## ERRATA

OBSERVATION OF TRIMUON PRODUCTION BY NEUTRINOS. B. C. Barish, J. F. Bartlett, A. Bodek, K. W. Brown, D. Buchholz, Y. K. Chu, F. Sciulli, E. Siskind, L. Stutte, E. Fisk, G. Krafczyk, D. Nease, and O. Fackler [Phys. Rev. Lett. 38, 577 (1977)].

Table I contains an entry, as printed, of 2 trimuon events produced by antineutrinos. The table should read 2 trimuon events for  $\nu$ , but 0 trimuon events for  $\overline{\nu}$ .

POSSIBILITY OF OBSERVING NONEXPONEN-TIAL DECAYS IN AUTOIONIZING STATES. Cleanthes A. Nicolaides and Donald R. Beck [Phys. Rev. Lett. 38, 683 (1977)].

The exact behavior of G(t) depends on the behavior of g(E) everywhere. The branch point  $I_1$  causes mathematical difficulties as to the exact value of g(E) at  $E = I_1$  and its neighborhood for a multielectron system. In our work on *narrow* autoionizing states, the emphasis is put on the pole only and g(E) is set equal to zero at  $E = I_1$ . Thus, Eq. (12) should be understood as G(z) = 1/2

 $(z - z_0)$  for Rez > 0, Imz < 0 and G(z) = 0 for  $E \le 0$ . The expansion of A(z) around E = 0 with the introduction of the phase-space weighting factor

$$\lim_{E \neq I_1} \Gamma(E) \sim \sqrt{E}$$

as was done by Goldberger and Watson (Ref. 5 of our paper), was just another approximation for G(z) which yielded a faster,  $t^{-3/2}$ , NED dependence for the amplitude. Different approximation for A(z), e.g., for cases of overlapping resonances or the presence of an external field, would presumably yield yet more expressions for NED.

FORMATION OF CONVECTIVE CELLS, ANOM-ALOUS DIFFUSION, AND STRONG PLASMA TURBULENCE DUE TO DRIFT INSTABILITIES. C. Z. Cheng and H. Okuda [Phys. Rev. Lett. <u>38</u>, 708 (1977)].

On page 710, last paragraph, in the third and tenth lines, the references to recent laboratory experiments in toroidal devices should read "Okabayashi and Arunasalam,<sup>4</sup>" instead of "Hamberger *et al.*<sup>3</sup>"