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<sup>3</sup>K. Neubeck *et al.*, *Phys. Rev. C* **10**, 320 (1974).

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<sup>9</sup>There is at present a sign ambiguity in  $\langle E1 \rangle$  since

the lifetime determines only  $\langle E1 \rangle^2$  and the theoretical sign is not believed to be reliable. Until this is resolved one can only infer the magnitude of  $\langle H_{\text{PNC}} \rangle$  from a measurement of  $P_\gamma$ .

<sup>10</sup>Calculated by Millener (Ref. 8). This is consistent with two experimental results—H. Grawe *et al.*, *Z. Phys.* **A280**, 271 (1977), and E. K. Warburton, J. W. Olness, and C. J. Lister, private communication.

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## Analyzing Power in Inclusive Proton-Nucleus Cross Sections

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This paper reports measurements of the analyzing power,  $A_y$ , in the production of both "backward" protons and forward (quasifree scattering) protons in the reaction  $p + A \rightarrow p + X$ , using 800-MeV polarized protons. For the backward protons the measurements show large negative  $A$ -dependent values of  $A_y$  at low momenta, changing to large positive values at high momenta; in the quasielastic region,  $A_y$  is large and positive, is smaller than  $A_y$  measured in hydrogen, and decreases with increasing  $A$ .

The measurement of inclusive cross sections in the reaction  $p + A \rightarrow p + X$ , in kinematic regions forbidden in  $p$ - $p$  reactions on free stationary protons, provides data that are sensitive to the high-momentum components in nuclei. Some of the general aspects of such studies appear in the recent literature<sup>1-7</sup> and several models<sup>1,3,8-10</sup> have been offered that attempt to account for the backward and, in particular, 180° production of high-momentum protons<sup>11</sup> and that differ greatly both in their assumed mechanisms and in their physical models of the nucleus.

Frankel and Woloshyn<sup>12</sup> (FW) have recently

pointed out how measurements of the analyzing power,<sup>13</sup>  $A_y$ , for the reaction  $p_{\text{polarized}} + A \rightarrow p + X$  are sensitive to and can distinguish between such models of high-momentum behavior within nuclei. For example, in the single-scattering model<sup>1,2</sup> shown in Fig. 1(a) the incoming proton of momentum  $\vec{p}$  lifts a target nucleon of virtual momentum  $\vec{k}$  on to the mass shell with observed momentum  $\vec{q}$ . FW estimate the analyzing power from the measured analyzing power<sup>14</sup> for  $p$ - $p$  interactions using the values of  $s$  and  $t$  appropriate to interactions with a bound nucleon of momentum  $\vec{k}$ . For backwardly detected protons

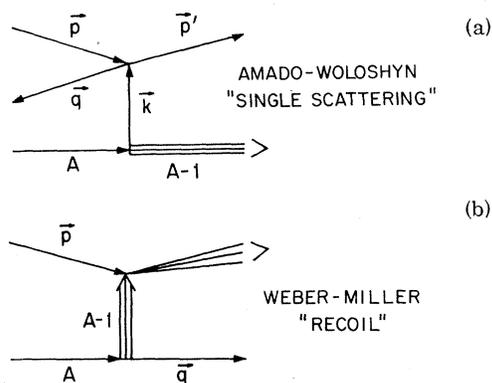


FIG. 1. (a) The single-scattering mechanism without final-state interactions; the struck nucleon is off shell. (b) The recoil mechanism without final-state interactions; the recoiling nucleus is off shell.

these analyzing powers are generally large and negative. In another example FW point out that the model of Weber and Miller<sup>8</sup> shown in Fig. 1(b), in which the incoming proton interacts with arbitrary recoiling off-shell fragments of the nucleus, gives zero analyzing power.

This experiment was undertaken expressly to obtain the first measurement<sup>15</sup> of  $A_y$  for backwardly emitted protons of high momenta and in part to test these predictions. It has also obtained data in the quasifree forward region. It utilized the polarized 800-MeV proton beam and the High-Resolution Spectrometer Facility (HRSF) at the Clinton P. Anderson Meson Physics Facility.

