

ly more complicated than for $l_n = 0$.

¹³This assumption is valid because q , which is greater than 10 fm for the cases considered here, is much larger than the range of V_{np} .

¹⁴J. D. Jackson, *Classical Electrodynamics* (Wiley, New York, 1962), p. 541.

¹⁵R. C. Johnson, Nucl. Phys. **A90**, 289 (1967).

¹⁶In the calculation the functions u_2 and u_0 were evaluated using the same approximations and parameters (see Ref. 6) which were used in the DWBA calculations. The values of q were calculated using Eq. (6) of Ref. 4. The kinetic energy was taken to be the average of the center-of-mass energies in the proton and deuteron channels.

Asymptotic Limit of $\sigma(e^+ + e^- \rightarrow \text{hadrons})/\sigma(e^+ + e^- \rightarrow \mu^+ + \mu^-)^\dagger$

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Combining the low-energy theorem by Crewther and by Chanowitz and Ellis on the anomalous vertex function for partially conserved dilation current and a new result from the most conventional vector-meson dominance, we predict the asymptotic limit of the ratio $\sigma(e^+ + e^- \rightarrow \text{hadrons})/\sigma(e^+ + e^- \rightarrow \mu^+ + \mu^-)$ to be $R \cong 16\pi^2/f_0^2 = 5.7 \pm 0.9$, which is in remarkable agreement with the preliminary data from the Cambridge Electron Accelerator e^+e^- colliding-beam experiment but in strong contradiction with the other predictions $R = \frac{2}{3}$ and $R = 2$ in the single- and three-triplet fractionally charged quark models.

One of the most intriguing experimental discoveries recently reported is the surprisingly large total cross section for hadron production which has been observed in the e^+e^- colliding-beam experiment at Cambridge Electron Accelerator (CEA).¹ The data have been analyzed in such a way that the ratio $R(q^2)$ of the observed total cross section for $e^+ + e^- \rightarrow \text{hadrons}$ to the theoretical one ($\cong 4\pi\alpha^2/3q^2$) for $e^+ + e^- \rightarrow \mu^+ + \mu^-$ is calculated to be 4.7 ± 1.1 at $q^2 = 16 \text{ GeV}^2$ and to be about 6.3 at $q^2 = 25 \text{ GeV}^2$, where q^2 is the square of the total c.m. energy of the colliding beams. A naive comparison of the CEA data with the Frascati

data² leads to the conclusion that $R(q^2)$ continues to rise as q^2 increases from above 1 up to 25 GeV^2 . On the other hand, many authors³ expect that $R(q^2)$ approaches a constant when q^2 goes up. It is now extremely interesting to ask what the asymptotic value for $R(q^2)$ is. The purpose of this short note is to report our new prediction for $R = \lim_{q^2 \rightarrow \infty} R(q^2)$, if it exists, based on the anomaly in partially conserved dilation current (PCDC) and the usual vector-meson dominance.

Let us start with the definition of the photon form factor for the trace of the stress-energy tensor,

$$\langle \gamma(\epsilon_1, k_1) | \theta_\lambda^\lambda(0) | \gamma(\epsilon_2, k_2) \rangle = -(\epsilon_1 \cdot \epsilon_2 k_1 \cdot k_2 - \epsilon_1 \cdot k_2 \epsilon_2 \cdot k_1) F((k_1 - k_2)^2). \quad (1)$$

Recently, Crewther⁴ and, independently, Chanowitz and Ellis⁵ have pointed out, on the canonical anomaly existing in the PCDC,⁶ that the coefficient of the anomalous term is completely determined by R . Furthermore, they have found that, because of the anomaly, $F(0)$ does not vanish and is given by

$$F(0) = -(e^2/12) \int dx dy x \cdot y \langle 0 | T(J_\mu(x) J^\mu(0) \theta_\lambda^\lambda(y)) | 0 \rangle = (e^2/6\pi^2) R. \quad (2)$$

