

Alvarado *et al.* Respond: It became clear to us that theoretical estimates of β_1 based on renormalization-group calculations^{1,2} give, directly or with the aid of well-known scaling relations, values in the range 0.81–0.88 for the Heisenberg model, 0.79–0.84 for the XY model, and 0.78–0.82 for the Ising model. So indeed the first data on β_1 for Ni(001) cannot apparently decide about Heisenberg XY or even Ising behavior. In the meantime experimental progress allowed more precise measurements of β_1 for another surface, Ni(110), with $\beta_1 = 0.77 \pm 0.02$ for $0.002 \leq |t| \leq 0.1$ and $E_K = 49$ eV.³ We point out that the bulk asymptotic regime in Ni appears to start at $|t| \lesssim 0.1$.^{4,5} It will be task of future experiments to determine β_1 for anisotropic systems and establish whether the surface behavior of such systems can be understood theoretically with the corresponding

bulk appropriate model.

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¹J. S. Reeve and A. J. Guttmann, Phys. Rev. Lett. **45**, 1581 (1980).

²H. W. Diehl and S. Dietrich, Z. Phys. B **42**, 65 (1981).

³S. F. Alvarado *et al.*, to be published.

⁴J. D. Cohen and T. R. Carver, Phys. Rev. B **15**, 5350 (1977).

⁵H. C. Binski, R. C. Reno, C. Hohenemser, R. Lyons, and C. Abeledo, Phys. Rev. B **6**, 4266 (1972).

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RESOLUTION OF THE EINSTEIN-PODOLSKY-ROSEN AND BELL PARADOXES. Itamar Pitowsky [Phys. Rev. Lett. **48**, 1299 (1982)].

On p. 1299, second column, last line, the expression should read

$$(2\pi \sin \theta)^{-1} M_{\theta}[\{x | s(x) = \frac{1}{2}\} \cap c(y, \theta)].$$

On p. 1301, first column, the unnumbered display formula should read

$$\frac{1}{n} \sum_{j=1}^n \chi_{A_j}(x) \xrightarrow{n \rightarrow \infty} \cos^2\left(\frac{\theta}{2}\right)$$

and the following line should read “ m_{θ} -almost everywhere on $c(y, \theta)$ ”.

CONFORMALLY INVARIANT QUANTIZATION OF THE LIOUVILLE THEORY. Thomas L. Curtright and Charles B. Thorn [Phys. Rev. Lett. **48**, 1309 (1982)].

Consequence (3) on p. 1311 should read as

follows:

(3) All eigenvalues of H are ≥ 0 because of (17a). Hence, among the energy eigenstates are special ones, $|E, 0\rangle$, which satisfy

$$H|E, 0\rangle = E|E, 0\rangle, \quad (19a)$$

$$L_n^{\pm}|E, 0\rangle = 0 \text{ for } n > 0. \quad (19b)$$

Otherwise, L_k 's could be used to obtain an eigenstate with negative energy. The variational argument, following Eq. (21), establishes that some of these special states have zero momentum.

EXPERIMENTAL DETECTION OF HOC^+ BY MICROWAVE SPECTROSCOPY. Christopher S. Gudeman and R. Claude Woods [Phys. Rev. Lett. **48**, 1344 (1982)].

The kilohertz digits of the HO^{13}C^+ frequency on page 1346 (column 2) were not printed. The complete frequency should be 85 752.714(15) MHz.