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Measurement of High-Energy Direct Photons in ψ Decays

G. S. Abrams, M. S. Alam, C. A. Blocker, A. M. Boyarski, M. Breidenbach, D. L. Burke, W. C. Carithers, W. Chinowsky, M. W. Coles, S. Cooper, W. E. Dieterle, J. B. Dillon, J. Dorenbosch, J. M. Dorfan, M. W. Eaton, G. J. Feldman, M. E. B. Franklin, G. Gidal, G. Goldhaber, G. Hanson, K. G. Hayes, T. Himel, D. G. Hitlin,^(a) R. J. Hollebeek, W. R. Innes, J. A. Jaros, P. Jenni, A. D. Johnson, J. A. Kadyk, A. J. Lankford, R. R. Larsen, V. Lüth, R. E. Millikan, M. E. Nelson, C. Y. Pang, J. F. Patrick, M. L. Perl, B. Richter, A. Roussarie, D. L. Scharre, R. H. Schindler, R. F. Schwitters,^(b) J. L. Siegrist, J. Strait, H. Taureg, M. Tonutti,^(c) G. H. Trilling, E. N. Vella, R. A. Vidal, I. Videau, J. M. Weiss, and H. Zaccane^(d)

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305, and Lawrence Berkeley Laboratory and Department of Physics, University of California, Berkeley, California 94720

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The inclusive γ and π^0 momentum distributions at the ψ have been measured. Using these data and estimates of η production, it is found that $(4.1 \pm 0.8)\%$ of ψ decays contain a direct photon with energy greater than 60% of the beam energy. The expected momentum distribution for direct photons calculated to lowest order in quantum chromodynamics is qualitatively different from that observed in the data.

First-order quantum chromodynamics (QCD) calculations predict that a significant fraction of the hadronic decays of heavy quark-antiquark 3S_1 resonances (such as the ψ) result in the production of direct γ 's (i.e., γ 's not coming from secondary decays of π^0 's or η 's).¹ We have measured the inclusive γ and π^0 momentum distributions at the ψ , and have made estimates of the η momentum distribution from the data. We ob-

serve γ production in excess of the expected contributions from π^0 and η decay.

The data were collected with the Mark II magnetic detector at the SLAC e^+e^- storage ring facility SPEAR at energies near the peak of the $\psi(3095)$ resonance. The detector has been described in detail elsewhere,² and only a brief description of the particle detection will be presented in this Letter.

Charged tracks are reconstructed from hits in the sixteen cylindrical drift-chamber layers which provide solid-angle coverage over 85% of 4π sr. The trigger requires two or more charged tracks, one of which must be within the central region of the drift chamber which covers 67% of 4π sr.

Photons are detected primarily in the eight lead-liquid argon (LA) shower counter modules which surround the solenoid, covering approximately 64% of 4π sr. The rms energy resolution for detected photons is given approximately by $\delta E/E = (0.12 \text{ GeV}^{1/2})/E^{1/2}$. The γ detection efficiency (including geometry) is approximately 30% at 0.2 GeV, 50% at 0.3 GeV, and 55% above 0.4 GeV.³

Neutral pions are reconstructed by combining pairs of γ 's each of which is required to have momentum greater than 150 MeV. Pairs with invariant mass between 0.075 and 0.200 GeV are considered to be π^0 candidates. The π^0 signal is extracted after subtraction of a background whose shape is a function of momentum. The background shape is obtained by combining real photons and "pseudophotons" from the same event. Pseudophotons are created in the analysis program by pretending that charged particles are π^0 's and allowing them to decay into pseudophotons.⁴ The background distribution is normalized to the data in the mass region between 0.3 and 1.0 GeV. Finally, a correction is made for tails of the π^0

which fall outside the specified mass cuts. The π^0 detection efficiency is 6% at 0.4 GeV, 18% at 0.6 GeV, and 30% at 1.0 GeV.