Backward data were accumulated at  $120^\circ$ ,  $101^\circ$ ,  $90^\circ$ , and  $75^\circ$  (lab) where large negative values of  $A_y$  were predicted by FW. Since rates fall off rapidly with increasing angle most of the data were taken at the smaller angles. Forward data were taken at  $19.3^\circ$ .

Time-of-flight measurements and  $dE/dx$  measurements allowed complete and background-free separation of protons, deuterons, and tritons.<sup>16</sup> The momentum acceptance of the spectrometer<sup>17</sup> was 2.8% (full width at half-maximum). The experiment was monitored on-line by sampling a large fraction of the events recorded on tape, the results at  $90^\circ$  suggesting measurements at  $75^\circ$ . The beam polarization was reversed automatically every two minutes and its polarization was monitored<sup>18</sup> continuously with a hydrogen polarimeter to correct each set of data. Although the polarimeter and the HRS have been well tested<sup>18</sup> we have as a further check re-

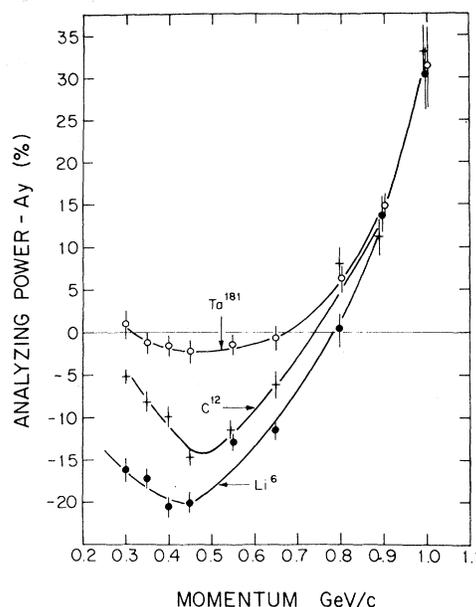


FIG. 2. Analyzing power (%) vs momentum (GeV/c) of detected proton for incident polarized 800-MeV protons; laboratory angle of  $75^\circ$ . Measurements were made at 0.05-GeV/c intervals. Points for different  $A$  are occasionally displaced so as not to obscure the statistical error bars in this and Figs. 3, 4, and 5.

measured the analyzing power for hydrogen using a (CH 1.1) scintillator target. We find  $A_y(\text{H}) = +0.47 \pm 0.02$  in agreement with the free-hydrogen value<sup>14</sup>  $A_y = +0.48 \pm 0.02$ . At this angle we also measured  $A_y$  for quasifree proton scattering over a wide range of momenta.

Figures 2, 3, 4, and 5 show our results for  $\text{Li}^6$ ,  $\text{C}^{12}$ , and  $\text{Ta}^{181}$  obtained for both thin and thick targets ranging from 40 mg/cm<sup>2</sup> to 700 mg/cm<sup>2</sup>. ( $A_y$  was observed to be thickness independent.) A few results are also shown for  $\text{Be}^9$ . Statistical errors are in all cases larger than the expected systematic errors. Not shown in Fig. 2 are our results for the very highest momenta which are statistically poorer than the plotted results but nevertheless show very large positive values of  $A_y$ . They are  $\text{Li}^6$  (1.1 GeV/c):  $+0.66 \pm 0.24$ ;  $\text{Ta}^{181}$  (1.1 GeV/c):  $+0.29 \pm 0.06$ ;  $\text{C}^{12}$  (1.2 GeV/c):  $+0.83 \pm 0.18$ ;  $\text{Ta}^{181}$  (1.2 GeV/c):  $+0.56 \pm 0.15$ .

We have also made the first measurement of  $A_y$  for the forward production of protons in nuclei at  $19.27^\circ$  at and near the "quasielastic" peak. Figure 6 shows these results for Li, C, and Ta.

