and (9) yield

$$\eta_2 = \tfrac{1}{18}[1 + \int_0^\infty dx\, x\hat{\psi}'(x)\hat{G}(x)]. \quad (11)$$

It is not hard to evaluate the expression for $\hat{G}(x)$ and the integral in Eq. (11) analytically in the limit where a tends to infinity. Then all dependence on the cutoff function $\beta((x/a)^{1/2})$ is trivially eliminated. To leading order in $1/a$ one obtains

$$\hat{G}(x) = [x(4+x)]^{-1/2} \times \ln\left|\frac{x(3+x)+(1+x)[x(4+x)]^{1/2}}{x-[x(4+x)]^{1/2}}\right|, \quad (12)$$

$$\hat{\psi}'(x) = -2/[1+x]^3. \quad (13)$$

The value of the integral in Eq. (11) is $\frac{2}{3}$, which implies

$$\eta = \tfrac{1}{54}\epsilon^2 + O(\epsilon^3), \quad (14)$$

in agreement with the result obtained previously by other techniques.[4,5]

For general values of the parameter a Eqs. (6) and (11) can be evaluated numerically. We have done this for two choices of the cutoff function: $\beta(k) = 2k^2$ and $\beta(k) = k^4$. In each case the result of Eq. (14) is obtained. Since there is nothing special about these choices of $\beta(k)$ we conclude that, as expected, η is independent of the redundant parameters present in Wilson's incomplete-integration renormalization-group approach.

*Work supported in part by the National Science Foundation through Grant No. GH-36882.

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ERRATUM

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The following papers whose content partially anticipates and overlaps that of our recent Letter have been brought to our attention: L. I. Gudzenko and S. I. Yakovlenko, Zh. Eksp. Teor. Fiz. 62, 1686 (1972) [Sov. Phys. JETP 35, 877 (1972)]; S. I. Yakovlenko, Zh. Eksp. Teor. Fiz. 64, 2020 (1973) [Sov. Phys. JETP 37, 1019 (1973)].