ERRATA

ION ACCELERATION IN STRONG ELECTROMAGNETIC INTERACTIONS WITH PLASMAS. A. Y. Wong and R. L. Stenzel [Phys. Rev. Lett. 34, 727 (1975)].

The following references should be added:

To Ref. 2, "G. A. Swartz, T. T. Reboul, G. D. Gordon, and H. W. Larber, Phys. Fluids 3, 973 (1960); this last reference detected ion accelerations in an expanding plasma much smaller than the electromagnetic wavelength."

To Ref. 4, "H. C. S. Hsuan, Phys. Rev. 172, 137 (1968); P. Koch and J. Albritton, Phys. Rev. Lett. 32, 1420 (1974)."

ASYMPTOTIC SU(4) IN THE l^+l^- ANNIHILATION OF NEW RESONANCES. E. Takasugi and S. Oneda [Phys. Rev. Lett. 34, 988 (1975)].

The following text was omitted from this Letter:

Note added.—The most general form of V_μ^{em} may be written as $V_\mu^{\text{em}} = V_\mu^3 + \frac{1}{3}\sqrt{3}\, V_\mu^8 - x(V_\mu^{15} - yV_\mu^0)$. If we take the usual fractional charge assignment of SU(3) quarks, $(Q_u, Q_d, Q_s) = (\frac{2}{3}, -\frac{1}{3}, -\frac{1}{3})$, the charge of the charmed quark is, in general, $Q_c = n - \frac{1}{3}$ in which case the charmed mesons $D_1(c\bar{d})$, $D_2(c\bar{u})$, and $F(c\bar{s})$ have the charge n, $n-1$, n, respectively, and $x = (\frac{2}{3})^{1/2}(n - \frac{1}{3})$, $y = \frac{1}{3}\sqrt{3}$. The requirement to suppress the strangeness-changing neutral currant restricts n to either 1 or -1.

In the case of SU(3) quarks of integral charge,[2] again there are two possibilities. However, the difference between the fractional and integral charge appears only through the coefficient of the single current V_μ^0. This change can be absorbed in our arbitrary parameter p in Eq. (9). Thus we consider two possibilities, $n=1$ and -1. In the "ideal" limit we obtain the sum rule *independent* of x and y, $(m_\rho \Gamma_\rho)^{1/2} = (m_\omega \Gamma_\omega)^{1/2} + (2m_\varphi \Gamma_\varphi)^{1/2}$. The $\Gamma(\varphi_c)$ depends on the choice of the value of x. The case $n=1$, i.e., $x = (\frac{2}{3})^{1/2}$, was discussed in the text. For $n=1$, i.e., $x = -2(\frac{2}{3})^{1/2}$, we obtain $(m_{\varphi_c} \Gamma_{\varphi_c})^{1/2} = (2m_\omega \Gamma_\omega)^{1/2} + 3(m_\varphi \Gamma_\varphi)^{1/2}$. Thus if we take $m_\omega \Gamma_\omega : m_\varphi \Gamma_\varphi = 1:2$ as in the text, we obtain $m_\omega \Gamma_\omega : m_\varphi \Gamma_\varphi : m_\rho \Gamma_\rho : m_{\varphi_c} \Gamma_{\varphi_c} = 1:2:9:32$. With $\Gamma_\omega = 0.76$ keV and $\Gamma_\varphi = 1.34$ keV, we obtain $\Gamma_{\varphi_c} \simeq 6.7$ keV from the above sum rule.

The present experimental value of $\Gamma_{\varphi_c} \simeq 5$ keV might favor the choice[13] $n=-1$, in which case the D and F form the D^{--}, D^- isodoublet and F^- isosinglet.

SEARCH FOR CHARMED-PARTICLE PRODUCTION IN 15-BeV/c π^+p INTERACTIONS. C. Baltay, C. V. Cautis, D. Cohen, S. Csorna, M. Kalelkar, D. Pisello, E. Schmidt, W. D. Smith, and N. Yeh [Phys. Rev. Lett. 34, 1118 (1975)].

The receipt date of this Letter was omitted. The manuscript was received 11 February 1975.