We have measured the inclusive γ and π^0 momentum distributions at the ψ using a sample of 280 000 observed ≥ 2 -prong hadron events. The trigger efficiency is measured by using a sample of events taken near the peak of the ψ '(3684). A sample of ψ events from the decay $\psi' - \psi\pi^+\pi^-$ is obtained by requiring the missing mass from observed pairs of oppositely charged pions to be consistent with the mass of the ψ , m_ψ .⁵ This sample of ψ events is identified purely from the π^+ and π^- and has no trigger bias arising from the ψ decay. As a function of γ or π^0 momentum, the trigger efficiency is calculated as the fraction of events which satisfy the trigger requirement after elimination of the recoiling π^+ and π^- from the event. The resulting trigger efficiency as a function of γ momentum p varies from 70% at $x = 0.1$ to approximately 40% near $x = 1.0$, where $x = 2p/m_\psi$. The trigger efficiency versus π^0 momentum varies from 70% at $x = 0.3$ to approximately 50% near $x = 1.0$. From the average trigger efficiency (based on the entire event sample), the number of produced ψ events corresponding to the 280 000 observed hadron events is calculated to be 435 000.

Figure 1 shows the inclusive γ momentum distribution, $(1/N_{\text{tot}})dN/dx$, as a function of x , where N_{tot} is the total number of produced ψ 's. The error bars represent the statistical errors only. Overall systematic errors are approximately $\pm 20\%$, roughly half of which are independent of x . Corrections for detection and trigger efficiency

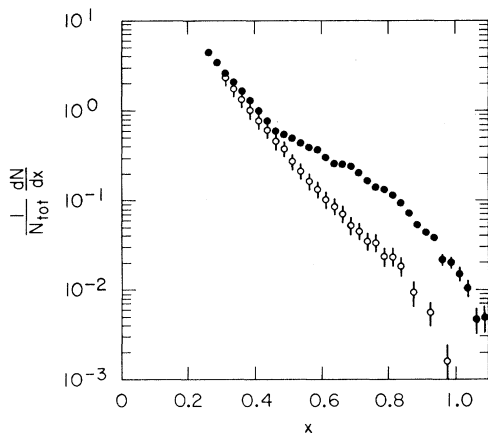


FIG. 1. Solid points: the inclusive γ momentum distribution $(1/N_{\text{tot}})dN/dx$. Open points: the expected γ momentum distribution based on measurements of π^0 and η production. The error bars on the measured γ distribution reflect only the statistical errors, whereas the error bars on the expected distribution include a $\pm 20\%$ systematic error.

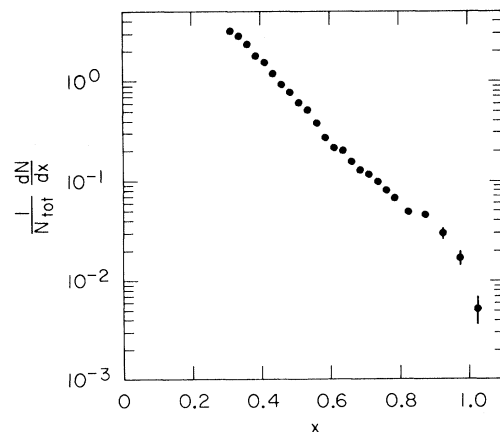


FIG. 2. Inclusive π^0 momentum distribution $(1/N_{\text{tot}}) \times dN/dx$, where $x = 2p/m_\psi$. Systematic errors of $\pm 30\%$ are not shown in the figure.

have been made. Background due to the e^+e^- final state in which one of the electrons radiates a γ is eliminated. Radiation in the initial state is negligible. The inclusive π^0 scaled momentum distribution is shown in Fig. 2. Systematic errors of $\pm 30\%$, much of which may be x dependent, are not shown in the figure.

From the measured π^0 distribution, we have made predictions for the expected γ momentum distribution. In addition to the contribution from π^0 decays, there is an additional contribution from η decays.⁶ To determine the η population, we have made least-squares fits to the background-subtracted $\gamma\gamma$ invariant-mass distributions in different momentum bins. The functional form used in the fits consists of a Gaussian, fixed at the η mass with width as determined by Monte Carlo calculation, over a linear background. We have measured the ratio

$$R(p) = B(\psi \rightarrow \eta + X)B(\eta \rightarrow \gamma\gamma) / B(\psi \rightarrow \pi^0 + X)$$

as a function of momentum. We find $R(p)$ to be less than 0.10 for all momentum bins except for $p > 1.2$ GeV, where we measure $R = 0.16 \pm 0.06$. From the measured value of R as a function of momentum, we estimate the contribution to γ production from η 's.

