ERRATA

MAGNETIC FIELD GENERATION BY DETONATION WAVES. Michael J. Frankel and Edward T. Toton [Phys. Rev. Lett. 43, 1814 (1979)].

Equation (2) should read

$$\frac{c^2}{4\pi}\frac{\partial}{\partial x}\left(\frac{1}{\sigma(x)}\frac{\partial B}{\partial x}\right)-\frac{\partial}{\partial x}[V(x)B] = \frac{ck_B}{eN_e}\left(\frac{\partial N_e}{\partial x}\frac{\partial T}{\partial z}-\frac{\partial N_e}{\partial z}\frac{\partial T}{\partial x}\right), \qquad (2)$$

and Eq. (5) should be written

$$S(x)=\frac{ck_B}{eN_e}\left(\frac{\partial N_e}{\partial x}\frac{\partial T}{\partial z}-\frac{\partial N_e}{\partial z}\frac{\partial T}{\partial x}\right), \qquad (5)$$

where derivatives with respect to y in the source term have been replaced by z derivatives, in conformity with the geometry of Fig. 1. The authors thank Dr. Frank Zerilli for pointing this out.

MECHANISM FOR THE DIFFERENCE IN LIFETIMES OF CHARGED AND NEUTRAL D MESONS. Myron Bander, D. Silverman, and A. Soni [Phys. Rev. Lett. 44, 7 (1980)].

With the normalization for f_D given in Eqs. (6) and (7) our formulas for the decay rate [Eqs. (8) and (9)] are too large by a factor of 2 whereas the numerical value of $(f_D/m_u)^2$, in Eq. (12), is too small by a factor of 2. The complete set of corrections can be accomplished by replacing f_D/m_u (or f_F/m_s) by $f_D/\sqrt{2}\,m_u$ (or $f_F/\sqrt{2}\,m_s$) in Eqs. (8)–(14).

We are grateful to Bob Cahn, Yung Kang, and Mahiko Suzuki for discussions and correspondence in this regard.

NARROW Σ-HYPERNUCLEAR STATES. A. Gal and C. B. Dover [Phys. Rev. Lett. 44, 379 (1980)].

Tables I and II should be recast as follows for clarity:

TABLE I. Number $n_c^{\pm}$ of $1p$ nucleons for the coherent excitation, $(1p)_N \to (1p)_\Sigma$, of Σ-hypernuclear $\frac{3}{2}^-$ states in $(K^-,\pi^\pm)$ reactions, respectively, at 0° on ^{7}Li and ^{9}Be for p_K = 720 MeV/c. The conversion width Γ relative to the nuclear-matter estimate Γ_{nm} is also shown.

Target nucleus	$^A_\Sigma Z$ structure	(I_N, I)	n_c^-	n_c^+	Γ/Γ_{nm}
^{7}Li	$\left(\{(5/9)^{1/2}S[2]+(4/9)^{1/2}D[2]\}\otimes {}^2p_\Sigma\right)_{^2p_{3/2}}$	(0,1)	3/2	0	0.9
		(1,0)	1/6	0	2.7
		(1,1)	0	0	2.0
		(1,2)	1/3	1	0.7
^{9}Be	$\left(\{(8/15)^{1/2}\,{}^1S[4]-(7/15)^{1/2}\,{}^1D[4]\}\otimes {}^2p_\Sigma\right)_{^2p_{3/2}}$ [a]	(0,1)	5/4	0	0.8
	$\left(\{(2/3)^{1/2}\,{}^{(2S_N+1)}P[3,1]+(1/3)^{1/2}\,{}^{(2S_N+1)}D[3,1]\}\otimes {}^2p_\Sigma\right)_{^2p_{3/2}}$ [b]	(0,1)	3/4	0	1.2
		(1,0)	1/3[c]	0	1.9[d] 2.0[e]
		(1,1)	0	0	1.6[d] 1.3[e]
		(1,2)	2/3[c]	2	1.0[d] 0.7[e]

[a] Lower peak.
[b] Upper peak; with $S_N = 1$ for $I_N = 0$ and $S_N = 0, 1$ for $I_N = 1$.
[c] Distrbuted according to $(2S_N+1)$ for $S_N = 0, 1$.
[d] $S_N = 0$.
[e] $S_N = 1$.