and post-doctoral fellowships. This work was supported by the U. S. Department of Energy under Contract No. W-7405-Eng-48.

<sup>1</sup>W. C. Stewart, Appl. Phys. Lett. <u>12</u>, 277 (1968); D. E. McCumber, J. Appl. Phys. <u>39</u>, 3113 (1968).

<sup>2</sup>B. D. Josephson, Phys. Lett. <u>1</u>, 251 (1962).

- $^{3}$ J. C. Gallop and B. W. Petley, J. Phys. E <u>9</u>, 417 (1976), have discussed this possibility.
- <sup>4</sup>K. K. Likharev and V. K. Semenov, Pis'ma Zh. Eksp. Teor. Fiz. <u>15</u>, 625 (1972) [JETP Lett. <u>15</u>, 442 (1972)].

<sup>5</sup>A. N. Vystavkin, V. N. Gubankov, L. S. Kuzmin,

K. K. Likharev, V. V. Migulin, and V. K. Semenov, Phys. Rev. Appl. 9, 79 (1974).

<sup>6</sup>M. J. Stephen, Phys. Rev. 182, 531 (1969).

<sup>7</sup>A. J. Dahm, A. Denenstein, D. N. Langenberg, W. H. Parker, D. Rogovin, and D. J. Scalapino, Phys. Rev. Lett. 22, 1416 (1969).

<sup>8</sup>H. B. Callen and T. A. Welton, Phys. Rev. <u>83</u>, 34 (1951).

<sup>9</sup>I. R. Senitzky, Phys. Rev. 119, 670 (1960).

<sup>10</sup>As  $\beta_c \rightarrow 0$ , the rolloff frequency of the noise, ~ 1/RC, increases, and the mean-square current noise available to the junction,  $\propto (1/RC)^2$ , eventually becomes so large that the noise-rounded critical current is reduced to zero. In the analytical discussion we choose  $0 < \beta_c << 1$ , while for the computer results we choose  $\beta_c \approx 0.1$ .

 $^{11}$ V. Ambegaokar and B. I. Halperin, Phys. Rev. Lett. 22, 1364 (1969).  $^{12}$ J. Kurkijärvi and V. Ambegaokar, Phys. Lett. <u>31A</u>,

<sup>12</sup>J. Kurkijärvi and V. Ambegaokar, Phys. Lett. <u>31A</u>, 314 (1970); R. F. Voss, to be published.

<sup>13</sup>Yu. M. Ivanchenko and L. A. Zil'berman, Zh. Eksp. Teor. Fiz. <u>55</u>, 2395 (1968) [Sov. Phys. JETP <u>28</u>, 272 (1969)].

<sup>14</sup>J. Kurkijärvi, in the Proceedings of the Second International Conference on SQUID's, Berlin, 6-9 May, 1980, to be published.

<sup>15</sup>A. O. Caldeira and A. J. Leggett, to be published.

## Ferroelectricity of Poly(Vinylidene Fluoride): Transition Temperature

P. Herchenröder,<sup>(a)</sup> Y. Segui,<sup>(b)</sup> D. Horne, and D. Y. Yoon *IBM Research Laboratory, San Jose, California* 95193 (Received 20 June 1980)

Temperature dependence of polarization of poly(vinylidene fluoride) (PVF<sub>2</sub>) has been measured under ac fields of 100-400 kV/cm at 50 Hz at temperatures ranging from 22 to 150 °C. It is found that PVF<sub>2</sub> is a ferroelectric with its transition temperature at ~140 °C. Above this temperature, which is about 30 °C below the melting point, no remnant polarization can be measured.

PACS numbers: 77.80.Bh

Since the discovery of large piezoelectricity and pyroelectricity in poly(vinylidene fluoride)  $(PVF_{2})$  films, there have been intense investigations on this polymer during the last ten years.<sup>1</sup> These investigations are mostly concerned with origin and characteristics of piezoelectricity and pyroelectricity of  $PVF_2$ . Recently, on the basis of ferroelectriclike hysteresis loops of polarization<sup>2,3</sup> and dipole orientation,<sup>4</sup> it has been suggested that PVF, is a ferroelectric rather than an electret as it is commonly known.<sup>1</sup> However, the transition temperature, which is an important characteristic of ferroelectric materials, has not been observed so far for PVF<sub>2</sub>. The purpose of this Letter is to present the first results on the temperature dependence of polarization of PVF<sub>2</sub>, which show the presence of a transition temperature above which remnant polarization is not observed.