The γ momentum distribution predicted from the measured π^0 and η momentum distributions is also shown in Fig. 1. The error bars include possible systematic errors ($\pm 22\%$) which affect the measured and predicted γ distributions differently, but do not include correlated errors. We observe an excess of produced γ 's over the prediction for $x \geq 0.5$.

Figure 3 shows the direct γ momentum distribution calculated by subtracting the expected distribution from π^0 's and η 's from the measured

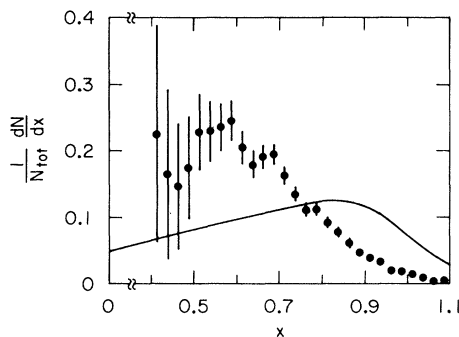


FIG. 3. Direct γ momentum distribution. Curve is the leading-order QCD prediction convoluted with the energy resolution.

distribution. The error bars include the statistical error and the systematic error in the difference between the measured and predicted distributions. The errors become small at large x because the π^0 contribution is small. For $x < 0.4$ the errors become too large to provide meaningful information, and only for $x > 0.5$ can a clear excess be demonstrated. In addition to the displayed error bars, there is a $\pm 17\%$ systematic error on the distribution due to uncertainties in the γ detection efficiency, the number of produced ψ events, the trigger efficiency, and the angular distribution of the direct γ 's as discussed below.

Figure 4 shows the distribution of observed photons with $x > 0.6$ as a function of $|\cos\theta|$, where θ is the polar angle of the γ with respect to the beam direction. The angular distribution has been fitted by the form $1 + \alpha_\gamma \cos^2\theta$ and gives $\alpha_\gamma = 0.14 \pm 0.12$. Approximately 25% of these events are background from π^0 or η decays. Analysis of the angular distribution of photons from observed π^0 decays shows this background to be consistent with an isotropic distribution. Correcting for this background gives $\alpha_\gamma = 0.18 \pm 0.18$ for the direct γ contribution. An overall correction of 3% is made to the data (which has been included in Fig. 3) to correct for the deviation of the observed angular distribution from isotropy. Integrating the direct γ momentum distribution from $x = 0.6$ to 1.0, we obtain an inclusive rate for direct photon production of $(4.1 \pm 0.8)\%$. This integrated cross section includes a correction (-6%) for feeddown from lower x due to the resolution of the LA.

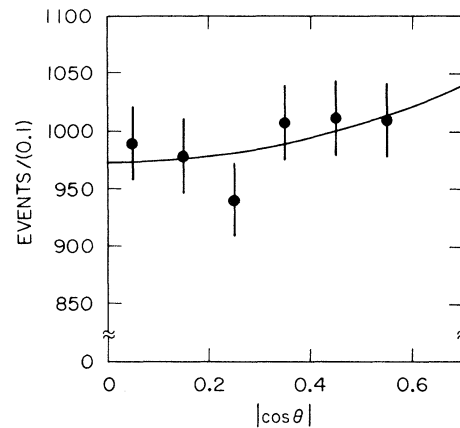


FIG. 4. Angular distribution of observed photons with $x > 0.6$ as a function of $|\cos\theta|$. Curve represents the best fit by the form $1 + \alpha_\gamma \cos^2\theta$.

As a check on possible systematic errors and trigger bias problems, this analysis has been repeated with two other independent methods. The first method makes use of photons which convert in the 0.06 radiation length of material preceding the drift chamber. The photons are reconstructed from measurement of the momenta of the resulting e^+ and e^- in the drift chamber. Neutral pions are reconstructed from combinations of a converted γ and a γ observed in the LA. As essentially all γ 's which are detected through pair conversion satisfy the trigger requirement, no trigger bias is introduced in this method. The data are corrected for the γ and π^0 detection efficiencies, which are calculated independently of those used in the LA. The second method makes use of the statistically independent sample of ψ events from the decay $\psi' \rightarrow \psi\pi^+\pi^-$. As described above, this sample is free of trigger bias. The resulting γ , π^0 , and direct γ distributions measured with these alternative methods of analysis are consistent with the results of the analysis described previously.

