ERRATA

MAGNETIC FIELD GENERATION BY DETONATION WAVES. Michael J. Frankel and Edward T. Toton [Phys. Rev. Lett. 43, 1814 (1979)].

Equation (2) should read

$$\frac{c^2}{4\pi}\frac{\partial}{\partial x}\left(\frac{1}{\sigma(x)}\frac{\partial B}{\partial x}\right)-\frac{\partial}{\partial x}[V(x)B] = \frac{ck_B}{eN_e}\left(\frac{\partial N_e}{\partial x}\frac{\partial T}{\partial z}-\frac{\partial N_e}{\partial z}\frac{\partial T}{\partial x}\right), \quad (2)$$

and Eq. (5) should be written

$$S(x)=\frac{ck_B}{eN_e}\left(\frac{\partial N_e}{\partial x}\frac{\partial T}{\partial z}-\frac{\partial N_e}{\partial z}\frac{\partial T}{\partial x}\right), \quad (5)$$

where derivatives with respect to y in the source term have been replaced by z derivatives, in conformity with the geometry of Fig. 1. The authors thank Dr. Frank Zerilli for pointing this out.

MECHANISM FOR THE DIFFERENCE IN LIFETIMES OF CHARGED AND NEUTRAL D MESONS. Myron Bander, D. Silverman, and A. Soni [Phys. Rev. Lett. 44, 7 (1980)].

With the normalization for f_D given in Eqs. (6) and (7) our formulas for the decay rate [Eqs. (8) and (9)] are too large by a factor of 2 whereas the numerical value of $(f_D/m_u)^2$, in Eq. (12), is too small by a factor of 2. The complete set of corrections can be accomplished by replacing f_D/m_u (or f_F/m_s) by $f_D/\sqrt{2}\,m_u$ (or $f_F/\sqrt{2}\,m_s$) in Eqs. (8)–(14).

We are grateful to Bob Cahn, Yung Kang, and Mahiko Suzuki for discussions and correspondence in this regard.

NARROW Σ-HYPERNUCLEAR STATES. A. Gal and C. B. Dover [Phys. Rev. Lett. 44, 379 (1980)].

Tables I and II should be recast as follows for clarity:

TABLE I. Number $n_c{}^{\pm}$ of $1p$ nucleons for the coherent excitation, $(1p)_N \rightarrow (1p)_\Sigma$, of Σ-hypernuclear $\frac{3}{2}^-$ states in $(K^-,\pi^\pm)$ reactions, respectively, at 0° on ^{7}Li and ^{9}Be for p_K = 720 MeV/c. The conversion width Γ relative to the nuclear-matter estimate Γ_{nm} is also shown.

Target nucleus	${}^A_\Sigma Z$ structure	(I_N, I)	$n_c{}^-$	$n_c{}^+$	Γ/Γ_{nm}
^{7}Li	$\left(\{(5/9)^{1/2}\,S[2]+(4/9)^{1/2}\,D[2]\}\otimes {}^2p_\Sigma\right)_{2p_{3/2}}$	(0,1)	3/2	0	0.9
		(1,0)	1/6	0	2.7
		(1,1)	0	0	2.0
		(1,2)	1/3	1	0.7
^{9}Be	$\left(\{(8/15)^{1/2}\,{}^1S[4]-(7/15)^{1/2}\,{}^1D[4]\}\otimes {}^2p_\Sigma\right)_{2p_{3/2}}$ [a]	(0,1)	5/4	0	0.8
	$\left(\{(2/3)^{1/2\,(2S_N+1)}P[3,1]+(1/3)^{1/2\,(2S_N+1)}D[3,1]\}\otimes {}^2p_\Sigma\right)_{2p_{3/2}}$ [b]	(0,1)	3/4	0	1.2
		(1,0)	1/3[c]	0	1.9[d] 2.0[e]
		(1,1)	0	0	1.6[d] 1.3[e]
		(1,2)	2/3[c]	2	1.0[d] 0.7[e]

[a] Lower peak.

[b] Upper peak; with $S_N = 1$ for $I_N = 0$ and $S_N = 0$, 1 for $I_N = 1$.

[c] Distrbuted according to $(2S_N + 1)$ for $S_N = 0, 1$.

[d] $S_N = 0$.

[e] $S_N = 1$.

TABLE II. Number $n_c^{\pm}$ of $1p$ nucleons involved in the coherent excitation of Σ-hypernuclear 0^+ states in $(K^-,\pi^{\pm})$ reactions, respectively, at 0° on ^{16}O for p_K = 720 MeV/c, and conversion width ratio Γ/Γ_{nm}. We show two representative cases, corresponding to weak or strong Σ spin-orbit coupling.

${}^{A}_{\Sigma}Z$ structure	I	n_c^-	n_c^+	Γ/Γ_{nm}
$(_N p_{3/2}^{-1} \otimes {}_\Sigma p_{3/2})_{0^+}$	1/2	4/3	0	1.3
	3/2	8/3	4	0.6
$(_N p_{3/2}^{-1} \otimes {}_\Sigma p_{1/2})_{0^+}$	1/2	2/3	0	1.2
	3/2	4/3	2	0.8
$^1S_0$	1/2	2	0	1.4
	3/2	4	6	0.3
$^3P_0$	1/2	0	0	1.1
	3/2	0	0	1.1

DYNAMIC CONFINEMENT FROM VELOCITY-DEPENDENT INTERACTIONS. M. King and F. Rohrlich [Phys. Rev. Lett. 44, 621 (1980)].

On page 622 in the top equation of the second column $(\vec{\xi}-\vec{\pi})^2$ should read $(\vec{\xi}\cdot\vec{\pi})^2$. On p. 624, Eq. (9b) $l-\frac{5}{2}$ should read $l+\frac{5}{2}$. In the equation preceding (9b) the argument of u_l should be $(\beta\xi)^2$ which corresponds to $t=(\beta\xi)^2$. In the last line of the left column on p. 624 "has been earlier" should read "has been noted earlier."

PHASE TRANSITIONS AND MAGNETIC MONOPOLE PRODUCTION IN THE VERY EARLY UNIVERSE. Alan H. Guth and S.-H. H. Tye [Phys. Rev. Lett. 44, 631 (1980)].

In the first paragraph, the time corresponding to $T=10^{17}$ GeV should be $t\sim 10^{-41}$ sec. Equation (2) should read

$$p(t)=\exp\left[-\frac{4\pi}{3}\int_0^t dt_1 R^3(t_1)\lambda(t_1)\left(\int_{t_1}^t dt_2\frac{v}{R(t_2)}\right)^3\right].$$

Footnote 12 should read the following: If the density of monopoles exceeds this bound, they would noticeably influence the cosmological deceleration parameter. See Ref. 7.