into a rod about 2 cm long and 4 mm in diameter, with rounded ends. The magnetic moment per gram,  $\sigma_{\varrho}$ , was measured by pulling the specimen from one search coil to another in a constant field, the two search coils being connected in series opposition to a ballistic galvanometer. Calibration was with nickel of high purity.

Measurements were made in increasing fields, after cooling in zero field to liquid helium temperature. Results are shown in Fig. 1. The initial points (circles) follow accurately the line for  $B=0$  ( $H=-4\pi\sigma_{\varrho}d$ , where d is the density, 8.9), and then begin to deviate at about 4000 to 5000 oersteds. The variations in the readings in fields from 5000 to 20000 oersteds reflect the wellknown irregular changes in magnetization resulting from changes in domain structure in the intermediate state, as observed by Schawlow et al.<sup>2</sup> and others. The general shape of the magnetization curve is that observed in a hard superconductor. Polishing, or annealing the specimen at 1100°C for several hours, made no essential change in the character of the curve.

When the field was decreased from its maxi-



FIG. 1. Magnetization of  $Nb<sub>3</sub>Sn$  as dependent on field strength, showing superconduction in entire specimen to about 5000 oersteds and superconduction in some parts of specimen to about 70 000 oersteds.

mum value (points marked with squares) some of the flux was frozen in, and irregularities were again observed.

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<sup>&</sup>lt;sup>1</sup>B. T. Matthias and T. H. Geballe, Phys. Rev. 95, 1435 (1954).

<sup>&</sup>lt;sup>2</sup>A. L. Schawlow, G. E. Devlin, and J. K. Hulm, Phys. Rev. 116, 626 (1959).