First-order QCD calculations¹ predict significant direct photon production at the ψ . The hadronic decay of the ψ must proceed via an intermediate state consisting of at least three color-octet gluons. One expects the dominant contribution to direct γ production to arise by replacing one of the outgoing gluon lines by a γ . From lowest-order perturbation theory, we have

$$B_\gamma = \Gamma(\psi \rightarrow \gamma gg) / \Gamma(\psi \rightarrow ggg) = (\alpha/\alpha_s) C (e_Q/e)^2,$$

where $C = 36/5$ is a color SU(3) factor, e_Q is the charge of the charmed quark, and α_s is the color fine-structure constant. For $\alpha_s = 0.18$,⁷ one calculates $B_\gamma = 13\%$ for all x . Equating the direct decay of the ψ into hadrons with the 3-gluon decay, and correcting for second-order electromagnetic decays, the lowest order QCD prediction for the branching ratio for direct γ decays from the ψ is 8%. The γ momentum spectrum is predicted to be roughly proportional to x and peak near $x = 1.0$. The lowest-order QCD prediction (convoluted with the γ energy resolution) is shown compared with the data in Fig. 3. Although the total rate of direct γ production approximately agrees with theoretical expectations (one expects 5% for $x > 0.6$), and the mean value of the coefficient of the $\cos^2\theta$ term in the angular distribution agrees with theory (one expects $\alpha_\gamma \approx 0.3$ for $x > 0.6$), the shape of the observed distribution is inconsistent with the prediction. However, it is expected that the observed distribu-

tion should be softer than the leading-order prediction since radiative effects and the mass of the final-state hadrons are not considered, but no calculation has been made. One calculation for the decay width of a heavy quark-antiquark bound state indicates that higher-order effects are of the same magnitude as lower-order terms,⁸ thus making it difficult to provide accurate theoretical predictions.

In conclusion, we have measured direct photon production at the ψ at large x . We find that $(4.1 \pm 0.8)\%$ of ψ decays produce a direct photon with $x > 0.6$.⁹ The observed rate is consistent with lowest-order QCD predictions, but the observed distribution is softer.

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^(a)Present address: California Institute of Technology, Pasadena, Cal. 91125.

^(b)Present address: Harvard University, Cambridge, Mass. 02138.

^(c)Present address: Universität Bonn, D-53 Bonn, Federal Republic of Germany.

^(d)Present address: Centre d'Etudes Nucléaires de Saclay, F-91190 Gif-sur-Yvette, France.

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³These efficiencies were obtained by Monte Carlo simulation of the electromagnetic shower development [see R. L. Ford and W. R. Nelson, SLAC Report No. 210, 1978 (unpublished)] in the Mark II detector assuming isotropic production of γ 's. These efficiencies agree with measured efficiencies obtained by studying the decays $\psi \rightarrow \pi^+\pi^-\pi^0$ and $\psi \rightarrow \pi^+\pi^-\pi^+\pi^-\pi^0$.

⁴It is assumed that the charged-particle and π^0 momentum distributions are similar. Resolution and efficiency effects are included in the generation of pseudophotons in order to simulate the observed photon spectrum.

⁵The background from accidental combinations falling in the ψ mass region is estimated by use of events in bands on either side of the peak and subtracted.

⁶The contribution from radiative decays of other secondary hadrons is estimated to be negligible compared with the overall uncertainty in the π^0 contribution. For instance, η' production equal to η production would change the estimated direct γ rate for $x > 0.6$ by

less than 3%.

⁷This value of α_s is calculated from the ratio of the leptonic to the hadronic width of the ψ [see T. Appelquist and H. D. Politzer, Phys. Rev. Lett. **34**, 43 (1975)]. The predicted direct γ contribution is included as part of the total width of the ψ used in the

calculation.

⁸R. Barbieri *et al.*, Nucl. Phys. **B154**, 535 (1979).

