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¹⁴We have simplified our algebra by eliminating certain physical constants. To gain a formalism corresponding to the Hamiltonian $H = -2J \sum_i [S_i^z S_{i+1}^z - \frac{1}{4} + \gamma(S_i^x S_{i+1}^x + S_i^y S_{i+1}^y)] - g\mu_B H_M \sum_i S_i^z$ with k , Boltzmann's constant, not equal to 1, make the transform $\Delta \rightarrow \gamma^{-1}$, $T \rightarrow k_B T / 2\gamma J$, and $H_0 \rightarrow g\mu_B H_M / 2\gamma J$ in all our formulas.

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²⁰For larger H_0 and $\Delta \geq 1$, except for an $O(T^0)$ neighborhood of $H_0 = 0$, $\Delta = 1$, our answer is Eq. (4) with the square-root term replaced by $-H_0/2$. E in this case is exponentially higher order than the integral. All results for $H_0 > O(T^2)$ have no structure and merge smoothly onto the $H_0 < O(T)$ results. Therefore, in all that follows we restrict $H_0 < O(T)$.

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NEW APPROACH TO PERTURBATION THEORY. Y. Aharonov and C. K. Au [Phys. Rev. Lett. **42**, 1582 (1979), and **43**, 176(E) (1979)].

Equations (12) and (40) in this paper should read as follows:

$$-\frac{1}{2}(g^2 - g')e^{-G} = (E - V_0 - \lambda V_1)e^{-G}; \quad (12)$$

$$F_i(x) \equiv 2\alpha_{i-1} [g_1 + g_0 \alpha_1 / (x - \alpha_0)] \exp(-2G_0) \\ + \sum_{m=2}^{i-1} \alpha_{i-m} \left[\sum_{j=0}^m g_j g_{m-j} - g_m' + 2E_m \right] (x - \alpha_0) \exp(-2G_0) - \sum_{j=1}^{i-1} g_j g_{i-j} (x - \alpha_0)^2 \exp(-2G_0). \quad (40)$$