If we assume vector-meson dominance, the left-hand side of Eq. (1) can be approximated by

$$(e^2/f_\rho^2) \langle \rho(\epsilon_1, k_1) | \theta_\lambda^\lambda(0) | \rho(\epsilon_2, k_2) \rangle + (\text{isoscalar terms}), \quad (3)$$

where f_ρ is the γ - ρ coupling constant ($f_\rho^2/4\pi = 2.2 \pm 0.3$). Now, the point is that we can evaluate the matrix element in (3) by the definition of the stress-energy tensor when $k_1 = k_2$, namely,

$$\langle \rho(\epsilon_1, k) | \theta_{00}(0) | \rho(\epsilon_2, k) \rangle = -2k_0 k_0 (\epsilon_1 \cdot \epsilon_2 - \epsilon_1 \cdot k \epsilon_2 \cdot k / k^2). \quad (4)$$

Thus we obtain the following approximate result:

$$F(0) \cong 2e^2/f_\rho^2 + (\text{isoscalar terms}). \quad (5)$$

If the electromagnetic current of hadrons $J_\mu(x)$ is a pure SU(3) octet operator,⁷ the isoscalar contribution is roughly $\frac{1}{3}$ the isovector (ρ) contribution. Therefore,

$$F(0) \cong \frac{4}{3}(2e^2/f_\rho^2). \quad (6)$$

Comparing these two results, (2) and (6), we finally predict the relation

$$R \cong 16\pi^2/f_\rho^2 = 5.7 \pm 0.9. \quad (7)$$

It is remarkable that this prediction is perfectly consistent with the CEA preliminary data for $q^2 = 25 \text{ GeV}^2$ as well as the rough estimate ($R \sim 6.5$)⁸ from the first evidence found at Frascati⁹ for hadron production by the two-photon process,¹⁰ but incompatible with the other predictions in the single- and three-triplet fractionally charged quark models. Notice that this prediction is independent of and different from those made by Bramon, Etim, and Greco¹¹ and by Sakurai,¹² based on an infinite series of vector-meson resonances or on the "new duality", since we have assumed the most conventional form of vector-meson dominance for *real photons*. It can not be stressed too strongly that *our application of the vector-meson dominance to the "softer" operator $T(J_\mu(x)J^\mu(0) \times \theta_\lambda^\lambda(y))$ for real photons is perfectly consistent with the PCDC anomaly and the nonvanishing constant R .*¹³ We have never applied vector-meson dominance directly to the "hard" operator $T(J_\mu(x)J_\nu(0))$. If one were to do this, one would reach the contradictory result $R = 0$ unless one adopts the "new duality".^{11,12} Notice also that the relation needs two corrections: One is due to the extrapolation of the matrix element from $k^2 = m_\rho^2$ to $k^2 = 0$, which has been estimated to be a few percent, and the other is due to vector states with negative charge conjugation other than ρ , ω , and ϕ , such as ρ' . The latter correction is supposed to be less than 10%. More detailed discussions on our result will be published elsewhere.⁸

Future e^+e^- colliding-beam experiments at Stanford Linear Accelerator Center and DESY will decide whether our prediction based on the PCDC anomaly and vector-meson dominance

agrees with the data.

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¹³It is worth mentioning here that the old application of vector-meson dominance to the decay of π^0 into 2γ is still consistent with the present data, independent of the PCAC anomaly. See M. Gell-Mann, D. Sharp, and W. G. Wagner, Phys. Rev. Lett. **8**, 261 (1962). By "softer" we mean that the dimension of θ_λ^λ is less than four. It may be more relevant here to state that, according to Weinberg's theorem [S. Weinberg, Phys. Rev. **118**, 838 (1960)], the Fourier transform of $T(J_\mu(x)J^\mu(0)\theta_\lambda^\lambda(y))$ increases more slowly (or decreases more rapidly) than that of $T(J_\mu(x)J_\nu(0))$ as k^2 , the square of the photon mass, goes up.