⁹Another experiment [M. T. Ronan *et al.*, Lawrence Berkeley Laboratory Report No. LBL-9256, 1979 (unpublished)] finds evidence of production of direct γ 's consistent in magnitude with this experiment.

Correlations between High-Momentum Mesons in $p + p \rightarrow \pi + \pi + X$ at $s^{1/2} = 63$ GeV

G. J. Bobbink, M. A. van Driel, F. C. Ern e, W. G. J. Langeveld, P. M. Kooijman,
J. C. Sens, and J. Timmer

National Institute for Nuclear and High Energy Physics, Amsterdam, The Netherlands

and

D. Favart and P. Leleux

Universit  Catholique de Louvain, B-1348 Louvain-la-Neuve, Belgium

and

M. M. Block, R. Campanini, and H. W. Ludwig

Northwestern University, Evanston, Illinois 60201

and

M. A. Botje and D. van Wezep

State University, Utrecht, The Netherlands

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In pp collisions at $\sqrt{s} = 62.3$ GeV where *each* proton fragments into at least one low- p_T , high- x pion, no significant correlations for $\pi^+\pi^+$, $\pi^+\pi^-$, and $\pi^-\pi^-$ are found, thus excluding quark exchange with Brodsky-Gunion counting rules as the dominant interaction mechanism. This, together with the experimental x dependence of single-particle spectra, suggests the introduction of a new counting rule, which is discussed.

We report a measurement of the process

$$p_1 + p_2 \rightarrow \pi_1 + \pi_2 + X$$

at the CERN intersecting storage rings (ISR) at $\sqrt{s} = 62.3$ GeV, in which π_1 (π_2) is emitted at small angle ($18 < \theta < 40$ mrad) with respect to p_1 (p_2). In this experiment we measure the invariant double-differential cross section

$$\sigma_{\text{INV}}(\pi_1, \pi_2) \equiv E_1 E_2 d^6\sigma / d^3p_1 d^3p_2, \quad (1)$$

for the combinations $\pi^+\pi^+$, $\pi^+\pi^-$, and $\pi^-\pi^-$, and the correlation coefficient R , defined as

$$R = \sigma_{\text{TOT}} \frac{\sigma_{\text{INV}}(\pi_1, \pi_2)}{\sigma_{\text{INV}}(\pi_1)\sigma_{\text{INV}}(\pi_2)} = \sigma_{\text{TOT}} \frac{N(\pi_1, \pi_2)}{N(\pi_1)N(\pi_2)}, \quad (2)$$

where $N(\pi_1, \pi_2)$ and $N(\pi_i)$ are the normalized double- and single-pion yields. R is a function of the variables x_1 , p_{T1} , x_2 , p_{T2} , φ_1 , φ_2 , and \sqrt{s} . The data cover the ranges $0.2 < x_1, x_2 < 0.95$, $0.2 < p_{T1}, p_{T2} < 1.2$ GeV/ c , and $\sqrt{s} = 62.3$ GeV. Assuming factorization in p_T and no φ dependence we have R

$= R(x_1, x_2)$ for data at fixed \sqrt{s} for each of the three combinations of charges.

The equipment consists of two small-angle spectrometers located on the downstream ends of one intersection of the CERN ISR; see Fig. 1. Each arm consists of a septum magnet (SM), trigger counters (S), five modules of drift chambers (DC), two hodoscopes (H1, H2), and three atmospheric  erenkov counters ( ).

The acceptance per arm is ~ 1 msr. The momentum resolution $dp/p \lesssim 1\%$. The first and the last  erenkov counters in each arm were filled with CO₂, the middle one with 50% N₂/50% He. The threshold momenta for $\pi/K/p$ in CO₂ and 50% N₂/50% He are 4.6/16.2/30.9 and 7.5/26.5/50.1 GeV/ c , respectively, resulting in unambiguous identification of pions in the range $6 \lesssim p \lesssim 26$ GeV/ c . In the range $26 < p < 30$ GeV/ c the contamination of the π^+ sample by protons and K^+ is calculated to be $< 2\%$ and $< 15\%$, respectively; the K^- and \bar{p} contamination in the π^- sample is neg-