PVF<sub>2</sub> samples used in these studies are capaci-

tor-grade films of 6  $\mu$ m in thickness obtained from Kureha Chemical Corporation. These films are biaxially oriented and contain both  $\alpha$ - and  $\beta$ phase crystalline regions of approximately equal amount.<sup>4</sup> These films were first annealed at 150 °C for about half an hour under slight tension in order to prevent shrinkage and structural changes that may occur at high temperatures. Film surfaces were then cleaned with acetone and aluminum electrodes were evaporated on both sides.

Electrical displacement D, electric field E, polarization P, and residual polarization  $P_r$  are related by the following well-known equations:

$$D = \epsilon_0 E + P, \tag{1}$$

$$P = \epsilon_0 \chi' E + P_r(E), \qquad (2)$$

where  $\epsilon_0$  is the free permitivity, and  $\chi'$  is the susceptibility at low field. In this paper, we are

mainly concerned with the hysteresis loops of residual polarization  $P_r(E)$  and the remnant polarization, i.e.,  $P_r(E=0)$ .

The electrical circuit used to measure the residual polarization is shown schematically in Fig. 1. It is a bridge system with the sample in one branch and adjustable impedances in the other. Through two amplifiers  $(A_s \text{ and } A_R)$  and one differential amplifier  $(A_p)$ , the net balance of the bridge can be recorded. A small field is first applied to the sample and adjustable impedances are chosen to equilibrate the bridge. These matching conditions are

$$C_R = C_S \text{ and } R_R/R_S = C_S/C_1, \tag{3}$$

where  $C_s$  and  $R_s$  represent the capacitance and resistance of the sample,  $C_R$  and  $R_R$  are adjustable capacitance and resistance, respectively, and  $C_1 (\gg C_s)$  is the measuring capacitance (see Fig. 1). Physically, this means that the  $R_{\kappa}-C_1$  branch cancels the dielectric loss, whereas the adjustable (air) capacitor  $C_R$  represents the capacitance of the sample at low field. Hence, the residual polarization  $P_{\star}(E)$  can be obtained by recording the differential output  $V_D$  as a function of the source voltage [for example, see Fig. 2(a)]. Measurements were made usually with a continuous sweep of the ac source. However, when the conductivity of the sample becomes large at high temperatures ( $\geq 80$  °C), a single pulse of the ac source was used in order to avoid the change of sample temperature due to excessive dielectric loss during measurements. Absolute calibration is carried out by equating the output  $V_R$  from air capacitor  $C_R$  to  $\epsilon' \epsilon_0 E$ , i.e.,

$$P_{r}(E) = (V_{D}/V_{R})\epsilon'\epsilon_{0}E, \qquad (4)$$



C<sub>1</sub> = 10<sup>5</sup> pF C<sub>S</sub>≤200 pF

FIG. 1. Schematic diagram of the circuit to measure polarization under ac field.

where  $\epsilon'$  is the relative permitivity at low field. The source and circuit were designed to supply voltages of maximum 1500 V (peak to peak) in the frequency range of 10 Hz to 20 kHz, and use of amplifiers enabled measurements of  $P_r$  as small as 0.5 nC/cm<sup>2</sup>. The frequency and field dependence of polarization hysteresis of PVF<sub>2</sub> will be published elsewhere.<sup>5</sup>

The residual polarization  $P_r(E)$  has been measured at temperatures ranging from 22 to  $150^{\circ}$ C. At lower temperatures, typical ferroelectric hysteresis loops such as the one shown in Fig. 2(a) are observed. At temperatures higher than ~140 °C, however, a double hysteresis loop with negligible remnant polarization  $P_r(0)$  appears [see Fig. 2(b)]. The double hysteresis loops look similar to those reported for conventional ferroelectrics, e.g., barium titanate<sup>6</sup> and ferric ammonium sulfate alum,<sup>7</sup> at temperatures above their Curie temperatures.

The temperature dependence of the remnant polarization  $P_r(0)$  measured at 50 Hz is shown in Fig. 3 for applied fields ranging from 100 to 400 kV/cm in amplitude. Remnant polarization is found to increase with temperature to reach a maximum at ~120 °C and then to fall down rapidly with further temperature increase to almost zero at ~140 °C. Beyond this temperature, no remnant polarization can be measured. Upon cooling down from 150 °C, the polarization maximum is reproduced, indicating that this temperature dependence of the polarization is reversible. Thus, these results show that PVF<sub>2</sub> undergoes a ferroelectric-to-nonferroelectric transition at ~140 °C



FIG. 2. Oscilloscope picture of residual polarization hysteresis of  $PVF_2$  under ac field of 250 kV/cm in amplitude at 50 Hz: (a) At 22 °C; (b) at 140 °C (single sweep).



