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Band-Gap Assignment in SnO₂ by Two-Photon Spectroscopy

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We present for the first time two-photon measurements on rutile $SnO₂$. The detaile polarization analysis yields Γ_3^* symmetry for the direct forbidden 1S exciton and the upper valence band. Our results clear up the long-standing controversy concerning the symmetry of the upper valence band.

As proposed by Loudon,¹ the direct forbidde transitions are specially suited to be studied by two-photon spectroscopy. Loudon calculated the two-photon absorption for $Cu₂O$, which was later studied by Pradère and co-workers.² Because of experimental difficulties these authors found only a weak structure, which they attribute to the 1S exciton. They were not able to measure the polarization dependence. We present for the first time a detailed polarization analysis of such a direct forbidden transition. Using this technique we are able to resolve the long-standing puzzle on the band-gap assignment in SnO₂ (D_{4h}^{14}) .

Though the early indication of SnO, being an indirect material^{3,4} could no longer be maintained after the excellent work of Nagasawa and Shion-'oya,⁵ the symmetry assignment of the direct gap remained as a controversy. Nagasawa and Shionoya were able to show that the band gap is direct forbidden at $\overline{K} = 0$. They offer three possibilities for the highest valence band $(\Gamma_2^+, \Gamma_3^+, \Gamma_4^+)$ assuming a Γ_1^{+} lowest conduction band, which is agreed on in later experimental^{6,7} as well as theoretical investigations.^{8,9} Agekyan¹⁰ analyzed the quadrupole transition to the 1S exciton and determined Γ_3^* symmetry for the top valence band. There are two band-structure calculations^{8,9} published which contradict each other as well as the assignment of Agekyan. Arlinghaus⁸ assigns the top valence band to be of Γ_1^+ symmetry where as Jacquemin⁹ comes up with a Γ_5^+ valence band. Jacquemin states that his assignment is compat-
ible with the experimental results of Agekyan.¹⁰ ible with the experimental results of Agekyan.

^A third band-structure calculation was very re-A third band-structure calculation was very re
cently presented by Robertson,¹¹ who calculate the highest valence band to be of Γ_3^* symmetry. In order to clear up this controversy, we present two-photon data on the lowest exciton. A complete polarization analysis of this transition allows an unambiguous assignment of the upper valence band to Γ_3^+ symmetry

Our basic experimental setup is described in Our basic experimental setup is described in detail by Fröhlich and Sondergeld.¹² To improv

FIG. 1. Two-photon spectra of 1S exciton in rutile SnO₂ at 4.5 K (resolution 0.15 meV). \parallel configuration $(0^{\circ}/0^{\circ})$, both polarization vectors $\bar{\epsilon}_1$ and $\bar{\epsilon}_2$ parallel to [100]; \perp configuration (0°/90°), $\bar{\epsilon}_1$ in [100] and $\bar{\epsilon}_2$ in [010] directions.

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FIG. 2. Two-photon spectra of 1S exciton in rutile SnO₂ at 4.5 K (resolution 5 meV). \parallel configuration (45°/ 45°), both polarization vectors $\vec{\epsilon}_1$ and $\vec{\epsilon}_2$ parallel to [110]; \perp configuration (45°/135°), $\vec{\epsilon}_1$ in [110] and $\vec{\epsilon}_2$ in [110] directions.

the signal-to-noise ratio and the spectral resolution we used for most of the measurements a stilben dye laser (0.15 meV resolution) pumped by a krypton laser instead of the standard xenon flashlamp plus monochromator (5 meV resolution). As a high-power laser we used a Raman-shifted (H, gas at 40 bars) Nd-doped yttrium-aluminumgarnet laser (0.6496 eV). The crystals¹³ were of high quality as indicated by the narrow linewidth of the 1S exciton.

In Figs. 1 and ² we present two examples of experimental spectra. All data presented here were taken with \bar{K} vectors parallel to the tetragonal axis. The high-resolution (0.15 meV) results of Fig. 1 were gained with the dye-laser setup whereas for the measurements in Fig. 2 we used the standard flash-lamp setup (5 meV resolution). From the high-resolution data of Fig. 1 we deduce for the position of the 1S exciton an energy of $(3.5630 \pm 3) \times 10^{-4}$ eV, which is in agreement with for the position of the 1S exciton an energy of $(3.5630 \pm 3) \times 10^{-4}$ eV, which is in agreement vone-photon data.^{5,10} The polarization analysi using the standard formalism¹⁴ allows an unambiguous assignment: From Fig. 1 we deduce the two possibilities of $\Gamma_1^{\;\;+}$ and $\Gamma_3^{\;\;+}$ as final state. The result of Fig. 2 excludes Γ_1^+ , thus leading The result of Fig. 2 excludes Γ_1 , thus leading
to the assignment of Γ_3^+ symmetry to the lowes exciton. Since the lowest conduction band is of ' ${\Gamma_1}^+$ symmetry we deduce for the upper valence

band Γ_3^* symmetry. The rather small absorp tion in the perpendicular configuration $(0^{\circ}/90^{\circ})$ in Fig. 1 and in the parallel configuration $(45\degree/45\degree)$ in Fig. 2 can be explained by the depolarization of the crystal. Additional measurements with circular polarized light confirm the above assignment.

Other compounds with a direct forbidden band gap like GeO₂, TiO₂, and Cu₂O are good candidates to be studied with this technique. The polarization analysis should help to clear up still existing uncertainties concerning the symmetry assignments not only of the uppermost valence band and lowest conduction band but of higherenergy transitions as well. Using this technique one should, e.g., be able to locate the second conduction band in $SnO₂$, which is supposedly of Γ_4 ⁺ symmetry.

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