One feature of the forward data is that the analyzing powers are large and positive as observed in the measurement of the forward-going proton in the interaction with free stationary protons

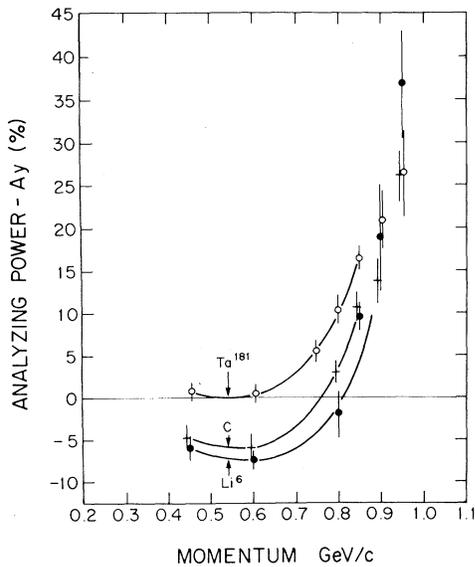


FIG. 3. Analyzing power (%) vs momentum (GeV/c); laboratory angle of 90°.

and, as one might expect, in quasifree scattering. However, it is important to observe that our values of  $A_y$  are appreciably smaller than for the free proton and are  $A$  dependent. One also sees in Fig. 6 that away from the quasifree peak where the internal momentum of the struck proton is not zero the analyzing powers fall rapidly. If  $A_y$  were calculated on the single-scattering hypothesis with the neglect of final-state interactions and at the appropriate  $s$  and  $t$  corresponding to a moving nucleon, one would not expect such a rapid falloff. Thus we have additional evidence from these analyzing-power measurements that  $A$ -dependent final-state interaction effects, known to be important in quasifree scattering, have a sensitive effect on the cross sections for polarized protons.<sup>19</sup>

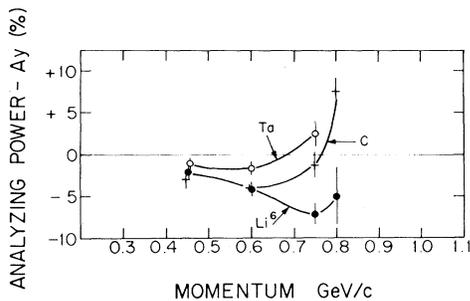


FIG. 4. Analyzing power (%) vs momentum (GeV/c); laboratory angle of 101°.

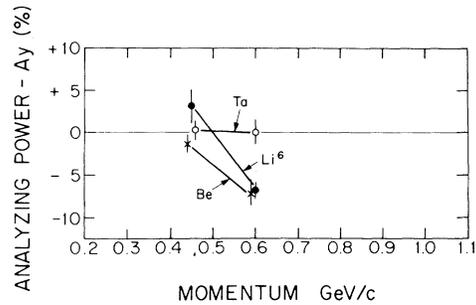


FIG. 5. Analyzing power (%) vs momentum (GeV/c); laboratory angle of 120°.

The backward data show these same features. Although Amado and Woloshyn<sup>3</sup> emphasized the crucial role of final-state interactions on the production of backward protons, the estimates of FW were made for the simple diagram of Fig. 1(a) neglecting final-state interactions. The low-momentum backward protons have an analyzing power of the predicted sign but not as large as predicted by FW. It would be important to extend these measurements to  $\text{He}^3, \text{He}^4$  to see if in the limit of small final-state interactions as  $A$  is decreased, the full FW predictions are approached. These data also show that as the energy of the backward proton becomes very low ( $\sim 45$  MeV), approaching the "evaporation" region,  $A_y$  seems to be approaching small values.

At the high momenta the analyzing powers change sign and become positive. Whether the sign reversal is an indication of more complex

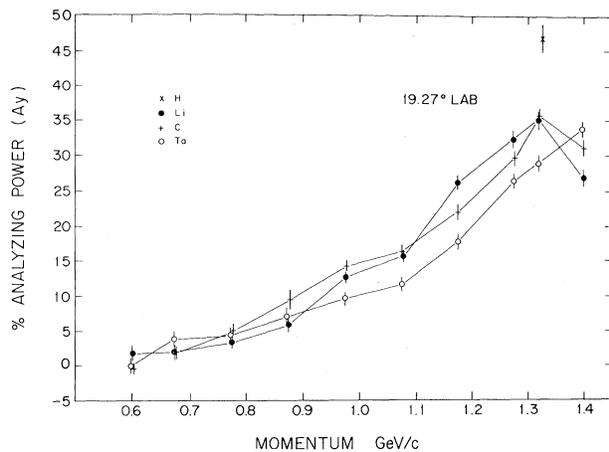


FIG. 6. Analyzing power (%) vs momentum (GeV/c); laboratory angle of 19.3°. The single point at 1.32 GeV/c shown as an X is our measured value for hydrogen which was obtained from the analyzing power we measured for a scintillator target (CH 1.1) and for carbon.