FIG. 3. Temperature dependence of remnant polarization  $P_r(0)$  of PVF<sub>2</sub> under ac field of 50 Hz. The amplitude of the field is indicated for each curve in the figure.

whereby it can no longer maintain its remnant polarization. This temperature is found to be about 30 °C below its melting point.

The rapid increase of remnant polarization with temperature observed for PVF<sub>2</sub> at temperatures lower than  $\sim 120$  °C (see Fig. 3) is rather unique when compared with conventional ferroelectrics.<sup>8</sup> For these materials remnant polarization usually remains constant or decreases slowly before falling abruptly to zero at the transition temperature. It is believed that this difference may arise from the much higher coercive field of PVF<sub>2</sub> compared to other ferroelectrics, ~100 kV/cm versus ~1  $kV/cm.^2$  That is, with the electric field used in these measurements, the condition of saturation is not reached. Under these conditions, the remnant polarization is likely to be determined by the dynamics of chain rotations under the influence of oscillating electric field. As temperature increases, the rate of chain rotation (orientation) increases, possibly because of increases in kink (defect) generation and propagation.<sup>9</sup> Hence, this will result in increasing remnant polarization with temperature. However, at higher temperatures, the crystal lattice energy becomes insufficient to contain the chain motions and, hence, to freeze the orientation of chain dipoles achieved

under the applied field. Then, the remanant polarization would vanish. It seems that the temperature dependence of the remnant polarization and the mechanism of ferroelectric-to-nonferroelectric transition of  $PVF_2$  are intimately related to the molecular motions in polymer crystallites.

In summary, it is demonstrated that  $PVF_2$  is a ferroelectric with its transition temperature located at ~140 °C, beyond which no remnant polarization can be measured. This temperature is found to be about 30 °C below the melting point. The detailed mechanism of this transition with regard to molecular motions in polymer crystallites has yet to be understood.

The authors wish to thank IBM World Trade of Germany and IBM World Trade of France for granting postdoctoral fellowships to P. Herchenröder and Y. Segui, respectively. Y. Segui also thanks the Centre National de la Recherche Scientifique, for allowing his sabbatical. The technical assistance of R. D. Diller is also acknowledged.

<sup>(a)</sup>Present address: Institut für Physikalische Chemie, Universität Mainz, Mainz, West Germany.

<sup>(b)</sup>Permanent address: Laboratoire de Genie Electrique, associé au Centre National de la Recherche Scientifique, Université Paul Sabatier, F-31077 Toulouse, France.

<sup>1</sup>See, for example, Y. Wada and R. Hayakawa, Jpn. J. Appl. Phys. <u>15</u>, 2041 (1976).

<sup>2</sup>M. Tamura, K. Ogasawara, N. Ono, and S. Hagiwara, J. Appl. Phys. <u>45</u>, 3768 (1974); M. Tamura, S. Hagi-wara, S. Matsumoto, and N. Ono, J. Appl. Phys. <u>48</u>, 513 (1977).

<sup>3</sup>R. G. Kepler, ACS Org. Coatings Plastics Chem. <u>38</u>, 278 (1978); J. E. Mckinney, G. T. Davis, and M. G. Broadhurst, J. Appl. Phys. <u>51</u>, 1676 (1980).

<sup>4</sup>D. Naegele and D. Y. Yoon, Appl. Phys. Lett. <u>33</u>, 132 (1978).

<sup>5</sup>P. Herchenröder, Y. Segui, D. Naegele, and D. Y. Yoon, to be published.

<sup>6</sup>W. J. Merz, Phys. Rev. 91, 513 (1953).

<sup>7</sup>F. Jona and G. Shirane, *Ferroelectric Crystals* (Pergamon, New York, 1962), p. 333.

<sup>8</sup>See, for example, J. C. Burfoot, *Ferroelectrics* (D. van Nostrand, London, 1967).

<sup>9</sup>P. L. Taylor, T. J. Sluckin, H. Dvey-Aharon, and A. J. Lewanski, Bull. Am. Phys. Soc. <u>24</u>, 415 (1979).



FIG. 2. Oscilloscope picture of residual polarization hysteresis of  $PVF_2$  under ac field of 250 kV/cm in amplitude at 50 Hz: (a) At 22 °C; (b) at 140 °C (single sweep).