mechanisms than have been proposed or whether the incorporation of final-state interaction effects in the single-scattering model can account for the observed analyzing powers, it will be important to extend such measurements over a wider range of angles, momenta, and  $A$ , since the data appear to be rich in new information concerning high-momentum behavior of nuclei.

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<sup>8</sup>H. S. Weber and L. D. Miller, Phys. Rev. C **16**, 726 (1977).

<sup>9</sup>T. Fujita, Phys. Rev. Lett. **39**, 174 (1977).

<sup>10</sup>E. A. Remler, "Backward Proton Production in Energetic Proton-Nucleus Collisions" (to be published).

<sup>11</sup>S. Frankel *et al.*, Phys. Rev. Lett. **36**, 642 (1976).

See also the data discussed in Ref. 2 and by S. Frankel

*et al.*, "Measurement of Inclusive Cross Sections  $p(0.8 \text{ GeV}) + A \rightarrow p + X$ , up to the  $\text{Li}^6$  Elastic Limit," Phys. Rev. C (to be published), and by S. Frankel *et al.*, "Detailed Application of Quasi-Two Body Scaling to  $p + A \rightarrow p + X$ ," Phys. Rev. C (to be published).

<sup>12</sup>S. Frankel and R. M. Woloshyn, Phys. Rev. C **16**, 1680 (1977).

<sup>13</sup>The analyzing power,  $A_y$ , is defined by the polarization-dependent counting rate  $N \cong 1 + A_y \vec{\sigma} \cdot \vec{n}$ , where  $\vec{\sigma}$  is the projectile polarization and  $\vec{n}$  is the unit vector in direction of  $\vec{p}_{in} \times \vec{p}_{out}$ ;  $\vec{p}_{out}$  is defined as the momentum of the forward proton in  $p$ - $p$  scattering.

<sup>14</sup>J. Bystricky, F. Lehar, and Z. Janout, Centre d'Etudes Nucléaires Report No. CEA-N 1574E, 1972 (unpublished).

<sup>15</sup>Our preliminary measurements of  $A_y$  for  $\text{Be}^9$  have previously been reported: S. Frankel, W. Frati, G. Blanpied, G. W. Hoffmann, H. A. Thiessen, O. Van Dyck, and C. Whitten, Univ. of Pennsylvania Research Report No. UPR-0050N, 1977 (unpublished). In this preliminary run we observed  $A_y = -0.085 \pm 0.028$  for  $\text{Be}^9$  at 610 MeV/c. Our new measurement for  $\text{Be}^9$  yields  $-0.064 \pm 0.009$  confirming our previous result and our early observation of a nonzero analyzing power which ruled out the Weber-Miller mechanism of Fig. 1(b). J. Källne *et al.*, Phys. Lett. **74B**, 170 (1978), have since found values of  $A_y$  generally consistent with zero. Their most precise data near 30 MeV ( $E_p = 516$  MeV) are, however, in the evaporation region and do not test the models in Fig. 1. Near 100 MeV their data are less precise  $A_y \cong 0 \pm 0.2$  and thus none of their data support their conclusion favoring the Weber-Miller mechanism.

<sup>16</sup>The detection apparatus is described in part by Frankel *et al.*, Ref. 11.

<sup>17</sup>The high-resolution capability of the HRS was not used for this experiment.

<sup>18</sup>G. W. Hoffmann *et al.*, Phys. Rev. Lett. **40**, 1256 (1978).

<sup>19</sup>D. M. Corley and N. S. Wall, Nucl. Phys. **A184**, 437 (1972).